## Erratum: Collinear factorization in wide-angle hadron pair production in $e^+e^-$ annihilation [Phys. Rev. D 100, 094014 (2019)]

E. Moffat, T.C. Rogers, N. Sato, and A. Signori<sup>®</sup>

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In our original publication, we presented analytical expressions for the short distance partonic cross sections for dihadron production in  $e^+e^-$  reaction in collinear factorization relevant for the large transverse momentum region of the exchange photon in the frame where the detected hadrons are back-to-back. After a new examination of the results, we found a mistake arising from the hadronic tensor decomposition. The full hadronic tensor is (e.g., [1])

$$W^{\mu\nu}(q, p_A, p_B) = \left(-g^{\mu\nu} + \frac{q^{\mu}q^{\nu}}{Q^2} - Z^{\mu}Z^{\nu}\right)W_T + Z^{\mu}Z^{\nu}W_L - (X^{\mu}Z^{\nu} + Z^{\mu}X^{\nu})W_{\Delta} + \left(-g^{\mu\nu} + \frac{q^{\mu}q^{\nu}}{Q^2} - 2X^{\mu}X^{\nu} - Z^{\mu}Z^{\nu}\right)W_{\Delta\Delta}.$$
(1)

In deriving the projection tensors, the paper dropped the contributions from  $W_{\Delta}$  and  $W_{\Delta\Delta}$ , which gives incorrect results. After implementing the necessary corrections, we have found that numerically the corrections are at most 3–4% at the very large  $q_{\rm T}$  and vanishes at small  $q_{\rm T}$ . Therefore, the phenomenological conclusions and the associated discussion in our manuscript are not significantly impacted by the mistake.

The relevant corrections are as follows:

- (i) Equation (12): This expression needs to be replaced by the above expression Eq. (1)
- (ii) In Eq. (14) the projector  $P_T^{\mu\nu}$  is now given by

$$P_T^{\mu\nu} = -\frac{1}{2} (g^{\mu\nu} + Z^{\mu} Z^{\nu}) \tag{2}$$

(iii) Equation (B2a) reads

$$\frac{\mathrm{d}\hat{\sigma}_{q\bar{q}}}{\mathrm{d}\hat{z}_{A}\mathrm{d}\hat{z}_{B}\mathrm{d}q_{\mathrm{T}}} = \frac{\mathrm{d}\hat{\sigma}_{\bar{q}q}}{\mathrm{d}\hat{z}_{A}\mathrm{d}\hat{z}_{B}\mathrm{d}q_{\mathrm{T}}} = F \frac{32(Q^{2} + q_{\mathrm{T}}^{2})^{2}(\hat{z}_{A}^{2} + \hat{z}_{B}^{2})}{(Q^{2}\hat{z}_{A} - Q^{2} + \hat{z}_{A}q_{\mathrm{T}}^{2})(Q^{2}\hat{z}_{B} - Q^{2} + \hat{z}_{B}q_{\mathrm{T}}^{2})}$$
(3)

(iv) Equation (B2b) reads

$$\frac{d\hat{\sigma}_{qg}}{d\hat{z}_{A}d\hat{z}_{B}dq_{T}} = \frac{d\hat{\sigma}_{\bar{q}g}}{d\hat{z}_{A}d\hat{z}_{B}dq_{T}} = F(-64Q^{4}\hat{z}_{A}^{2} - 64Q^{4}\hat{z}_{A}\hat{z}_{B} + 128Q^{4}\hat{z}_{A} - 32Q^{4}\hat{z}_{B}^{2} + 128Q^{4}\hat{z}_{B} - 128Q^{4} - 32q_{T}^{4}(2\hat{z}_{A}^{2} + 2\hat{z}_{A}\hat{z}_{B} + \hat{z}_{B}^{2}) \\
- 32q_{T}^{2}(4Q^{2}\hat{z}_{A}^{2} + 4Q^{2}\hat{z}_{A}\hat{z}_{B} - 4Q^{2}\hat{z}_{A} + 2Q^{2}\hat{z}_{B}^{2} - 4Q^{2}\hat{z}_{B})/(Q^{2}(\hat{z}_{A} - 1) + \hat{z}_{A}q_{T}^{2})((Q^{2} + q_{T}^{2})(\hat{z}_{A} + \hat{z}_{B}) - Q^{2})$$
(4)

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(v) Equation (B2c) reads

$$\frac{d\hat{\sigma}_{g\bar{q}}}{d\hat{z}_{A}d\hat{z}_{B}dq_{T}} = \frac{d\hat{\sigma}_{gq}}{d\hat{z}_{A}d\hat{z}_{B}dq_{T}} = F(-32Q^{4}\hat{z}_{A}^{2} - 64Q^{4}\hat{z}_{A}\hat{z}_{B} + 128Q^{4}\hat{z}_{A} - 64Q^{4}\hat{z}_{B}^{2} + 128Q^{4}\hat{z}_{B} - 128Q^{4} \\ - 32q_{T}^{4}(\hat{z}_{A}^{2} + 2\hat{z}_{A}\hat{z}_{B} + 2\hat{z}_{B}^{2}) - 32q_{T}^{2}(2Q^{2}\hat{z}_{A}^{2} + 4Q^{2}\hat{z}_{A}\hat{z}_{B} - 4Q^{2}\hat{z}_{A} + 4Q^{2}\hat{z}_{B}^{2} - 4Q^{2}\hat{z}_{B}))/ \\ (Q^{2}(\hat{z}_{B} - 1) + \hat{z}_{B}q_{T}^{2})((Q^{2} + q_{T}^{2})(\hat{z}_{A} + \hat{z}_{B}) - Q^{2}) \tag{5}$$

where

$$F = \delta_{+}(k_{C}^{2}) \frac{\alpha_{\rm em}^{2} e_{q}^{2} \alpha_{s} \hat{z}_{A} \hat{z}_{B} q_{\rm T} (Q^{2} + q_{\rm T}^{2})^{2}}{6Q^{6}}$$
(6)

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[1] J. C. Collins, Foundations of Perturbative QCD (Cambridge University Press, Cambridge, England, 2011).