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Monitoring therapeutic proton beams with LGAD silicon detectors

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ABSTRACT: The University and the National Institute for Nuclear Physics of Torino are 15 developing LGAD-based prototypes for beam monitoring in proton therapy. The direct 16 measurement of single beam particles could overcome some features of currently used ionization 17 18 chambers, such as slow charge collection and reduced sensitivity, which limit the implementation of advanced delivery techniques (e.g. rescanning). LGAD strip sensors have been designed and 19 produced by Bruno Kessler Foundation (FBK, Trento) specifically for this project. A counter 20 prototype to directly count individual protons at clinical fluence rates $(10^6-10^{10} \text{ protons/cm}^2 \cdot \text{s})$ 21 and a telescope system to measure the beam energy with time-of-flight (TOF) techniques are 22 described. Tests of LGAD silicon strip sensors performed on synchrotron and cyclotron beams of 23 therapeutic centers, using a pin-hole ionization chamber for the independent measurement of the 24 particle flux, already showed the possibility to keep the counting error < 1 % up to a beam fluence 25 rate of few 10^8 protons/cm² s. The ongoing tests of counting sensors readout by a dedicated fast 26 Charge Sensitive Amplifier chip are reported. The telescope system, made of two sensors at a 27 distance up to 95 cm, allows measuring the beam energy in the clinical range (70-230 MeV) with 28 a maximum deviation of 310 keV in respect to the nominal one, with an uncertainty of 500 keV, 29 thus achieving the prescribed clinical accuracy of 1 mm in the range in water. 30

31 KEYWORDS: particle therapy, beam monitors, silicon detector

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45 **1. Intro**

Within the MoVeIT project, the University and the National Institute for Nuclear Physics of Torino are proposing the direct measurement of particles using LGAD-based innovative detectors for monitoring clinical proton beams [1,2]. This approach has the potential to overcome the limits of current beam monitors, such as slow charge collection times, thus supporting the future implementation of advanced and fast delivery techniques (e.g. fast rescanning modalities) in particle therapy. This contribution reports about the status of the project.

52 2. Materials and Methods

53 2.1 Thin LGAD silicon sensors

54 Ultra Fast Silicon Detectors (UFSD) are based on Low Gain Avalanche Detectors (LGAD) design, characterized by small thicknesses (typically ~50 µm) and controlled low gain (~10-30), 55 56 providing an enhanced signal in thin detectors with the same noise level of traditional silicon 57 sensors of similar geometry, with very short time duration (ns), allowing particle counting, and excellent time resolution (ps), allowing beam energy measurement through TOF techniques [3]. 58 59 Thin planar UFSD prototypes were specifically designed and produced for the MoVeIT project by Bruno Kessler Foundation (FBK, Trento, Italy). Strip segmentation has been chosen for all 60 structures (reported in Table 1) in order to reduce the expected particle rate per channel and the 61 Table 1. Characteristics of the MoVeIT sensors for counting and timing purposes.

	Number of strips	Strip length	Strip area	Active thickness
Counting sensor (short)	20	1.5 cm	2.2 mm^2	60 µm
Counting sensor (long)	30	3.0 cm	2.4 mm^2	60 µm
Counting sensor (final)	144	2.6 cm	3.0 mm^2	50 µm
Timing sensor	11	4.0 mm	2.2 mm^2	55 µm

sensor capacitance. The final counting sensor, whose sensitive area ($\sim 2.7x2.7 \text{ cm}^2$) was defined to cover the clinical proton beam cross section characterized by a FHWM of ~ 1 cm at the isocentre, is described in more details in the contribution of [4].

65 2.2 Counting

A dedicated fast Charge Sensitive Amplifier chip (ABACUS [5], Fig. 1a and 1b) with an area of $2x5 \text{ mm}^2$ and 24 channels has been designed to discriminate the expected signal pulses in a wide charge range (3-150 fC, corresponding to the energy released in silicon by clinical protons of 60-230 MeV) with a maximum dead-time of 10 ns to minimize pile-up counting inefficiencies at clinical fluence rates (10^6 - 10^{10} protons/cm² s)

70 at clinical fluence rates $(10^6-10^{10} \text{ protons/cm}^2 \cdot \text{s})$.



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Figure 1. a. Test-board mounting a counting long strip readout by an Abacus chip, zoomed in **b**. **c.** Telescope system prototype for TOF measurements.

The outputs of the chip are sent to an FPGA for the pulse counting, and the FPGA is connected to a computer where a dedicated LabVIEW program controls and integrates the FPGA input-outputs, displays online the counting rate from each strip and stores the useful data. The same FPGA is also used to initialize the chip by setting the local DACs of the chip channels.

76 **2.3 Timing**

77 Timing sensors (Table 1), thinned down to a total thickness of 70 μ m to minimize the beam perturbation, are readout by dedicated boards with two stages of amplification optimized for time 78 79 measurements with large signals (3-150 fC charge dynamic range) at high fluxes ($10^9 \text{ p/cm}^2 \cdot \text{s}$). 80 Signals generated by incoming protons are first amplified by the front-end electronics and then readout by a 16+1 channels digitizer (5 GS/s, 12 bits resolution) controlled by a PC with an 80 81 MB/s optical link. A prototype of telescope made of two timing sensors aligned along the beam 82 direction has been built (Fig. 1c). The first sensor is kept fixed at the isocenter of the beam 83 84 distribution system, whereas the second sensor can be moved in the transversal plane (x-y) for the alignment of the telescope with the beam trajectory, and along the longitudinal direction (z axis) 85 to vary the distance between the two sensors from 300 mm to 950 mm. An optical encoder (0.01 86 87 μ m resolution) provides the measurement of the position displacements of the second sensor. Starting from the measured time differences of coincident protons crossing the two detectors for 88 different beam energies and distances between the sensors, the mean kinetic energy of the beam 89 90 can be determined and finally the range in water, which is the clinically relevant parameter [6]. A 91 self-calibration method, where only the measured mean time differences and the relative displacements are used as input parameters, has been implemented to remove the systematic 92

errors due to the experimental setup, independently from any a priori knowledge of the beam
 parameters (the method is object of the Italian Patent Application No. 102021000025190).

95 3. Results



Figure 2. *Left.* Signal amplitude distribution obtained from 1 chip channel fitted with convolution of Landau and Gaussian (red curves) for 4 beam energies. *Right.* Most Probable Values (MPV) vs energies.

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97 Tests of LGAD silicon sensors performed on synchrotron and cyclotron beams of therapeutic 98 centers, using a pin-hole ionization chamber for the independent measurement of the particle flux, 99 showed the possibility to keep the counting error < 1% up to a beam fluence rates of few 10^8 100 $p/cm^2 \cdot s$ [2]. Algorithms based on logical combinations of signals from two independent detector channels under the same radiation field are being developed for count-loss correction at fluence 101 rates up to 10^{10} p/cm² s [7]. To characterize the counting prototype (sensor + asic), threshold scans 102 (rate as a function of the threshold voltage set in the leading-edge discriminator of each channel) 103 have been acquired with different proton beam energies (70-228 MeV) at the Trento Proton 104 Therapy Center. The signal amplitude distribution of one channel, obtained through the discrete 105 106 derivative of the threshold scan, is shown in Fig. 2. The Most Probable Values, obtained by fitting the amplitude distributions with the convolution of a Landau and a Gaussian, vs the beam energy 107



Figure 3. Time difference vs distance for two distances (600 and 950 mm) between the sensors of the telescope and 7 beam energies. The legend shows the nominal energy (E_N , retrieved from the Bragg peak positions measured by CNAO using a Peakfinder water column with a maximum deviation within ±0.15 mm) and the measured energy (E_M).

reproduce the mean energy loss per unit path length described by the Bethe Bloch formula, asexpected (Fig. 2).

For TOF measurements, the self-calibration method allowed defining the time offsets 110 between different combinations of strips in sensor 1 and sensor 2. These time-offsets were used 111 112 to correct the time-of-arrival differences of protons crossing sensor 1 and sensor 2 of the telescope, for different beam energies and 2 distances between the sensors (Fig. 3) at CNAO 113 (Pavia). A maximum deviation of 310 keV between nominal and measured energies, obtained 114 115 from the linear interpolation of time differences vs distance, have been observed with a uncertainty of 0.5 MeV on the measured energy, thus achieving the prescribed clinical accuracy 116 117 of 1 mm in the range in water (Fig. 3).

118 **4. Conclusion**

This contribution reports about the status of the development of two prototypes, based on the UFSD technology, for online monitoring of clinical proton beams: the first one to directly count individual protons and the second to measure the beam energy with time-of-flight (TOF) techniques. The final counting prototype, featuring a 2.7x2.7 cm² sensor with 144 strips and its custom readout will be tested in the next months, while the promising results of the telescope system suggests its future use in beam quality control procedures.

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