

## Search for Scalar Leptoquarks Produced via $\tau$ -Lepton–Quark Scattering in $pp$ Collisions at $\sqrt{s} = 13$ TeV

A. Hayrapetyan *et al.*<sup>\*</sup>  
(CMS Collaboration)



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The first search for scalar leptoquarks produced in  $\tau$ -lepton–quark collisions is presented. It is based on a set of proton-proton collision data recorded with the CMS detector at the LHC at a center-of-mass energy of 13 TeV corresponding to an integrated luminosity of  $138 \text{ fb}^{-1}$ . The reconstructed final state consists of a jet, significant missing transverse momentum, and a  $\tau$  lepton reconstructed through its hadronic or leptonic decays. Limits are set on the product of the leptoquark production cross section and branching fraction and interpreted as exclusions in the plane of the leptoquark mass and the leptoquark– $\tau$ -quark coupling strength.

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Leptoquarks (LQs) are hypothetical color-triplet bosons carrying both baryon and lepton numbers and having fractional electric charge [1–6]. Mechanisms involving LQs coupled to third-generation quarks and leptons could explain the deviations from standard model predictions observed in several measurements of B meson decays, as explained in Refs. [7–10] and references therein.

At the CERN LHC, LQs can be produced singly through quark-gluon fusion, in pairs through gluon fusion or quark-antiquark annihilation, as part of the nonresonant production of two leptons in the  $t$ -channel, or from lepton-quark collisions. While the ATLAS and CMS Collaborations have performed searches for LQs targeting the first three production mechanisms [11–17], the last production mode has never been explored at the LHC. Recent theoretical progress in the determination of the lepton and photon density functions in the proton [18], based on the LUX approach [19,20], has shown that a significant LQ production cross section can be expected in proton-proton ( $pp$ ) collisions [21–24].

Lepton-induced production of LQs proceeds via the collision of a lepton and a quark, where the lepton is produced through a photon via quantum fluctuations in the proton. A leading-order Feynman diagram for this process is shown in Fig. 1. After the decay of the LQ, the final state consists of a high transverse momentum ( $p_T$ ) centrally produced lepton from the LQ decay, a same-flavor opposite-sign forward lepton from the photon decay, and a high- $p_T$  centrally produced jet from the LQ decay. The forward

lepton in this production mode usually has  $p_T$  below reconstruction threshold and differs from other production modes, which typically have two high- $p_T$  leptons. This Letter presents the first search for scalar LQs produced in  $\tau$ -lepton–quark collisions. For the strongest experimental sensitivity, we consider couplings to light-flavor ( $u$ ,  $d$ , and  $s$ ) and  $b$  quarks. The search is based on  $pp$  collision data at  $\sqrt{s} = 13$  TeV collected with the CMS detector in 2016–2018 corresponding to an integrated luminosity of  $138 \text{ fb}^{-1}$ . Three different decay modes of the  $\tau$  lepton are studied, resulting in the  $\tau_h + \text{jet}$ ,  $e + \text{jet}$ , and  $\mu + \text{jet}$  final states, where  $\tau_h$  denotes a  $\tau$  lepton decaying hadronically. Tabulated results are provided in the HEPData record for this analysis [25].

The CMS apparatus [26] is a multipurpose, nearly hermetic detector, designed to trigger on [27,28] and identify electrons, muons, photons, and both charged and neutral hadrons [29–31]. A global “particle-flow” (PF) algorithm [32] aims to reconstruct all individual particles in an event, combining information provided by the all-silicon inner tracker and by the crystal electromagnetic, and a brass and scintillator hadron calorimeters, operating inside a 3.8 T superconducting solenoid, with data from the gas-ionization muon detectors embedded in the flux-return yoke outside the solenoid. The reconstructed particles are used to build  $\tau$  leptons, jets, and missing transverse momentum ( $\vec{p}_T^{\text{miss}}$ ) [33–36].

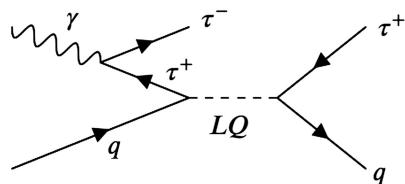


FIG. 1. Feynman diagram of the lepton-induced LQ production.

<sup>\*</sup>Full author list given at the end of the Letter.

The signal and some of the background contributions are estimated using simulations. Signal samples for scalar LQs in the  $s$ -channel single LQ production are generated at next-to-leading-order (NLO) with POWHEG interfaced with HERWIG [37] for the showering, using the models and LUXlep lepton parton distribution functions (PDFs) described in Refs. [18,21]. The LUXlep PDF set results from the combination of the NNPDF3.1luxQED set [38] and the lepton PDFs of Ref. [18], which are obtained with a framework similar to that used to determine the PDFs of the photon in the proton using electron-proton scattering data [19,20]. The signal samples are generated for  $b\tau$  ( $u\tau$ ) couplings, for LQ masses  $m_{\text{LQ}}$  between 0.6 and 2.0 (0.6 and 3.0) TeV, an absolute electric charge of  $2/3e$  ( $1/3e$ ), and a Yukawa coupling at the LQ-lepton–quark vertex,  $\lambda$ , between 0.5 and 3.0 (0.2 and 2.0). Charge conjugate pairs are simulated together. The analysis acceptance and efficiency, determined from simulation, are similar for couplings to any light-flavor quark, and the samples with LQ- $u$ - $\tau$  couplings are also used to extract results on LQ- $d$ - $\tau$  and LQ- $s$ - $\tau$  couplings. Results on LQ- $c$ - $\tau$  couplings cannot be extracted from the same simulation because events with such couplings have a different efficiency for the requirement on the identification of jets originating from  $b$  quarks, which is used to build analysis categories as described later.

The cross sections for all coupling hypotheses are computed at NLO [23,24]. The MadGraph5\_aMC@NLO 2.6.5 event generator [39] is used to generate events originating from  $Z + \text{jets}$  and  $W + \text{jets}$  processes. They are simulated at NLO with the FxFx jet matching and merging [40]. The vector boson  $p_T$  distribution is corrected to match a calculation at next-to-NLO (NNLO) for strong interactions and NLO with the next-to-leading logarithmic Sudakov approximation for electroweak interactions [41]. The MadGraph5\_aMC@NLO generator is also used for the simulation of diboson production, while POWHEG 2.0 [42–46] is used for  $t\bar{t}$  and single top quark production. The generators are interfaced with PYTHIA 8.240 [47] to model the parton showering, fragmentation, and hadronization, as well as the decay of the  $\tau$  leptons. The PYTHIA parameters affecting the description of the underlying event are set to the CP5 tune [48]. The NNPDF3.1 PDF set [49–51] at NNLO precision is used for background simulations. Additional  $pp$  interactions per bunch crossing are added to the simulated samples with the frequency distribution matching that observed in data. Generated events are processed through a GEANT4 [52] simulation of the CMS detector.

Electrons are reconstructed from energy deposits in the calorimeters and tracks in the tracking system, and identified with a cut-based discriminant [29]. Muons are reconstructed from tracks and hits in the tracker and muon systems [30,53]. Jets are clustered from PF candidates using the anti- $k_T$  FASTJET algorithm with a distance parameter  $R$  of 0.4 [54,55]. Their energy is corrected on an event-by-event basis [56]. Jets originating from  $b$  quarks

are identified with the medium working point of the DeepJet algorithm [57,58]. The decay products of  $\tau$  leptons decaying hadronically are reconstructed with the hadrons-plus-strips algorithm [33]. Quark and gluon jets, electrons, and muons misidentified as  $\tau_h$  candidates are reduced with deep neural network discriminants [36]. The tight working point is used to separate  $\tau_h$  candidates from jets; its efficiency is about 75% for  $\tau_h$  with  $p_T > 100$  GeV [36]. The loosest working point, used in the background estimation procedure, has an efficiency above 98%. The vector  $\vec{p}_T^{\text{miss}}$  is defined as the projection onto the plane perpendicular to the beam axis of the negative vector momentum sum of all reconstructed PF objects in an event. Its magnitude is referred to as  $p_T^{\text{miss}}$ .

The selection of events in the  $\tau_h + \text{jet}$  final state relies either on a single  $\tau_h$  trigger with online thresholds ranging from 140 to 180 GeV, depending on the data-taking year, or on a trigger requiring the scalar  $p_T$  sum of jets in the event with  $p_T > 40$  GeV to be above 900–1050 GeV. Events are selected in the  $e + \text{jet}$  final state using a single isolated electron trigger with online thresholds in the range of 27–35 GeV complemented with a photon trigger for  $p_T > 175$ –200 GeV. In the  $\mu + \text{jet}$  final state, single muon triggers with online thresholds between 24 and 50 GeV are used. Depending on whether the jet is tagged as originating from a  $b$  quark [57,58], the event is classified to be in a “btags” or “no-btags” category. A significant  $\vec{p}_T^{\text{miss}}$  satisfying  $|\Delta\phi(\ell, \vec{p}_T^{\text{miss}})| < 0.2$ –0.3 is required, where  $\ell$  stands for the visible decay products of the  $\tau$  lepton, as expected from the presence of one or more neutrinos in the  $\tau$  lepton decay. The off-line selection criteria for the lepton, leading jet, and  $\vec{p}_T^{\text{miss}}$  are presented in Table I. They differ slightly for final states with leptonically and hadronically decaying  $\tau$  leptons because of the different background contributions, triggers, and fractions of visible  $\tau$  lepton momentum. To select the LQ production mechanism targeted in this analysis, events are discarded if a second well-identified and isolated electron, muon, or  $\tau_h$  candidate with  $p_T > 50$  GeV and absolute pseudorapidity,  $|\eta|$ , less than 2.1 is found. This additional lepton veto removes about 5% of the signal

TABLE I. Selection criteria per final state.

Variable	$\tau_h + \text{jet}$	$e + \text{jet}$	$\mu + \text{jet}$
$p_T^\ell$ (GeV)	$> 200$	$> 100$	$> 100$
$ \eta^\ell $	$< 2.1$	$< 2.1$	$< 2.1$
$p_T^{\text{jet}}$ (GeV)	$> 300$	$> 200$	$> 200$
$ \eta^{\text{jet}} $	$< 2.4$	$< 2.4$	$< 2.4$
$p_T^{\text{miss}}$ (GeV)	$> 100$	$> 150$	$> 150$
$p_T(\vec{\ell} + \vec{p}_T^{\text{miss}})$ (GeV)	$> 100$	$> 100$	$> 100$
$ \Delta\phi(\ell, \vec{p}_T^{\text{miss}}) $ (radians)	$< 0.3$	$< 0.2$	$< 0.2$
$\Delta R(\ell, \text{jet})$	$> 0.5$	$> 0.5$	$> 0.5$

events in simulation. With this veto, other LQ production modes contribute less than 1% to the signal yield and are neglected.

As an estimation of  $m_{\text{LQ}}$  at reconstruction level, we calculate the collinear mass, defined as  $m_{\text{coll}} = m_{\text{vis}}(\tau, \text{jet})/\sqrt{x_{\text{vis}}}$ , where  $m_{\text{vis}}(\tau, \text{jet})$  is the invariant mass of the visible  $\tau$  decay products and the jet,  $x_{\text{vis}} = p_{\text{T}}^{\text{vis}}(\tau)/[p_{\text{T}}^{\text{vis}}(\tau) + p_{\text{T}}^{\text{invis}}(\tau)]$ , and  $p_{\text{T}}^{\text{invis}}(\tau)$  is the component of  $\vec{p}_{\text{T}}^{\text{miss}}$  in the direction of the visible  $\tau$  decay products. This calculation assumes that the neutrinos from  $\tau$  lepton decays are the only source of  $\vec{p}_{\text{T}}^{\text{miss}}$  and have the same  $\eta$  as the visible  $\tau$  decay products. The experimental resolution on  $m_{\text{coll}}$  is about 5%–10%, depending on the final state and  $m_{\text{LQ}}$ , and for  $\lambda \lesssim 1$  the natural width of the signal is negligible by comparison. In simulated events, the  $m_{\text{coll}}$  distribution peaks at the generated  $m_{\text{LQ}}$  values.

Background events with a prompt lepton in the final state ( $W + \text{jets}$ , Drell-Yan, diboson, single top quark, and  $t\bar{t}$  production events) are estimated from simulation and normalized to their theoretical cross sections [59–64], with the exception of the normalization of  $W + \text{jets}$  events in the btag categories. Because of the lack of a precise prediction for the cross section of  $W + \text{jets}$  in association with heavy-flavor jets, the latter is extracted from data in control regions (CRs) in the  $e + \text{jet}$  and  $\mu + \text{jet}$  final states with  $0.2 < |\Delta\phi(\ell, \vec{p}_{\text{T}}^{\text{miss}})| < 0.4$ , where  $\phi$  is the azimuthal angle in radians. The variable  $N_{\text{jets}}$ , defined as the number of jets with  $p_{\text{T}} > 30 \text{ GeV}$  and  $|\eta| < 2.4$ , and separated by  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} > 0.5$  from the selected lepton, is used as an observable to separate  $W + \text{jets}$  from  $t\bar{t}$  events, which on average have a higher number of jets. After the maximum likelihood fit including the signal regions (SRs), described later, the cross section for the  $W + \text{jets}$  background in the btag categories is measured to be  $1.04^{+0.07}_{-0.06}$  times the prediction from simulation, using the inclusive  $W + \text{jets}$  cross section at NNLO [59].

The backgrounds with a jet misidentified as a  $\tau_h$  candidate (an electron or muon), dominated by QCD multijet events, constitute about 50% (10%) of the total background and are estimated from data. The probability for jets selected with the loosest working point of the  $\tau_h$  discriminator to pass the tighter identification requirements in the signal region, called the misidentification factor (MF), is measured as a function of  $p_{\text{T}}(\tau_h)$ ,  $N_{\text{jets}}$ ,  $|\eta|$ , and jet  $p_{\text{T}}$ . It is in the range 0.10–0.25. This measurement is performed in a CR with a selection identical to that in the signal region, except that the  $p_{\text{T}}^{\text{miss}}$  requirement is inverted. A multiplicative correction for the extrapolation to higher  $p_{\text{T}}^{\text{miss}}$  values in the SR is applied. The MF for electrons is measured as a function of the electron  $p_{\text{T}}$  because the trigger requirements are  $p_{\text{T}}$ -dependent, while the muon MF is measured as a function of  $p_{\text{T}}^{\text{miss}}$ , since the relative fraction of QCD multijet and  $W + \text{jets}$  backgrounds varies with  $p_{\text{T}}^{\text{miss}}$ . The observed  $p_{\text{T}}^{\text{miss}}$  linear dependence is extrapolated

to the SR for  $150 < p_{\text{T}}^{\text{miss}} < 300 \text{ GeV}$  and is assumed constant above 300 GeV. Events that pass the SR selection with the exception of the  $\tau_h$ , electron, or muon identification, depending on the final state, are reweighted using the MFs to estimate the background with misidentified jets. Contributions from prompt  $\tau_h$  candidates, electrons, or muons in this region are estimated from simulation and subtracted to avoid double counting events containing genuine leptons.

After the application of the selection criteria shown in Table I, boosted decision trees (BDTs) are trained with the TMVA package [65] for each final state to improve the separation between the signal and background. The input variables, chosen to have a limited dependence on  $m_{\text{LQ}}$ , are the following:  $|\Delta\phi|$  between pairs of analysis objects (electrons, muons,  $\tau_h$ , jets,  $\vec{p}_{\text{T}}^{\text{miss}}$ ), the ratio of the objects'  $p_{\text{T}}$  to  $m_{\text{coll}}$  or to each other's  $p_{\text{T}}$ , the  $\Delta R$  separation between the jet and the reconstructed  $\tau$  lepton, and  $N_{\text{jets}}$ . The BDTs are trained with a mixture of all signal samples with various  $m_{\text{LQ}}$  and  $\lambda$ , against background events coming from  $W + \text{jets}$  production and either processes with misidentified jets in the  $\tau_h + \text{jet}$  channel, or  $t\bar{t}$  events in the  $e + \text{jet}$  and  $\mu + \text{jet}$  channels. The BDT training is performed once for each final state, and is used in the btag and no-btag categories, for all data-taking years and  $m_{\text{LQ}}$  hypotheses. The BDT output distribution is verified to be well predicted in a validation region where the requirement on  $|\Delta\phi(\ell, \vec{p}_{\text{T}}^{\text{miss}})|$  is inverted.

Events in the no-btag (btag) category are further split into 4 (3) subcategories on the basis of their BDT output. In the btag subcategories, events with  $N_{\text{jets}} > 2$  are vetoed to reduce the contribution from the  $t\bar{t}$  background. The discriminating observable is  $m_{\text{coll}}$  in all subcategories and final states.

Uncertainties in the reconstruction, identification, and isolation of  $\tau_h$  candidates (electrons, muons) are determined via the “tag-and-probe” method [66] and can be as high as 15% (2%, 2%). The  $b$  tagging uncertainties for heavy-flavor jets and mistagging uncertainties for light-flavor quark and gluon jets are included with partial correlations between the data-taking years. They are in the range 6%–9% in the btag category for LQs coupled to  $b$  quarks. The uncertainty in the efficiency of the  $\tau_h$  (electron, muon) trigger ranges up to 11% (2%, 2%). The uncertainty in the energy scale of  $\tau_h$  candidates is below the percent level for all  $\tau_h$  decay modes, and is propagated to the  $m_{\text{coll}}$  distributions. The impact of uncertainties in the energy scale and resolution of jets, and in the measurement of  $\vec{p}_{\text{T}}^{\text{miss}}$ , is smaller by comparison. The uncertainties in the energy scale of electrons and muons are negligible as compared with the uncertainties mentioned above.

The uncertainties in the  $t\bar{t}$ , diboson, and single top quark production cross sections are 5.2%, 2.5%, and 3.7%, respectively [61–64]. The uncertainties in the NNLO cross sections of the  $Z + \text{jets}$  and  $W + \text{jets}$  backgrounds are 2%

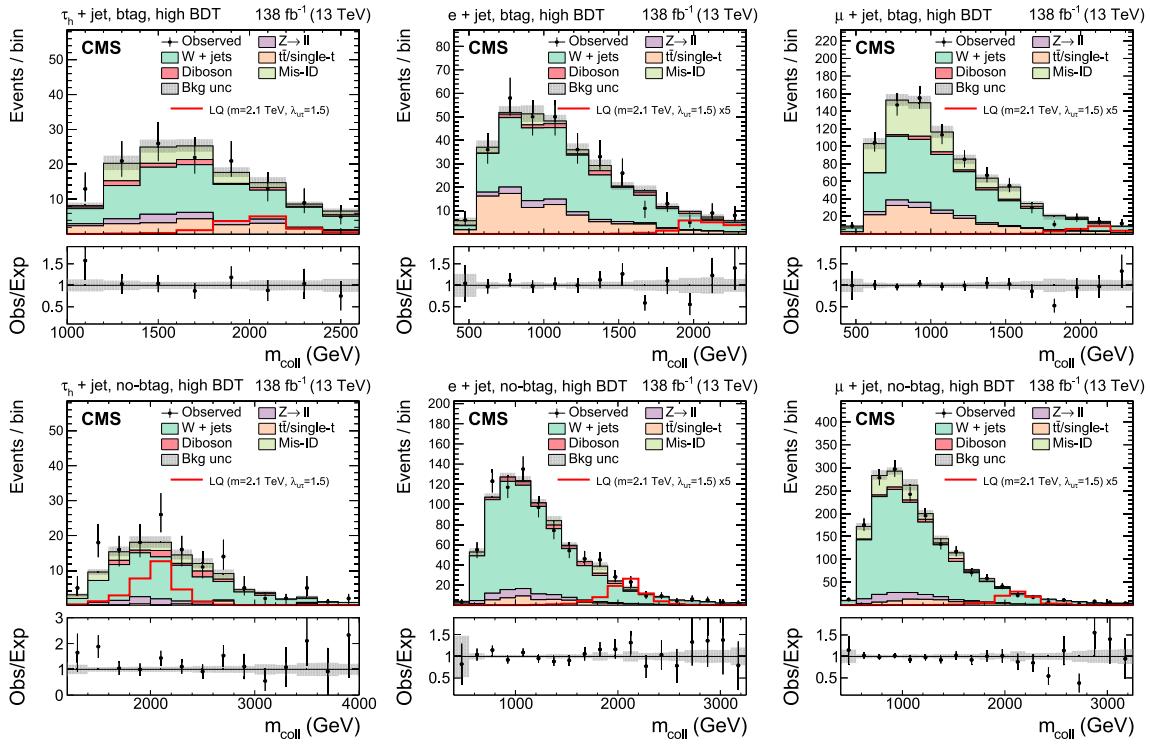


FIG. 2. Observed and expected distributions of  $m_{\text{coll}}$  in the  $\tau_h + \text{jet}$  (left),  $e + \text{jet}$  (center), and  $\mu + \text{jet}$  (right) channels for the btag (upper) and no-btag (lower) subcategories with the BDT requirements selecting the most signal-like events. The bands include statistical and systematic uncertainties. The background distributions are the results of the maximum likelihood fit. For clarity, the signal distributions in the  $e + \text{jet}$  and  $\mu + \text{jet}$  final states are multiplied by a factor of 5.

and 3%, respectively [59,60]. In the btag category, the  $Z + \text{jets}$  background normalization uncertainty is increased to 20% on the basis of the agreement observed in a CR with 2  $\tau_h$  candidates and 1  $b$  tagged jet, while the normalization of the  $W + \text{jets}$  background is left floating in the fit, resulting in an uncertainty of 7% as detailed above. The acceptance uncertainties in the renormalization and factorization scales, PDFs, and parton showering for these simulated processes are also included. For the signal, uncertainties in the renormalization and factorization scales, and in the PDFs are included. They affect the  $m_{\text{coll}}$  distributions and have a normalization effect in the ranges 3%–4% (4%–6%) (scales) and 2%–3% (2%–4%) (PDFs), respectively, for LQs with  $ut$  ( $bt$ ) couplings, in the mass ranges considered in the analysis [23]. The uncertainty in the integrated luminosity is 1.6% [67–69].

Different uncertainties affecting the  $m_{\text{coll}}$  distributions are considered for the background with jets misidentified as  $\tau_h$  candidates: 10% uncertainty per detector region (barrel or end caps), 10% uncertainty per  $N_{\text{jets}}$  bin, 30% uncertainty for events with  $p_T(\tau_h) > 600$  GeV, 10% uncertainty in the leading jet  $p_T$  correction, and the  $p_T^{\text{miss}}$  correction uncertainty described earlier. The uncertainty in the MFs for electrons and muons originates from the limited number

of events in the measurement and from the choice of observables and selection criteria. It results in a 20% normalization uncertainty for backgrounds with misidentified electrons or muons, uncorrelated between the data-taking years. In the  $\mu + \text{jet}$  channel, an uncertainty in the  $p_T^{\text{miss}}$  dependence is included.

A maximum likelihood fit is performed with the signal normalization as a free parameter, and the systematic uncertainties described above as nuisance parameters. The  $m_{\text{coll}}$  distributions in the different subcategories and final states are fitted simultaneously, together with the  $N_{\text{jets}}$  distributions in the CRs that control the normalization of the background from  $W + \text{jets}$  with a  $b$  tagged jet. Those with the highest BDT output requirement, selecting the most signal-like events, are shown in Fig. 2. The btag subcategory with the highest BDT output requirement contains about 70% (41%, 56%) of the LQ events with  $m_{\text{LQ}} = 2$  TeV and  $\lambda_{bt} = 1.0$  entering the btag category, and about 9.8% (5.6%, 12%) of the total background in the  $\tau_h + \text{jet}$  ( $e + \text{jet}$ ,  $\mu + \text{jet}$ ) final state. The no-btag subcategory with the highest BDT output requirement contains about 55% (30%, 33%) of the LQ events with  $m_{\text{LQ}} = 3$  TeV and  $\lambda_{ut} = 1.0$  entering the no-btag category, and about 0.86% (1.1%, 1.9%) of the total background in the no-btag category in the  $\tau_h + \text{jet}$  ( $e + \text{jet}$ ,  $\mu + \text{jet}$ ) final state.

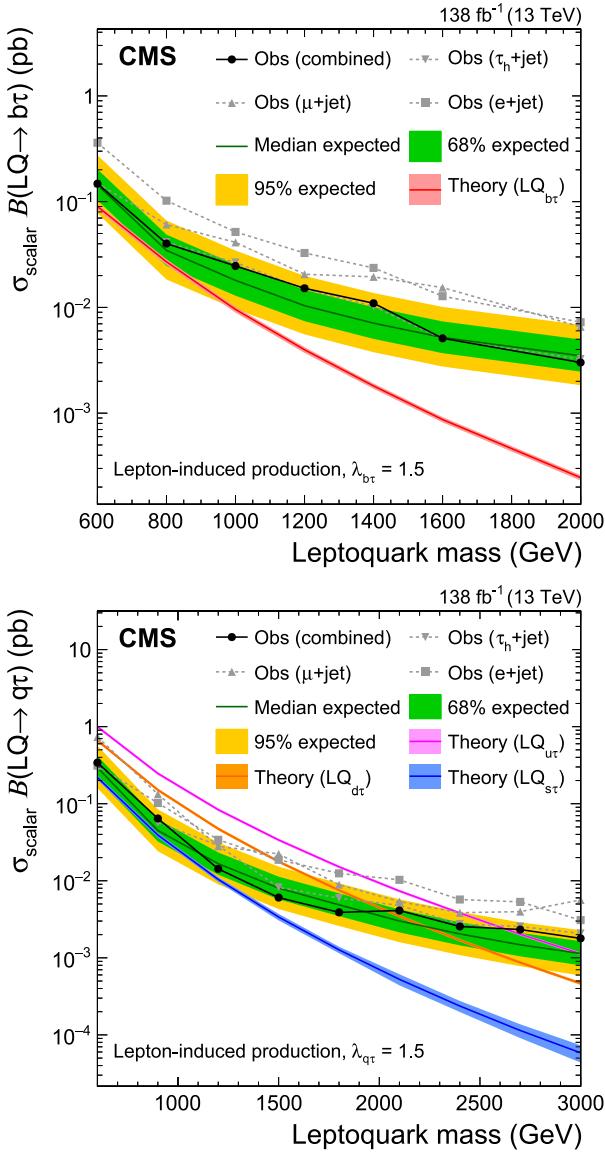


FIG. 3. Expected and observed upper limits at 95% C.L. on the product of the scalar lepton-induced LQ production cross section and the branching fraction for a LQ coupled to  $b$  quarks and  $\tau$  leptons (left), or to light-flavor quarks and  $\tau$  leptons (right), using  $\lambda = 1.5$ . The theoretical cross sections correspond to the calculations of Refs. [23,24]. The inner (green) band and the outer (yellow) band indicate the regions containing 68% and 95%, respectively, of the distribution of limits expected under the background-only hypothesis. The filled circles show the observed limits for the combination of final states, while the other markers indicate the observed results per final state.

No statistically significant excess above the standard model backgrounds is observed. Upper limits at 95% confidence level (C.L.) are set on the product of the LQ production cross section and the branching fraction for

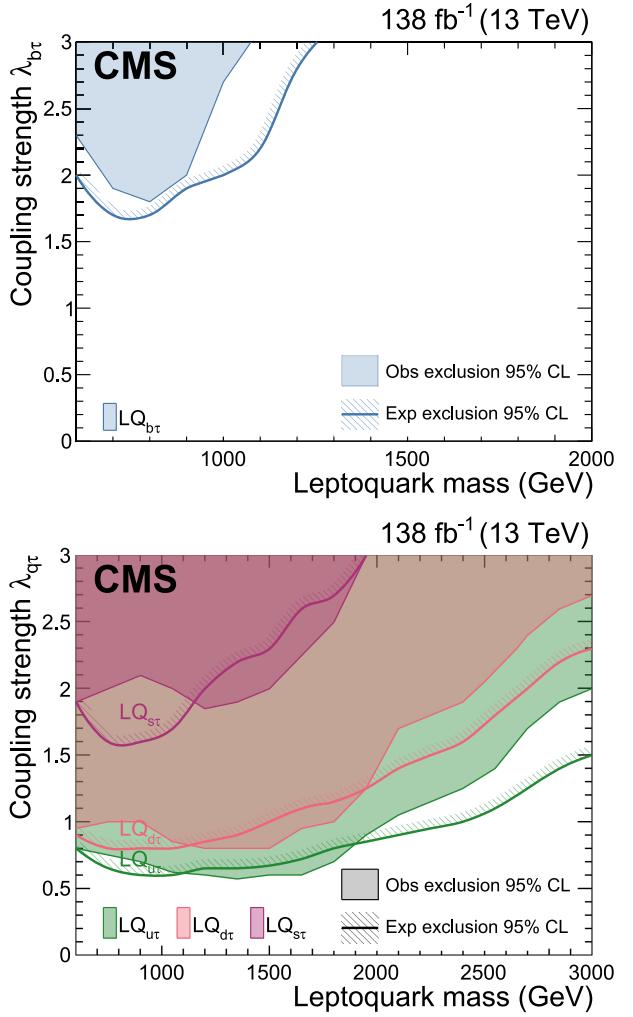


FIG. 4. Upper limit at 95% C.L. on the coupling strength  $\lambda$  of a scalar LQ to  $b$  quarks and  $\tau$  leptons (upper), and to light-flavor quarks and  $\tau$  leptons (lower). Regions above the hatched lines correspond to the expected exclusions.

different coupling hypotheses, using the C.L.<sub>s</sub> method [70,71] in the asymptotic approximation [72]. Because of the trigger requirements, the signal acceptance in the  $\tau_h + \text{jet}$  final state is low for  $m_{\text{LQ}} = 600$  GeV and the limits for this mass point are derived from the distributions in the  $e + \text{jet}$  and  $\mu + \text{jet}$  final states only. The 95% C.L. limits are in the range 0.34–0.0018 pb (0.15–0.0030 pb) for masses between 0.6 and 3.0 TeV (0.6 and 2.0 TeV) for  $\lambda_{ut} = 1.5$  ( $\lambda_{bt} = 1.5$ ), as shown in Fig. 3. The limits are translated into exclusion regions in the  $m_{\text{LQ}}-\lambda$  plane, as shown in Fig. 4, assuming the branching fraction of the LQs to a quark and a  $\tau$  lepton to be 100% for the considered quark flavor. The observed limits on LQs coupling to  $b$  quarks and  $\tau$  leptons extend existing constraints from searches in other production modes at high  $m_{\text{LQ}}$  [12]. Leptoquarks

coupling to light-flavor quarks and  $\tau$  leptons can be excluded at masses above the previously existing limit of 1.3 TeV [17], for  $\lambda_{ut}$  ( $\lambda_{dt}$ ,  $\lambda_{st}$ ) above 0.6 (0.8, 1.9).

In summary, a search for leptoquarks produced in lepton-quark collisions and coupled to  $\tau$  leptons has been performed for the first time, using data collected with the CMS detector in 2016–2018. These limits are complementary to those set using other production modes at high mass and coupling values for  $b\tau$  couplings, while the limits on the couplings of leptoquarks to light-flavor quarks extend the mass range excluded by previous searches in other production modes.

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Myllymäki<sup>12</sup>, M. m. Rantanen<sup>12</sup>, H. Siikonen<sup>12</sup>, E. Tuominen<sup>12</sup>, J. Tuominen<sup>12</sup>, P. Luukka<sup>12</sup>, H. Petrow<sup>12</sup>, T. Tuuva,<sup>36,a</sup> M. Besancon<sup>12</sup>, F. Couderc<sup>12</sup>, M. Dejardin<sup>12</sup>, D. Denegri,<sup>37</sup> J. L. Faure,<sup>37</sup> F. Ferri<sup>12</sup>, S. Ganjour<sup>12</sup>, P. Gras<sup>12</sup>, G. Hamel de Monchenault<sup>12</sup>, V. Lohezic<sup>12</sup>, J. Malcles<sup>12</sup>, J. Rander,<sup>37</sup> A. Rosowsky<sup>12</sup>, M. Ö. Sahin<sup>12</sup>, A. Savoy-Navarro<sup>12</sup>, P. Simkina<sup>12</sup>, M. Titov<sup>12</sup>, C. Baldenegro Barrera<sup>12</sup>, F. Beaudette<sup>12</sup>, A. Buchot Perraguin<sup>12</sup>, P. Busson<sup>12</sup>, A. Cappati<sup>12</sup>, C. Charlote<sup>12</sup>, F. Damas<sup>12</sup>, O. Davignon<sup>12</sup>, A. De Wit<sup>12</sup>, G. Falmagne<sup>12</sup>, B. A. Fontana Santos Alves<sup>12</sup>, S. Ghosh<sup>12</sup>, A. Gilbert<sup>12</sup>, R. Granier de Cassagnac<sup>12</sup>, A. Hakimi<sup>12</sup>, B. Harikrishnan<sup>12</sup>, L. Kalipoliti<sup>12</sup>, G. Liu<sup>12</sup>, J. Motta<sup>12</sup>, M. Nguyen<sup>12</sup>, C. Ochando<sup>12</sup>, L. Portales<sup>12</sup>, R. Salerno<sup>12</sup>, U. Sarkar<sup>12</sup>, J. B. Sauvan<sup>12</sup>, Y. Sirois<sup>12</sup>, A. Tarabini<sup>12</sup>, E. Vernazza<sup>12</sup>, A. Zabi<sup>12</sup>, A. Zghiche<sup>12</sup>, J.-L. Agram<sup>12</sup>, J. Andrea<sup>12</sup>, D. Apparau<sup>12</sup>, D. Bloch<sup>12</sup>, J.-M. Brom<sup>12</sup>, E. C. Chabert<sup>12</sup>, C. Collard<sup>12</sup>, S. Falke<sup>12</sup>, U. Goerlach<sup>12</sup>, C. Grimault,<sup>39</sup> R. Haeberle,<sup>39</sup> A.-C. Le Bihan<sup>12</sup>, M. A. Sessini<sup>12</sup>, P. Van Hove<sup>12</sup>, S. Beauceron<sup>12</sup>, B. Blancon<sup>12</sup>, G. Boudoul<sup>12</sup>, N. Chanon<sup>12</sup>, J. Choi<sup>12</sup>, D. Contardo<sup>12</sup>, P. Depasse<sup>12</sup>, C. Dozen<sup>12</sup>, H. El Mamouni,<sup>40</sup> J. Fay<sup>12</sup>, S. Gascon<sup>12</sup>, M. Gouzevitch<sup>12</sup>, C. Greenberg,<sup>40</sup> G. Grenier<sup>12</sup>, B. Ille<sup>12</sup>, I. B. Laktineh,<sup>40</sup> M. Lethuillier<sup>12</sup>, L. Mirabito,<sup>40</sup> S. Perries,<sup>40</sup> A. Purohit<sup>12</sup>, M. Vander Donckt<sup>12</sup>, P. Verdier<sup>12</sup>, J. Xiao<sup>12</sup>, A. Khvedelidze<sup>12</sup>, I. Lomidze<sup>12</sup>, Z. Tsamalaidze<sup>12</sup>, V. Botta<sup>12</sup>, L. Feld<sup>12</sup>, K. Klein<sup>12</sup>, M. Lipinski<sup>12</sup>, D. Meuser<sup>12</sup>, A. Pauls<sup>12</sup>, N. Röwert<sup>12</sup>, M. Teroerde<sup>12</sup>, S. Diekmann<sup>12</sup>, A. Dodonova<sup>12</sup>, N. Eich<sup>12</sup>, D. Eliseev<sup>12</sup>, F. Engelke<sup>12</sup>, M. Erdmann<sup>12</sup>, P. Fackeldey<sup>12</sup>, B. Fischer<sup>12</sup>, T. Hebbeker<sup>12</sup>, K. Hoepfner<sup>12</sup>, F. Ivone<sup>12</sup>, A. Jung<sup>12</sup>, M. y. Lee<sup>12</sup>, L. Mastrolorenzo,<sup>43</sup> M. Merschmeyer<sup>12</sup>, A. Meyer<sup>12</sup>, S. Mukherjee<sup>12</sup>, D. Noll<sup>12</sup>, A. Novak<sup>12</sup>, F. Nowotny,<sup>43</sup> A. Pozdnyakov<sup>12</sup>, Y. Rath,<sup>43</sup> W. Redjeb<sup>12</sup>, F. Rehm,<sup>43</sup> H. Reithler<sup>12</sup>, V. Sarkisovi<sup>12</sup>, A. Schmidt<sup>12</sup>, S. C. Schuler,<sup>43</sup> A. Sharma<sup>12</sup>, A. Stein<sup>12</sup>, F. Torres Da Silva De Araujo<sup>12</sup>, L. Vigilante,<sup>43</sup> S. Wiedenbeck<sup>12</sup>, S. Zaleski,<sup>43</sup> C. Dziwok<sup>12</sup>, G. Flügge<sup>12</sup>, W. Haj Ahmad<sup>12</sup>, T. Kress<sup>12</sup>, A. Nowack<sup>12</sup>, O. Pooth<sup>12</sup>, A. Stahl<sup>12</sup>, T. Ziemons<sup>12</sup>, A. Zotz<sup>12</sup>, H. Aarup Petersen<sup>12</sup>, M. Aldaya Martin<sup>12</sup>, J. Alimena<sup>12</sup>, S. Amoroso,<sup>45</sup> Y. An<sup>12</sup>, S. Baxter<sup>12</sup>, M. Bayatmakou<sup>12</sup>, H. Becerril Gonzalez<sup>12</sup>, O. Behnke<sup>12</sup>, A. Belvedere<sup>12</sup>, S. Bhattacharya<sup>12</sup>, F. Blekman<sup>12</sup>, K. Borras<sup>12</sup>, D. Brunner<sup>12</sup>, A. Campbell<sup>12</sup>, A. Cardini<sup>12</sup>, C. Cheng,<sup>45</sup> F. Colombina<sup>12</sup>, S. Consuegra Rodríguez<sup>12</sup>, G. Correia Silva<sup>12</sup>, M. De Silva<sup>12</sup>, G. Eckerlin,<sup>45</sup> D. Eckstein<sup>12</sup>, L. I. Estevez Banos<sup>12</sup>, O. Filatov<sup>12</sup>, E. Gallo<sup>12</sup>, A. Geiser<sup>12</sup>, A. Giraldi<sup>12</sup>, G. Greau,<sup>45</sup> V. Guglielmi<sup>12</sup>, M. Guthoff<sup>12</sup>, A. Hinzmann<sup>12</sup>, A. Jafari<sup>12</sup>, L. Jeppe<sup>12</sup>, N. Z. Jomhari<sup>12</sup>, B. Kaech<sup>12</sup>, M. Kasemann<sup>12</sup>, H. Kaveh<sup>12</sup>, C. Kleinwort<sup>12</sup>, R. Kogler<sup>12</sup>, M. Komm<sup>12</sup>, D. Krücker<sup>12</sup>, W. Lange,<sup>45</sup> D. Leyva Pernia<sup>12</sup>, K. Lipka<sup>12</sup>, W. Lohmann<sup>12</sup>, R. Mankel<sup>12</sup>, I.-A. Melzer-Pellmann<sup>12</sup>, M. Mendizabal Morentin<sup>12</sup>, J. Metwally,<sup>45</sup> A. B. Meyer<sup>12</sup>, G. Milella<sup>12</sup>, A. Mussgiller<sup>12</sup>, A. Nürnberg<sup>12</sup>, Y. Otarid,<sup>45</sup> D. Pérez Adán<sup>12</sup>, E. Ranken<sup>12</sup>, A. Raspereza<sup>12</sup>, B. Ribeiro Lopes<sup>12</sup>, J. Rübenach,<sup>45</sup> A. Saggio<sup>12</sup>, M. Scham<sup>12</sup>, S. Schnake<sup>12</sup>, P. Schütze<sup>12</sup>, C. Schwanenberger<sup>12</sup>, D. Selivanova<sup>12</sup>, M. Shchedrolosiev<sup>12</sup>, R. E. Sosa Ricardo<sup>12</sup>, L. P. Sreelatha Pramod<sup>12</sup>, D. Stafford,<sup>45</sup> F. Vazzoler<sup>12</sup>, A. Ventura Barroso<sup>12</sup>, R. Walsh<sup>12</sup>, Q. Wang<sup>12</sup>, Y. Wen<sup>12</sup>, K. Wichmann,<sup>45</sup> L. Wiens<sup>12</sup>, C. Wissing<sup>12</sup>, S. Wuchterl<sup>12</sup>, Y. Yang<sup>12</sup>, A. Zimmermann Castro Santos<sup>12</sup>, A. Albrecht<sup>12</sup>, S. Albrecht<sup>12</sup>, M. Antonello<sup>12</sup>, S. Bein<sup>12</sup>, L. Benato<sup>12</sup>, M. Bonanomi<sup>12</sup>, P. Connor<sup>12</sup>, M. Eich,<sup>46</sup> K. El Morabit<sup>12</sup>, Y. Fischer<sup>12</sup>, A. Fröhlich,<sup>46</sup> C. Garbers<sup>12</sup>, E. Garutti<sup>12</sup>, A. Grohsjean<sup>12</sup>, M. Hajheidari,<sup>46</sup> J. Haller<sup>12</sup>, H. R. Jabusch<sup>12</sup>, G. Kasieczka<sup>12</sup>, P. Keicher,<sup>46</sup> R. Klanner<sup>12</sup>, W. Korcari<sup>12</sup>, T. Kramer<sup>12</sup>, V. Kutzner<sup>12</sup>, F. Labe<sup>12</sup>, J. Lange<sup>12</sup>, A. Lobanov<sup>12</sup>, C. Matthies<sup>12</sup>, A. Mehta<sup>12</sup>, L. Moureaux<sup>12</sup>, M. Mrowietz,<sup>46</sup> A. Nigamova<sup>12</sup>, Y. Nissan,<sup>46</sup> A. Paasch<sup>12</sup>, K. J. Pena Rodriguez<sup>12</sup>, T. Quadfasel<sup>12</sup>, B. Raciti<sup>12</sup>, M. Rieger<sup>12</sup>, D. Savoie<sup>12</sup>, J. Schindler<sup>12</sup>, P. Schleper<sup>12</sup>, M. Schröder<sup>12</sup>, J. Schwandt<sup>12</sup>, M. Sommerhalder<sup>12</sup>, H. Stadie<sup>12</sup>, G. Steinbrück<sup>12</sup>, A. Tews,<sup>46</sup> M. Wolf<sup>12</sup>, S. Brommer<sup>12</sup>, M. Burkart,<sup>47</sup> E. Butz<sup>12</sup>, T. Chwalek<sup>12</sup>, A. Dierlamm<sup>12</sup>, A. Droll,<sup>47</sup> N. Faltermann<sup>12</sup>, M. Giffels<sup>12</sup>, A. Gottmann<sup>12</sup>, F. Hartmann<sup>12</sup>, R. Hofsaess<sup>12</sup>, M. Horzela<sup>12</sup>, U. Husemann<sup>12</sup>, M. Klute<sup>12</sup>, R. Koppenhöfer<sup>12</sup>, M. Link,<sup>47</sup> A. Lintuluoto<sup>12</sup>, S. Maier<sup>12</sup>, S. Mitra<sup>12</sup>, M. Mormile<sup>12</sup>, Th. Müller<sup>12</sup>, M. Neukum,<sup>47</sup> M. Oh<sup>12</sup>, G. Quast<sup>12</sup>, K. Rabbertz<sup>12</sup>, B. Regnery<sup>12</sup>, N. Shadskiy<sup>12</sup>, I. Shvetsov<sup>12</sup>, H. J. Simonis<sup>12</sup>, N. Trevisani<sup>12</sup>,

- R. Ulrich<sup>47</sup> J. van der Linden<sup>47</sup> R. F. Von Cube<sup>47</sup> M. Wassmer<sup>47</sup> S. Wieland<sup>47</sup> F. Wittig<sup>47</sup> R. Wolf<sup>47</sup>  
 S. Wunsch<sup>47</sup> X. Zuo<sup>47</sup> G. Anagnostou<sup>48</sup> P. Assiouras<sup>48</sup> G. Daskalakis<sup>48</sup> A. Kyriakis<sup>48</sup> A. Papadopoulos,<sup>48,ff</sup>  
 A. Stakia<sup>48</sup> P. Kontaxakis<sup>49</sup> G. Melachroinos<sup>49</sup> A. Panagiotou<sup>49</sup> I. Papavergou<sup>49</sup> I. Paraskevas<sup>49</sup> N. Saoulidou<sup>49</sup>  
 K. Theofilatos<sup>49</sup> E. Tziaferi<sup>49</sup> K. Vellidis<sup>49</sup> I. Zisopoulos<sup>49</sup> G. Bakas<sup>50</sup> T. Chatzistavrou<sup>50</sup> G. Karapostoli<sup>50</sup>  
 K. Kousouris<sup>50</sup> I. Papakrivopoulos<sup>50</sup> E. Siamarkou<sup>50</sup> G. Tsipolitis<sup>50</sup> A. Zacharopoulou<sup>50</sup> K. Adamidis<sup>51</sup>  
 I. Bestintzanos<sup>51</sup> I. Evangelou<sup>51</sup> C. Foudas<sup>51</sup> P. Gianneios<sup>51</sup> C. Kamtsikis<sup>51</sup> P. Katsoulis<sup>51</sup> P. Kokkas<sup>51</sup>  
 P. G. Kosmoglou Kioseoglou<sup>51</sup> N. Manthos<sup>51</sup> I. Papadopoulos<sup>51</sup> J. Strologas<sup>51</sup> M. Csand<sup>52</sup> K. Farkas<sup>52</sup>  
 M. M. A. Gadallah<sup>52,gg</sup> . Kadlecik<sup>52</sup> P. Major<sup>52</sup> K. Mandal<sup>52</sup> G. Psztor<sup>52</sup> A. J. Radl<sup>52,hh</sup> G. I. Veres<sup>52</sup>  
 M. Bartok<sup>53,ii</sup> C. Hajdu<sup>53</sup> D. Horvath<sup>53,jj,kk</sup> F. Sikler<sup>53</sup> V. Veszpremi<sup>53</sup> P. Raics<sup>54</sup> B. Ujvari<sup>54,ll</sup> G. Zilizi<sup>54</sup>  
 G. Bencze<sup>55</sup> S. Czellar<sup>55</sup> J. Karancsi<sup>55,ii</sup> J. Molnar<sup>55</sup> Z. Szillasi<sup>55</sup> T. Csorgo<sup>56,hh</sup> F. Nemes<sup>56,hh</sup> T. Novak<sup>56</sup>  
 J. Babbar<sup>57</sup> S. Bansal<sup>57</sup> S. B. Beri<sup>57</sup> V. Bhatnagar<sup>57</sup> G. Chaudhary<sup>57</sup> S. Chauhan<sup>57</sup> N. Dhingra<sup>57,mm</sup>  
 R. Gupta<sup>57</sup> A. Kaur<sup>57</sup> A. Kaur<sup>57</sup> H. Kaur<sup>57</sup> M. Kaur<sup>57</sup> S. Kumar<sup>57</sup> M. Meena<sup>57</sup> K. Sandeep<sup>57</sup>  
 T. Sheokand<sup>57</sup> J. B. Singh<sup>57</sup> A. Singla<sup>57</sup> A. Ahmed<sup>58</sup> A. Bhardwaj<sup>58</sup> A. Chhetri<sup>58</sup> B. C. Choudhary<sup>58</sup>  
 A. Kumar<sup>58</sup> M. Naimuddin<sup>58</sup> K. Ranjan<sup>58</sup> S. Saumya<sup>58</sup> S. Acharya<sup>59,nn</sup> S. Baradia<sup>59</sup> S. Barman<sup>59,oo</sup>  
 S. Bhattacharya<sup>59</sup> D. Bhowmik<sup>59</sup> S. Dutta<sup>59</sup> S. Dutta<sup>59</sup> B. Gomber<sup>59,nn</sup> P. Palit<sup>59</sup> G. Saha<sup>59</sup> B. Sahu<sup>59,nn</sup>  
 S. Sarkar<sup>59</sup> M. M. Ameen<sup>60</sup> P. K. Behera<sup>60</sup> S. C. Behera<sup>60</sup> S. Chatterjee<sup>60</sup> P. Jana<sup>60</sup> P. Kalbhor<sup>60</sup>  
 J. R. Komaragiri<sup>60,pp</sup> D. Kumar<sup>60,pp</sup> L. Panwar<sup>60,pp</sup> R. Pradhan<sup>60</sup> P. R. Pujahari<sup>60</sup> N. R. Saha<sup>60</sup> A. Sharma<sup>60</sup>  
 A. K. Sikdar<sup>60</sup> S. Verma<sup>60</sup> T. Aziz<sup>61</sup> I. Das<sup>61</sup> S. Dugad<sup>61</sup> M. Kumar<sup>61</sup> G. B. Mohanty<sup>61</sup> P. Suryadevara,<sup>61</sup>  
 A. Bala<sup>62</sup> S. Banerjee<sup>62</sup> R. M. Chatterjee<sup>62</sup> M. Guchait<sup>62</sup> Sh. Jain<sup>62</sup> S. Karmakar<sup>62</sup> S. Kumar<sup>62</sup>  
 G. Majumder<sup>62</sup> K. Mazumdar<sup>62</sup> S. Mukherjee<sup>62</sup> S. Parolia<sup>62</sup> A. Thachayath<sup>62</sup> S. Bahinipati<sup>63,qq</sup> A. K. Das,<sup>63</sup>  
 C. Kar<sup>63</sup> D. Maity<sup>63,rr</sup> P. Mal<sup>63</sup> T. Mishra<sup>63</sup> V. K. Muraleedharan Nair Bindhu<sup>63,rr</sup> K. Naskar<sup>63,rr</sup>  
 A. Nayak<sup>63,rr</sup> P. Sadangi<sup>63</sup> P. Saha<sup>63</sup> S. K. Swain<sup>63</sup> S. Varghese<sup>63,rr</sup> D. Vats<sup>63,rr</sup> A. Alpana<sup>64</sup> S. Dube<sup>64</sup>  
 B. Kansal<sup>64</sup> A. Laha<sup>64</sup> A. Rastogi<sup>64</sup> S. Sharma<sup>64</sup> H. Bakhshiansohi<sup>65,ss</sup> E. Khazaie<sup>65,tt</sup> M. Zeinali<sup>65,uu</sup>  
 S. Chenarani<sup>66,vv</sup> S. M. Etesami<sup>66</sup> M. Khakzad<sup>66</sup> M. Mohammadi Najafabadi<sup>66</sup> M. Grunewald<sup>67</sup>  
 M. Abbrescia<sup>68a,68b</sup> R. Aly<sup>68a,68c,ww</sup> A. Colaleo<sup>68a,68b</sup> D. Creanza<sup>68a,68c</sup> B. D' Anzi<sup>68a,68b</sup> N. De Filippis<sup>68a,68c</sup>  
 M. De Palma<sup>68a,68b</sup> A. Di Florio<sup>68a,68c</sup> W. Elmetenawee<sup>68a,68b</sup> L. Fiore<sup>68a</sup> G. Iaselli<sup>68a,68c</sup> G. Maggi<sup>68a,68c</sup>  
 M. Maggi<sup>68a</sup> I. Margjeka<sup>68a,68b</sup> V. Mastrapasqua<sup>68a,68b</sup> S. My<sup>68a,68b</sup> S. Nuzzo<sup>68a,68b</sup> A. Pellecchia<sup>68a,68b</sup>  
 A. Pompili<sup>68a,68b</sup> G. Pugliese<sup>68a,68c</sup> R. Radogna<sup>68a</sup> G. Ramirez-Sanchez<sup>68a,68c</sup> D. Ramos<sup>68a</sup> A. Ranieri<sup>68a</sup>  
 L. Silvestris<sup>68a</sup> F. M. Simone<sup>68a,68b</sup> . Szbilir<sup>68a</sup> A. Stamerra<sup>68a</sup> R. Venditti<sup>68a</sup> P. Verwilligen<sup>68a</sup>  
 A. Zaza<sup>68a,68b</sup> G. Abbiendi<sup>69a</sup> C. Battilana<sup>69a,69b</sup> L. Borgonovi<sup>69a</sup> R. Campanini<sup>69a,69b</sup> P. Capiluppi<sup>69a,69b</sup>  
 A. Castro<sup>69a,69b</sup> F. R. Cavallo<sup>69a</sup> M. Cuffiani<sup>69a,69b</sup> G. M. Dallavalle<sup>69a</sup> T. Diotalevi<sup>69a,69b</sup> F. Fabbri<sup>69a</sup>  
 A. Fanfani<sup>69a,69b</sup> D. Fasanella<sup>69a,69b</sup> P. Giacomelli<sup>69a</sup> L. Giommi<sup>69a,69b</sup> C. Grandi<sup>69a</sup> L. Guiducci<sup>69a,69b</sup>  
 S. Lo Meo<sup>69a,xx</sup> L. Lunerti<sup>69a,69b</sup> S. Marcellini<sup>69a</sup> G. Masetti<sup>69a</sup> F. L. Navarria<sup>69a,69b</sup> A. Perrotta<sup>69a</sup>  
 F. Primavera<sup>69a,69b</sup> A. M. Rossi<sup>69a,69b</sup> T. Rovelli<sup>69a,69b</sup> G. P. Siroli<sup>69a,69b</sup> S. Costa<sup>70a,70b,yy</sup> A. Di Mattia<sup>70a</sup>  
 R. Potenza<sup>70a,70b</sup> A. Tricomi<sup>70a,70b,yy</sup> C. Tuve<sup>70a,70b</sup> G. Barbagli<sup>71a</sup> G. Bardelli<sup>71a,71b</sup> B. Camaiani<sup>71a,71b</sup>  
 A. Cassese<sup>71a</sup> R. Ceccarelli<sup>71a</sup> V. Ciulli<sup>71a,71b</sup> C. Civinini<sup>71a</sup> R. D'Alessandro<sup>71a,71b</sup> E. Focardi<sup>71a,71b</sup>  
 T. Kello<sup>71a</sup> G. Latino<sup>71a,71b</sup> P. Lenzi<sup>71a,71b</sup> M. Lizzo<sup>71a</sup> M. Meschini<sup>71a</sup> S. Paoletti<sup>71a</sup> A. Papanastassiou,<sup>71a,71b</sup>  
 G. Sguazzoni<sup>71a</sup> L. Viliani<sup>71a</sup> L. Benussi<sup>72</sup> S. Bianco<sup>72</sup> S. Meola<sup>72,zz</sup> D. Piccolo<sup>72</sup> P. Chatagnon<sup>73a</sup>  
 F. Ferro<sup>73a</sup> E. Robutti<sup>73a</sup> S. Tosi<sup>73a,73b</sup> A. Benaglia<sup>74a</sup> G. Boldrini<sup>74a,74b</sup> F. Brivio<sup>74a</sup> F. Cetorelli<sup>74a</sup>  
 F. De Guio<sup>74a,74b</sup> M. E. Dinardo<sup>74a,74b</sup> P. Dini<sup>74a</sup> S. Gennai<sup>74a</sup> R. Gerosa<sup>74a,74b</sup> A. Ghezzi<sup>74a,74b</sup>  
 P. Govoni<sup>74a,74b</sup> L. Guzzi<sup>74a</sup> M. T. Lucchini<sup>74a,74b</sup> M. Malberti<sup>74a</sup> S. Malvezzi<sup>74a</sup> A. Massironi<sup>74a</sup>  
 D. Menasce<sup>74a</sup> L. Moroni<sup>74a</sup> M. Paganoni<sup>74a,74b</sup> D. Pedrini<sup>74a</sup> B. S. Pinolini<sup>74a</sup> S. Ragazzi<sup>74a,74b</sup>  
 T. Tabarelli de Fatis<sup>74a,74b</sup> D. Zuolo<sup>74a</sup> S. Buontempo<sup>75a</sup> A. Cagnotta<sup>75a,75b</sup> F. Carnevali,<sup>75a,75b</sup> N. Cavallo<sup>75a,75c</sup>  
 A. De Iorio<sup>75a,75b</sup> F. Fabozzi<sup>75a,75c</sup> A. O. M. Iorio<sup>75a,75b</sup> L. Lista<sup>75a,75b,aaa</sup> P. Paolucci<sup>75a,ff</sup> B. Rossi<sup>75a</sup>  
 C. Sciacca<sup>75a,75b</sup> R. Ardino<sup>76a</sup> P. Azz<sup>76a</sup> N. Bacchetta<sup>76a,bbb</sup> D. Bisello<sup>76a,76b</sup> P. Bortignon<sup>76a</sup>  
 A. Bragagnolo<sup>76a,76b</sup> R. Carlin<sup>76a,76b</sup> P. Checchia<sup>76a</sup> T. Dorigo<sup>76a</sup> S. Fantinel<sup>76a</sup> F. Fanzago<sup>76a</sup>  
 U. Gasparini<sup>76a,76b</sup> G. Grossi<sup>76a,ccc</sup> L. Layer<sup>76a,76b</sup> E. Lusiani<sup>76a</sup> M. Margoni<sup>76a,76b</sup> M. Migliorini<sup>76a,76b</sup>  
 J. Pazzini<sup>76a,76b</sup> P. Ronchese<sup>76a,76b</sup> R. Rossin<sup>76a,76b</sup> F. Simonetto<sup>76a,76b</sup> G. Strong<sup>76a</sup> M. Tosi<sup>76a,76b</sup>  
 A. Triossi<sup>76a,76b</sup> S. Ventura<sup>76a</sup> H. Yasar<sup>76a,76b</sup> M. Zanetti<sup>76a,76b</sup> P. Zotto<sup>76a,76b</sup> A. Zucchetta<sup>76a,76b</sup>

- G. Zumerle<sup>76a,76b</sup> S. Abu Zeid<sup>77a,ddd</sup> C. Aimè<sup>77a,77b</sup> A. Braghieri<sup>77a</sup> S. Calzaferri<sup>77a,77b</sup> D. Fiorina<sup>77a,77b</sup>  
 P. Montagna<sup>77a,77b</sup> V. Re<sup>77a</sup> C. Riccardi<sup>77a,77b</sup> P. Salvini<sup>77a</sup> I. Vai<sup>77a,77b</sup> P. Vitulo<sup>77a,77b</sup> S. Ajmal<sup>78a,78b</sup>  
 P. Asenov<sup>78a,eee</sup> G. M. Bilei<sup>78a</sup> D. Ciangottini<sup>78a,78b</sup> L. Fanò<sup>78a,78b</sup> M. Magherini<sup>78a,78b</sup> G. Mantovani,<sup>78a,78b</sup>  
 V. Mariani<sup>78a,78b</sup> M. Menichelli<sup>78a</sup> F. Moscatelli<sup>78a,eee</sup> A. Piccinelli<sup>78a,78b</sup> M. Presilla<sup>78a,78b</sup> A. Rossi<sup>78a,78b</sup>  
 A. Santocchia<sup>78a,78b</sup> D. Spiga<sup>78a</sup> T. Tedeschi<sup>78a,78b</sup> P. Azzurri<sup>79a</sup> G. Bagliesi<sup>79a</sup> R. Bhattacharya<sup>79a</sup>  
 L. Bianchini<sup>79a,79b</sup> T. Boccali<sup>79a</sup> E. Bossini<sup>79a</sup> D. Bruschini<sup>79a,79c</sup> R. Castaldi<sup>79a</sup> M. A. Ciocci<sup>79a,79b</sup>  
 M. Cipriani<sup>79a,79b</sup> V. D'Amante<sup>79a,79d</sup> R. Dell'Orso<sup>79a</sup> S. Donato<sup>79a</sup> A. Giassi<sup>79a</sup> F. Ligabue<sup>79a,79c</sup>  
 D. Matos Figueiredo<sup>79a</sup> A. Messineo<sup>79a,79b</sup> M. Musich<sup>79a,79b</sup> F. Palla<sup>79a</sup> A. Rizzi<sup>79a,79b</sup> G. Rolandi<sup>79a,79c</sup>  
 S. Roy Chowdhury<sup>79a</sup> T. Sarkar<sup>79a</sup> A. Scribano<sup>79a</sup> P. Spagnolo<sup>79a</sup> R. Tenchini<sup>79a,79b</sup> G. Tonelli<sup>79a,79b</sup>  
 N. Turini<sup>79a,79d</sup> A. Venturi<sup>79a</sup> P. G. Verdini<sup>79a</sup> P. Barria<sup>80a</sup> M. Campana<sup>80a,80b</sup> F. Cavallari<sup>80a</sup>  
 L. Cunqueiro Mendez<sup>80a,80b</sup> D. Del Re<sup>80a,80b</sup> E. Di Marco<sup>80a</sup> M. Diemoz<sup>80a</sup> F. Errico<sup>80a,80b</sup> E. Longo<sup>80a,80b</sup>  
 P. Meridiani<sup>80a</sup> J. Mijuskovic<sup>80a,80b</sup> G. Organtini<sup>80a,80b</sup> F. Pandolfi<sup>80a</sup> R. Paramatti<sup>80a,80b</sup> C. Quaranta<sup>80a,80b</sup>  
 S. Rahatlou<sup>80a,80b</sup> C. Rovelli<sup>80a</sup> F. Santanastasio<sup>80a,80b</sup> L. Soffi<sup>80a</sup> N. Amapane<sup>81a,81b</sup> R. Arcidiacono<sup>81a,81c</sup>  
 S. Argiro<sup>81a,81b</sup> M. Arneodo<sup>81a,81c</sup> N. Bartosik<sup>81a</sup> R. Bellan<sup>81a,81b</sup> A. Bellora<sup>81a,81b</sup> C. Biino<sup>81a</sup> N. Cartiglia<sup>81a</sup>  
 M. Costa<sup>81a,81b</sup> R. Covarelli<sup>81a,81b</sup> N. Demaria<sup>81a</sup> L. Finco<sup>81a</sup> M. Grippo<sup>81a,81b</sup> B. Kiani<sup>81a,81b</sup> F. Legger<sup>81a</sup>  
 F. Luongo<sup>81a,81b</sup> C. Mariotti<sup>81a</sup> S. Maselli<sup>81a</sup> A. Mecca<sup>81a,81b</sup> E. Migliore<sup>81a,81b</sup> M. Monteno<sup>81a</sup>  
 R. Mulargia<sup>81a</sup> M. M. Obertino<sup>81a,81b</sup> G. Ortona<sup>81a</sup> L. Pacher<sup>81a,81b</sup> N. Pastrone<sup>81a</sup> M. Pelliccioni<sup>81a</sup>  
 M. Ruspa<sup>81a,81c</sup> F. Siviero<sup>81a,81b</sup> V. Sola<sup>81a,81b</sup> A. Solano<sup>81a,81b</sup> D. Soldi<sup>81a,81b</sup> A. Staiano<sup>81a</sup> C. Tarricone<sup>81a,81b</sup>  
 M. Tornago<sup>81a,81b</sup> D. Trocino<sup>81a</sup> G. Umoret<sup>81a,81b</sup> E. Vlasov<sup>81a,81b</sup> S. Belforte<sup>82a</sup> V. Candelise<sup>82a,82b</sup>  
 M. Casarsa<sup>82a</sup> F. Cossutti<sup>82a</sup> K. De Leo<sup>82a,82b</sup> G. Della Ricca<sup>82a,82b</sup> S. Dogra<sup>83</sup> J. Hong<sup>83</sup> C. Huh<sup>83</sup>  
 B. Kim<sup>83</sup> D. H. Kim<sup>83</sup> J. Kim<sup>83</sup> H. Lee<sup>83</sup> S. W. Lee<sup>83</sup> C. S. Moon<sup>83</sup> Y. D. Oh<sup>83</sup> M. S. Ryu<sup>83</sup> S. Sekmen<sup>83</sup>  
 Y. C. Yang<sup>83</sup> G. Bak<sup>84</sup> P. Gwak<sup>84</sup> H. Kim<sup>84</sup> D. H. Moon<sup>84</sup> E. Asilar<sup>85</sup> D. Kim<sup>85</sup> T. J. Kim<sup>85</sup> J. A. Merlin,<sup>85</sup>  
 J. Park<sup>85</sup> S. Choi<sup>86</sup> S. Han<sup>86</sup> B. Hong<sup>86</sup> K. Lee<sup>86</sup> K. S. Lee<sup>86</sup> S. Lee<sup>86</sup> J. Park<sup>86</sup> S. K. Park<sup>86</sup> J. Yoo<sup>86</sup>  
 J. Goh<sup>87</sup> H. S. Kim<sup>88</sup> Y. Kim<sup>88</sup> S. Lee<sup>88</sup> J. Almond<sup>89</sup> J. H. Bhyun<sup>89</sup> J. Choi<sup>89</sup> W. Jun<sup>89</sup> J. Kim<sup>89</sup> J. S. Kim,<sup>89</sup>  
 S. Ko<sup>89</sup> H. Kwon<sup>89</sup> H. Lee<sup>89</sup> J. Lee<sup>89</sup> B. H. Oh<sup>89</sup> S. B. Oh<sup>89</sup> H. Seo<sup>89</sup> U. K. Yang<sup>89</sup> I. Yoon<sup>89</sup>  
 W. Jang<sup>90</sup> D. Y. Kang<sup>90</sup> Y. Kang<sup>90</sup> S. Kim<sup>90</sup> B. Ko<sup>90</sup> J. S. H. Lee<sup>90</sup> Y. Lee<sup>90</sup> I. C. Park<sup>90</sup> Y. Roh,<sup>90</sup>  
 I. J. Watson<sup>90</sup> S. Yang<sup>90</sup> S. Ha<sup>91</sup> H. D. Yoo<sup>91</sup> M. Choi<sup>92</sup> M. R. Kim<sup>92</sup> H. Lee<sup>92</sup> Y. Lee<sup>92</sup> I. Yu<sup>92</sup>  
 T. Beyrouthy<sup>93</sup> Y. Maghrbi<sup>93</sup> K. Dreimanis<sup>94</sup> A. Gaile<sup>94</sup> G. Pikurs<sup>94</sup> A. Potrebko<sup>94</sup> M. Seidel<sup>94</sup>  
 V. Veckalns<sup>94,fff</sup> N. R. Strautnieks<sup>95</sup> M. Ambrozas<sup>96</sup> A. Juodagalvis<sup>96</sup> A. Rinkevicius<sup>96</sup> G. Tamulaitis<sup>96</sup>  
 N. Bin Norjoharuddeen<sup>97</sup> I. Yusuff<sup>97,egg</sup> Z. Zolkapli<sup>97</sup> J. F. Benitez<sup>98</sup> A. Castaneda Hernandez<sup>98</sup>  
 H. A. Encinas Acosta<sup>98</sup> L. G. Gallegos Maríñez<sup>98</sup> M. León Coello<sup>98</sup> J. A. Murillo Quijada<sup>98</sup> A. Sehrawat<sup>98</sup>  
 L. Valencia Palomo<sup>98</sup> G. Ayala<sup>99</sup> H. Castilla-Valdez<sup>99</sup> E. De La Cruz-Burelo<sup>99</sup> I. Heredia-De La Cruz<sup>99,hhh</sup>  
 R. Lopez-Fernandez<sup>99</sup> C. A. Mondragon Herrera<sup>99</sup> A. Sánchez Hernández<sup>99</sup> C. Oropeza Barrera<sup>100</sup>  
 M. Ramírez García<sup>100</sup> I. Bautista<sup>101</sup> I. Pedraza<sup>101</sup> H. A. Salazar Ibarguen<sup>101</sup> C. Uribe Estrada<sup>101</sup> I. Bubanja,<sup>102</sup>  
 N. Raicevic<sup>102</sup> P. H. Butler<sup>103</sup> A. Ahmad<sup>104</sup> M. I. Asghar<sup>104</sup> A. Awais<sup>104</sup> M. I. M. Awan,<sup>104</sup> H. R. Hoorani<sup>104</sup>  
 W. A. Khan<sup>104</sup> V. Avati<sup>105</sup> L. Grzanka<sup>105</sup> M. Malawski<sup>105</sup> H. Bialkowska<sup>106</sup> M. Bluj<sup>106</sup> B. Boimska<sup>106</sup>  
 M. Górski<sup>106</sup> M. Kazana<sup>106</sup> M. Szleper<sup>106</sup> P. Zalewski<sup>106</sup> K. Bunkowski<sup>107</sup> K. Doroba<sup>107</sup> A. Kalinowski<sup>107</sup>  
 M. Konecki<sup>107</sup> J. Krolikowski<sup>107</sup> A. Muhammad<sup>107</sup> M. Araujo<sup>108</sup> D. Bastos<sup>108</sup> C. Beirão Da Cruz E Silva<sup>108</sup>  
 A. Boletti<sup>108</sup> M. Bozzo<sup>108</sup> P. Faccioli<sup>108</sup> M. Gallinaro<sup>108</sup> J. Hollar<sup>108</sup> N. Leonardo<sup>108</sup> T. Niknejad<sup>108</sup>  
 A. Petrilli<sup>108</sup> M. Pisano<sup>108</sup> J. Seixas<sup>108</sup> J. Varela<sup>108</sup> J. W. Wulff<sup>108</sup> P. Adzic<sup>109</sup> P. Milenovic<sup>109</sup>  
 M. Dordevic<sup>110</sup> J. Milosevic<sup>110</sup> V. Rekovic<sup>110</sup> M. Aguilar-Benitez<sup>111</sup> J. Alcaraz Maestre<sup>111</sup> Cristina F. Bedoya<sup>111</sup>  
 M. Cepeda<sup>111</sup> M. Cerrada<sup>111</sup> N. Colino<sup>111</sup> B. De La Cruz<sup>111</sup> A. Delgado Peris<sup>111</sup> D. Fernández Del Val<sup>111</sup>  
 J. P. Fernández Ramos<sup>111</sup> J. Flix<sup>111</sup> M. C. Fouz<sup>111</sup> O. Gonzalez Lopez<sup>111</sup> S. Goy Lopez<sup>111</sup> J. M. Hernandez<sup>111</sup>  
 M. I. Josa<sup>111</sup> J. León Holgado<sup>111</sup> D. Moran<sup>111</sup> C. M. Morcillo Perez<sup>111</sup> Á. Navarro Tobar<sup>111</sup>  
 C. Perez Dengra<sup>111</sup> A. Pérez-Calero Yzquierdo<sup>111</sup> J. Puerta Pelayo<sup>111</sup> I. Redondo<sup>111</sup> D. D. Redondo Ferrero<sup>111</sup>  
 L. Romero,<sup>111</sup> S. Sánchez Navas<sup>111</sup> L. Urda Gómez<sup>111</sup> J. Vazquez Escobar<sup>111</sup> C. Willmott,<sup>111</sup> J. F. de Trocóniz<sup>112</sup>  
 B. Alvarez Gonzalez<sup>113</sup> J. Cuevas<sup>113</sup> J. Fernandez Menendez<sup>113</sup> S. Folgueras<sup>113</sup> I. Gonzalez Caballero<sup>113</sup>  
 J. R. González Fernández<sup>113</sup> E. Palencia Cortezon<sup>113</sup> C. Ramón Álvarez<sup>113</sup> V. Rodríguez Bouza<sup>113</sup>  
 A. Soto Rodríguez<sup>113</sup> A. Trapote<sup>113</sup> C. Vico Villalba<sup>113</sup> P. Vischia<sup>113</sup> S. Bhowmik<sup>114</sup> S. Blanco Fernández<sup>114</sup>

- J. A. Brochero Cifuentes<sup>114</sup> I. J. Cabrillo<sup>114</sup> A. Calderon<sup>114</sup> J. Duarte Campderros<sup>114</sup> M. Fernandez<sup>114</sup>  
C. Fernandez Madrazo<sup>114</sup> G. Gomez<sup>114</sup> C. Lasosa Garcia<sup>114</sup> C. Martinez Rivero<sup>114</sup>
- P. Martinez Ruiz del Arbol<sup>114</sup> F. Matorras<sup>114</sup> P. Matorras Cuevas<sup>114</sup> E. Navarrete Ramos<sup>114</sup> J. Piedra Gomez<sup>114</sup>  
L. Scodellaro<sup>114</sup> I. Vila<sup>114</sup> J. M. Vizan Garcia<sup>114</sup> M. K. Jayananda<sup>115</sup> B. Kailasapathy<sup>115,iii</sup>
- D. U. J. Sonnadara<sup>115</sup> D. D. C. Wickramarathna<sup>115</sup> W. G. D. Dharmaratna<sup>116</sup> K. Liyanage<sup>116</sup> N. Perera<sup>116</sup>  
N. Wickramage<sup>116</sup> D. Abbaneo<sup>117</sup> C. Amendola<sup>117</sup> E. Auffray<sup>117</sup> G. Auizinger<sup>117</sup> J. Baechler<sup>117</sup> D. Barney<sup>117</sup>  
A. Bermudez Martínez<sup>117</sup> M. Bianco<sup>117</sup> B. Bilin<sup>117</sup> A. A. Bin Anuar<sup>117</sup> A. Bocci<sup>117</sup> E. Brondolin<sup>117</sup>  
C. Caillol<sup>117</sup> T. Camporesi<sup>117</sup> G. Cerminara<sup>117</sup> N. Chernyavskaya<sup>117</sup> D. d'Enterria<sup>117</sup> A. Dabrowski<sup>117</sup>  
A. David<sup>117</sup> A. De Roeck<sup>117</sup> M. M. Defranchis<sup>117</sup> M. Deile<sup>117</sup> M. Dobson<sup>117</sup> F. Fallavollita,<sup>117,iii</sup>  
L. Forthomme<sup>117</sup> G. Franzoni<sup>117</sup> W. Funk<sup>117</sup> S. Giani<sup>117</sup> D. Gigi<sup>117</sup> K. Gill<sup>117</sup> F. Glege<sup>117</sup> L. Gouskos<sup>117</sup>  
M. Haranko<sup>117</sup> J. Hegeman<sup>117</sup> B. Huber<sup>117</sup> V. Innocente<sup>117</sup> T. James<sup>117</sup> P. Janot<sup>117</sup> J. Kieseler<sup>117</sup>  
S. Laurila<sup>117</sup> P. Lecoq<sup>117</sup> E. Leutgeb<sup>117</sup> C. Lourenço<sup>117</sup> B. Maier<sup>117</sup> L. Malgeri<sup>117</sup> M. Mannelli<sup>117</sup>  
A. C. Marini<sup>117</sup> M. Matthewman<sup>117</sup> F. Meijers<sup>117</sup> S. Mersi<sup>117</sup> E. Meschi<sup>117</sup> V. Milosevic<sup>117</sup> F. Moortgat<sup>117</sup>  
M. Mulders<sup>117</sup> S. Orfanelli<sup>117</sup> F. Pantaleo<sup>117</sup> M. Peruzzi<sup>117</sup> G. Petrussciani<sup>117</sup> A. Pfeiffer<sup>117</sup> M. Pierini<sup>117</sup>  
D. Piparo<sup>117</sup> H. Qu<sup>117</sup> D. Rabady<sup>117</sup> G. Reales Gutierrez,<sup>117</sup> M. Rovere<sup>117</sup> H. Sakulin<sup>117</sup> S. Scarfi<sup>117</sup>  
M. Selvaggi<sup>117</sup> A. Sharma<sup>117</sup> K. Shchelina<sup>117</sup> P. Silva<sup>117</sup> P. Spicas<sup>117,kkk</sup> A. G. Stahl Leiton<sup>117</sup> A. Steen<sup>117</sup>  
S. Summers<sup>117</sup> D. Treille<sup>117</sup> P. Tropea<sup>117</sup> A. Tsirou,<sup>117</sup> D. Walter<sup>117</sup> J. Wanczyk<sup>117,III</sup> K. A. Wozniak<sup>117,mmm</sup>  
P. Zehetner<sup>117</sup> P. Zejdl<sup>117</sup> W. D. Zeuner,<sup>117</sup> T. Bevilacqua<sup>118,nnn</sup> L. Caminada<sup>118,nnn</sup> A. Ebrahimi<sup>118</sup>  
W. Erdmann<sup>118</sup> R. Horisberger<sup>118</sup> Q. Ingram<sup>118</sup> H. C. Kaestli<sup>118</sup> D. Kotlinski<sup>118</sup> C. Lange<sup>118</sup>  
M. Missiroli<sup>118,nnn</sup> L. Noehte<sup>118,nnn</sup> T. Rohe<sup>118</sup> T. K. Arrestad<sup>119</sup> K. Androsov<sup>119,III</sup> M. Backhaus<sup>119</sup>  
A. Calandri<sup>119</sup> C. Cazzaniga<sup>119</sup> K. Datta<sup>119</sup> A. De Cosa<sup>119</sup> G. Dissertori<sup>119</sup> M. Dittmar,<sup>119</sup> M. Donegà<sup>119</sup>  
F. Eble<sup>119</sup> M. Galli<sup>119</sup> K. Gedea<sup>119</sup> F. Glessgen<sup>119</sup> C. Grab<sup>119</sup> D. Hits<sup>119</sup> W. Lustermann<sup>119</sup> A.-M. Lyon<sup>119</sup>  
R. A. Manzoni<sup>119</sup> M. Marchegiani<sup>119</sup> L. Marchese<sup>119</sup> C. Martin Perez<sup>119</sup> A. Mascellani<sup>119,III</sup>  
F. Nessi-Tedaldi<sup>119</sup> F. Pauss<sup>119</sup> V. Perovic<sup>119</sup> S. Pigazzini<sup>119</sup> M. G. Ratti<sup>119</sup> M. Reichmann<sup>119</sup> C. Reissel<sup>119</sup>  
T. Reitenspiess<sup>119</sup> B. Ristic<sup>119</sup> F. Riti<sup>119</sup> D. Ruini,<sup>119</sup> D. A. Sanz Becerra<sup>119</sup> R. Seidita<sup>119</sup> J. Steggemann<sup>119,III</sup>  
D. Valsecchi<sup>119</sup> R. Wallny<sup>119</sup> C. Amsler<sup>120,ooo</sup> P. Bärtschi<sup>120</sup> C. Botta<sup>120</sup> D. Brzhechko,<sup>120</sup> M. F. Canelli<sup>120</sup>  
K. Cormier<sup>120</sup> R. Del Burgo,<sup>120</sup> J. K. Heikkilä<sup>120</sup> M. Huwiler<sup>120</sup> W. Jin<sup>120</sup> A. Jofrehei<sup>120</sup> B. Kilminster<sup>120</sup>  
S. Leontsinis<sup>120</sup> S. P. Liechti<sup>120</sup> A. Macchiolo<sup>120</sup> P. Meiring<sup>120</sup> V. M. Mikuni<sup>120</sup> U. Molinatti<sup>120</sup>  
I. Neutelings<sup>120</sup> A. Reimers<sup>120</sup> P. Robmann,<sup>120</sup> S. Sanchez Cruz<sup>120</sup> K. Schweiger<sup>120</sup> M. Senger<sup>120</sup>  
Y. Takahashi<sup>120</sup> R. Tramontano<sup>120</sup> C. Adloff,<sup>121,ppp</sup> C. M. Kuo,<sup>121</sup> W. Lin,<sup>121</sup> P. K. Rout<sup>121</sup> P. C. Tiwari<sup>121,pp</sup>  
S. S. Yu<sup>121</sup> L. Ceard,<sup>122</sup> Y. Chao<sup>122</sup> K. F. Chen<sup>122</sup> P. s. Chen,<sup>122</sup> Z. g. Chen,<sup>122</sup> W.-S. Hou<sup>122</sup> T. h. Hsu,<sup>122</sup>  
Y. w. Kao,<sup>122</sup> R. Khurana,<sup>122</sup> G. Kole<sup>122</sup> Y. y. Li<sup>122</sup> R.-S. Lu<sup>122</sup> E. Paganis<sup>122</sup> A. Psallidas,<sup>122</sup> X. f. Su,<sup>122</sup>  
J. Thomas-Wilsker<sup>122</sup> L. s. Tsai,<sup>122</sup> H. y. Wu,<sup>122</sup> E. Yazgan<sup>122</sup> C. Asawatangtrakuldee<sup>123</sup> N. Srimanobhas<sup>123</sup>  
V. Wachirapusanand<sup>123</sup> D. Agyel<sup>124</sup> F. Boran<sup>124</sup> Z. S. Demiroglu<sup>124</sup> F. Dolek<sup>124</sup> I. Dumanoglu<sup>124,qqq</sup>  
E. Eskut<sup>124</sup> Y. Guler<sup>124,rrr</sup> E. Gurpinar Guler<sup>124,rrr</sup> C. Isik<sup>124</sup> O. Kara,<sup>124</sup> A. Kayis Topaksu<sup>124</sup> U. Kiminsu<sup>124</sup>  
G. Onengut<sup>124</sup> K. Ozdemir<sup>124,sss</sup> A. Polatoz<sup>124</sup> B. Tali<sup>124,ttt</sup> U. G. Tok<sup>124</sup> S. Turkcapar<sup>124</sup> E. Uslan<sup>124</sup>  
I. S. Zorbakir<sup>124</sup> M. Yalvac<sup>125,uuu</sup> B. Akgun<sup>126</sup> I. O. Atakisi<sup>126</sup> E. Gürmez<sup>126</sup> M. Kaya<sup>126,vvv</sup> O. Kaya<sup>126,www</sup>  
S. Tekten<sup>126,xxx</sup> A. Cakir<sup>127</sup> K. Cankocak<sup>127,qqq,yyy</sup> Y. Komurcu<sup>127</sup> S. Sen<sup>127,zzz</sup> O. Aydilek<sup>128</sup> S. Cerci<sup>128,ttt</sup>  
V. Epshteyn<sup>128</sup> B. Hacisahinoglu<sup>128</sup> I. Hos<sup>128,aaaa</sup> B. Isildak<sup>128,bbbb</sup> B. Kaynak<sup>128</sup> S. Ozkorucuklu<sup>128</sup>  
O. Potok<sup>128</sup> H. Sert<sup>128</sup> C. Simsek<sup>128</sup> D. Sunar Cerci<sup>128,ttt</sup> C. Zorbilmez<sup>128</sup> A. Boyaryntsev<sup>129</sup> B. Grynyov<sup>129</sup>  
L. Levchuk<sup>130</sup> D. Anthony<sup>131</sup> J. J. Brooke<sup>131</sup> A. Bancock<sup>131</sup> F. Bury<sup>131</sup> E. Clement<sup>131</sup> D. Cussans<sup>131</sup>  
H. Flacher<sup>131</sup> M. Glowacki,<sup>131</sup> J. Goldstein<sup>131</sup> H. F. Heath<sup>131</sup> L. Kreczko<sup>131</sup> B. Krikler<sup>131</sup> S. Paramesvaran<sup>131</sup>  
S. Seif El Nasr-Storey,<sup>131</sup> V. J. Smith<sup>131</sup> N. Stylianou<sup>131,cccc</sup> K. Walkingshaw Pass,<sup>131</sup> R. White<sup>131</sup> A. H. Ball,<sup>132</sup>  
K. W. Bell<sup>132</sup> A. Belyaev<sup>132,dddd</sup> C. Brew<sup>132</sup> R. M. Brown<sup>132</sup> D. J. A. Cockerill<sup>132</sup> C. Cooke<sup>132</sup> K. V. Ellis,<sup>132</sup>  
K. Harder<sup>132</sup> S. Harper<sup>132</sup> M.-L. Holmberg<sup>132,eeee</sup> J. Linacre<sup>132</sup> K. Manolopoulos,<sup>132</sup> D. M. Newbold<sup>132</sup>  
E. Olaiya,<sup>132</sup> D. Petyt<sup>132</sup> T. Reis<sup>132</sup> G. Salvi<sup>132</sup> T. Schuh,<sup>132</sup> C. H. Shepherd-Themistocleous<sup>132</sup> I. R. Tomalin<sup>132</sup>  
T. Williams<sup>132</sup> R. Bainbridge<sup>133</sup> P. Bloch<sup>133</sup> C. E. Brown<sup>133</sup> O. Buchmuller,<sup>133</sup> V. Cacchio,<sup>133</sup>  
C. A. Carrillo Montoya<sup>133</sup> G. S. Chahal<sup>133,ffff</sup> D. Colling<sup>133</sup> J. S. Dancu,<sup>133</sup> P. Dauncey<sup>133</sup> G. Davies<sup>133</sup>  
J. Davies,<sup>133</sup> M. Della Negra<sup>133</sup> S. Fayer,<sup>133</sup> G. Fedi<sup>133</sup> G. Hall<sup>133</sup> M. H. Hassanshahi<sup>133</sup> A. Howard,<sup>133</sup> G. Iles<sup>133</sup>

- M. Knight<sup>133</sup>, J. Langford<sup>133</sup>, L. Lyons<sup>133</sup>, A.-M. Magnan<sup>133</sup>, S. Malik,<sup>133</sup> A. Martelli<sup>133</sup>, M. Mieskolainen<sup>133</sup>, J. Nash<sup>133,gggg</sup>, M. Pesaresi,<sup>133</sup> B. C. Radburn-Smith<sup>133</sup>, A. Richards,<sup>133</sup> A. Rose<sup>133</sup>, C. Seez<sup>133</sup>, R. Shukla<sup>133</sup>, A. Tapper<sup>133</sup>, K. Uchida<sup>133</sup>, G. P. Uttley<sup>133</sup>, L. H. Vage,<sup>133</sup> T. Virdee<sup>133,ff</sup>, M. Vojinovic<sup>133</sup>, N. Wardle<sup>133</sup>, D. Winterbottom<sup>133</sup>, K. Coldham,<sup>134</sup> J. E. Cole<sup>134</sup>, A. Khan,<sup>134</sup> P. Kyberd<sup>134</sup>, I. D. Reid<sup>134</sup>, S. Abdullin<sup>135</sup>, A. Brinkerhoff<sup>135</sup>, B. Caraway<sup>135</sup>, J. Dittmann<sup>135</sup>, K. Hatakeyama<sup>135</sup>, J. Hiltbrand<sup>135</sup>, A. R. Kanuganti<sup>135</sup>, B. McMaster<sup>135</sup>, M. Saunders<sup>135</sup>, S. Sawant<sup>135</sup>, C. Sutantawibul<sup>135</sup>, M. Toms<sup>135,q</sup>, J. Wilson<sup>135</sup>, R. Bartek<sup>136</sup>, A. Dominguez<sup>136</sup>, C. Huerta Escamilla,<sup>136</sup> A. E. Simsek<sup>136</sup>, R. Uniyal<sup>136</sup>, A. M. Vargas Hernandez<sup>136</sup>, R. Chudasama<sup>137</sup>, S. I. Cooper<sup>137</sup>, S. V. Gleyzer<sup>137</sup>, C. U. Perez<sup>137</sup>, P. Rumerio<sup>137,hhh</sup>, E. Usai<sup>137</sup>, C. West<sup>137</sup>, R. Yi<sup>137</sup>, A. Akpinar<sup>138</sup>, A. Albert<sup>138</sup>, D. Arcaro<sup>138</sup>, C. Cosby<sup>138</sup>, Z. Demiragli<sup>138</sup>, C. Erice<sup>138</sup>, E. Fontanesi<sup>138</sup>, D. Gastler<sup>138</sup>, S. Jeon<sup>138</sup>, J. Rohlfs<sup>138</sup>, K. Salyer<sup>138</sup>, D. Sperka<sup>138</sup>, D. Spitzbart<sup>138</sup>, I. Suarez<sup>138</sup>, A. Tsatsos<sup>138</sup>, S. Yuan<sup>138</sup>, G. Benelli<sup>139</sup>, X. Coubez,<sup>139,aa</sup> D. Cutts<sup>139</sup>, M. Hadley<sup>139</sup>, U. Heintz<sup>139</sup>, J. M. Hogan<sup>139,iii</sup>, T. Kwon<sup>139</sup>, G. Landsberg<sup>139</sup>, K. T. Lau<sup>139</sup>, D. Li<sup>139</sup>, J. Luo<sup>139</sup>, S. Mondal<sup>139</sup>, M. Narain<sup>139,a</sup>, N. Pervan<sup>139</sup>, S. Sagir<sup>139,jjj</sup>, F. Simpson<sup>139</sup>, M. Stamenkovic<sup>139</sup>, W. Y. Wong,<sup>139</sup>, X. Yan<sup>139</sup>, W. Zhang,<sup>139</sup>, S. Abbott<sup>140</sup>, J. Bonilla<sup>140</sup>, C. Brainerd<sup>140</sup>, R. Breedon<sup>140</sup>, M. Calderon De La Barca Sanchez<sup>140</sup>, M. Chertok<sup>140</sup>, M. Citron<sup>140</sup>, J. Conway<sup>140</sup>, P. T. Cox<sup>140</sup>, R. Erbacher<sup>140</sup>, F. Jensen<sup>140</sup>, O. Kukral<sup>140</sup>, G. Mocellin<sup>140</sup>, M. Mulhearn<sup>140</sup>, D. Pellett<sup>140</sup>, W. Wei,<sup>140</sup>, Y. Yao<sup>140</sup>, F. Zhang<sup>140</sup>, M. Bachtis<sup>141</sup>, R. Cousins<sup>141</sup>, A. Datta<sup>141</sup>, J. Hauser<sup>141</sup>, M. Ignatenko<sup>141</sup>, M. A. Iqbal<sup>141</sup>, T. Lam<sup>141</sup>, E. Manca<sup>141</sup>, W. A. Nash<sup>141</sup>, D. Saltzberg<sup>141</sup>, B. Stone<sup>141</sup>, V. Valuev<sup>141</sup>, R. Clare<sup>142</sup>, M. Gordon,<sup>142</sup>, G. Hanson<sup>142</sup>, W. Si<sup>142</sup>, S. Wimpenny<sup>142,a</sup>, J. G. Branson<sup>143</sup>, S. Cittolin<sup>143</sup>, S. Cooperstein<sup>143</sup>, D. Diaz<sup>143</sup>, J. Duarte<sup>143</sup>, L. Giannini<sup>143</sup>, J. Guiang<sup>143</sup>, R. Kansal<sup>143</sup>, V. Krutelyov<sup>143</sup>, R. Lee<sup>143</sup>, J. Letts<sup>143</sup>, M. Masciovecchio<sup>143</sup>, F. Mokhtar<sup>143</sup>, M. Pieri<sup>143</sup>, M. Quinnan<sup>143</sup>, B. V. Sathia Narayanan<sup>143</sup>, V. Sharma<sup>143</sup>, M. Tadel<sup>143</sup>, E. Vourliotis<sup>143</sup>, F. Würthwein<sup>143</sup>, Y. Xiang<sup>143</sup>, A. Yagil<sup>143</sup>, A. Barzdukas<sup>144</sup>, L. Brennan,<sup>144</sup>, C. Campagnari<sup>144</sup>, G. Collura<sup>144</sup>, A. Dorsett<sup>144</sup>, J. Incandela<sup>144</sup>, M. Kilpatrick<sup>144</sup>, J. Kim<sup>144</sup>, A. J. Li<sup>144</sup>, P. Masterson<sup>144</sup>, H. Mei<sup>144</sup>, M. Oshiro<sup>144</sup>, J. Richman<sup>144</sup>, U. Sarica<sup>144</sup>, R. Schmitz<sup>144</sup>, F. Setti<sup>144</sup>, J. Sheplock<sup>144</sup>, D. Stuart<sup>144</sup>, S. Wang<sup>144</sup>, A. Bornheim<sup>145</sup>, O. Cerri,<sup>145</sup> A. Latorre,<sup>145</sup> J. M. Lawhorn<sup>145</sup>, J. Mao<sup>145</sup>, H. B. Newman<sup>145</sup>, T. Q. Nguyen<sup>145</sup>, M. Spiropulu<sup>145</sup>, J. R. Vlimant<sup>145</sup>, C. Wang<sup>145</sup>, S. Xie<sup>145</sup>, R. Y. Zhu<sup>145</sup>, J. Alison<sup>146</sup>, S. An<sup>146</sup>, M. B. Andrews<sup>146</sup>, P. Bryant<sup>146</sup>, V. Dutta<sup>146</sup>, T. Ferguson<sup>146</sup>, A. Harilal<sup>146</sup>, C. Liu<sup>146</sup>, T. Mudholkar<sup>146</sup>, S. Murthy<sup>146</sup>, M. Paulini<sup>146</sup>, A. Roberts<sup>146</sup>, A. Sanchez<sup>146</sup>, W. Terrill<sup>146</sup>, J. P. Cumalat<sup>147</sup>, W. T. Ford<sup>147</sup>, A. Hassani<sup>147</sup>, G. Karathanasis<sup>147</sup>, E. MacDonald,<sup>147</sup> N. Manganelli<sup>147</sup>, F. Marin<sup>147</sup>, A. Perloff<sup>147</sup>, C. Savard<sup>147</sup>, N. Schonbeck<sup>147</sup>, K. Stenson<sup>147</sup>, K. A. Ulmer<sup>147</sup>, S. R. Wagner<sup>147</sup>, N. Zipper<sup>147</sup>, J. Alexander<sup>148</sup>, S. Bright-Thonney<sup>148</sup>, X. Chen<sup>148</sup>, D. J. Cranshaw<sup>148</sup>, J. Fan<sup>148</sup>, X. Fan<sup>148</sup>, D. Gadkari<sup>148</sup>, S. Hogan<sup>148</sup>, J. Monroy<sup>148</sup>, J. R. Patterson<sup>148</sup>, J. Reichert<sup>148</sup>, M. Reid<sup>148</sup>, A. Ryd<sup>148</sup>, J. Thom<sup>148</sup>, P. Wittich<sup>148</sup>, R. Zou<sup>148</sup>, M. Albrow<sup>149</sup>, M. Alyari<sup>149</sup>, O. Amram<sup>149</sup>, G. Apollinari<sup>149</sup>, A. Apresyan<sup>149</sup>, L. A. T. Bauerdtick<sup>149</sup>, D. Berry<sup>149</sup>, J. Berryhill<sup>149</sup>, P. C. Bhat<sup>149</sup>, K. Burkett<sup>149</sup>, J. N. Butler<sup>149</sup>, A. Canepa<sup>149</sup>, G. B. Cerati<sup>149</sup>, H. W. K. Cheung<sup>149</sup>, F. Chlebana<sup>149</sup>, G. Cummings<sup>149</sup>, J. Dickinson<sup>149</sup>, I. Dutta<sup>149</sup>, V. D. Elvira<sup>149</sup>, Y. Feng<sup>149</sup>, J. Freeman<sup>149</sup>, A. Gandrakota<sup>149</sup>, Z. Gece<sup>149</sup>, L. Gray<sup>149</sup>, D. Green,<sup>149</sup> A. Grummer<sup>149</sup>, S. Grünendahl<sup>149</sup>, D. Guerrero<sup>149</sup>, O. Gutsche<sup>149</sup>, R. M. Harris<sup>149</sup>, R. Heller<sup>149</sup>, T. C. Herwig<sup>149</sup>, J. Hirschauer<sup>149</sup>, L. Horyn<sup>149</sup>, B. Jayatilaka<sup>149</sup>, S. Jindariani<sup>149</sup>, M. Johnson<sup>149</sup>, U. Joshi<sup>149</sup>, T. Klijnsma<sup>149</sup>, B. Klima<sup>149</sup>, K. H. M. Kwok<sup>149</sup>, S. Lammel<sup>149</sup>, D. Lincoln<sup>149</sup>, R. Lipton<sup>149</sup>, T. Liu<sup>149</sup>, C. Madrid<sup>149</sup>, K. Maeshima<sup>149</sup>, C. Mantilla<sup>149</sup>, D. Mason<sup>149</sup>, P. McBride<sup>149</sup>, P. Merkel<sup>149</sup>, S. Mrenna<sup>149</sup>, S. Nahn<sup>149</sup>, J. Ngadiuba<sup>149</sup>, D. Noonan<sup>149</sup>, V. Papadimitriou<sup>149</sup>, N. Pastika<sup>149</sup>, K. Pedro<sup>149</sup>, C. Pena<sup>149,kkkk</sup>, F. Ravera<sup>149</sup>, A. Reinsvold Hall<sup>149,III</sup>, L. Ristori<sup>149</sup>, E. Sexton-Kennedy<sup>149</sup>, N. Smith<sup>149</sup>, A. Soha<sup>149</sup>, L. Spiegel<sup>149</sup>, S. Stoynev<sup>149</sup>, L. Taylor<sup>149</sup>, S. Tkaczyk<sup>149</sup>, N. V. Tran<sup>149</sup>, L. Uplegger<sup>149</sup>, E. W. Vaandering<sup>149</sup>, I. Zoi<sup>149</sup>, C. Aruta<sup>150</sup>, P. Avery<sup>150</sup>, D. Bourilkov<sup>150</sup>, L. Cadamuro<sup>150</sup>, P. Chang<sup>150</sup>, V. Cherepanov<sup>150</sup>, R. D. Field,<sup>150</sup>, E. Koenig<sup>150</sup>, M. Kolosova<sup>150</sup>, J. Konigsberg<sup>150</sup>, A. Korytov<sup>150</sup>, K. H. Lo,<sup>150</sup> K. Matchev<sup>150</sup>, N. Menendez<sup>150</sup>, G. Mitselmakher<sup>150</sup>, K. Mohrman<sup>150</sup>, A. Muthirakalayil Madhu<sup>150</sup>, N. Rawal<sup>150</sup>, D. Rosenzweig<sup>150</sup>, S. Rosenzweig<sup>150</sup>, K. Shi<sup>150</sup>, J. Wang<sup>150</sup>, T. Adams<sup>151</sup>, A. Al Kadhim<sup>151</sup>, A. Askew<sup>151</sup>, N. Bower<sup>151</sup>, R. Habibullah<sup>151</sup>, V. Hagopian<sup>151</sup>, R. Hashmi<sup>151</sup>, R. S. Kim<sup>151</sup>, S. Kim<sup>151</sup>, T. Kolberg<sup>151</sup>, G. Martinez,<sup>151</sup>, H. Prosper<sup>151</sup>, P. R. Prova,<sup>151</sup>, O. Viazlo<sup>151</sup>, M. Wulansatiti<sup>151</sup>, R. Yohay<sup>151</sup>, J. Zhang,<sup>151</sup>, B. Alsufyani,<sup>152</sup>, M. M. Baarmand<sup>152</sup>, S. Butalla<sup>152</sup>, T. Elkafrawy<sup>152,ddd</sup>, M. Hohlmann<sup>152</sup>, R. Kumar Verma<sup>152</sup>

- M. Rahmani,<sup>152</sup> M. R. Adams,<sup>153</sup> C. Bennett,<sup>153</sup> R. Cavanaugh,<sup>153</sup> S. Dittmer,<sup>153</sup> R. Escobar Franco,<sup>153</sup>  
O. Evdokimov,<sup>153</sup> C. E. Gerber,<sup>153</sup> D. J. Hofman,<sup>153</sup> J. h. Lee,<sup>153</sup> D. S. Lemos,<sup>153</sup> A. H. Merrit,<sup>153</sup> C. Mills,<sup>153</sup>  
S. Nanda,<sup>153</sup> G. Oh,<sup>153</sup> B. Ozek,<sup>153</sup> D. Pilipovic,<sup>153</sup> T. Roy,<sup>153</sup> S. Rudrabhatla,<sup>153</sup> M. B. Tonjes,<sup>153</sup>  
N. Varelas,<sup>153</sup> X. Wang,<sup>153</sup> Z. Ye,<sup>153</sup> J. Yoo,<sup>153</sup> M. Alhusseini,<sup>154</sup> D. Blend,<sup>154</sup> K. Dilisiz,<sup>154,mmmm</sup> L. Emediato,<sup>154</sup>  
G. Karaman,<sup>154</sup> O. K. Köseyan,<sup>154</sup> J.-P. Merlo,<sup>154</sup> A. Mestvirishvili,<sup>154,nnnn</sup> J. Nachtman,<sup>154</sup> O. Neogi,<sup>154</sup>  
H. Ogul,<sup>154,oooo</sup> Y. Onel,<sup>154</sup> A. Penzo,<sup>154</sup> C. Snyder,<sup>154</sup> E. Tiras,<sup>154,pppp</sup> B. Blumenfeld,<sup>155</sup> L. Corcodilos,<sup>155</sup>  
J. Davis,<sup>155</sup> A. V. Gritsan,<sup>155</sup> L. Kang,<sup>155</sup> S. Kyriacou,<sup>155</sup> P. Maksimovic,<sup>155</sup> M. Roguljic,<sup>155</sup> J. Roskes,<sup>155</sup>  
S. Sekhar,<sup>155</sup> M. Swartz,<sup>155</sup> T. Á. Vámi,<sup>155</sup> A. Abreu,<sup>156</sup> L. F. Alcerro Alcerro,<sup>156</sup> J. Anguiano,<sup>156</sup> P. Baringer,<sup>156</sup>  
A. Bean,<sup>156</sup> Z. Flowers,<sup>156</sup> D. Grove,<sup>156</sup> J. King,<sup>156</sup> G. Krintiras,<sup>156</sup> M. Lazarovits,<sup>156</sup> C. Le Mahieu,<sup>156</sup>  
C. Lindsey,<sup>156</sup> J. Marquez,<sup>156</sup> N. Minafra,<sup>156</sup> M. Murray,<sup>156</sup> M. Nickel,<sup>156</sup> M. Pitt,<sup>156</sup> S. Popescu,<sup>156,qqqq</sup>  
C. Rogan,<sup>156</sup> C. Royon,<sup>156</sup> R. Salvatico,<sup>156</sup> S. Sanders,<sup>156</sup> C. Smith,<sup>156</sup> Q. Wang,<sup>156</sup> G. Wilson,<sup>156</sup>  
B. Allmond,<sup>157</sup> A. Ivanov,<sup>157</sup> K. Kaadze,<sup>157</sup> A. Kalogeropoulos,<sup>157</sup> D. Kim,<sup>157</sup> Y. Maravin,<sup>157</sup> K. Nam,<sup>157</sup>  
J. Natoli,<sup>157</sup> D. Roy,<sup>157</sup> G. Sorrentino,<sup>157</sup> F. Rebassoo,<sup>158</sup> D. Wright,<sup>158</sup> E. Adams,<sup>159</sup> A. Baden,<sup>159</sup> O. Baron,<sup>159</sup>  
A. Belloni,<sup>159</sup> A. Bethani,<sup>159</sup> Y. M. Chen,<sup>159</sup> S. C. Eno,<sup>159</sup> N. J. Hadley,<sup>159</sup> S. Jabeen,<sup>159</sup> R. G. Kellogg,<sup>159</sup>  
T. Koeth,<sup>159</sup> Y. Lai,<sup>159</sup> S. Lascio,<sup>159</sup> A. C. Mignerey,<sup>159</sup> S. Nabili,<sup>159</sup> C. Palmer,<sup>159</sup> C. Papageorgakis,<sup>159</sup>  
M. M. Paranjpe,<sup>159</sup> L. Wang,<sup>159</sup> K. Wong,<sup>159</sup> J. Bendavid,<sup>160</sup> W. Busza,<sup>160</sup> I. A. Cali,<sup>160</sup> Y. Chen,<sup>160</sup>  
M. D'Alfonso,<sup>160</sup> J. Eysermans,<sup>160</sup> C. Freer,<sup>160</sup> G. Gomez-Ceballos,<sup>160</sup> M. Goncharov,<sup>160</sup> P. Harris,<sup>160</sup> D. Hoang,<sup>160</sup>  
D. Kovalskyi,<sup>160</sup> J. Krupa,<sup>160</sup> L. Lavezzo,<sup>160</sup> Y.-J. Lee,<sup>160</sup> K. Long,<sup>160</sup> C. Mironov,<sup>160</sup> C. Paus,<sup>160</sup>  
D. Rankin,<sup>160</sup> C. Roland,<sup>160</sup> G. Roland,<sup>160</sup> S. Rothman,<sup>160</sup> Z. Shi,<sup>160</sup> G. S. F. Stephans,<sup>160</sup> J. Wang,<sup>160</sup>  
Z. Wang,<sup>160</sup> B. Wyslouch,<sup>160</sup> T. J. Yang,<sup>160</sup> B. Crossman,<sup>161</sup> B. M. Joshi,<sup>161</sup> C. Kapsiak,<sup>161</sup> M. Krohn,<sup>161</sup>  
D. Mahon,<sup>161</sup> J. Mans,<sup>161</sup> B. Marzocchi,<sup>161</sup> S. Pandey,<sup>161</sup> M. Revering,<sup>161</sup> R. Rusack,<sup>161</sup> R. Saradhy,<sup>161</sup>  
N. Schroeder,<sup>161</sup> N. Strobbe,<sup>161</sup> M. A. Wadud,<sup>161</sup> L. M. Cremaldi,<sup>162</sup> K. Bloom,<sup>163</sup> M. Bryson,<sup>163</sup> D. R. Claes,<sup>163</sup>  
C. Fangmeier,<sup>163</sup> F. Golf,<sup>163</sup> G. Haza,<sup>163</sup> J. Hossain,<sup>163</sup> C. Joo,<sup>163</sup> I. Kravchenko,<sup>163</sup> I. Reed,<sup>163</sup> J. E. Siado,<sup>163</sup>  
W. Tabb,<sup>163</sup> A. Vagnerini,<sup>163</sup> A. Wightman,<sup>163</sup> F. Yan,<sup>163</sup> D. Yu,<sup>163</sup> A. G. Zecchinelli,<sup>163</sup> G. Agarwal,<sup>164</sup>  
H. Bandyopadhyay,<sup>164</sup> L. Hay,<sup>164</sup> I. Iashvili,<sup>164</sup> A. Kharchilava,<sup>164</sup> C. McLean,<sup>164</sup> M. Morris,<sup>164</sup> D. Nguyen,<sup>164</sup>  
S. Rappoccio,<sup>164</sup> H. Rejeb Sfar,<sup>164</sup> A. Williams,<sup>164</sup> G. Alverson,<sup>165</sup> E. Barberis,<sup>165</sup> Y. Haddad,<sup>165</sup> Y. Han,<sup>165</sup>  
A. Krishna,<sup>165</sup> J. Li,<sup>165</sup> M. Lu,<sup>165</sup> G. Madigan,<sup>165</sup> R. McCarthy,<sup>165</sup> D. M. Morse,<sup>165</sup> V. Nguyen,<sup>165</sup>  
T. Orimoto,<sup>165</sup> A. Parker,<sup>165</sup> L. Skinnari,<sup>165</sup> A. Tishelman-Charny,<sup>165</sup> B. Wang,<sup>165</sup> D. Wood,<sup>165</sup>  
S. Bhattacharya,<sup>166</sup> J. Bueghly,<sup>166</sup> Z. Chen,<sup>166</sup> K. A. Hahn,<sup>166</sup> Y. Liu,<sup>166</sup> Y. Miao,<sup>166</sup> D. G. Monk,<sup>166</sup>  
M. H. Schmitt,<sup>166</sup> A. Taliercio,<sup>166</sup> M. Velasco,<sup>166</sup> R. Band,<sup>167</sup> R. Bucci,<sup>167</sup> S. Castells,<sup>167</sup> M. Cremonesi,<sup>167</sup>  
A. Das,<sup>167</sup> R. Goldouzian,<sup>167</sup> M. Hildreth,<sup>167</sup> K. W. Ho,<sup>167</sup> K. Hurtado Anampa,<sup>167</sup> C. Jessop,<sup>167</sup> K. Lannon,<sup>167</sup>  
J. Lawrence,<sup>167</sup> N. Loukas,<sup>167</sup> L. Lutton,<sup>167</sup> J. Mariano,<sup>167</sup> N. Marinelli,<sup>167</sup> I. McAlister,<sup>167</sup> T. McCauley,<sup>167</sup>  
C. McGrady,<sup>167</sup> C. Moore,<sup>167</sup> Y. Musienko,<sup>167,q</sup> H. Nelson,<sup>167</sup> M. Osherson,<sup>167</sup> R. Ruchti,<sup>167</sup> A. Townsend,<sup>167</sup>  
M. Wayne,<sup>167</sup> H. Yockey,<sup>167</sup> M. Zarucki,<sup>167</sup> L. Zygalas,<sup>167</sup> A. Basnet,<sup>168</sup> B. Bylsma,<sup>168</sup> M. Carrigan,<sup>168</sup>  
L. S. Durkin,<sup>168</sup> C. Hill,<sup>168</sup> M. Joyce,<sup>168</sup> A. Lesauvage,<sup>168</sup> M. Nunez Ornelas,<sup>168</sup> K. Wei,<sup>168</sup> B. L. Winer,<sup>168</sup>  
B. R. Yates,<sup>168</sup> F. M. Addesa,<sup>169</sup> H. Bouchamaoui,<sup>169</sup> P. Das,<sup>169</sup> G. Dezoort,<sup>169</sup> P. Elmer,<sup>169</sup> A. Frankenthal,<sup>169</sup>  
B. Greenberg,<sup>169</sup> N. Haubrich,<sup>169</sup> S. Higginbotham,<sup>169</sup> G. Kopp,<sup>169</sup> S. Kwan,<sup>169</sup> D. Lange,<sup>169</sup> A. Loeliger,<sup>169</sup>  
D. Marlow,<sup>169</sup> I. Ojalvo,<sup>169</sup> J. Olsen,<sup>169</sup> A. Shevelev,<sup>169</sup> D. Stickland,<sup>169</sup> C. Tully,<sup>169</sup> S. Malik,<sup>170</sup>  
A. S. Bakshi,<sup>171</sup> V. E. Barnes,<sup>171</sup> S. Chandra,<sup>171</sup> R. Chawla,<sup>171</sup> S. Das,<sup>171</sup> A. Gu,<sup>171</sup> L. Gutay,<sup>171</sup> M. Jones,<sup>171</sup>  
A. W. Jung,<sup>171</sup> D. Kondratyev,<sup>171</sup> A. M. Koshy,<sup>171</sup> M. Liu,<sup>171</sup> G. Negro,<sup>171</sup> N. Neumeister,<sup>171</sup> G. Paspalaki,<sup>171</sup>  
S. Piperov,<sup>171</sup> V. Scheurer,<sup>171</sup> J. F. Schulte,<sup>171</sup> M. Stojanovic,<sup>171</sup> J. Thieman,<sup>171</sup> A. K. Virdi,<sup>171</sup> F. Wang,<sup>171</sup>  
W. Xie,<sup>171</sup> J. Dolen,<sup>172</sup> N. Parashar,<sup>172</sup> A. Pathak,<sup>172</sup> D. Acosta,<sup>173</sup> A. Baty,<sup>173</sup> T. Carnahan,<sup>173</sup> S. Dildick,<sup>173</sup>  
K. M. Ecklund,<sup>173</sup> P. J. Fernández Manteca,<sup>173</sup> S. Freed,<sup>173</sup> P. Gardner,<sup>173</sup> F. J. M. Geurts,<sup>173</sup> A. Kumar,<sup>173</sup> W. Li,<sup>173</sup>  
O. Miguel Colin,<sup>173</sup> B. P. Padley,<sup>173</sup> R. Redjimi,<sup>173</sup> J. Rotter,<sup>173</sup> E. Yigitbasi,<sup>173</sup> Y. Zhang,<sup>173</sup> A. Bodek,<sup>174</sup>  
P. de Barbaro,<sup>174</sup> R. Demina,<sup>174</sup> J. L. Dulemba,<sup>174</sup> C. Fallon,<sup>174</sup> A. Garcia-Bellido,<sup>174</sup> O. Hindrichs,<sup>174</sup>  
A. Khukhunaishvili,<sup>174</sup> P. Parygin,<sup>174,q</sup> E. Popova,<sup>174,q</sup> R. Taus,<sup>174</sup> G. P. Van Onsem,<sup>174</sup> K. Goulianatos,<sup>175</sup>  
B. Chiarito,<sup>176</sup> J. P. Chou,<sup>176</sup> Y. Gershtein,<sup>176</sup> E. Halkiadakis,<sup>176</sup> A. Hart,<sup>176</sup> M. Heindl,<sup>176</sup> D. Jaroslawski,<sup>176</sup>  
O. Karacheban,<sup>176,dd</sup> I. Laflotte,<sup>176</sup> A. Lath,<sup>176</sup> R. Montalvo,<sup>176</sup> K. Nash,<sup>176</sup> H. Routray,<sup>176</sup> S. Salur,<sup>176</sup>  
S. Schnetzer,<sup>176</sup> S. Somalwar,<sup>176</sup> R. Stone,<sup>176</sup> S. A. Thayil,<sup>176</sup> S. Thomas,<sup>176</sup> J. Vora,<sup>176</sup> H. Wang,<sup>176</sup> H. Acharya,<sup>177</sup>

- D. Ally<sup>177</sup>, A. G. Delannoy<sup>177</sup>, S. Fiorendi<sup>177</sup>, T. Holmes<sup>177</sup>, N. Karunaratna<sup>177</sup>, L. Lee<sup>177</sup>, E. Nibigira<sup>177</sup>, S. Spanier<sup>177</sup>, D. Aebi<sup>178</sup>, M. Ahmad<sup>178</sup>, O. Bouhalil<sup>178,rrr</sup>, M. Dalchenko<sup>178</sup>, R. Eusebi<sup>178</sup>, J. Gilmore<sup>178</sup>, T. Huang<sup>178</sup>, T. Kamon<sup>178,ssss</sup>, H. Kim<sup>178</sup>, S. Luo<sup>178</sup>, S. Malhotra<sup>178</sup>, R. Mueller<sup>178</sup>, D. Overton<sup>178</sup>, D. Rathjens<sup>178</sup>, A. Safonov<sup>178</sup>, N. Akchurin<sup>179</sup>, J. Damgov<sup>179</sup>, V. Hegde<sup>179</sup>, A. Hussain<sup>179</sup>, Y. Kazhykarim,<sup>179</sup>, K. Lamichhane<sup>179</sup>, S. W. Lee<sup>179</sup>, A. Mankel<sup>179</sup>, T. Mengke<sup>179</sup>, S. Muthumuni<sup>179</sup>, T. Peltola<sup>179</sup>, I. Volobouev<sup>179</sup>, A. Whitbeck<sup>179</sup>, E. Appelt<sup>180</sup>, S. Greene<sup>180</sup>, A. Gurrola<sup>180</sup>, W. Johns<sup>180</sup>, R. Kunnavalkam Elayavalli<sup>180</sup>, A. Melo<sup>180</sup>, F. Romeo<sup>180</sup>, P. Sheldon<sup>180</sup>, S. Tuo<sup>180</sup>, J. Velkovska<sup>180</sup>, J. Viinikainen<sup>180</sup>, B. Cardwell<sup>181</sup>, B. Cox<sup>181</sup>, J. Hakala<sup>181</sup>, R. Hirosky<sup>181</sup>, A. Ledovskoy<sup>181</sup>, A. Li<sup>181</sup>, C. Neu<sup>181</sup>, C. E. Perez Lara<sup>181</sup>, P. E. Karchin<sup>182</sup>, A. Aravind,<sup>183</sup>, S. Banerjee<sup>183</sup>, K. Black<sup>183</sup>, T. Bose<sup>183</sup>, S. Dasu<sup>183</sup>, I. De Bruyn<sup>183</sup>, P. Everaerts<sup>183</sup>, C. Galloni,<sup>183</sup>, H. He<sup>183</sup>, M. Herndon<sup>183</sup>, A. Herve<sup>183</sup>, C. K. Koraka<sup>183</sup>, A. Lanaro,<sup>183</sup>, R. Loveless<sup>183</sup>, J. Madhusudanan Sreekala<sup>183</sup>, A. Mallampalli<sup>183</sup>, A. Mohammadi<sup>183</sup>, S. Mondal,<sup>183</sup>, G. Parida<sup>183</sup>, D. Pinna,<sup>183</sup>, A. Savin,<sup>183</sup>, V. Shang<sup>183</sup>, V. Sharma<sup>183</sup>, W. H. Smith<sup>183</sup>, D. Teague,<sup>183</sup>, H. F. Tsoi<sup>183</sup>, W. Vetens<sup>183</sup>, A. Warden<sup>183</sup>, S. Afanasiev<sup>184</sup>, V. Andreev<sup>184</sup>, Yu. Andreev<sup>184</sup>, T. Aushev<sup>184</sup>, M. Azarkin<sup>184</sup>, A. Babaev<sup>184</sup>, A. Belyaev<sup>184</sup>, V. Blinov,<sup>184,q</sup>, E. Boos<sup>184</sup>, V. Borshch<sup>184</sup>, D. Budkouski<sup>184</sup>, V. Bunichev<sup>184</sup>, M. Chadeeva<sup>184,q</sup>, V. Chekhovsky,<sup>184</sup>, R. Chistov<sup>184,q</sup>, A. Dermenev<sup>184</sup>, T. Dimova<sup>184,q</sup>, D. Druzhkin<sup>184,ttt</sup>, M. Dubinin<sup>184,kkkk</sup>, L. Dudko<sup>184</sup>, A. Ershov<sup>184</sup>, G. Gavrilov<sup>184</sup>, V. Gavrilov<sup>184</sup>, S. Gninenco<sup>184</sup>, V. Golovtcov<sup>184</sup>, N. Golubev<sup>184</sup>, I. Golutvin<sup>184</sup>, I. Gorbunov<sup>184</sup>, A. Gribushin<sup>184</sup>, Y. Ivanov<sup>184</sup>, V. Kachanov<sup>184</sup>, L. Kardapoltsev<sup>184,q</sup>, V. Karjavine<sup>184</sup>, A. Karneyeu<sup>184</sup>, V. Kim<sup>184,q</sup>, M. Kirakosyan,<sup>184</sup>, D. Kirpichnikov<sup>184</sup>, M. Kirsanov<sup>184</sup>, V. Klyukhin<sup>184</sup>, O. Kodolova<sup>184,uuuu</sup>, D. Konstantinov<sup>184</sup>, V. Korenkov<sup>184</sup>, A. Kozyrev<sup>184,q</sup>, N. Krasnikov<sup>184</sup>, A. Lanev<sup>184</sup>, P. Levchenko<sup>184,vvv</sup>, N. Lychkovskaya<sup>184</sup>, V. Makarenko<sup>184</sup>, A. Malakhov<sup>184</sup>, V. Matveev<sup>184,q</sup>, V. Murzin<sup>184</sup>, A. Nikitenko<sup>184,wwww,uuuu</sup>, S. Obraztsov<sup>184</sup>, V. Oreshkin<sup>184</sup>, V. Palichik<sup>184</sup>, V. Perelygin<sup>184</sup>, M. Perfilov,<sup>184</sup>, S. Petrushanko<sup>184</sup>, S. Polikarpov<sup>184,q</sup>, V. Popov,<sup>184</sup>, O. Radchenko<sup>184,q</sup>, M. Savina<sup>184</sup>, V. Savrin<sup>184</sup>, V. Shalaev<sup>184</sup>, S. Shmatov<sup>184</sup>, S. Shulha<sup>184</sup>, Y. Skovpen<sup>184,q</sup>, S. Slabospitskii<sup>184</sup>, V. Smirnov<sup>184</sup>, D. Sosnov<sup>184</sup>, V. Sulimov<sup>184</sup>, E. Tcherniaev<sup>184</sup>, A. Terkulov<sup>184</sup>, O. Teryaev<sup>184</sup>, I. Tlisova<sup>184</sup>, A. Toropin<sup>184</sup>, L. Uvarov<sup>184</sup>, A. Uzunian<sup>184</sup>, A. Vorobyev,<sup>184,a</sup>, N. Voytishin<sup>184</sup>, B. S. Yuldashev,<sup>184,xxxx</sup>, A. Zarubin<sup>184</sup>, I. Zhizhin<sup>184</sup>, and A. Zhokin<sup>184</sup>

(CMS Collaboration)

<sup>1</sup>*Yerevan Physics Institute, Yerevan, Armenia*<sup>2</sup>*Institut für Hochenergiephysik, Vienna, Austria*<sup>3</sup>*Universiteit Antwerpen, Antwerpen, Belgium*<sup>4</sup>*Vrije Universiteit Brussel, Brussel, Belgium*<sup>5</sup>*Université Libre de Bruxelles, Bruxelles, Belgium*<sup>6</sup>*Ghent University, Ghent, Belgium*<sup>7</sup>*Université Catholique de Louvain, Louvain-la-Neuve, Belgium*<sup>8</sup>*Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil*<sup>9</sup>*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*<sup>10</sup>*Universidade Estadual Paulista, Universidade Federal do ABC, São Paulo, Brazil*<sup>11</sup>*Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria*<sup>12</sup>*University of Sofia, Sofia, Bulgaria*<sup>13</sup>*Instituto De Alta Investigación, Universidad de Tarapacá, Casilla 7 D, Arica, Chile*<sup>14</sup>*Beihang University, Beijing, China*<sup>15</sup>*Department of Physics, Tsinghua University, Beijing, China*<sup>16</sup>*Institute of High Energy Physics, Beijing, China*<sup>17</sup>*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*<sup>18</sup>*Sun Yat-Sen University, Guangzhou, China*<sup>19</sup>*University of Science and Technology of China, Hefei, China*<sup>20</sup>*Institute of Modern Physics and Key Laboratory of Nuclear Physics and**Ion-beam Application (MOE) - Fudan University, Shanghai, China*<sup>21</sup>*Zhejiang University, Hangzhou, Zhejiang, China*<sup>22</sup>*Universidad de Los Andes, Bogota, Colombia*<sup>23</sup>*Universidad de Antioquia, Medellin, Colombia*

- <sup>24</sup>University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia  
<sup>25</sup>University of Split, Faculty of Science, Split, Croatia  
<sup>26</sup>Institute Rudjer Boskovic, Zagreb, Croatia  
<sup>27</sup>University of Cyprus, Nicosia, Cyprus  
<sup>28</sup>Charles University, Prague, Czech Republic  
<sup>29</sup>Escuela Politecnica Nacional, Quito, Ecuador  
<sup>30</sup>Universidad San Francisco de Quito, Quito, Ecuador  
<sup>31</sup>Academy of Scientific Research and Technology of the Arab Republic of Egypt,  
Egyptian Network of High Energy Physics, Cairo, Egypt  
<sup>32</sup>Center for High Energy Physics (CHEP-FU), Fayoum University, El-Fayoum, Egypt  
<sup>33</sup>National Institute of Chemical Physics and Biophysics, Tallinn, Estonia  
<sup>34</sup>Department of Physics, University of Helsinki, Helsinki, Finland  
<sup>35</sup>Helsinki Institute of Physics, Helsinki, Finland  
<sup>36</sup>Lappeenranta-Lahti University of Technology, Lappeenranta, Finland  
<sup>37</sup>IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France  
<sup>38</sup>Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France  
<sup>39</sup>Université de Strasbourg, CNRS, IPHC UMR 7178, Strasbourg, France  
<sup>40</sup>Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France  
<sup>41</sup>Georgian Technical University, Tbilisi, Georgia  
<sup>42</sup>RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany  
<sup>43</sup>RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany  
<sup>44</sup>RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany  
<sup>45</sup>Deutsches Elektronen-Synchrotron, Hamburg, Germany  
<sup>46</sup>University of Hamburg, Hamburg, Germany  
<sup>47</sup>Karlsruhe Institut fuer Technologie, Karlsruhe, Germany  
<sup>48</sup>Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece  
<sup>49</sup>National and Kapodistrian University of Athens, Athens, Greece  
<sup>50</sup>National Technical University of Athens, Athens, Greece  
<sup>51</sup>University of Ioánnina, Ioánnina, Greece  
<sup>52</sup>MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary  
<sup>53</sup>Wigner Research Centre for Physics, Budapest, Hungary  
<sup>54</sup>Faculty of Informatics, University of Debrecen, Debrecen, Hungary  
<sup>55</sup>Institute of Nuclear Research ATOMKI, Debrecen, Hungary  
<sup>56</sup>Karoly Robert Campus, MATE Institute of Technology, Gyongyos, Hungary  
<sup>57</sup>Panjab University, Chandigarh, India  
<sup>58</sup>University of Delhi, Delhi, India  
<sup>59</sup>Saha Institute of Nuclear Physics, HBNI, Kolkata, India  
<sup>60</sup>Indian Institute of Technology Madras, Madras, India  
<sup>61</sup>Tata Institute of Fundamental Research-A, Mumbai, India  
<sup>62</sup>Tata Institute of Fundamental Research-B, Mumbai, India  
<sup>63</sup>National Institute of Science Education and Research, An OCC of Homi Bhabha National Institute, Bhubaneswar, Odisha, India  
<sup>64</sup>Indian Institute of Science Education and Research (IISER), Pune, India  
<sup>65</sup>Isfahan University of Technology, Isfahan, Iran  
<sup>66</sup>Institute for Research in Fundamental Sciences (IPM), Tehran, Iran  
<sup>67</sup>University College Dublin, Dublin, Ireland  
<sup>68a</sup>INFN Sezione di Bari, Bari, Italy  
<sup>68b</sup>Università di Bari, Bari, Italy  
<sup>68c</sup>Politecnico di Bari, Bari, Italy  
<sup>69a</sup>INFN Sezione di Bologna, Bologna, Italy  
<sup>69b</sup>Università di Bologna, Bologna, Italy  
<sup>70a</sup>INFN Sezione di Catania, Catania, Italy  
<sup>70b</sup>Università di Catania, Catania, Italy  
<sup>71a</sup>INFN Sezione di Firenze, Firenze, Italy  
<sup>71b</sup>Università di Firenze, Firenze, Italy  
<sup>72</sup>INFN Laboratori Nazionali di Frascati, Frascati, Italy  
<sup>73a</sup>INFN Sezione di Genova, Genova, Italy  
<sup>73b</sup>Università di Genova, Genova, Italy  
<sup>74a</sup>INFN Sezione di Milano-Bicocca, Milano, Italy  
<sup>74b</sup>Università di Milano-Bicocca, Milano, Italy  
<sup>75a</sup>INFN Sezione di Napoli, Napoli, Italy

- <sup>75b</sup>Università di Napoli 'Federico II', Napoli, Italy  
<sup>75c</sup>Università della Basilicata, Potenza, Italy  
<sup>75d</sup>Università G. Marconi, Roma, Italy  
<sup>76a</sup>INFN Sezione di Padova, Padova, Italy  
<sup>76b</sup>Università di Padova, Padova, Italy  
<sup>76c</sup>Università di Trento, Trento, Italy  
<sup>77a</sup>INFN Sezione di Pavia, Pavia, Italy  
<sup>77b</sup>Università di Pavia, Pavia, Italy  
<sup>78a</sup>INFN Sezione di Perugia, Perugia, Italy  
<sup>78b</sup>Università di Perugia, Perugia, Italy  
<sup>79a</sup>INFN Sezione di Pisa, Pisa, Italy  
<sup>79b</sup>Università di Pisa, Pisa, Italy  
<sup>79c</sup>Scuola Normale Superiore di Pisa, Pisa, Italy  
<sup>79d</sup>Università di Siena, Siena, Italy  
<sup>80a</sup>INFN Sezione di Roma, Roma, Italy  
<sup>80b</sup>Sapienza Università di Roma, Roma, Italy  
<sup>81a</sup>INFN Sezione di Torino, Torino, Italy  
<sup>81b</sup>Università di Torino, Torino, Italy  
<sup>81c</sup>Università del Piemonte Orientale, Novara, Italy  
<sup>82a</sup>INFN Sezione di Trieste, Trieste, Italy  
<sup>82b</sup>Università di Trieste, Trieste, Italy  
<sup>83</sup>Kyungpook National University, Daegu, Korea  
<sup>84</sup>Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea  
<sup>85</sup>Hanyang University, Seoul, Korea  
<sup>86</sup>Korea University, Seoul, Korea  
<sup>87</sup>Kyung Hee University, Department of Physics, Seoul, Korea  
<sup>88</sup>Sejong University, Seoul, Korea  
<sup>89</sup>Seoul National University, Seoul, Korea  
<sup>90</sup>University of Seoul, Seoul, Korea  
<sup>91</sup>Yonsei University, Department of Physics, Seoul, Korea  
<sup>92</sup>Sungkyunkwan University, Suwon, Korea  
<sup>93</sup>College of Engineering and Technology, American University of the Middle East (AUM), Dasman, Kuwait  
<sup>94</sup>Riga Technical University, Riga, Latvia  
<sup>95</sup>University of Latvia (LU), Riga, Latvia  
<sup>96</sup>Vilnius University, Vilnius, Lithuania  
<sup>97</sup>National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia  
<sup>98</sup>Universidad de Sonora (UNISON), Hermosillo, Mexico  
<sup>99</sup>Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico  
<sup>100</sup>Universidad Iberoamericana, Mexico City, Mexico  
<sup>101</sup>Benemerita Universidad Autonoma de Puebla, Puebla, Mexico  
<sup>102</sup>University of Montenegro, Podgorica, Montenegro  
<sup>103</sup>University of Canterbury, Christchurch, New Zealand  
<sup>104</sup>National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan  
<sup>105</sup>AGH University of Science and Technology Faculty of Computer Science, Electronics and Telecommunications, Krakow, Poland  
<sup>106</sup>National Centre for Nuclear Research, Swierk, Poland  
<sup>107</sup>Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland  
<sup>108</sup>Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal  
<sup>109</sup>Faculty of Physics, University of Belgrade, Belgrade, Serbia  
<sup>110</sup>VINCA Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia  
<sup>111</sup>Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain  
<sup>112</sup>Universidad Autónoma de Madrid, Madrid, Spain  
<sup>113</sup>Universidad de Oviedo, Instituto Universitario de Ciencias y Tecnologías Espaciales de Asturias (ICTEA), Oviedo, Spain  
<sup>114</sup>Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain  
<sup>115</sup>University of Colombo, Colombo, Sri Lanka  
<sup>116</sup>University of Ruhuna, Department of Physics, Matara, Sri Lanka  
<sup>117</sup>CERN, European Organization for Nuclear Research, Geneva, Switzerland  
<sup>118</sup>Paul Scherrer Institut, Villigen, Switzerland  
<sup>119</sup>ETH Zurich - Institute for Particle Physics and Astrophysics (IPA), Zurich, Switzerland  
<sup>120</sup>Universität Zürich, Zurich, Switzerland  
<sup>121</sup>National Central University, Chung-Li, Taiwan

- <sup>122</sup>National Taiwan University (NTU), Taipei, Taiwan  
<sup>123</sup>Chulalongkorn University, Faculty of Science, Department of Physics, Bangkok, Thailand  
<sup>124</sup>Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey  
<sup>125</sup>Middle East Technical University, Physics Department, Ankara, Turkey  
<sup>126</sup>Bogazici University, Istanbul, Turkey  
<sup>127</sup>Istanbul Technical University, Istanbul, Turkey  
<sup>128</sup>Istanbul University, Istanbul, Turkey  
<sup>129</sup>Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkiv, Ukraine  
<sup>130</sup>National Science Centre, Kharkiv Institute of Physics and Technology, Kharkiv, Ukraine  
<sup>131</sup>University of Bristol, Bristol, United Kingdom  
<sup>132</sup>Rutherford Appleton Laboratory, Didcot, United Kingdom  
<sup>133</sup>Imperial College, London, United Kingdom  
<sup>134</sup>Brunel University, Uxbridge, United Kingdom  
<sup>135</sup>Baylor University, Waco, Texas, USA  
<sup>136</sup>Catholic University of America, Washington, DC, USA  
<sup>137</sup>The University of Alabama, Tuscaloosa, Alabama, USA  
<sup>138</sup>Boston University, Boston, Massachusetts, USA  
<sup>139</sup>Brown University, Providence, Rhode Island, USA  
<sup>140</sup>University of California, Davis, Davis, California, USA  
<sup>141</sup>University of California, Los Angeles, California, USA  
<sup>142</sup>University of California, Riverside, Riverside, California, USA  
<sup>143</sup>University of California, San Diego, La Jolla, California, USA  
<sup>144</sup>University of California, Santa Barbara - Department of Physics, Santa Barbara, California, USA  
<sup>145</sup>California Institute of Technology, Pasadena, California, USA  
<sup>146</sup>Carnegie Mellon University, Pittsburgh, Pennsylvania, USA  
<sup>147</sup>University of Colorado Boulder, Boulder, Colorado, USA  
<sup>148</sup>Cornell University, Ithaca, New York, USA  
<sup>149</sup>Fermi National Accelerator Laboratory, Batavia, Illinois, USA  
<sup>150</sup>University of Florida, Gainesville, Florida, USA  
<sup>151</sup>Florida State University, Tallahassee, Florida, USA  
<sup>152</sup>Florida Institute of Technology, Melbourne, Florida, USA  
<sup>153</sup>University of Illinois at Chicago (UIC), Chicago, Illinois, USA  
<sup>154</sup>The University of Iowa, Iowa City, Iowa, USA  
<sup>155</sup>Johns Hopkins University, Baltimore, Maryland, USA  
<sup>156</sup>The University of Kansas, Lawrence, Kansas, USA  
<sup>157</sup>Kansas State University, Manhattan, Kansas, USA  
<sup>158</sup>Lawrence Livermore National Laboratory, Livermore, California, USA  
<sup>159</sup>University of Maryland, College Park, Maryland, USA  
<sup>160</sup>Massachusetts Institute of Technology, Cambridge, Massachusetts, USA  
<sup>161</sup>University of Minnesota, Minneapolis, Minnesota, USA  
<sup>162</sup>University of Mississippi, Oxford, Mississippi, USA  
<sup>163</sup>University of Nebraska-Lincoln, Lincoln, Nebraska, USA  
<sup>164</sup>State University of New York at Buffalo, Buffalo, New York, USA  
<sup>165</sup>Northeastern University, Boston, Massachusetts, USA  
<sup>166</sup>Northwestern University, Evanston, Illinois, USA  
<sup>167</sup>University of Notre Dame, Notre Dame, Indiana, USA  
<sup>168</sup>The Ohio State University, Columbus, Ohio, USA  
<sup>169</sup>Princeton University, Princeton, New Jersey, USA  
<sup>170</sup>University of Puerto Rico, Mayaguez, Puerto Rico, USA  
<sup>171</sup>Purdue University, West Lafayette, Indiana, USA  
<sup>172</sup>Purdue University Northwest, Hammond, Indiana, USA  
<sup>173</sup>Rice University, Houston, Texas, USA  
<sup>174</sup>University of Rochester, Rochester, New York, USA  
<sup>175</sup>The Rockefeller University, New York, New York, USA  
<sup>176</sup>Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA  
<sup>177</sup>University of Tennessee, Knoxville, Tennessee, USA  
<sup>178</sup>Texas A&M University, College Station, Texas, USA  
<sup>179</sup>Texas Tech University, Lubbock, Texas, USA  
<sup>180</sup>Vanderbilt University, Nashville, Tennessee, USA  
<sup>181</sup>University of Virginia, Charlottesville, Virginia, USA

<sup>182</sup>*Wayne State University, Detroit, Michigan, USA*<sup>183</sup>*University of Wisconsin - Madison, Madison, Wisconsin, USA*<sup>184</sup>*An institute or international laboratory covered by a cooperation agreement with CERN*<sup>a</sup>Deceased.<sup>b</sup>Also at Yerevan State University, Yerevan, Armenia.<sup>c</sup>Also at TU Wien, Vienna, Austria.<sup>d</sup>Also at Institute of Basic and Applied Sciences, Faculty of Engineering, Arab Academy for Science, Technology and Maritime Transport, Alexandria, Egypt.<sup>e</sup>Also at Ghent University, Ghent, Belgium.<sup>f</sup>Also at Universidade Estadual de Campinas, Campinas, Brazil.<sup>g</sup>Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil.<sup>h</sup>Also at UFMS, Nova Andradina, Brazil.<sup>i</sup>Also at Nanjing Normal University, Nanjing, China.<sup>j</sup>Also at Henan Normal University, Xinxiang, China.<sup>k</sup>Also at The University of Iowa, Iowa City, Iowa, USA.<sup>l</sup>Also at University of Chinese Academy of Sciences, Beijing, China.<sup>m</sup>Also at China Center of Advanced Science and Technology, Beijing, China.<sup>n</sup>Also at University of Chinese Academy of Sciences, Beijing, China.<sup>o</sup>Also at China Spallation Neutron Source, Guangdong, China.<sup>p</sup>Also at Université Libre de Bruxelles, Bruxelles, Belgium.<sup>q</sup>Also at Another institute or international laboratory covered by a cooperation agreement with CERN.<sup>r</sup>Also at Suez University, Suez, Egypt.<sup>s</sup>Also at British University in Egypt, Cairo, Egypt.<sup>t</sup>Also at Birla Institute of Technology, Mesra, Mesra, India.<sup>u</sup>Also at Purdue University, West Lafayette, Indiana, USA.<sup>v</sup>Also at Université de Haute Alsace, Mulhouse, France.<sup>w</sup>Also at Department of Physics, Tsinghua University, Beijing, China.<sup>x</sup>Also at The University of the State of Amazonas, Manaus, Brazil.<sup>y</sup>Also at Erzincan Binali Yildirim University, Erzincan, Turkey.<sup>z</sup>Also at University of Hamburg, Hamburg, Germany.<sup>aa</sup>Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany.<sup>bb</sup>Also at Isfahan University of Technology, Isfahan, Iran.<sup>cc</sup>Also at Bergische University Wuppertal (BUW), Wuppertal, Germany.<sup>dd</sup>Also at Brandenburg University of Technology, Cottbus, Germany.<sup>ee</sup>Also at Forschungszentrum Jülich, Juelich, Germany.<sup>ff</sup>Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.<sup>gg</sup>Also at Physics Department, Faculty of Science, Assiut University, Assiut, Egypt.<sup>hh</sup>Also at Wigner Research Centre for Physics, Budapest, Hungary.<sup>ii</sup>Also at Institute of Physics, University of Debrecen, Debrecen, Hungary.<sup>jj</sup>Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.<sup>kk</sup>Also at Universitatea Babes-Bolyai—Facultatea de Fizica, Cluj-Napoca, Romania.<sup>ll</sup>Also at Faculty of Informatics, University of Debrecen, Debrecen, Hungary.<sup>mm</sup>Also at Punjab Agricultural University, Ludhiana, India.<sup>nn</sup>Also at University of Hyderabad, Hyderabad, India.<sup>oo</sup>Also at University of Visva-Bharati, Santiniketan, India.<sup>pp</sup>Also at Indian Institute of Science (IISc), Bangalore, India.<sup>qq</sup>Also at IIT Bhubaneswar, Bhubaneswar, India.<sup>rr</sup>Also at Institute of Physics, Bhubaneswar, India.<sup>ss</sup>Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany.<sup>tt</sup>Also at Department of Physics, Isfahan University of Technology, Isfahan, Iran.<sup>uu</sup>Also at Sharif University of Technology, Tehran, Iran.<sup>vv</sup>Also at Department of Physics, University of Science and Technology of Mazandaran, Behshahr, Iran.<sup>ww</sup>Also at Helwan University, Cairo, Egypt.<sup>xx</sup>Also at Italian National Agency for New Technologies, Energy and Sustainable Economic Development, Bologna, Italy.<sup>yy</sup>Also at Centro Siciliano di Fisica Nucleare e di Struttura Della Materia, Catania, Italy.<sup>zz</sup>Also at Università degli Studi Guglielmo Marconi, Roma, Italy.<sup>aaa</sup>Also at Scuola Superiore Meridionale, Università di Napoli “Federico II”, Napoli, Italy.<sup>bbb</sup>Also at Fermi National Accelerator Laboratory, Batavia, Illinois, USA.<sup>ccc</sup>Also at Università di Napoli “Federico II”, Napoli, Italy.

- <sup>ddd</sup> Also at Ain Shams University, Cairo, Egypt.
- <sup>eee</sup> Also at Consiglio Nazionale delle Ricerche—Istituto Officina dei Materiali, Perugia, Italy.
- <sup>fff</sup> Also at Riga Technical University, Riga, Latvia.
- <sup>ggg</sup> Also at Department of Applied Physics, Faculty of Science and Technology, Universiti Kebangsaan Malaysia, Bangi, Malaysia.
- <sup>hhh</sup> Also at Consejo Nacional de Ciencia y Tecnología, Mexico City, Mexico.
- <sup>iii</sup> Also at Trincomalee Campus, Eastern University, Sri Lanka, Nilaveli, Sri Lanka.
- <sup>jjj</sup> Also at INFN Sezione di Pavia, Università di Pavia, Pavia, Italy.
- <sup>kkk</sup> Also at National and Kapodistrian University of Athens, Athens, Greece.
- <sup>lll</sup> Also at Ecole Polytechnique Fédérale Lausanne, Lausanne, Switzerland.
- <sup>mmm</sup> Also at University of Vienna Faculty of Computer Science, Vienna, Austria.
- <sup>nnn</sup> Also at Universität Zürich, Zurich, Switzerland.
- <sup>ooo</sup> Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria.
- <sup>ppp</sup> Also at Laboratoire d'Annecy-le-Vieux de Physique des Particules, IN2P3-CNRS, Annecy-le-Vieux, France.
- <sup>qqq</sup> Also at Near East University, Research Center of Experimental Health Science, Mersin, Turkey.
- <sup>rrr</sup> Also at Konya Technical University, Konya, Turkey.
- <sup>sss</sup> Also at Izmir Bakircay University, Izmir, Turkey.
- <sup>ttt</sup> Also at Adiyaman University, Adiyaman, Turkey.
- <sup>uuu</sup> Also at Bozok Üniversitesi Rektörlüğü, Yozgat, Turkey.
- <sup>vvv</sup> Also at Marmara University, Istanbul, Turkey.
- <sup>www</sup> Also at Milli Savunma University, Istanbul, Turkey.
- <sup>xxx</sup> Also at Kafkas University, Kars, Turkey.
- <sup>yyy</sup> Also at stanbul Okan University, Istanbul, Turkey.
- <sup>zzz</sup> Also at Hacettepe University, Ankara, Turkey.
- <sup>aaaa</sup> Also at Istanbul University—Cerrahpaşa, Faculty of Engineering, Istanbul, Turkey.
- <sup>bbbb</sup> Also at Yildiz Technical University, Istanbul, Turkey.
- <sup>cccc</sup> Also at Vrije Universiteit Brussel, Brussel, Belgium.
- <sup>dddd</sup> Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
- <sup>eeee</sup> Also at University of Bristol, Bristol, United Kingdom.
- <sup>ffff</sup> Also at IPPP Durham University, Durham, United Kingdom.
- <sup>gggg</sup> Also at Monash University, Faculty of Science, Clayton, Australia.
- <sup>hhhh</sup> Also at Università di Torino, Torino, Italy.
- <sup>iiii</sup> Also at Bethel University, St. Paul, Minnesota, USA.
- <sup>jjjj</sup> Also at Karamanoğlu Mehmetbey University, Karaman, Turkey.
- <sup>kkkk</sup> Also at California Institute of Technology, Pasadena, California, USA.
- <sup>llll</sup> Also at United States Naval Academy, Annapolis, Maryland, USA.
- <sup>mmmm</sup> Also at Bingöl University, Bingöl, Turkey.
- <sup>nnnn</sup> Also at Georgian Technical University, Tbilisi, Georgia.
- <sup>oooo</sup> Also at Sinop University, Sinop, Turkey.
- <sup>pppp</sup> Also at Erciyes University, Kayseri, Turkey.
- <sup>qqqq</sup> Also at Horia Hulubei National Institute of Physics and Nuclear Engineering (IFIN-HH), Bucharest, Romania.
- <sup>rrrr</sup> Also at Texas A&M University at Qatar, Doha, Qatar.
- <sup>ssss</sup> Also at Kyungpook National University, Daegu, Korea.
- <sup>tttt</sup> Also at Universiteit Antwerpen, Antwerpen, Belgium.
- <sup>uuuu</sup> Also at Yerevan Physics Institute, Yerevan, Armenia.
- <sup>vvvv</sup> Also at Northeastern University, Boston, Massachusetts, USA.
- <sup>wwww</sup> Also at Imperial College, London, United Kingdom.
- <sup>xxxx</sup> Also at Institute of Nuclear Physics of the Uzbekistan Academy of Sciences, Tashkent, Uzbekistan.