

AperTO - Archivio Istituzionale Open Access dell'Università di Torino

On the radius of spatial analyticity for semilinear symmetric hyperbolic systems

This is the author's manuscript

Original Citation:

Availability:

This version is available <http://hdl.handle.net/2318/142296> since 2017-05-19T15:08:56Z

Published version:

DOI:10.1016/j.jde.2014.01.020

Terms of use:

Open Access

Anyone can freely access the full text of works made available as "Open Access". Works made available under a Creative Commons license can be used according to the terms and conditions of said license. Use of all other works requires consent of the right holder (author or publisher) if not exempted from copyright protection by the applicable law.

(Article begins on next page)

ON THE RADIUS OF SPATIAL ANALYTICITY FOR SEMILINEAR SYMMETRIC HYPERBOLIC SYSTEMS

MARCO CAPPIELLO, PIERO D'ANCONA AND FABIO NICOLA

ABSTRACT. We study the problem of propagation of analytic regularity for semilinear symmetric hyperbolic systems. We adopt a global perspective and we prove that if the initial datum extends to a holomorphic function in a strip of radius (=width) ε_0 , the same happens for the solution $u(t, \cdot)$ for a certain radius $\varepsilon(t)$, as long as the solution exists. Our focus is on precise lower bounds on the spatial radius of analyticity $\varepsilon(t)$ as t grows.

We also get similar results for the Schrödinger equation with a real-analytic electromagnetic potential.

1. INTRODUCTION

The main concern in this paper is the propagation of the analytic regularity for systems of semilinear evolution equations. Local propagation of the analyticity for nonlinear partial differential equations has been studied in several papers, see for instance [1, 3, 4, 18, 19], starting from classical results of existence and uniqueness. In our paper we focus our attention on systems of semilinear equations with coefficients defined for $(t, x) \in \overline{\mathbb{R}}_+ \times \mathbb{R}^d$, hence globally defined in the space variables. Motivation for this type of study comes from the global analytic theory of nonlinear evolution PDE started with the paper by Kato and Masuda [28] and based on the proof of the existence of global analytic solutions in the space variables for the related Cauchy problem. After [28], many authors proved results of this type for several models as the (generalized) Korteweg-de Vries equation, the Euler equations, the Benjamin-Ono equation, the nonlinear Schrödinger equation, see for instance [5, 8, 20, 21, 22, 23, 24, 25, 26]. A parallel study in an elliptic setting devoted in particular to the analyticity of travelling waves has been developed in [6, 10, 11, 13, 14]. Besides the existence and uniqueness of an analytic solution, also the estimate of the radius of analyticity represents an interesting issue. Although there exist results in the literature treating particular models, cf. [7, 9, 29], we are not aware of similar results for general semilinear hyperbolic systems on \mathbb{R}^d . In the present paper we consider systems with initial data analytic on \mathbb{R}^d and admitting a holomorphic extension in a strip of the form $\{x + iy : |y| < \varepsilon_0\}$ for

2000 *Mathematics Subject Classification.* 35L45, 35A20, 35S10, 35Q35.

Key words and phrases. Hyperbolic systems, holomorphic extension, radius of analyticity.

some $\varepsilon_0 > 0$. We assume moreover the existence of a smooth (C^∞) solution of the Cauchy problem. Under these conditions, we prove that the solution is indeed analytic for every t in a strip of the form $\{x + iy \in \mathbb{C}^d : |y| < \varepsilon(t)\}$ and we give a precise estimate of lower bounds for $\varepsilon(t)$ as t grows. The class of systems to which our results apply includes semilinear symmetric hyperbolic systems but also strictly hyperbolic systems, see Remark 5.4, and other types of equations like the Schrödinger equation with electromagnetic potential (cf. [17]). For the sake of simplicity, in this Introduction we shall consider the special case of differential operators; we refer to Theorem 2.1 below for a more general statement involving pseudodifferential operators.

Consider a system of the form

$$Lu = \mathcal{N}[u]$$

in $\overline{\mathbb{R}}_+ \times \mathbb{R}^d \ni (t, x)$, $u = (u_1, \dots, u_n)$, with

$$(1.1) \quad L = \partial_t + iA_0(i\nabla_x) + \sum_{j=1}^d A_j(t, x)\partial_j + B(t, x),$$

where $A_0(i\nabla_x)$ is any formally self-adjoint Fourier multiplier (i.e. the symbol $A_0(\xi)$ is a smooth Hermitian $n \times n$ matrix), and $A_j(t, x)$ are $n \times n$ Hermitian matrices with entries in $C(\overline{\mathbb{R}}_+; C_b^\infty(\mathbb{R}^d))$, analytic with respect to x , satisfying

$$(1.2) \quad \sup_{(t,x) \in \overline{\mathbb{R}}_+ \times \mathbb{R}^d} \|\partial_x^\alpha A_j(t, x)\| \leq C^{|\alpha|+1} \alpha!, \quad \alpha \in \mathbb{N}^d$$

for some constant $C > 0$ (here $C_b^\infty(\mathbb{R}^d)$ denotes the space of all C^∞ functions on \mathbb{R}^d bounded with all their derivatives). We assume moreover that $B(t, x)$ is a $n \times n$ matrix with entries in $C(\overline{\mathbb{R}}_+; C_b^\infty(\mathbb{R}^d))$, satisfying

$$(1.3) \quad \sup_{(t,x) \in \overline{\mathbb{R}}_+ \times \mathbb{R}^d} \|\partial_x^\alpha B(t, x)\| \leq C^{|\alpha|+1} \alpha!, \quad \alpha \in \mathbb{N}^d.$$

The components of the non-linearity $\mathcal{N}[u]$ are assumed to be polynomials of degree $p \in \mathbb{N}$, $p \geq 2$, in the components of u with analytic coefficients, namely:

$$(1.4) \quad \mathcal{N}[u]_k = \sum_{\gamma \in \mathbb{N}^n: 2 \leq |\gamma| \leq p} g_{k,\gamma}(t, x) u^\gamma, \quad k = 1, \dots, n,$$

with $g_{k,\gamma}(t, x)$ continuous, satisfying

$$(1.5) \quad \|\partial_x^\alpha g_{k,\gamma}\|_{L^\infty(\overline{\mathbb{R}}_+ \times \mathbb{R}^d)} \leq M^{|\alpha|+1} \alpha!, \quad \alpha \in \mathbb{N}^d$$

for some constant $M > 0$.

When $A_0 = 0$ in (1.1) we get what is known as a *symmetric hyperbolic system* with bounded coefficients. However, it is important to allow the skew-adjoint term $iA_0(i\nabla_x)$ too, in view of the applications to Schrödinger equation below.

Incidentally, we observe that the form of this term is admittedly very special, but the reader will note from the proof that allowing more general skew-adjoint operators entails challenging issues.

We study the above system in the space of uniformly analytic functions, defined as follows.

Let \mathcal{A} be the space of functions f satisfying the estimates

$$\|\partial^\alpha f\|_{L^\infty} \leq C^{|\alpha|+1} \alpha!, \quad \alpha \in \mathbb{N}^d,$$

for some constant $C > 0$. Every function $f \in \mathcal{A}$ extends to a holomorphic function $f(x + iy)$ in the strip $\{x + iy \in \mathbb{C}^d : |y| < C^{-1}\}$ and we say that it has radius of analyticity $\geq C^{-1}$. For technical reasons it is convenient to work with an equivalent definition, where the L^∞ norm is replaced by the norm $\|\cdot\|_s$ in the Sobolev space $H^s(\mathbb{R}^d)$ for some $s > d/2$, fixed once and for all.

By Sobolev embedding theorem it is easy to prove that $\mathcal{A} = \cup_{\varepsilon > 0} \mathcal{A}(\varepsilon)$, where $\mathcal{A}(\varepsilon)$ is given by the smooth functions $f \in L^2(\mathbb{R}^d)$ such that $\sup_N E_N^\varepsilon[f] < \infty$, with

$$E_N^\varepsilon[f] := \sum_{|\alpha|=N} \frac{\varepsilon^{|\alpha|-1}}{\alpha!} \|\partial^\alpha f\|_s, \quad N \geq 1.$$

We have then the following result.

Theorem 1.1. *There exists a positive number $\bar{\varepsilon}_0 > 0$, depending only on the equation, such that for all $0 < \varepsilon_0 \leq \bar{\varepsilon}_0$ the following result holds true.*

Let $0 \in I \subset [0, +\infty)$ be a time interval. Let $u_0 \in \mathcal{A}(\varepsilon_0)$, and assume that there exists $s_0 \geq 0$ such that the Cauchy problem

$$\begin{cases} Lu = \mathcal{N}[u] & (t, x) \in I \times \mathbb{R}^d, \\ u(0, x) = u_0(x) \end{cases}$$

admits a solution $u \in C^1(I, H^{s'}(\mathbb{R}^d)) \cap C^0(I, H^{s'+1}(\mathbb{R}^d))$ for every $s' \geq s_0$. Then $u(t, \cdot) \in \mathcal{A}(\varepsilon(t))$, for every $t \in I$, with

$$(1.6) \quad \varepsilon(t) = \varepsilon_0 \exp\left(-A \int_0^t (1 + \|u(\sigma)\|_{L^\infty}^{p-1}) d\sigma\right),$$

for some constant $A > 0$ depending only on $\|u_0\|_s + \sup_{N \geq 1} E_N^{\varepsilon_0}[u_0]$. When the equation is linear ($p = 1$), the constant A depends only on the equation and it is independent of u_0 .

It will follow from the proof that we can take $A \simeq (\|u_0\|_s + \sup_{N \geq 1} E_N^{\varepsilon_0}[u_0])^{p-1}$.

The threshold $\bar{\varepsilon}_0$ corresponds, broadly speaking, to the radius of analyticity of the coefficients A_j , B and $g_{k,\gamma}$ in the equation (see (1.2), (1.3) and (1.5)). Hence $\mathcal{A}(\bar{\varepsilon}_0)$ is in some sense the larger space preserved by the equation.

Roughly speaking the phenomenon of shrinking of the radius of spatial analyticity is due to the combination of two effects: for linear equations the solution may

already compress more and more as t grows; this produces a growth of the spatial derivatives and causes the exponential decay in (1.6). In the presence of a non-linearity this growth may be amplified (if the initial datum is not too small), and this causes the presence of the integral term in (1.6). We refer to Section 5 for two model examples which illustrate these facts. Some variants are also discussed there, together with speculations on other lower bounds for $\varepsilon(t)$, as well as the issue of the dependence of the constants on the initial datum.

As anticipated, the above result applies also to the Schrödinger equation with real analytic electromagnetic potential. Namely, we have the following result.

Corollary 1.2. *Let $A(t, x) = (a_1(t, x), \dots, a_d(t, x))$ and $V(t, x)$ be the magnetic and the electric potential, respectively, with $a_j(t, x)$ and $V(t, x)$ in $C(\overline{\mathbb{R}}_+; C_b^\infty(\mathbb{R}^d))$, with $a_j(t, x)$ real-valued, and satisfying*

$$\|\partial_x^\alpha a_j\|_{L^\infty(\mathbb{R}_+ \times \mathbb{R}^d)} \leq C^{|\alpha|+1} \alpha!, \quad \|\partial_x^\alpha V\|_{L^\infty(\mathbb{R}_+ \times \mathbb{R}^d)} \leq C^{|\alpha|+1} \alpha!, \quad \alpha \in \mathbb{N}^d.$$

Consider the magnetic Laplacian

$$\Delta_A = \sum_{j=1}^d (\partial_j + ia_j)^2 = \Delta + 2iA \cdot \nabla + i \operatorname{div} A - \sum_{j=1}^d a_j^2$$

and the equation

$$(1.7) \quad Lu = \partial_t u - i\Delta_A u - iV(t, x)u = \mathcal{N}[u],$$

where the non-linearity $\mathcal{N}[u]$ has the same form as in (1.4) and (1.5) (without the subscript k , because we are now considering a scalar case).

Then the result of Theorem 1.1 holds true for the equation (1.7).

In fact, it is sufficient to apply Theorem 1.1 with $A_0(i\nabla_x) = -\Delta$, $A_j(t, x) = 2a_j(t, x)$ (which are assumed real-valued, hence Hermitian as 1×1 matrices) and $B = \operatorname{div} A + i \sum_{j=1}^d a_j^2 + iV$ in (1.1).

The core of the proof of Theorem 1.1 and of the more general Theorem 2.1 below consists in proving a suitable analytic energy estimate for the solution of the Cauchy problem. In future works we plan to prove similar results for more general nonlinear terms, possibly treating the case of quasi-linear or fully nonlinear systems, and also to improve the above lower bound for the radius of analyticity for special classes of equations (e.g. with constant coefficients). We also recall that in some recent papers [14, 15, 16] we proved holomorphic extensions of the solutions of elliptic equations in conical subsets of \mathbb{C}^d of the form $\{x + iy \in \mathbb{C}^d : |y| < \varepsilon(1 + |x|)\}$, $\varepsilon > 0$, improving some of the results of [10, 11, 13] on the analyticity in a strip for travelling waves and stationary Schrödinger equations. An interesting issue would be to investigate the same problem for the corresponding evolution equations and to estimate the decay of the width of the sector as the time variable grows.

2. STATEMENT AND PROOF OF THE MAIN RESULT

Let us now consider a more general system of the form

$$Lu = \mathcal{N}[u]$$

in $\overline{\mathbb{R}}_+ \times \mathbb{R}^d \ni (t, x)$, $u = (u_1, \dots, u_n)$, with

$$(2.1) \quad L = \partial_t + iA_0(D_x) + A(t, x, D_x),$$

where $A_0(D_x)$ is any formally self-adjoint Fourier multiplier (i.e. the symbol $A_0(\xi)$ is a smooth Hermitian $n \times n$ matrix), and $A(t, x, D_x)$ is a system of pseudodifferential operators, i.e.

$$A(t, x, D_x)u = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ix\xi} A(t, x, \xi) \hat{u}(\xi) d\xi$$

whose symbol $A(t, x, \xi)$ is a $n \times n$ matrix satisfying the following conditions:

$$\langle \xi \rangle^{|\beta|-1} \partial_x^\alpha \partial_\xi^\beta A(t, x, \xi) \text{ has entries in } C(\overline{\mathbb{R}}_+; C_b^0(\mathbb{R}^d \times \mathbb{R}^d)) \quad \forall \alpha, \beta \in \mathbb{N}^d$$

with bounds

$$(2.2) \quad \|\partial_x^\alpha \partial_\xi^\beta A(t, x, \xi)\| \leq C_\beta C^{|\alpha|+1} \alpha! \langle \xi \rangle^{1-|\beta|}, \quad \alpha, \beta \in \mathbb{N}^d$$

and moreover

$$(2.3) \quad \sup_{t \in \overline{\mathbb{R}}_+; x, \xi \in \mathbb{R}^d} \|A(t, x, \xi) + A(t, x, \xi)^*\| < \infty.$$

It is immediate to see that the simpler case considered in the Introduction, namely the operator L in (1.1), (1.2), (1.3) is in fact a special case.

The non-linearity $\mathcal{N}[u]$ is the same considered in the Introduction, therefore is given by (1.4), with coefficients $g_{k,\gamma}(t, x)$ satisfying (1.5).

Theorem 2.1. *Let L be given by (2.1), (2.2), (2.3) and $\mathcal{N}[u]$ as in (1.4), (1.5). Then the same result as in Theorem 1.1 holds true.*

Theorem 2.1 will follow if we prove that $E_N^{\varepsilon(t)}[u(t)]$, $N \geq 1$, is bounded, with $\varepsilon(t)$ as in (1.6). Namely, we will prove by induction that for some constant $C_0 > 0$ we have

$$(2.4) \quad E_N^{\varepsilon(t)}[u(t)] \leq \Phi(t) := C_0 \exp\left(C_0 \int_0^t (1 + \|u(\sigma)\|_{L^\infty}^{p-1}) d\sigma\right),$$

for every $N \geq 1$, if $\varepsilon_0 \leq \bar{\varepsilon}_0$, with $\bar{\varepsilon}_0$ small enough, and the constant A in (1.6) is large enough.

Under the hypothesis of Theorem 2.1 we have the following classical energy estimates

$$(2.5) \quad \|v(t)\|_s \leq C e^{Ct} \|v(0)\|_s + C e^{Ct} \int_0^t e^{-C\sigma} \|Lv(\sigma)\|_s d\sigma$$

for every function $v \in C^1(I, H^s) \cap C^0(I, H^{s+1})$ and $t \in I$, where $I \subset [0, +\infty)$ is a time interval containing 0, and the constant $C > 0$ depends on L and s , see [30, Proposition 6.1]. Usually this estimate is proved without the term $iA_0(i\nabla_x)$ in (1.1), but since it does not depend on x and is skew-adjoint, it does not affect the classical proof (in particular the fact that it commutes with spatial derivatives is useful to extend the L^2 -energy estimates to the H^s version); e.g. the proof of [32, Proposition 2.1.2] can be repeated almost verbatim for our operator L in (2.1).

Now, we apply the above estimate to $v = \partial^\alpha u$, where u solves $Lu = \mathcal{N}[u]$, $u(0) = u_0$. We get

$$\|\partial^\alpha u(t)\|_s \leq Ce^{Ct} \left(\|\partial^\alpha u_0\|_s + \int_0^t e^{-C\sigma} (\|[L, \partial^\alpha]u(\sigma)\|_s + \|\partial^\alpha \mathcal{N}[u(\sigma)]\|_s) d\sigma \right).$$

We multiply by $\varepsilon(t)^{|\alpha|-1}/\alpha!$, $|\alpha| = N \geq 1$, and we obtain, since $\varepsilon(t) \leq \varepsilon(0) = \varepsilon_0$,

(2.6)

$$\begin{aligned} e^{-Ct} E_N^{\varepsilon(t)}[u(t)] &\leq CE_N^{\varepsilon_0}[u_0] + C \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} \sum_{|\alpha|=N} \frac{\varepsilon(\sigma)^{|\alpha|-1}}{\alpha!} \|[L, \partial^\alpha]u(\sigma)\|_s d\sigma \\ &\quad + C \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} E_N^{\varepsilon(\sigma)}[\mathcal{N}[u(\sigma)]] d\sigma. \end{aligned}$$

Now, we can choose the constant C_0 in (2.4) such that

$$(2.7) \quad C_0 > C \max\{4 \sup_N E_N^{\varepsilon_0}[u_0], 2\}$$

where C is the constant arising in (2.5). Observe that by the condition (2.7) we have $C < C_0$ and $C \sup_N E_N^{\varepsilon_0}[u_0] < C_0/4$ so that the first term in the right-hand side of (2.6) is then dominated by $\frac{1}{4}e^{-Ct}\Phi(t)$. We will then prove that (2.4) holds for every $N \geq 1$ possibly for a bigger constant C_0 (independent of N). More precisely, we will prove directly the case $N = 1$ whereas for $N > 1$ we shall argue by induction supposing that (2.4) holds for subscripts $< N$. Under this hypothesis we will verify in the next sections the following estimates on the two integrals involved in the right-hand side of (2.6).

Proposition 2.2. *There exist constants $C' > 0$ and $\bar{\varepsilon}_0 > 0$, depending only on the equation, such that, for every $\varepsilon_0 \leq \bar{\varepsilon}_0$ and every $A \geq 1$ in (1.6), and for every sufficiently large constant C_0 in (2.4) we have*

$$(2.8) \quad \begin{aligned} \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} \sum_{|\alpha|=N} \frac{\varepsilon(\sigma)^{|\alpha|-1}}{\alpha!} \|[L, \partial^\alpha]u(\sigma)\|_s d\sigma \\ \leq \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} C' N e^{-C\sigma} E_N^{\varepsilon(\sigma)}[u(\sigma)] d\sigma + \frac{1}{4} e^{-Ct} \Phi(t) \end{aligned}$$

where $\Phi(t)$ is defined in (2.4).

Proposition 2.3. *There exist constants $C' > 0$ and $\bar{\varepsilon}_0 > 0$, depending only on the equation, such that, for every $\varepsilon_0 \leq \bar{\varepsilon}_0$ and for every sufficiently large constants A in (1.6) and C_0 in (2.4), we have*

$$(2.9) \quad \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} E_N^{\varepsilon(\sigma)} [\mathcal{N}[u(\sigma)]] d\sigma \\ \leq \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} C' \|u(\sigma)\|_{L^\infty}^{p-1} e^{-C\sigma} E_N^{\varepsilon(\sigma)} [u(\sigma)] d\sigma + \frac{1}{4} e^{-Ct} \Phi(t)$$

where $\Phi(t)$ is defined in (2.4).

Assume that Propositions 2.2 and 2.3 hold true and set $\Psi_N(t) = e^{-Ct} E_N^{\varepsilon(t)} [u(t)]$. We can continue the computation in (2.6) as

$$(2.10) \quad \Psi_N(t) \leq \frac{3}{4} e^{-Ct} \Phi(t) + \varepsilon(t)^{N-1} \int_0^t \frac{\Psi_N(\sigma)}{\varepsilon(\sigma)^{N-1}} C' (N + \|u(\sigma)\|_{L^\infty}^{p-1}) d\sigma.$$

We now use the following form of Gronwall inequality, which can be deduced easily from the classical one (see e.g. [32, Lemma 2.1.3]).

Lemma 2.4. *Let $0 \leq g, \psi, a \in L_{loc}^\infty([0, T])$, with $a > 0$, and $0 \leq h \in L_{loc}^1([0, T])$, and*

$$\psi(t) \leq g(t) + a(t) \int_0^t h(\sigma) \frac{\psi(\sigma)}{a(\sigma)} d\sigma.$$

Then, with $H(t) := \int_0^t h(\sigma) d\sigma$,

$$\psi(t) \leq g(t) + a(t) e^{H(t)} \int_0^t e^{-H(\sigma)} h(\sigma) \frac{g(\sigma)}{a(\sigma)} d\sigma.$$

We apply Lemma 2.4 to (2.10) with $\psi(t) = \Psi_N(t)$, $g(t) = \frac{3}{4} e^{-Ct} \Phi(t)$, $h(\sigma) = C' (N + \|u(\sigma)\|_{L^\infty}^{p-1})$, $a(t) = \varepsilon(t)^{N-1}$.

From (2.10) we therefore get

$$\Psi_N(t) \leq \frac{3}{4} e^{-Ct} \Phi(t) + \tilde{\Phi}_N(t)$$

with

$$\tilde{\Phi}_N(t) = \frac{3}{4} \varepsilon(t)^{N-1} e^{C'Nt + C' \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1}} \int_0^t e^{-C'N\sigma - C' \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} C' (N + \|u(\sigma)\|_{L^\infty}^{p-1}) \\ \times \frac{e^{-C\sigma} \Phi(\sigma)}{\varepsilon(\sigma)^{N-1}} d\sigma.$$

We want to prove that

$$(2.11) \quad \tilde{\Phi}_N(t) \leq \frac{1}{4} e^{-Ct} \Phi(t),$$

if the constants A in (1.6) and C_0 in (2.4) are large enough, so that we get $\Psi_N(t) \leq e^{-Ct} \Phi(t)$ and the proof of (2.4) is concluded.

To this end we use the explicit expression of $\varepsilon(t)$ in (1.6); we have

$$\begin{aligned} \tilde{\Phi}_N(t) &= \frac{3}{4} C_0 C' N e^{(-A(N-1)+C'N)t+(C'-A(N-1)) \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma} \\ &\quad \times \int_0^t e^{(-C'N+A(N-1)-C+C_0)\sigma+(-C'+C_0+A(N-1)) \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} d\sigma \\ &\quad + \frac{3}{4} C_0 C' e^{(-A(N-1)+C'N)t+(C'-A(N-1)) \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma} \\ &\quad \times \int_0^t e^{(-C'N+A(N-1)-C+C_0)\sigma+(-C'+C_0+A(N-1)) \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma. \end{aligned}$$

In the first integral we estimate

$$e^{(-C'+C_0+A(N-1)) \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} \leq e^{(-C'+C_0+A(N-1)) \int_0^t \|u(\tau)\|_{L^\infty}^{p-1} d\tau},$$

which holds if $C_0 > C'$, whereas in the second integral we use

$$e^{(-C'N+A(N-1)-C+C_0)\sigma} \leq e^{(-C'N+A(N-1)-C+C_0)t},$$

which holds if $C_0 > C' + C$, $A > 2C'$. We get

$$\tilde{\Phi}_N(t) \leq \frac{3}{4} \left[\frac{C'N}{-C'N + A(N-1) - C + C_0} + \frac{C'}{-C' + C_0 + A(N-1)} \right] e^{-Ct} \Phi(t)$$

We conclude by observing that for $N = 1$ the expression in parenthesis is $< 1/3$ if C_0 is large enough, while the same holds for every $N \geq 2$, if A is large enough. This shows that (2.11) holds true and concludes the proof. \square

3. PROOF OF PROPOSITION 2.2

We have, for $k = 1, \dots, n$,

$$([L, \partial^\alpha]u)_k = - \sum_{h=1}^n \sum_{0 \neq \gamma \leq \alpha} \binom{\alpha}{\gamma} (\partial^\gamma A)(t, x, D_x)_{k,h} \partial^{\alpha-\gamma} u_h,$$

In view of (2.2) the symbols $\partial_x^\gamma A(t, x, \xi)_{k,h}$ belong to the Hörmander's classes $S_{1,0}^1$ and therefore the corresponding operators are bounded $H^{s+1} \rightarrow H^s$ with operator norm dominated by a seminorm of their symbol in that symbol class, cf. ([27,

Chapter XVIII]). Hence, using $\|f\|_s + \sum_{j=1}^d \|\partial_j f\|_s$ as an equivalent norm in H^{s+1} we get

$$\begin{aligned} \frac{\varepsilon^{|\alpha|-1}}{\alpha!} \|[L, \partial^\alpha]u\|_s &\leq C_1 \sum_{j=1}^d \sum_{0 \neq \gamma \leq \alpha} (C_1 \varepsilon)^{|\gamma|-1} |\alpha - \gamma + e_j| \varepsilon^{|\alpha-\gamma|} \frac{\|\partial^{\alpha-\gamma+e_j} u\|_s}{(\alpha - \gamma + e_j)!} \\ &\quad + C_1 \sum_{\substack{0 \neq \gamma \leq \alpha \\ |\gamma| \leq |\alpha|-1}} (C_1 \varepsilon)^{|\gamma|} \varepsilon^{|\alpha-\gamma|-1} \frac{\|\partial^{\alpha-\gamma} u\|_s}{(\alpha - \gamma)!} + (C_1 \varepsilon)^{|\alpha|} \|u\|_s \end{aligned}$$

where e_j denotes the j -th element of the canonical basis of \mathbb{R}^d . Hence we get

$$\sum_{|\alpha|=N} \frac{\varepsilon^{|\alpha|-1}}{\alpha!} \|[L, \partial^\alpha]u\|_s \leq C_2 \sum_{1 \leq |\tilde{\alpha}| \leq N} \sum_{|\gamma|=N-|\tilde{\alpha}|+1} (C_2 \varepsilon)^{N-|\tilde{\alpha}|} |\tilde{\alpha}| \frac{\|\partial^{\tilde{\alpha}} u\|_s}{\tilde{\alpha}!} + (C_2 \varepsilon)^N \|u\|_s.$$

Since the number of multi-indices γ of length $N - |\tilde{\alpha}| + 1$ is dominated by $2^{N-|\tilde{\alpha}|+d}$ we obtain

$$\sum_{|\alpha|=N} \frac{\varepsilon^{|\alpha|-1}}{\alpha!} \|[L, \partial^\alpha]u\|_s d\sigma \leq C' \sum_{j=1}^N (C' \varepsilon)^{N-j} j E_j^\varepsilon[u] + (C_2 \varepsilon)^N \|u\|_s.$$

The above constants C_1, C_2, C' depend on the operator L but are independent of ε, N . Hence we got

$$\begin{aligned} \sum_{|\alpha|=N} \frac{\varepsilon(\sigma)^{|\alpha|-1}}{\alpha!} \|[L, \partial^\alpha]u(\sigma)\|_s &\leq C' N E_N^{\varepsilon(\sigma)}[u(\sigma)] \\ &\quad + C' \sum_{j=1}^{N-1} (C' \varepsilon(\sigma))^{N-j} j E_j^{\varepsilon(\sigma)}[u(\sigma)] + (C_2 \varepsilon(\sigma))^N \|u(\sigma)\|_s. \end{aligned}$$

Substituting in the left hand side of (2.8) we see that it is sufficient to prove that

$$(3.1) \quad \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} (C_2 \varepsilon(\sigma))^N \|u(\sigma)\|_s d\sigma \leq \frac{1}{8} e^{-Ct} \Phi(t) \quad \text{for } N \geq 1,$$

and

$$(3.2) \quad \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} C' \sum_{j=1}^{N-1} (C' \varepsilon(\sigma))^{N-j} j E_j^{\varepsilon(\sigma)}[u(\sigma)] d\sigma \leq \frac{1}{8} e^{-Ct} \Phi(t) \quad \text{for } N \geq 2,$$

if $\varepsilon(0) = \varepsilon_0 \leq \bar{\varepsilon}_0$ with $\bar{\varepsilon}_0$ small enough, and the constant C_0 in (2.4) is large enough.

Let us prove estimate (3.1). We apply the energy estimate (2.5) with $v = u$. Since $Lu = \mathcal{N}[u]$ has the form in (1.4), (1.5), with the aid of the Schauder's

estimates $\|fg\|_s \lesssim \|f\|_s \|g\|_{L^\infty}$ one obtains

$$e^{-C\sigma} \|u(\sigma)\|_s \leq C \|u_0\|_s + C_3 \int_0^\sigma (1 + \|u(\tau)\|_{L^\infty}^{p-1}) e^{-C\tau} \|u(\tau)\|_s d\tau$$

for a suitable constant $C_3 > 0$. By an application of Lemma 2.4 with $a(t) \equiv 1$, $\psi(t) = e^{-Ct} \|u(t)\|_s$, $g(t) = C \|u_0\|_s$ and $h(t) = C_3(1 + \|u\|_{L^\infty}^{p-1})$, we get

$$(3.3) \quad e^{-C\sigma} \|u(\sigma)\|_s \leq C \|u_0\|_s e^{C_3 \int_0^\sigma (1 + \|u(\tau)\|_{L^\infty}^{p-1}) d\tau},$$

which gives, using $\varepsilon(t)^{N-1}/\varepsilon(\sigma)^{N-1} \leq 1$ and $\varepsilon(\sigma) \leq \varepsilon_0 e^{-A\sigma}$,

$$\begin{aligned} & \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} (C_2 \varepsilon(\sigma))^N \|u(\sigma)\|_s d\sigma \\ & \leq C \|u_0\|_s e^{C_3 \int_0^t (1 + \|u(\sigma)\|_{L^\infty}^{p-1}) d\sigma} \int_0^t (C_2 \varepsilon_0)^N e^{-AN\sigma} d\sigma \\ & \leq \frac{C(C_2 \varepsilon_0)^N}{AN} \|u_0\|_s e^{C_3 \int_0^t (1 + \|u(\sigma)\|_{L^\infty}^{p-1}) d\sigma} \leq \frac{1}{8} e^{-Ct} \Phi(t) \end{aligned}$$

if the constant C_0 in (2.4) satisfies $C_0 \geq C + C_3$ and $C_0 \geq C \|u_0\|_s$ (hence these constraints on C_0 have to be added to those in (2.7)), $A \geq 1$ and $\varepsilon_0 \leq \bar{\varepsilon}_0 := 1/(8C_2)$. This proves (3.1).

Let us now prove (3.2) for $N \geq 2$. We can use then the inductive hypothesis (2.4) for subscripts $j < N$ and the estimates $\varepsilon(t)^{N-1}/\varepsilon(\sigma)^{N-1} \leq e^{-A(t-\sigma)(N-1)}$, as well as $\varepsilon(t) \leq \varepsilon_0 e^{-At}$. Then the left-hand side of (3.2) is

$$\begin{aligned} & \leq C' C_0 e^{C_0 \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma} \sum_{j=1}^{N-1} (C' \varepsilon_0)^{N-j} \int_0^t j e^{(A(j-1) + C_0 - C)\sigma - At(N-1)} d\sigma \\ & \leq e^{-Ct} \Phi(t) \sum_{j=1}^{N-1} (C' \varepsilon_0)^{N-j} \frac{C' j}{A(j-1) + C_0 - C}. \end{aligned}$$

The above sum is $< 1/8$ for all $\varepsilon_0 \leq \bar{\varepsilon}_0$ if $\bar{\varepsilon}_0$ is small enough, depending on C and C' (uniformly with respect to $C_0 > 2C$, cf. (2.7), and $A \geq 1$). \square

4. PROOF OF PROPOSITION 2.3

To simplify the notation and without loss of generality we shall prove Proposition 2.3 in the scalar case $n = 1$, and for a nonlinear term $\mathcal{N}[u] = g(t, x)u^p$, where $p \in \mathbb{N}$, $p \geq 2$ and $g(t, x)$ is a smooth function defined on $\overline{\mathbb{R}}_+ \times \mathbb{R}$ and satisfying the

estimate (1.5). Let us first consider the case $N = 1$. We have

$$\begin{aligned} E_1^c[\mathcal{N}[u]] &= E_1[\mathcal{N}[u]] = \sum_{j=1}^d \|\partial_{x_j}(g(t, x)u^p)\|_s \\ &\leq C_1 \sum_{j=1}^d \|\partial_{x_j}g\|_{L^\infty(\overline{\mathbb{R}_+ \times \mathbb{R}})} \cdot \|u\|_s^p + C_1 \|g\|_{L^\infty(\overline{\mathbb{R}_+ \times \mathbb{R}})} \cdot \|u\|_{L^\infty}^{p-1} \cdot E_1[u] \\ &\leq C'(\|u\|_s^p + \|u\|_{L^\infty}^{p-1} \cdot E_1[u]). \end{aligned}$$

To prove (2.9) for $N = 1$ we are then reduced to prove that

$$C' \int_0^t e^{-C\sigma} \|u(\sigma)\|_s^p d\sigma \leq \frac{1}{4} e^{-Ct} \Phi(t).$$

Using the estimate (3.3) and assuming $C_0 > p(C_3 + C)$ we obtain that

$$\begin{aligned} \int_0^t e^{-C\sigma} \|u(\sigma)\|_s^p d\sigma &\leq (C\|u_0\|_s)^p \int_0^t e^{(C(p-1)+C_3p)\sigma} e^{C_3p \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} \\ &\leq (C\|u_0\|_s)^p e^{(C_0-C)t} \int_0^t e^{-(C_0-Cp-C_3p)\sigma} e^{C_3p \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} d\sigma \\ &\leq \frac{(C\|u_0\|_s)^p}{C_0} e^{-Ct} \Phi(t) \int_0^t e^{-(C_0-Cp-C_3p)\sigma} d\sigma \\ &= \frac{(C\|u_0\|_s)^p}{C_0(C_0 - Cp - C_3p)} e^{-Ct} \Phi(t) (1 - e^{-(C_0-Cp-C_3p)t}) \\ &\leq \frac{(C\|u_0\|_s)^p}{C_0(C_0 - Cp - C_3p)} e^{-Ct} \Phi(t). \end{aligned}$$

Then one can choose C_0 large enough to have $\frac{(C\|u_0\|_s)^p}{C_0(C_0 - Cp - C_3p)} \leq \frac{1}{4}$. The proposition is then proved in the case $N = 1$.

For $N \geq 2$, using Leibniz formula we have

$$\|\partial_x^\alpha \mathcal{N}[u]\|_s \leq \|g(t, \cdot)u^{p-1} \partial_x^\alpha u\|_s + \sum_{\substack{\gamma + \alpha_1 + \dots + \alpha_p = \alpha \\ |\alpha_j| < |\alpha| \forall j}} \frac{\alpha!}{\gamma! \alpha_1! \dots \alpha_p!} \|\partial_x^\gamma g(t, \cdot) \partial_x^{\alpha_1} u \dots \partial_x^{\alpha_p} u\|_s.$$

Then by Schauder's lemma and using the condition (1.5) we obtain

$$\begin{aligned}
E_N^\varepsilon[\mathcal{N}[u]] &= \sum_{|\alpha|=N} \frac{\varepsilon^{|\alpha|-1}}{\alpha!} \|\partial_x^\alpha \mathcal{N}[u]\|_s \leq C \|g\|_{L^\infty(\bar{\mathbb{R}}_+ \times \mathbb{R})} \cdot \|u\|_{L^\infty}^{p-1} \cdot E_N^\varepsilon[u] \\
&+ C \sum_{|\alpha|=N} \frac{\varepsilon^{|\alpha|-1}}{\alpha!} \sum_{\substack{\gamma+\alpha_1+\dots+\alpha_p=\alpha \\ |\alpha_j|<|\alpha|\forall j}} \frac{\alpha!}{\gamma! \alpha_1! \dots \alpha_p!} \|\partial_x^\gamma g(t, \cdot)\|_{L^\infty(\bar{\mathbb{R}}_+ \times \mathbb{R})} \cdot \|\partial_x^{\alpha_1} u\|_s \dots \|\partial_x^{\alpha_p} u\|_s \\
&\leq CM \|u\|_{L^\infty}^{p-1} \cdot E_N^\varepsilon[u] + C' \varepsilon^{p-1} \sum_{|\alpha|=N} \sum_{\substack{\alpha_1+\dots+\alpha_p=\alpha \\ |\alpha_j|<|\alpha|\forall j}} (M\varepsilon)^{|\gamma|} \prod_{j=1}^p \frac{\varepsilon^{|\alpha_j|-1} \|\partial_x^{\alpha_j} u\|_s}{\alpha_j!} \\
&\leq C' \|u\|_{L^\infty}^{p-1} \cdot E_N^\varepsilon[u] + C' \varepsilon \left(\sup_{1 \leq j \leq N-1} E_j^\varepsilon[u] \right)^p,
\end{aligned}$$

where we used the fact that $p-1 \geq 1$ and we assume $\varepsilon \leq \min\{1, M^{-1}\}$, M being the constant appearing in (1.5). Then we obtain

$$\begin{aligned}
&\int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} e^{-C\sigma} E_N^{\varepsilon(\sigma)}[\mathcal{N}[u(\sigma)]] d\sigma \\
&\leq C' \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} \|u(\sigma)\|_{L^\infty}^{p-1} e^{-C\sigma} E_N^{\varepsilon(\sigma)}[u(\sigma)] d\sigma \\
&\quad + C' \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} \varepsilon(\sigma) e^{-C\sigma} \left(\sup_{1 \leq j \leq N-1} E_j^{\varepsilon(\sigma)}[u(\sigma)] \right)^p d\sigma
\end{aligned}$$

if $\varepsilon \leq \bar{\varepsilon}_0 := \min\{1, M^{-1}\}$.

To conclude the proof, it is then sufficient to prove that

$$(4.1) \quad \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} \varepsilon(\sigma) e^{-C\sigma} \left(\sup_{1 \leq j \leq N-1} E_j^{\varepsilon(\sigma)}[u(\sigma)] \right)^p d\sigma \leq \frac{1}{4} e^{-Ct} \Phi(t).$$

First of all we observe that since the function $\varepsilon(t)$ is decreasing, we have $\frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} \leq 1$ for $0 \leq \sigma \leq t$. Moreover, writing

$$e^{-C\sigma} \left(\sup_{1 \leq j \leq N-1} E_j^{\varepsilon(\sigma)}[u(\sigma)] \right)^p = e^{(p-1)C\sigma} \left(\sup_{1 \leq j \leq N-1} e^{-C\sigma} E_j^{\varepsilon(\sigma)}[u(\sigma)] \right)^p$$

and using the inductive assumption (2.4), we obtain that

$$\left(\sup_{1 \leq j \leq N-1} E_j^{\varepsilon(\sigma)}[u(\sigma)] \right)^p \leq C_0^p e^{Cp \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma}.$$

Hence, using the explicit expression of $\varepsilon(t)$ in (1.6) and choosing $A > C_0(p-1) \geq C(p-1)$, we get

$$\begin{aligned}
& \int_0^t \frac{\varepsilon(t)^{N-1}}{\varepsilon(\sigma)^{N-1}} \varepsilon(\sigma) e^{-C\sigma} \left(\sup_{1 \leq j \leq N-1} E_j^{\varepsilon(\sigma)}[u(\sigma)] \right)^p d\sigma \\
& \leq C_0^p \varepsilon_0 \int_0^t e^{((p-1)C-A)\sigma - (A-C_0p) \int_0^\sigma \|u(\tau)\|_{L^\infty}^{p-1} d\tau} d\sigma \\
& \leq C_0^p \varepsilon_0 e^{C_0 \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma} \cdot \int_0^t e^{-(A-(p-1)C)\sigma} d\sigma \\
& = \frac{C_0^p \varepsilon_0}{A - (p-1)C} (1 - e^{-(A-(p-1)C)t}) e^{C_0 \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma} \\
& \leq \frac{C_0^p \varepsilon_0}{A - (p-1)C} e^{-Ct} \Phi(t) \leq \frac{1}{4} e^{-Ct} \Phi(t)
\end{aligned}$$

choosing A sufficiently large. This concludes the proof. \square

5. EXAMPLES AND REMARKS

The following two examples make clear the expression for $\varepsilon(t)$ given in (1.6).

Example 5.1. Consider first the linear Cauchy problem (hence $p = 1$) in $\overline{\mathbb{R}}_+ \times \mathbb{R}$ given by

$$\begin{cases} \partial_t u - x \partial_x u = 0 \\ u(0, x) = u_0(x) = (1 + x^2)^{-1}. \end{cases}$$

The initial datum extends to a holomorphic function $u_0(x + iy)$ in the strip $|y| < 1$, whereas the solution, given by $u(t, x) = (1 + x^2 e^{2t})^{-1}$, extends, for every fixed t , to a holomorphic function in the strip $|y| < e^{-t}$. Hence the linear part of the equation is already responsible of the exponential decay in (1.6). For equations with constant coefficients (even non-linear) one instead expects algebraic decay for the radius of analyticity (cf. e.g. [9]). We plan to present a systematic study of this special case in a future work.

Example 5.2. Fix now $p \geq 2$, $p \in \mathbb{N}$, and consider instead the non-linear Cauchy problem

$$\begin{cases} \partial_t u = \frac{1}{p-1} u^p \\ u(0, x) = u_0(x) = (1 + x^2)^{-1}. \end{cases}$$

The maximal solution is defined in $[0, 1) \times \mathbb{R}$ and reads

$$u(t, x) = ((1 + x^2)^{p-1} - t)^{-1/(p-1)}.$$

The initial datum is the same as before. Now the solution extends to a holomorphic function in the strip $|y| < \sqrt{1 - t^{1/(p-1)}} \sim (p-1)^{-1/2} \sqrt{1-t}$ as $t \rightarrow 1$. This agrees with the formula (1.6) with $A = 1/2$; in fact, for $0 \leq t < 1$,

$$\exp\left(-\frac{1}{2} \int_0^t \|u(\sigma)\|_{L^\infty}^{p-1} d\sigma\right) = \exp\left(-\frac{1}{2} \int_0^t (1-\sigma)^{-1} d\sigma\right) = \sqrt{1-t}.$$

Remark 5.3. A careful inspection of the proof shows that the conclusions of Theorems 1.1 and 2.1 hold with a radius of analyticity

$$\varepsilon(t) = \frac{1}{B} \exp\left(-A \int_0^t (1 + \|u(\sigma)\|_{L^\infty}^{p-1}) d\sigma\right)$$

for a constant $A > 0$ depending only on the equation, and a constant B depending on the initial datum; to be precise $B \simeq (\|u_0\|_s + \sup_{N \geq 1} E_N^{\varepsilon_0}[u_0])^{p-1}$.

Remark 5.4. We observe that the condition (2.3) is used in the proof of Theorem 2.1 only to prove the energy estimate (2.5). Hence Theorem 2.1 still holds for every system admitting a smooth solution satisfying an estimate of the form (2.5). In particular, we can apply our result to a strictly hyperbolic system of the form

$$L'u = \mathcal{N}[u]$$

$\mathcal{N}[u]$ is of the form (1.4), (1.5) and

$$L' = \partial_t + A(t, x, D_x)$$

for some $n \times n$ matrix $A(t, x, \xi)$ satisfying the condition (2.2) and admitting distinct purely imaginary eigenvalues $i\lambda_j(t, x, \xi)$, $j = 1, \dots, n$, with

$$\lambda_1(t, x, \xi) < \lambda_2(t, x, \xi) < \dots < \lambda_n(t, x, \xi), \quad (t, x, \xi) \in \overline{\mathbb{R}}_+ \times \mathbb{R}^d \times \mathbb{R}^d.$$

It is well known that under these assumptions for s sufficiently large the Cauchy problem

$$\begin{cases} L'u = \mathcal{N}[u] \\ u(0, x) = u_0 \end{cases}$$

admits a solution $u \in C^1(I, H^s) \cap C^0(I, H^{s+1})$ defined for $(t, x) \in I \times \mathbb{R}^d$ for some time interval I and the energy estimate (2.5) holds true. Hence the arguments in the proof of Theorem 2.1 can be applied to prove the analyticity of the solution and the estimate of the radius of analyticity.

REFERENCES

- [1] S. Alinhac and G. Metivier, *Propagation de l'analyticit  des solutions de syst mes hyperboliques non-lin aires*, Invent. Math. **75** (1984), 189–204.
- [2] P. Albert, J. L. Bona and J.-C. Saut, *Model equations for waves in stratified fluids*, Proc. R. Soc. Lond. A **453** (1997), 1233–1260.

- [3] M.S. Baouendi and C. Goulaouic, *Cauchy problems for analytic pseudodifferential operators*, Comm. Partial Diff. Eqs. **1** (1976), 135–189.
- [4] M.S. Baouendi and C. Goulaouic, *Remarks on the abstract form of nonlinear Cauchy Kowalevsky theorems*, Comm. Partial Diff. Eqs. **2** (1977), 1151–1162.
- [5] T. B. Benjamin, J. L. Bona and D. K. Bose, *Solitary wave solutions of nonlinear problems*, Phil. Trans. R. Soc. Lond. A **331** (1990), 195–244.
- [6] H. A. Biagioni and T. Gramchev, *Fractional derivative estimates in Gevrey classes, global regularity and decay for solutions to semilinear equations in \mathbb{R}^n* , J. Differential Equations, **194** (2003), 140–165.
- [7] M. Bjørkavag and H. Kalisch, *Radius of analyticity and exponential convergence for spectral projections of the generalized KdV equation*, Commun. Nonlinear Sci. Numer. Simul. **15** (2010), no. 4, 869–880.
- [8] J. Bona and Z. Grujic', *Spatial analyticity properties of nonlinear waves*. Dedicated to Jim Douglas, Jr. on the occasion of his 75th birthday. Math. Models Methods Appl. Sci., **13** (2003), 345–360.
- [9] J. Bona, Z. Grujic' and H. Kalisch *Algebraic lower bounds for the uniform radius of spatial analyticity for the generalized KdV equation*, Ann. Inst. H. Poincaré Anal. Non Linéaire, **22** (2005), 783–797.
- [10] J. Bona and Y. Li, *Analyticity of solitary-wave solutions of model equations for long waves*. SIAM J. Math. Anal. **27** (1996), n. 3, 725–737.
- [11] J. Bona and Y. Li, *Decay and analyticity of solitary waves*, J. Math. Pures Appl., **76** (1997), 377–430.
- [12] J. Bona and F.B. Weissler, *Pole dynamics of interacting solitons and blowup of complex-valued solutions of KdV*, Nonlinearity, **22** (2009), 311–349.
- [13] M. Cappiello, T. Gramchev and L. Rodino, *Semilinear pseudo-differential equations and travelling waves*, Fields Institute Communications, **52** (2007), 213–238.
- [14] M. Cappiello and F. Nicola, *Holomorphic extension of solutions of semilinear elliptic equations*, Nonlinear Anal. **74** (2011), 2663–2681.
- [15] M. Cappiello and F. Nicola, *Regularity and decay of solutions of nonlinear harmonic oscillators*, Adv. Math. **229** (2012), 1266–1299.
- [16] M. Cappiello and F. Nicola, *On the holomorphic extension of solutions of elliptic pseudodifferential equations*, Pliska Stud. Math. Bulgar. **21** (2012), 113–126.
- [17] P. D'Ancona, L.Fanelli, L.Vega and N.Visciglia, *Endpoint Strichartz estimates for the magnetic Schrödinger equation*, J. Funct. Anal. **258** (2010), 3227–3240.
- [18] P. D'Ancona and S. Spagnolo, *Small analytic solutions to nonlinear weakly hyperbolic systems*, Ann. Scu. Norm. Sup. Pisa IV.22 (1995), 469–491.
- [19] P. D'Ancona and S. Spagnolo, *On the life span of the analytic solutions to quasilinear weakly hyperbolic equations*, Indiana Univ. Math. J. **40** (1991), 71–99.
- [20] J. Gorsky, A. A. Himonas, C. Holliman and G. Petronilho, *The Cauchy problem for a periodic higher order KdV equation in analytic Gevrey spaces*, J. Math. Anal. Appl., to appear.
- [21] Z. Grujic' and H. Kalisch, *Local well-posedness of the generalized Korteweg-de Vries equation in spaces of analytic functions*, Diff. Int. Equations **15** (2002), 1325–1334.
- [22] H. Hannah, A. A. Himonas and G. Petronilho, *Anisotropic Gevrey regularity for MKdV on the circle*, Discrete Contin. Dyn. Systems, Suppl. 2011, 634–642.
- [23] N. Hayashi, *Global existence of small analytic solutions to nonlinear Schrödinger equations*, Duke Math. J. **60** (1990), 717–727.

- [24] N. Hayashi, *Analyticity of solutions of the Korteweg-de Vries equation*, SIAM J. Math. Anal. **22** (1991) no. 6, 1738–1743.
- [25] N. Hayashi and K. Kato, *Global existence of small analytic solutions to Schrödinger equations with quadratic nonlinearity*, Comm. Partial Differential Equations **22** (1997), 773–798.
- [26] A. A. Himonas and G. Petronilho, *Analytic well-posedness of periodic gKdV*, J. Differ. Equations **253** (2012) no. 11, 3101–3112.
- [27] L. Hörmander, *The analysis of linear partial differential operators, III*, Springer-Verlag, 1985.
- [28] T. Kato and K. Masuda, *Nonlinear evolution equations and analyticity. I*, Ann. Inst. H. Poincaré Anal. Non Linéaire **3** (1986), 455–467.
- [29] I. Kukavica and V. Vicol, *On the radius of analyticity of solutions to the three-dimensional Euler equations*, Proc. Amer. Math. Soc. **137** (2009) no. 2, 669–677.
- [30] S. Mizohata, *The theory of partial differential equations*, Cambridge University Press, Cambridge, 1973.
- [31] F. Nicola and L. Rodino, *Global Pseudodifferential Calculus on Euclidean Spaces*, Birkhäuser, Basel, 2010.
- [32] J. Rauch, *Hyperbolic Partial Differential Equations and Geometric Optics*, GSM Series **133**, Amer. Math. Soc., 2012.

DIPARTIMENTO DI MATEMATICA, UNIVERSITÀ DEGLI STUDI DI TORINO, VIA CARLO ALBERTO 10, 10123 TORINO, ITALY

SAPIENZA – UNIVERSITÀ DI ROMA, DIPARTIMENTO DI MATEMATICA, PIAZZALE A. MORO 2, I-00185 ROMA, ITALY

DIPARTIMENTO DI SCIENZE MATEMATICHE, POLITECNICO DI TORINO, CORSO DUCA DEGLI ABRUZZI 24, 10129 TORINO, ITALY

E-mail address: marco.cappiello@unito.it

E-mail address: dancona@mat.uniroma1.it

E-mail address: fabio.nicola@polito.it