

# Observation of $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$ and study of $D^+ \rightarrow \pi^+\pi^-\ell^+\nu_\ell$ decay dynamics

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Using  $2.93 \text{ fb}^{-1}$  of  $e^+e^-$  collision data collected with the BESIII detector at the center-of-mass energy of 3.773 GeV, we investigate the semileptonic decays  $D^+ \rightarrow \pi^+\pi^-\ell^+\nu_\ell$  ( $\ell = e$  and  $\mu$ ). The  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  decay is observed for the first time. By analyzing simultaneously the differential decay rates of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  and  $D^+ \rightarrow f_0(500)e^+\nu_e$  in different  $\ell^+\nu_\ell$  four-momentum transfer intervals, the product of the relevant hadronic form factor  $f_+^{f_0}(0)$  and the magnitude of the  $c \rightarrow d$  Cabibbo-Kobayashi-Maskawa matrix element  $|V_{cd}|$  is determined to be  $f_+^{f_0}(0)|V_{cd}| = 0.143 \pm 0.014_{\text{stat}} \pm 0.011_{\text{syst}}$  for the first time. With the input of  $|V_{cd}|$  from the global fit in the standard model, we determine  $f_+^{f_0}(0) = 0.63 \pm 0.06_{\text{stat}} \pm 0.05_{\text{syst}}$ . The absolute branching fractions of  $D^+ \rightarrow f_0(500)_{(\pi^+\pi^-)}\mu^+\nu_\mu$  and  $D^+ \rightarrow \rho^0_{(\pi^+\pi^-)}\mu^+\nu_\mu$  are determined as  $(0.72 \pm 0.13_{\text{stat}} \pm 0.08_{\text{syst}}) \times 10^{-3}$  and  $(1.64 \pm 0.13_{\text{stat}} \pm 0.10_{\text{syst}}) \times 10^{-3}$ . Combining these results with those of previous BESIII measurements on their semi-electronic counterparts from the same data sample, we test lepton flavor universality by measuring the branching fraction ratios  $\mathcal{B}_{D^+ \rightarrow \rho^0\mu^+\nu_\mu}/\mathcal{B}_{D^+ \rightarrow \rho^0e^+\nu_e} = 0.88 \pm 0.10$  and  $\mathcal{B}_{D^+ \rightarrow f_0(500)\mu^+\nu_\mu}/\mathcal{B}_{D^+ \rightarrow f_0(500)e^+\nu_e} = 1.14 \pm 0.26$ , which are compatible with the standard model expectation.

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## I. MOTIVATION

A series of ground breaking discoveries of the  $XYZ$  states [1], implying tetraquark composition, at the turn of this century and the subsequent observation of pentaquarks  $P_c$  [2,3] have revolutionized hadron spectroscopy. These related topics have been the focus of much experimental and theoretical interest in recent years. Quantum chromodynamics (QCD) has been established for half a century as the fundamental theory of the strong interaction, whose development has benefited from such spectroscopy studies since the beginning. However, it is surprising that the nature of the light scalar mesons, which constitute the spectrum of lowest mass states, is still under debate after almost seven decades [1,4,5]. Notably, they are located near the  $S$ -wave thresholds of two-body final states, such as  $KK$  and  $\pi\pi$ . The ongoing puzzle behind these phenomena can be ascribed to the nonperturbative dynamics at low energies and quark confinement.

The lightest scalar meson  $f_0(500)$ , traditionally known as the  $\sigma$  meson, was not well-established for many years, with consensus on its mass and width only being reached in the past decade or so [4]. The list of references is huge and more information can be found in the Review of Particle Properties [1], and recent reviews [4,6]. Sharing the same spin-parity quantum numbers with the QCD vacuum, this meson plays a significant role in the dynamical generation of mass through the spontaneous breaking of QCD chiral symmetry, and consequently the confinement of quarks. Clarifying the nature of the  $f_0(500)$  could shed light on these issues. However, a  $q\bar{q}$  configuration of the  $f_0(500)$  and other light scalar mesons forming the SU(3) nonet in the naive quark model cannot explain their inverted mass hierarchy. There is still the possibility that these are mixtures of  $q\bar{q}$  states [7]. Other interpretations include diquark-antidiquark states [8], meson-meson bound states [9], and even more complicated scenarios [10].

The form factor (FF) of semileptonic (SL)  $D^+$  transition to the  $f_0(500)$  not only provides new information about the nature of the  $f_0(500)$  but also helps to understand the dynamics of SL charmed-meson decays by testing different nonperturbative theoretical methods. Predictions exist of the branching fractions (BFs) [11–13] and the FF [14–16] of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  ( $\ell = e$  and  $\mu$ ). Additionally, scrutinizing the BF ratios of SL  $D$  decays allows for probes of lepton flavor universality (LFU), and therefore tests the standard model (SM) in the charm sector [17].

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To date, the  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  decay has not yet been observed, and only the FOCUS, E791, and E687 experiments have performed studies of  $D^+ \rightarrow \rho^0\mu^+\nu_\mu$  [18–20]; while the BESIII experiment has observed  $D^+ \rightarrow f_0(500)e^+\nu_e$  [21]. This paper reports the observation and BF measurement of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$ , studies of the dynamics of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ , the absolute measurement of the BF of  $D^+ \rightarrow \rho^0\mu^+\nu_\mu$ , and comprehensive LFU tests with  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$  and  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ , for the first time. These analyses make use of  $2.93 \text{ fb}^{-1}$  of  $e^+e^-$  collision data taken at the center-of-mass energy of 3.773 GeV. Throughout this paper, charge conjugate channels are always implied. Due to slightly different event selection criteria with Ref. [21] as mentioned later, the BFs of  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$ ,  $D^+ \rightarrow f_0(500)e^+\nu_e$ , and  $D^+ \rightarrow f_0(500)e^+\nu_e$  obtained in this work offer independent checks; while the hadronic FF of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  is measured for the first time.

## II. BESIII DETECTOR AND MONTE CARLO SIMULATIONS

The BESIII [22] detector is a magnetic spectrometer located at the Beijing Electron Positron Collider (BEPCII) [23]. More details about the BESIII detector are described in Ref. [22]. Monte Carlo (MC) simulated data samples, produced with a GEANT4-based [24] software package including the geometric description of the BESIII detector [25] and the detector response, are used to determine the detection efficiencies and to estimate the background contributions. The simulation includes the beam-energy spread and initial-state radiation in the  $e^+e^-$  annihilations modeled with the generator KKMC [26]. The inclusive MC sample consists of the production of  $D\bar{D}$  pairs with consideration of quantum coherence for all neutral  $D$  modes, the non- $D\bar{D}$  decays of the  $\psi(3770)$ , the initial-state radiation production of the  $J/\psi$  and  $\psi(3686)$  states, and the continuum processes. The known decay modes are modeled with EVTGEN [27] using the BFs taken from the PDG [1], and the remaining unknown decays from the charmonium states are modeled with LUNDCHARM [28]. Final-state radiation from charged final state particles is incorporated with the PHOTOS package [29]. The SL decays  $D^+ \rightarrow \pi^+\pi^-\ell^+\nu_\ell$  are simulated according to the results of a previous BESIII amplitude analysis measurement in the  $e$  channel [21].

## III. SINGLE-TAG $\bar{D}$ CANDIDATES

A detailed description of the selection criteria for charged and neutral particles is provided in Refs. [30–32]. The tagged  $D^-$  mesons are reconstructed in six hadronic final states  $K^+\pi^-\pi^-$ ,  $K^+\pi^-\pi^-\pi^0$ ,  $K_S^0\pi^-$ ,  $K_S^0\pi^-\pi^0$ ,  $K_S^0\pi^-\pi^+\pi^-$  and  $K^+K^-\pi^-$ . The  $D^-$  signals are

separated from backgrounds with two kinematic variables;  $\Delta E \equiv E_{D^-} - E_{\text{beam}}$  and  $M_{\text{BC}} \equiv \sqrt{E_{\text{beam}}^2 - |\vec{p}_{D^-}|^2}$ , where  $E_{D^-}$  and  $\vec{p}_{D^-}$  are the reconstructed energy and momentum of the  $D^-$  candidate in the  $e^+e^-$  center-of-mass frame, and  $E_{\text{beam}}$  is the beam energy. If there are multiple candidates per tag mode, the one with the minimum  $|\Delta E|$  is kept. For each tag mode, the yield is extracted by fitting the  $M_{\text{BC}}$  distribution as in Refs. [30–32]. The total yield of tagged  $D^-$  mesons is  $N_{\text{tag}}^{\text{tot}} = (1522.5 \pm 2.2_{\text{stat}}) \times 10^3$ .

## IV. DOUBLE-TAG EVENTS

After a tagged candidate is selected, the  $D^+ \rightarrow \pi^+\pi^-\ell^+\nu_\ell$  decays are reconstructed recoiling against the tagged  $D^-$ . This requires a  $\ell^+$  candidate and a  $\pi^+\pi^-$  pair in the signal side. The  $\pi^\pm$  and  $e^+$  are identified with the same criteria as Ref. [31]. Particle identification (PID) for the  $\mu^+$  considers the measurements of the specific ionization energy loss by the main drift chamber (MDC), the flight time by the time-of-flight system, and the energy deposited in the electromagnetic calorimeter (EMC). Based on these measurements, we calculate the combined confidence levels for positron ( $CL_e$ ), muon ( $CL_\mu$ ), pion ( $CL_\pi$ ), and kaon ( $CL_K$ ) hypotheses for each charged track. The charged track satisfying  $CL_\mu > 0.001$ ,  $CL_\mu > CL_e$ ,  $CL_\mu > CL_\pi$ ,  $CL_\mu > CL_K$ , and  $E_{\text{EMC}} \in (0.09, 0.31) \text{ GeV}$  is assigned as  $\mu^+$  candidate, where  $E_{\text{EMC}}$  is its energy deposited in the EMC. The two requirements  $CL_\mu > CL_\pi$  and  $E_{\text{EMC}}$  together suppress about 75% of background at the cost of 50% of signal. To suppress the backgrounds from the hadronic  $D$  decays, the maximum energy of any photon that is not used in the tag selection ( $E_{\text{extra}\gamma}^{\text{max}}$ ) is required to be less than 0.25 GeV, and no additional charged track ( $N_{\text{char}}^{\text{extra}}$ ) is allowed.

In the selection of  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$ , the energy deposited in the EMC of  $\pi^+$  is required to be less than 0.8 times its reconstructed momentum in the MDC, to minimize the misidentified background from positrons. To reject the background events associated with  $K_S^0 \rightarrow \pi_{\pi \rightarrow \mu}^+\pi^-$ , where  $\pi_{\pi \rightarrow \mu}^+$  denotes a track of pion identified as a muon candidate, the invariant mass of the  $\mu^+\pi^-$  combination ( $M_{\mu^+\pi^-}$ ) is required to be outside the interval (0.446, 0.518) GeV/ $c^2$ . To reduce the peaking backgrounds of  $D^+ \rightarrow \bar{K}^0\pi^+(\pi^0)$  and  $D^+ \rightarrow \pi^+\pi^+\pi^-(\pi^0)$ , the invariant mass of the  $\pi^+\pi^-\mu^+$  combination ( $M_{\pi^+\pi^-\mu^+}$ ) is required to be less than 1.56 GeV/ $c^2$ . To further suppress backgrounds associated with additional neutral pions, it is required that there is no extra  $\pi^0$  besides that used in the SL selection ( $N_{\pi^0}^{\text{extra}}$ ), which suppresses about of 36% background at the cost of 4% of signal.

To help separate the SL signals from backgrounds, we define a kinematic quantity of  $U_{\text{miss}} \equiv E_{\text{miss}} - |\vec{p}_{\text{miss}}|$ .

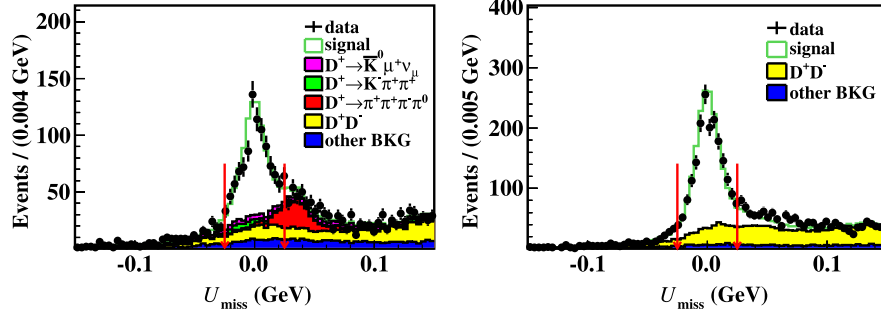


FIG. 1. The  $U_{\text{miss}}$  distributions of the accepted candidates for (left)  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$  and (right)  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$ . The red arrows denote the signal region. The events entering the distributions are required to lie outside  $M_{\pi^+\pi^-} \in (0.461, 0.533)$   $\text{GeV}/c^2$  for  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$ , and  $M_{\pi^+\pi^-} \in (0.42, 0.56)$   $\text{GeV}/c^2$  for  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$ .

Here,  $E_{\text{miss}} \equiv E_{\text{beam}} - E_{\pi^+\pi^-} - E_{\ell^+}$  and  $\vec{p}_{\text{miss}} \equiv \vec{p}_{D^+} - \vec{p}_{\pi^+\pi^-} - \vec{p}_{\ell^+}$  are the missing energy and momentum of the SL event in the  $e^+e^-$  center-of-mass frame, respectively, in which  $E_{\pi^+\pi^-}(\ell^+)$  and  $\vec{p}_{\pi^+\pi^-}(\ell^+)$  are the energy and momentum of the  $\pi^+\pi^-(\ell^+)$  candidates. The  $U_{\text{miss}}$  resolution is improved using  $\vec{p}_{D^+} \equiv -\hat{p}_{D^-} \cdot \sqrt{E_{\text{beam}}^2 - m_{D^+}^2}$ , where  $\hat{p}_{D^-}$  is the unit vector in the momentum direction of the tag  $D^-$  and  $m_{D^+}$  is the nominal  $D^+$  mass [1]. The  $U_{\text{miss}}$  distribution of the correctly reconstructed signals is expected to peak around zero, as shown in Fig. 1. The events within  $U_{\text{miss}} \in (-0.025, 0.025)$   $\text{GeV}$  are retained for further analysis. This requirement retains 83% of the signal and removes 74% of the all backgrounds for  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$ , in which the major peaking background is from  $D^+ \rightarrow \pi^+\pi^-\pi^+\pi^0$ ; while retains 75% of the signal and removes 82% of the all backgrounds for  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$ .

## V. BRANCHING FRACTIONS

### A. Signal efficiencies and branching fractions

Figure 2 shows the  $M_{\pi^+\pi^-}$  distributions of the accepted SL events in data. According to the previous study of  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$  [21], we ignore the nonresonant component except for  $f_0(500)$ ,  $\rho^0$ , and  $\omega$ . In this case, the signal

yields from different components are determined by an unbinned maximum likelihood fit to these distributions. In the fit, the shapes of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ ,  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$ ,  $D^+ \rightarrow \omega\ell^+\nu_\ell$  are derived from individual signal MC samples where the Bugg function [33], Gounaris-Sakurai (GS) function [34] and GS  $\times$  RBW function [35] are used to describe the  $f_0(500)$ ,  $\rho^0$  and  $\omega$  resonances, respectively. Here RBW is a relativistic Breit-Wigner function with a constant width [35]. The background shapes are derived from the inclusive MC sample. MC studies show that the peaking backgrounds around 0.498  $\text{GeV}/c^2$  are mainly from the decays  $D^+ \rightarrow K_S^0\ell^+\nu_\ell$  and some hadronic  $D$  decays containing  $K_S^0$ . In the fit, the peaking background component containing  $K_S^0$  has been forced to have the same normalization as the smoother background derived from the inclusive MC sample. The amplitude analysis of  $D \rightarrow \pi^+\pi^-e^+\nu_e$  in Ref. [21] shows that the interference between  $D^+ \rightarrow f_0(500)e^+\nu_e$  and  $D^+ \rightarrow \rho^0e^+\nu_e$  is negligible. With the same data sample, we take this conclusion into account to simplify the analysis. The yields of  $D^+ \rightarrow \omega\ell^+\nu_\ell$  are fixed based on the BF previously measured by BESIII [21], while the yields of other components are free parameters. The fit results on the  $M_{\pi^+\pi^-}$  distributions are also shown in Fig. 2. From these fits, we obtain the yields of  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$  and  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  decays [36].

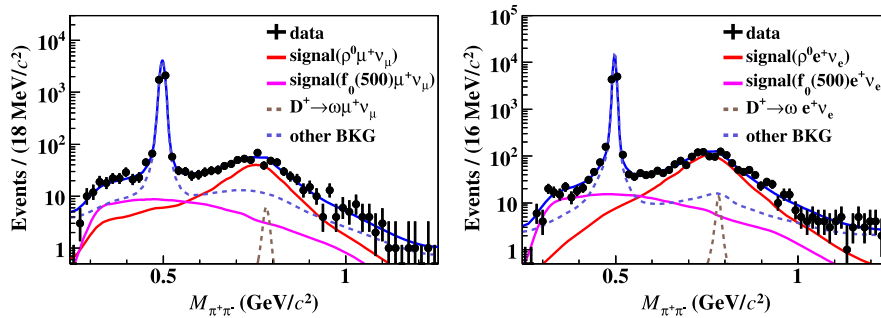


FIG. 2. The  $M_{\pi^+\pi^-}$  distributions of the candidates for (left)  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$  and (right)  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$  with the fit results overlaid (blue solid line). The contributions of the signal and various background components are indicated, where ‘other BKG’ includes all other components.

TABLE I. The signal yields in data, the statistical significances, the signal efficiencies and the absolute BF of different signal modes. The BFs of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  include the BF of  $f_0(500) \rightarrow \pi^+\pi^-$ . The statistical uncertainties of the  $\mu$  channels are worse than those of the  $e$  channels by a factor of two mainly due to stricter selection criteria and much higher backgrounds.

Signal mode	$N_{\text{obs}}$	$S$ ( $\sigma$ )	$\epsilon_{\text{sig}}$ (%)	$\mathcal{B}_{\text{sig}}(\times 10^{-3})$
$f_0(500)\mu^+\nu_\mu$	$209 \pm 38$	5.9	$18.93 \pm 0.13$	$0.72 \pm 0.13$
$\rho^0\mu^+\nu_\mu$	$496 \pm 38$	>10	$19.86 \pm 0.13$	$1.64 \pm 0.13$
$f_0(500)e^+\nu_e$	$412 \pm 43$	>10	$44.76 \pm 0.25$	$0.60 \pm 0.06$
$\rho^0e^+\nu_e$	$1237 \pm 47$	>10	$44.12 \pm 0.25$	$1.84 \pm 0.07$

The BF of the SL decay is determined by

$$\mathcal{B}_{\text{sig}} = N_{\text{obs}} / (N_{\text{tag}}^{\text{tot}} \cdot \epsilon_{\text{sig}}), \quad (1)$$

where  $N_{\text{tag}}^{\text{tot}}$  and  $N_{\text{obs}}$  are the tag and SL signal yields in the data sample, and  $\epsilon_{\text{sig}} = \sum_i [\epsilon_{\text{tag, sig}}^i \cdot N_{\text{tag}}^i] / (\epsilon_{\text{tag}}^i \cdot N_{\text{tag}}^{\text{tot}})$  is the efficiency of detecting the SL decay in the presence of the tag  $D^-$  meson. Here,  $i$  denotes the tag mode, and  $\epsilon_{\text{tag}}$

and  $\epsilon_{\text{tag, sig}}$  are the efficiencies of selecting the tag and simultaneously selecting the tag and signal candidates, respectively. Table I summarizes the SL yields in data ( $N_{\text{obs}}$ ), the statistical significances ( $S$ ), the signal efficiencies ( $\epsilon_{\text{sig}}$ ), and the obtained BFs ( $\mathcal{B}_{\text{sig}}$ ). For each signal decay, the statistical significance is calculated according to  $\sqrt{-2\ln(\mathcal{L}_0/\mathcal{L}_{\text{max}})}$ , where  $\mathcal{L}_{\text{max}}$  and  $\mathcal{L}_0$  are the maximum likelihoods of the fits with and without including the signal component, respectively. The reliability of the efficiency determinations has been verified by the comparison of the distributions of momenta and  $\cos\theta$  of the  $\pi^+$ ,  $\pi^-$ , and  $\ell^+$  of the selected  $D^+ \rightarrow \pi^+\pi^-\ell^+\nu_\ell$  candidate events between data and MC simulation. As shown in Fig. 3, a good consistency between data and MC simulation for  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$  is observed. Similar comparison has also shown good data-MC consistence of these variables for  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$  in Ref. [21].

## B. Systematic uncertainties in branching fractions

The systematic uncertainties in the BF measurements are discussed below. The uncertainty in the total yield of tag

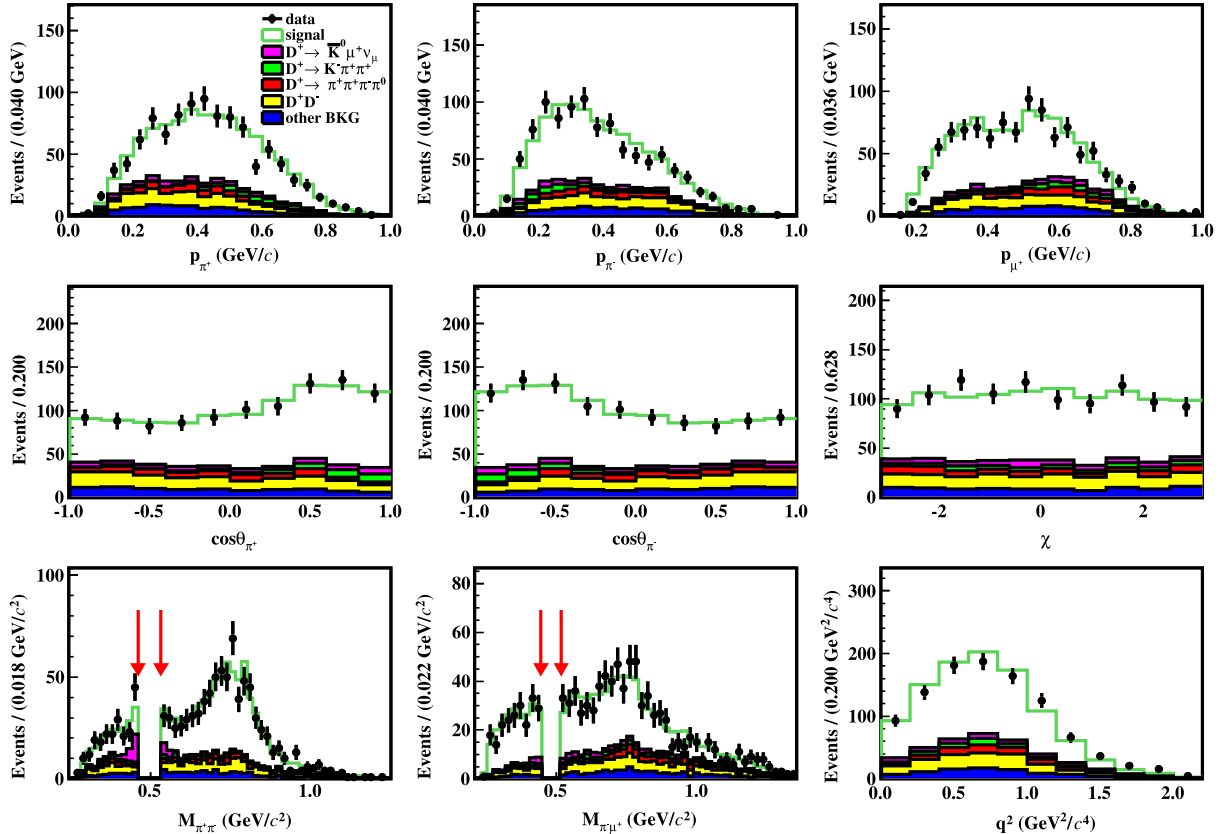


FIG. 3. The distributions of (top row) momenta and (middle row)  $\cos\theta$  of  $\pi^+$ ,  $\pi^-$  and  $\mu^+$ , (bottom row) invariant-mass distributions of  $\pi^+\pi^-$ ,  $\pi^-\mu^+$  and  $q^2$  distribution from the accepted candidates of  $D^+ \rightarrow \pi^+\pi^-\mu^+\nu_\mu$  for data and inclusive MC sample.  $q^2$  is the mass square of the  $\mu^+\nu_\mu$  system;  $\cos\theta_{\pi^+}$  ( $\cos\theta_{\pi^-}$ ), the helicity angle is the angle between the  $\pi^+$  ( $\pi^-$ ) and the mother decay plane in the  $\pi^+\pi^-$  rest frame;  $\chi$  is the coplanarity angle between the  $\pi^+\pi^-$  and  $\mu^+\nu_\mu$  decay planes. To enhance the signal distributions, an additional requirement of  $M_{\pi^+\pi^-} \in (0.461, 0.533)$  GeV/ $c^2$  has been imposed in addition to the baseline selection criteria.

TABLE II. Relative systematic uncertainties (in %) in the measurements of the BFs of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  and  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$ .

Source	$D^+ \rightarrow \rho^0\mu^+\nu_\mu$	$D^+ \rightarrow f_0(500)\mu^+\nu_\mu$	$D^+ \rightarrow \rho^0e^+\nu_e$	$D^+ \rightarrow f_0(500)e^+\nu_e$
$N_{\text{tag}}^{\text{tot}}$	0.5	0.5	0.5	0.5
$\pi^\pm$ tracking	0.4	0.4	0.4	0.4
$\pi^\pm$ PID	0.6	0.6	0.6	0.6
$\ell^+$ tracking	0.2	0.2	0.2	0.2
$\ell^+$ PID	0.5	0.5	0.5	0.5
$U_{\text{miss}}$ requirement	0.5	0.5	0.5	0.5
$E_{\text{extray}}^{\text{max}}$ & $N_{\text{char}}^{\text{extra}}$	1.3	1.3	1.3	1.3
$M_{\pi^-\mu^+}$ requirement	0.2	0.2	...	...
$M_{\pi^+\pi^-\mu^+}$ requirement	0.5	0.5	...	...
$M_{\pi^+\pi^-}$ fit	4.9	10.2	2.8	7.4
MC sample size	0.6	0.7	0.6	0.6
MC model	3.6	3.6	0.9	0.9
Total	6.4	11.0	3.5	7.7

$D^-$  mesons was previously assigned as 0.5% in Refs. [30–32]. The systematic uncertainty in the tracking (PID) efficiency per  $\pi^\pm$  is assigned as 0.2% (0.3%), by analyzing the  $D\bar{D}$  hadronic events [37]. The systematic uncertainties due to the tracking (PID) efficiencies per  $\mu^+$  or  $e^+$  are assigned as 0.2% (0.5%) by using the  $e^+e^- \rightarrow \gamma\mu^+\mu^-$  and  $e^+e^- \rightarrow \gamma e^+e^-$  control samples, respectively, where the data-MC differences of the two-dimensional (momentum and  $\cos\theta$ ) distributions of the control samples are re-weighted to match those of the SL signal decays. The uncertainty associated with the combined  $E_{\text{extray}}^{\text{max}}$ ,  $N_{\text{char}}^{\text{extra}}$  (and  $N_{\pi^0}^{\text{extra}}$ ) requirements is taken to be 1.3% by analyzing hadronic  $D^+D^-$  candidate events. The systematic uncertainty due to the  $U_{\text{miss}}$  requirement is estimated to be 0.5% by analyzing the control sample of  $D^+ \rightarrow K_S^0\ell^+\nu_\ell$ .

To estimate the uncertainties due to the  $M_{\mu^+\pi^-}$  requirements, we remeasure the BFs while varying the  $M_{\mu^+\pi^-}$  requirement in the range from (0.456, 0.508)  $\text{GeV}/c^2$  to (0.436, 0.528)  $\text{GeV}/c^2$ , with a step of 4  $\text{MeV}/c^2$  corresponding to the approximate  $K_S^0$  mass resolution. To estimate the uncertainty due to the  $M_{\pi^+\pi^-\mu^+}$  requirement, we remeasure the BFs while varying the  $M_{\pi^+\pi^-\mu^+}$  condition from 1.4  $\text{GeV}/c^2$  to 1.6  $\text{GeV}/c^2$ , with a step size of 20  $\text{MeV}/c^2$ . Fitting the measurements obtained from this procedure with a linear function yields an average BF from this exercise. The differences between the average BF and the baseline value are assigned as the systematic uncertainties, which are 0.2% and 0.5% for the  $M_{\mu^+\pi^-}$  and  $M_{\pi^+\pi^-\mu^+}$  requirements, respectively.

The uncertainty associated with the  $M_{\pi^+\pi^-}$  fit is estimated by examining the BF changes when using an RBW signal description with Gaussian smearing of the MC-simulated signal shape, varying the amplitude analysis model parameters in their uncertainties 600 times randomly, varying the assumed BFs of the peaking background components

by  $\pm 1\sigma$ , and varying the smooth parameters of other background contributions [36]. The uncertainties arising from the finite sample size of the signal MC are 0.7%, 0.6% and 0.6% for  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$ ,  $D^+ \rightarrow \rho^0\mu^+\nu_\mu$ , and  $D^+ \rightarrow f_0(500)e^+\nu_e$ , respectively. The uncertainties due to the signal MC model for these three decays are 3.6%, 3.6% and 0.9%, respectively, which are the differences of the baseline signal efficiencies and those obtained with the signal MC samples after varying the values of the input FFs by  $\pm 1\sigma$ .

For each signal decay, adding these uncertainties in quadrature yields the total systematic uncertainty, which is 11.0% for  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$ , 6.4% for  $D^+ \rightarrow \rho^0\mu^+\nu_\mu$ , 7.7% for  $D^+ \rightarrow f_0(500)e^+\nu_e$ , and 3.5% for  $D^+ \rightarrow \rho^0e^+\nu_e$ . Table II summarizes the sources of the systematic uncertainties in the measurements of the BFs of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  and  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$ .

### C. Results and discussion

Finally, the BFs of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  with  $f_0(500) \rightarrow \pi^+\pi^-$  and  $D^+ \rightarrow \rho^0\mu^+\nu_\mu$  are determined to be  $\mathcal{B}_{D^+ \rightarrow f_0(500)\mu^+\nu_\mu} \times \mathcal{B}_{f_0(500) \rightarrow \pi^+\pi^-} = (0.72 \pm 0.13_{\text{stat}} \pm 0.08_{\text{syst}}) \times 10^{-3}$  and  $\mathcal{B}_{D^+ \rightarrow \rho^0\mu^+\nu_\mu} = (1.64 \pm 0.13_{\text{stat}} \pm 0.10_{\text{syst}}) \times 10^{-3}$ . As cross checks, we have also determined the BFs of  $D^+ \rightarrow f_0(500)e^+\nu_e$ ,  $f_0(500) \rightarrow \pi^+\pi^-$  and  $D^+ \rightarrow \rho^0e^+\nu_e$  to be  $\mathcal{B}_{D^+ \rightarrow f_0(500)e^+\nu_e} \times \mathcal{B}_{f_0(500) \rightarrow \pi^+\pi^-} = (0.60 \pm 0.06_{\text{stat}} \pm 0.05_{\text{syst}}) \times 10^{-3}$  and  $\mathcal{B}_{D^+ \rightarrow \rho^0e^+\nu_e} = (1.84 \pm 0.07_{\text{stat}} \pm 0.06_{\text{syst}}) \times 10^{-3}$ . They are consistent with those from the amplitude analysis of  $D^+ \rightarrow \pi^+\pi^-e^+\nu_e$  [21]. Combining the  $\mathcal{B}_{D^+ \rightarrow h\mu^+\nu_\mu}$  ( $h = f_0(500)$  or  $\rho^0$ ) obtained in this work and the  $\mathcal{B}_{D^+ \rightarrow he^+\nu_e}$  from a previous BESIII analysis [21] gives the BF ratios  $\mathcal{B}_{D^+ \rightarrow \rho^0\mu^+\nu_\mu} / \mathcal{B}_{D^+ \rightarrow \rho^0e^+\nu_e} = 0.88 \pm 0.10$  and  $\mathcal{B}_{D^+ \rightarrow f_0(500)\mu^+\nu_\mu} / \mathcal{B}_{D^+ \rightarrow f_0(500)e^+\nu_e} = 1.14 \pm 0.26$ , which are compatible with the predictions of 0.96 [38,39] and

0.89 ~ 0.91 [11,12], respectively, made under the assumption of lepton universality.

## VI. HADRONIC FORM FACTORS

### A. Hadronic form factors

We study the  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  decay dynamics, by dividing the SL candidate events into three  $q^2$  intervals: (0.0, 0.5), (0.5, 1.0), and (1.0, 2.2)  $\text{GeV}^2/c^4$ , where  $q^2$  is the four-momentum transfer square of the  $\ell^+\nu_\ell$  system. A least- $\chi^2$  fit is performed to the measured ( $\Delta\Gamma_{\text{msr}}^i$ ) and theoretically expected ( $\Delta\Gamma_{\text{th}}^i$ ) partial decay rates in the  $i$ th  $q^2$  interval. The  $\Delta\Gamma_{\text{msr}}^i$  is determined by

$$\Delta\Gamma_{\text{msr}}^i = \sum_j^3 (\varepsilon^{-1})_{ij} N_{\text{obs}}^j / (\tau_{D^+} \cdot N_{\text{tag}}^{\text{tot}}), \quad (2)$$

where  $N_{\text{obs}}^j$  is the SL signal yield in the  $j$ th  $q^2$  interval in data,  $\tau_{D^+}$  is the  $D^+$  lifetime, and  $\varepsilon$  is the efficiency matrix. Figure 4 shows the results of the fits to the  $M_{\pi^+\pi^-}$  distributions in the reconstructed  $q^2$  bins, using the same fitting method as the data. Table III gives the weighted efficiency matrices of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ , which have been averaged over all six tag modes. Table IV presents the numbers of reconstructed events in data  $N_{\text{obs}}^i$  obtained by fitting the  $M_{\pi^+\pi^-}$  distributions, the numbers of produced events  $N_{\text{pro}}^i$ , and the measured partial decay rates  $\Delta\Gamma_{\text{msr}}^i$  in various  $q^2$  intervals for  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ .

TABLE III. The efficiency matrices of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  averaged over six tag modes, where  $\varepsilon_{ij}$  represents the efficiency of events produced in the  $j$ th  $q^2$  interval and reconstructed in the  $i$ th  $q^2$  interval.

$\varepsilon_{ij}$ (%)	$D^+ \rightarrow f_0(500)\mu^+\nu_\mu$			$D^+ \rightarrow f_0(500)e^+\nu_e$		
	1	2	3	1	2	3
1	15.18	0.83	0.41	45.33	0.96	0.02
2	1.31	19.01	0.98	0.45	42.45	1.10
3	1.05	0.89	17.17	0.02	0.55	40.11

The  $\Delta\Gamma_{\text{th}}^i$  is expressed as [40,41]

$$\Delta\Gamma_{\text{th}}^i = \int_{q_{\text{min}}^2}^{q_{\text{max}}^2(i)} \int_{s_{\text{min}}}^{s_{\text{max}}} \frac{G_F^2 |V_{cd}|^2}{192\pi^4 m_{D^+}^3} \left(1 - \frac{m_\ell^2}{q^2}\right)^2 \times \left(1 + \frac{m_\ell^2}{2q^2}\right) \lambda^{\frac{3}{2}}(m_{D^+}^2, s, q^2) |f_+^0(q^2)|^2 P(s) ds dq^2, \quad (3)$$

where  $G_F$  is the Fermi constant,  $|V_{cd}|$  is the magnitude of the Cabibbo-Kobayashi-Maskawa matrix element governing  $c \rightarrow d$  transition,  $m_{D^+}$  is the mass of  $D^+$ ,  $m_\ell$  is the mass of lepton,  $f_0$  denotes the  $f_0(500)$  meson,  $s$  is the mass-squared of the  $\pi^+\pi^-$  pair,  $\lambda(x, y, z) = x^2 + y^2 + z^2 - 2xy - 2xz - 2yz$ , and the  $P(s)$  is formed with the Bugg resonant lineshape.

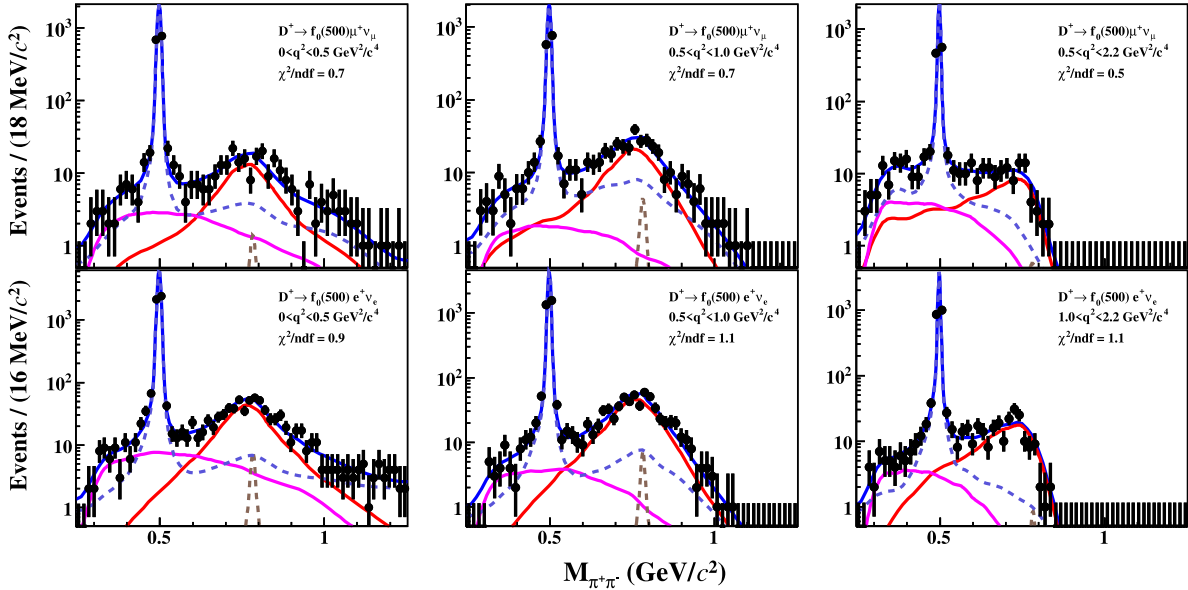


FIG. 4. Fits to the  $M_{\pi^+\pi^-}$  distributions of the  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  candidate events. Data are shown as dots with error bars. The blue solid curves are the fit results; the red solid line and pink solid lines denote the SL signals of  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$  and  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ , respectively; the black dotted curves are the peaking background of  $D^+ \rightarrow \omega\ell^+\nu_\ell$ ; and the blue dotted curves are the fitted combinatorial backgrounds.

TABLE IV. The partial decay rates of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  in various  $q^2$  intervals.

$i$	$D^+ \rightarrow f_0(500)\mu^+\nu_\mu$			$D^+ \rightarrow f_0(500)e^+\nu_e$		
	1	2	3	1	2	3
$q^2(\text{GeV}^2/c^4)$	(0.0, 0.5)	(0.5, 1.0)	(1.0, 2.2)	(0.0, 0.5)	(0.5, 1.0)	(1.0, 2.2)
$N_{\text{obs}}^i$	$74 \pm 21$	$45 \pm 22$	$72 \pm 22$	$221 \pm 31$	$95 \pm 21$	$63 \pm 16$
$N_{\text{pro}}^i$	$469 \pm 136$	$187 \pm 115$	$380 \pm 131$	$483 \pm 68$	$215 \pm 49$	$154 \pm 39$
$\Delta\Gamma_{\text{msr}}^i(\text{ns}^{-1})$	$0.298 \pm 0.087$	$0.119 \pm 0.073$	$0.242 \pm 0.083$	$0.307 \pm 0.043$	$0.137 \pm 0.031$	$0.098 \pm 0.025$

In this analysis, the hadronic FF ( $f_+^{f_0}(q^2)$ ) is parametrized by a two-parameter series expansion [42] due to very limited statistics. The series expansion is given as [42]

$$f_+(t) = \frac{1}{P(t)\Phi(t, t_0)} a_0(t_0) \left( 1 + \sum_{k=1}^{\infty} r_k(t_0) [z(t, t_0)]^k \right), \quad (4)$$

where

$$\begin{aligned} z(t, t_0) &= \frac{\sqrt{t_+ - t} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - t} + \sqrt{t_+ - t_0}}, \\ t_{\pm} &= (m_D \pm m_{f_0})^2, \\ t_0 &= t_+ (1 - \sqrt{1 - t_-/t_+}). \end{aligned} \quad (5)$$

The  $m_{f_0}$ , mass of  $f_0(500)$  meson, is taken as  $(538 \pm 12)$  MeV/ $c^2$  [43]. The function  $P(t) = z(t, m_{D^*}^2)$  and  $\Phi$  is given by

$$\begin{aligned} \Phi(t, t_0) &= \sqrt{\frac{1}{24\pi\chi_V} \left( \frac{t_+ - t}{t_+ - t_0} \right)^{1/4}} (\sqrt{t_+ - t} + \sqrt{t_+})^{-5} \\ &\times (\sqrt{t_+ - t} + \sqrt{t_+ - t_0}) (\sqrt{t_+ - t} + \sqrt{t_+ - t_-})^{3/2} \\ &\times (t_+ - t)^{3/4}, \end{aligned} \quad (6)$$

where  $\chi_V$  can be obtained from dispersion relations using perturbative QCD and depends on the ratio of the  $d$ -quark mass over  $c$ -quark mass,  $u = m_d/m_c$  [44]. At the leading order, with  $u = 0$ ,

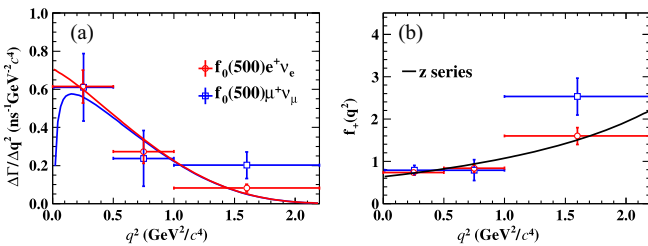


FIG. 5. (a) The partial decay rates of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  and (b) projections to  $f_+^{f_0}(q^2)$  with fit results overlaid, where the dots with error bars are data and the curves are the fit results.

$$\chi_V = \frac{3}{32\pi^2 m_c^2}. \quad (7)$$

Due to limited sample size in this analysis, we select the  $k = 1$  case, which gives

$$f_+(t) = \frac{1}{P(t)\Phi(t, t_0)} \frac{f_+(0)P(0)\Phi(0, t_0)}{1 + r_1(t_0)z(0, t_0)} (1 + r_1(t_0)z(t, t_0)). \quad (8)$$

A simultaneous fit is performed on the differential decay rates of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  and  $D^+ \rightarrow f_0(500)e^+\nu_e$ , in which the two decays are constrained to have the same parameters for the hadronic FF. The fit results and the projections on  $f_+^{f_0}(q^2)$  are shown in Fig. 5. The baseline fit parameters are taken from the fit performed with the combined statistical and systematic covariance matrix ( $\rho_{ij}^{\text{stat}}$  and  $\rho_{ij}^{\text{sys}}$ ), and their statistical uncertainties are taken from the fit only with the statistical covariance matrix. Tables V and VI summarize the statistical and systematic covariance matrices ( $\rho_{ij}^{\text{stat}}$  and  $\rho_{ij}^{\text{sys}}$ ) for the measured partial decay rates of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  in different  $q^2$  intervals, respectively. The upper-left  $3 \times 3$  matrix represents the result of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  channel, the lower-right  $3 \times 3$  matrix represents the result of  $D^+ \rightarrow f_0(500)e^+\nu_e$  channel, and the upper-right  $3 \times 3$  matrix represents the correlation between the two channels in different  $q^2$  intervals.

TABLE V. Statistical matrices for the measured partial decay rates of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  in different  $q^2$  intervals.

$\rho_{ij}^{\text{stat}}$	$D^+ \rightarrow f_0(500)\mu^+\nu_\mu$			$D^+ \rightarrow f_0(500)e^+\nu_e$		
	1	2	3	1	2	3
1	1.000	-0.121	-0.081	0.000	0.000	0.000
2		1.000	-0.094	0.000	0.000	0.000
3			1.000	0.000	0.000	0.000
1				1.000	-0.030	-0.000
2					1.000	-0.038
3						1.000

TABLE VI. Systematic density matrices of the measured partial decay rates of  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  in different  $q^2$  intervals.

$\rho_{ij}^{\text{syst}}$	$D^+ \rightarrow f_0(500)\mu^+\nu_\mu$			$D^+ \rightarrow f_0(500)e^+\nu_e$		
	1	2	3	1	2	3
1	1.000	0.985	0.679	0.002	0.002	0.003
2		1.000	0.611	0.002	0.002	0.002
3			1.000	0.004	0.004	0.005
1				1.000	0.994	0.879
2					1.000	0.880
3						1.000

For each parameter, the systematic uncertainty is obtained by calculating the quadratic difference of the uncertainties between these two fits. The systematic uncertainties related to  $\Delta\Gamma_{\text{msr}}^i$  is estimated to be 4.4% by the method in Ref. [45], which is the same as the measurement content of the systematic uncertainties of BF's as in Table II. The uncertainty associated with  $\Delta\Gamma_{\text{th}}^i$  is assigned as 3.0% by changing the parameters of Bugg resonant lineshape [33]. The total systematic uncertainty is taken to be the square root of the quadratic sum of these contributions.

### B. Results and discussion

From the fit, we determine  $f_+^{f_0}(0)|V_{cd}| = 0.143 \pm 0.014_{\text{stat}} \pm 0.011_{\text{syst}}$  for the first time, by analyzing simultaneously the partial decay rates of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  and  $D^+ \rightarrow f_0(500)e^+\nu_e$ . Using the value of  $|V_{cd}| = 0.22486 \pm 0.00067$  given in Ref. [1] yields  $f_+^{f_0}(0) = 0.63 \pm 0.06_{\text{stat}} \pm 0.05_{\text{syst}}$ . We obtain the parameter of hadronic FF  $r_1 = -5.8 \pm 2.4 \pm 0.5$  and the correlation coefficient of the two parameters is 0.83. The fit quality is  $\chi^2/\text{ndf} = 5.3/4$ , where  $\text{ndf}$  is the number of degrees of freedom. The obtained hadronic FF is consistent with the prediction given by QCD light-cone sum rules (LCSR) [14], covariant conned quark model (CCQM) [15] and chiral perturbation theory (chPT) [16], as shown in Fig. 6.

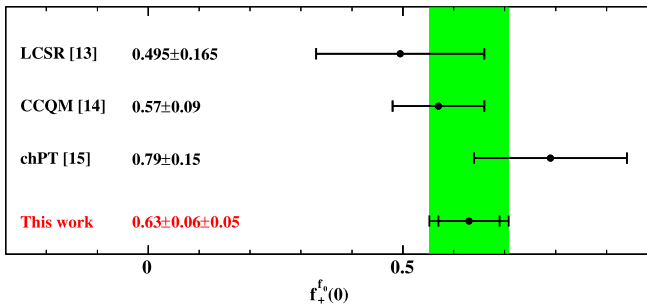


FIG. 6. The comparison of  $f_+^{f_0}(0)$  obtained in this work with theoretical calculations. The hadronic FF's are shown as dots with error bars.

The obtained results are helpful for deepening and understanding of the nature of the  $f_0(500)$  meson.

### VII. SUMMARY

In summary, by analyzing  $2.93 \text{ fb}^{-1}$  of  $e^+e^-$  collision data taken at  $\sqrt{s} = 3.773 \text{ GeV}$  with the BESIII detector, we report the first observation of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  with statistical significance of  $5.9\sigma$ . The absolute BF's of  $D^+ \rightarrow \rho^0\ell^+\nu_\ell$  and  $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$  are measured. As cross checks, we have also presented the BF's of  $D^+ \rightarrow f_0(500)e^+\nu_e$  with  $f_0(500) \rightarrow \pi^+\pi^-$  and  $D^+ \rightarrow \rho^0e^+\nu_e$  with  $\rho^0 \rightarrow \pi^+\pi^-$ , which are in agreement with the previous BESIII results [21]. From simultaneous fit to the partial decay rates of  $D^+ \rightarrow f_0(500)\mu^+\nu_\mu$  and  $D^+ \rightarrow f_0(500)e^+\nu_e$ , the product of  $f_+^{f_0}(0)|V_{cd}|$  is determined. Furthermore, taking the value of  $|V_{cd}|$  from a standard model fit [1] as input, the hadronic FF at zero momentum transfer squared  $f_+^{f_0}(0)$  is determined. The measured hadronic FF is important to test the various theoretical calculations [14–16].

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