TRANSPORT SEMIGROUP ASSOCIATED TO POSITIVE BOUNDARY CONDITIONS OF UNIT NORM: A DYSON-PHILLIPS APPROACH

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ABSTRACT. We revisit our study of general transport operator with general force field and general invariant measure by considering, in the L^1 setting, the linear transport operator $\mathcal{T}_H$ associated to a linear and positive boundary operator H of unit norm. It is known that in this case an extension of $\mathcal{T}_H$ generates a substochastic (i.e. positive contraction) C_0-semigroup $(V_H(t))_{t \geqslant 0}$. We show here that $(V_H(t))_{t \geqslant 0}$ is the smallest substochastic C_0-semigroup with the above mentioned property and provides a representation of $(V_H(t))_{t \geqslant 0}$ as the sum of an expansion series similar to Dyson-Phillips series. We develop an honesty theory for such boundary perturbations that allows to consider the honesty of trajectories on subintervals $J \subseteq [0, \infty)$. New necessary and sufficient conditions for a trajectory to be honest are given in terms of the aforementioned series expansion.

1. **Introduction.** We investigate here the well-posedness (in the sense of semigroup theory) in $L^1(\mathbf{\Omega}, \mathrm{d}\mu)$ of the general transport equation

$$\partial_t f(\mathbf{x}, t) + \mathscr{F}(\mathbf{x}) \cdot \nabla_{\mathbf{x}} f(\mathbf{x}, t) = 0 \qquad (\mathbf{x} \in \mathbf{\Omega}, \ t > 0), \tag{1.1a}$$

supplemented by the abstract boundary condition

$$f_{|\Gamma_-}(\mathbf{y}, t) = H(f_{|\Gamma_+})(\mathbf{y}, t), \qquad (\mathbf{y} \in \Gamma_-, t > 0), \tag{1.1b}$$

and the initial condition

$$f(\mathbf{x}, 0) = f_0(\mathbf{x}), \qquad (\mathbf{x} \in \mathbf{\Omega}). \tag{1.1c}$$

Here $\mathbf{\Omega}$ is a sufficiently smooth open subset of $\mathbb{R}^N$ endowed with a positive Radon measure μ, $\Gamma_{\pm}$ are suitable boundaries of the phase space and the field $\mathscr{F}$ is globally Lipschitz and divergence free with respect to μ, in the sense that μ is a measure invariant by the (globally defined) flow associated to $\mathscr{F}$. Our main concern here is

the influence of the boundary conditions (1.1b) and we treat here the delicate case of a boundary operator

$$H \,:\, L^1_+ \to L^1_-$$

which is linear, positive, bounded ($L^1_\pm$ being suitable trace spaces corresponding to the boundaries $\Gamma_\pm$, see Section 2 for details) and of unit norm

$$\|H\|_{\mathscr{B}(L^1_+, L^1_-)} = \sup_{f \in L^1_+, \|f\|_{L^1_+} = 1} \|Hf\|_{L^1_-} = 1. \tag{1.2}$$

Our motivation for studying such a problem is the study of kinetic equation of Vlasov-type for which the phase space $\mathbf{\Omega}$ is a cylindrical domain $\mathbf{\Omega} = D \times \mathbb{R}^N \subset \mathbb{R}^{2N}$ (D being a sufficiently smooth open subset of $\mathbb{R}^N$) and the field $\mathscr{F}$ is given by

$$\mathscr{F}(\mathbf{x}) = (v, \mathbf{F}(x, v)) \quad \text{for any} \quad \mathbf{x} = (x, v) \in \mathbf{\Omega} \tag{1.3}$$

where $\mathbf{F} : \mathbf{\Omega} \to \mathbb{R}^N$ is a time independent force field. The simplest (but already very rich) example of such a kinetic equation is the so-called *free-streaming* equation for which $\mathbf{F} = 0$. Boundary conditions in such kinetic equations are usually modeled by a boundary operator H which relates the incoming and outgoing boundary fluxes of particles; the form of this operator depends on the gas-surface interaction (see [13] for more details on such a topic).

The mathematical study of the aforementioned problem has already a long story starting from the seminal paper [10] who considered the case in which μ is the Lebesgue measure and the so-called 'no re-entry' boundary conditions (i.e. $H = 0$ in (1.1b)). More general fields and boundary conditions (but still mostly associated with the Lebesgue measure) have been considered in [11]. The free-streaming case (i.e. $\mathscr{F}(x, v) = (v, 0)$) received much more attention, starting from [20], where the free streaming transport operator associated to different boundary operators H is deeply investigated (see also [17] for general boundary conditions). Recently, transport operators associated to general external fields and general measures, with general bounded boundary conditions have been thoroughly investigated by the authors in collaboration with J. Banasiak in a series of papers [1, 2, 3] that contain both a generalization of the theory developed in the free streaming case and some new results. Summarizing the known results on this topic, one can say that the transport operator associated to H, which shall be denoted by $\mathcal{T}_H$ (see Section 2 for a precise definition), is the generator of a strongly continuous semigroup when the boundary operator H is a contraction (and also for some very peculiar multiplying boundary conditions, [20, 17, 2, 3, 12]).

A very interesting and important case, both from the mathematical and physical point of view, arises whenever H is a *positive boundary operator of unit norm* (see (1.2)). In such a case, one can not state *a priori* that $\mathcal{T}_H$ generates a C_0-semigroup in $L^1(\mathbf{\Omega}, \mathrm{d}\mu)$. Nevertheless, since for each $r \in [0, 1)$ the operator $H_r := rH$ is a strict contraction, the transport operator $\mathcal{T}_{H_r}$ associated to H_r does generate a C_0-semigroup $(V_r(t))_{t \geqslant 0}$. These semigroups are *substochastic, i. e. they are positive contraction semigroups* and one can show that the strong limit $\lim_{r \nearrow 1-} V_r(t) := V_H(t)$ exists and defines a C_0-semigroup in $L^1(\mathbf{\Omega}, \mathrm{d}\mu)$. Its generator $\mathcal{A}$ is then an extension of $\mathcal{T}_H$ and a natural question is to recognize if $\mathcal{A} = \overline{\mathcal{T}_H}$ or not. For conservative conditions, i.e. if

$$\|Hf\|_{L^1_-} = \|f\|_{L^1_+} \qquad \forall f \in L^1_+ \ f \geqslant 0$$

it is known that the semigroup $(V_H(t))_{t \geqslant 0}$ is conservative if and only if $\mathcal{A} = \overline{\mathcal{T}_H}$. On the contrary, whenever $\mathcal{A} \supsetneq \overline{\mathcal{T}_H}$ a mass loss occurs, i.e. there exists nonnegative f such that $\|V_H(t)f\| < \|f\|$ for some $t > 0$.

As first observed in [5], such a problem is very similar to what occurs in the so-called substochastic theory of additive perturbations of semigroups, (see the monograph [9]), where one is faced with the following problem: let $(T, D(T))$ be the generator of a substochastic C_0-semigroup $(G_T(t))_{t \geqslant 0}$ in $X = L^1(\Sigma, \mathrm{d}\nu)$ (where (Σ, ν) is a given measure space) and let $(B, \mathscr{D}(B))$ be a non-negative linear operator in X such that $\mathscr{D}(T) \subseteq \mathscr{D}(B)$ and $\int_\Sigma (T + B) f \mathrm{d}\nu \leqslant 0$ for all $f \in \mathscr{D}(T)_+ = \mathscr{D}(T) \cap X_+$. Then for any $0 < r < 1$ operator $(T + rB, \mathscr{D}(T))$ generates a C_0- semigroup $(G_r(t))_{t \geqslant 0}$. These semigroups are such that the strong limit $\lim_{r \nearrow 1-} G_r(t) := G_K(t)$ exists and the family $(G_K(t))_{t \geqslant 0}$ is a C_0-semigroup generated by an extension K of $T + B$. In the context of additive perturbations of substochastic semigroups a complete characterization of K is given; it is shown that $(G_K(t))_{t \geqslant 0}$ is the smallest (in the lattice sense) C_0-semigroup generated by an extension of $T + B$. Moreover $G_K(t)$ can be written as the sum of a strongly convergent series of linear positive operators (Dyson-Phillips expansion series) and a satisfying honesty theory, dealing with the mass carried by individual trajectories, has been developed [9, 6, 19]. Such a honesty theory for additive perturbation has been based mainly on the so-called *resolvent approach* (i.e. on the study of the resolvent of $(\lambda - K)^{-1}$) and such a resolvent approach has been applied to the boundary perturbation case in [5, 18]. Recently a new approach to honesty has been proposed, based now on the semigroup approach and the fine properties of the Dyson-Phillips iterated [6]. Such an approach is equivalent to the resolvent one but its main interest lies in the fact that it is robust enough to be applied to other kind of problems in which the resolvent approach is more difficult to handle (e.g. for non-autonomous families [7] or integrated semigroups [8]).

In the present paper we want to recognize that a fully similar study can be carried out for the operator $\mathcal{T}_H$. Notice that several results concerning the transport operator $\mathcal{T}_H$ and the semigroup $(V_H(t))_{t \geqslant 0}$ are already available in the literature. A complete description of $\mathcal{A}$ is given in [2] where it is shown that $\mathcal{A}$ is an extension of $\mathcal{T}_H$; the study of conservative boundary conditions has been performed, in the free-streaming case, in [5] and, for general force fields, in [2]. The general case of boundary operators with unit norm has been handled with in [18] where a detailed *honesty theory* has been performed. Nevertheless the obtained results are not so satisfying as those obtained in the substochastic theory of additive perturbations of semigroups. In particular the question of whether $(V_H(t))_{t \geqslant 0}$ is the *smallest* substochastic C_0-semigroup generated by an extension of $\mathcal{T}_H$ remains open and the honesty theory performed in [2, 18] is based on the resolvent approach only.

The objective of the present paper is to fill this blank. In particular, the main novelty of the paper lies in the following:

i) First, we prove that indeed the semigroup $(V_H(t))_{t \geqslant 0}$ is the smallest (in the lattice sense) substochastic C_0-semigroup generated by an extension of $\mathcal{T}_H$. This is the object of Theorem 3.3.

ii) Second, and more important, we develop a 'semigroup approach' to the honesty theory of boundary perturbations, exploiting the recent results in [3] which allow to provide a characterization of the semigroup $(V_H(t))_{t \geqslant 0}$ as an expansion series, similar to the Dyson-Phillips arising in the additive perturbation case

(see Theorem 3.5). While the resolvent approach allows to establish necessary and sufficient conditions for a trajectory to be honest (i. e. honest on $[0, \infty)$) the new semigroup approach allows to establish more general necessary and sufficient conditions for a trajectory to be honest on a subinterval $J \subseteq [0, \infty)$ (see Theorem 4.3). Such honesty of trajectory on subinterval is reminiscent from [19] and allows to get a more precise picture of the eventual "mass loss" in the conservative case. We strongly believe that such a semigroup approach has its own interest and that, as it occurs for additive perturbation [7], it could hopefully be extended to deal with non-autonomous problems.

To be more precise, the contents of the paper are as follows. In Section 2 we introduce the necessary notation and define the transport operator $\mathcal{T}_H$. This section is mainly taken from the recent contributions [2, 3]. In Section 3 we establish the most important properties of the semigroup $(V_H(t))_{t \geqslant 0}$ and its generator, in particular showing that $(V_H(t))_{t \geqslant 0}$ is the smallest substochastic C_0-semigroup generated by an extension of $\mathcal{T}_H$. In Section 4 we develop the honesty theory for boundary perturbations, introducing first useful functionals and defining then the concept of honesty of trajectories on subintervals $J \subseteq [0, \infty)$. We obtain also necessary and sufficient conditions for the honesty in the spirit of [2, Section 6] and [18] not only using the usual resolvent approach but also using the series approach introduced in [6]. In Section 6 two well-known examples are revisited using our new approach, that allows us to deduce new interesting properties.

2. **Preliminaries.** In the present section, we introduce the general mathematical framework we shall consider in the sequel. The material from this section is mainly taken from [1, 2] and we refer to these two contributions for further properties of abstract transport operators. We begin with the rigorous definition of the transport operator $\mathcal{T}_H$ associated to bounded boundary operator H.

2.1. **Definition of the transport operator $\mathcal{T}_H$.** In this paper we consider transport operators associated to general external fields and general measures, according to the theory developed in two recent contributions [1],[2]. More precisely, given a smooth open subset $\mathbf{\Omega}$ of $\mathbb{R}^N$, we consider a time independent globally Lipschitz vector field $\mathscr{F} : \mathbb{R}^N \to \mathbb{R}^N$ so that, for any $\mathbf{x} \in \mathbf{\Omega}$, the Cauchy problem

$$\frac{\mathrm{d}\mathbf{X}}{\mathrm{d}t}(t) = \mathscr{F}(\mathbf{X}(t)), \qquad \forall t \in \mathbb{R}; \qquad \mathbf{X}(0) = \mathbf{x} \in \mathbf{\Omega} \tag{2.1}$$

admits a unique global solution

$$(\mathbf{x}, t) \in \mathbf{\Omega} \times \mathbb{R} \longmapsto \Phi(\mathbf{x}, t) \in \mathbb{R}^N,$$

that allows to define a flow $(T_t)_{t \in \mathbb{R}}$ given by $T_t = \Phi(\cdot, t)$. As in [1], we assume that there exists a Radon measure μ over $\mathbb{R}^N$ which is invariant under the flow $(T_t)_{t \in \mathbb{R}}$, i.e.

$$\mu(T_t A) = \mu(A) \text{ for any measurable subset } A \subset \mathbb{R}^N \text{ and any } t \in \mathbb{R}. \tag{2.2}$$

Of course, solutions to (2.1) do not necessarily belong to $\mathbf{\Omega}$ for all times, leading to the definition of stay times of the characteristic curves in $\mathbf{\Omega}$: for any $\mathbf{x} \in \mathbf{\Omega}$, define

$$\tau_{\pm}(\mathbf{x}) = \inf\{s > 0 \,; \Phi(\mathbf{x}, \pm s) \notin \mathbf{\Omega}\}, \tag{2.3}$$

with the convention that inf $\varnothing = \infty$. This allows to represent $\boldsymbol{\Omega}$ as $\boldsymbol{\Omega} = \boldsymbol{\Omega}_\pm \cup \boldsymbol{\Omega}_{\pm\infty}$ where

$$\boldsymbol{\Omega}_\pm = \{\mathbf{x} \in \boldsymbol{\Omega}\,;\, \tau_\pm(\mathbf{x}) < \infty\}, \qquad \text{and} \quad \boldsymbol{\Omega}_{\pm\infty} = \{\mathbf{x} \in \boldsymbol{\Omega}\,;\, \tau_\pm(\mathbf{x}) = \infty\}.$$

Moreover, we define the *incoming* and *outgoing* boundaries as

$$\Gamma_\pm := \{\mathbf{y} \in \partial\boldsymbol{\Omega}\,;\, \exists \mathbf{x} \in \boldsymbol{\Omega},\, \tau_\pm(\mathbf{x}) < \infty \text{ and } \mathbf{y} = \Phi(\mathbf{x}, \pm\tau_\pm(\mathbf{x}))\,\}. \qquad (2.4)$$

The definition of the stay time $\tau_\pm$ extends then to $\Gamma_\pm$ by setting simply $\tau_\pm(\mathbf{y}) = 0$ and $\tau_\mp(\mathbf{y}) = \tau_+(\mathbf{x}) + \tau_-(\mathbf{x})$ for any $\mathbf{y} \in \Gamma_\pm$ with $\mathbf{y} = \Phi(\mathbf{x}, \pm\tau_\pm(\mathbf{x}))$. Notice that, with the above definition, $\tau_\mp(\mathbf{y})$ is well defined (i.e. the definition does not depend on the choice of $\mathbf{x} \in \boldsymbol{\Omega}_\pm$) and $\tau_\mp(\mathbf{y})$ is nothing but the length of the characteristic curves having $\mathbf{y}$ as its left (respectively right) end-point. We finally set

$$\Gamma_{\pm\infty} = \{\mathbf{y} \in \Gamma_\pm\,;\, \tau_\mp(\mathbf{y}) = \infty\}.$$

With such notations, one can prove (see [1, Section 2]) the existence of unique positive Borel measures $\mu_\pm$ on $\Gamma_\pm$ such that the measure μ on $\boldsymbol{\Omega}_\pm$ is identified to the product measure of $\mu_\pm$ with the Lebesgue measure on $\mathbb{R}$ (see [1, Proposition 2.10]). The construction of such measures allow to define the trace spaces

$$L^1_\pm = L^1(\Gamma_\pm, \mathrm{d}\mu_\pm)$$

with the usual norm. In the Banach space

$$X := L^1(\boldsymbol{\Omega}, \mathrm{d}\mu)$$

endowed with its usual norm, we can define the maximal *transport operator* as follows (see [1, Theorem 3.6])

Definition 2.1. Given $f \in L^1(\boldsymbol{\Omega}, \mathrm{d}\mu)$, f belongs to the domain $\mathscr{D}(\mathcal{T}_{\max})$ of $\mathcal{T}_{\max}$ if and only if there exists $g \in L^1(\boldsymbol{\Omega}, \mathrm{d}\mu)$ and a *representative* $f^\sharp$ of f (i.e. $f^\sharp(x) = f(x)$ for μ-a.e. $x \in \boldsymbol{\Omega}$) such that, for μ-almost every $\mathbf{x} \in \boldsymbol{\Omega}$ and any $-\tau_-(\mathbf{x}) < t_1 \leqslant t_2 < \tau_+(\mathbf{x})$ one has

$$f^\sharp(\Phi(\mathbf{x}, t_1)) - f^\sharp(\Phi(\mathbf{x}, t_2)) = \int_{t_1}^{t_2} g(\Phi(\mathbf{x}, s))\mathrm{d}s. \qquad (2.5)$$

In this case, we set $\mathcal{T}_{\max} f = g$.

Remark 1. Notice that the above operator $\mathcal{T}_{\max}$ is *well-defined*, i.e. $\mathcal{T}_{\max} f$ does not depend on the representative $f^\sharp$. Finally, we wish to emphasize the fact that the domain $\mathscr{D}(\mathcal{T}_{\max})$ is precisely the set of functions $f \in L^1(\boldsymbol{\Omega}, \mathrm{d}\mu)$ that admit a representative which is absolutely continuous along almost any characteristic curve.

With the above definition, each function $f \in \mathscr{D}(\mathcal{T}_{\max})$ is such that the limits

$$\mathsf{B}^+ f(\mathbf{y}) := \lim_{s \to 0+} f^\sharp(\Phi(\mathbf{y}, -s)) \qquad \text{and} \qquad \mathsf{B}^- f(\mathbf{y}) := \lim_{s \to 0+} f^\sharp(\Phi(\mathbf{y}, s))$$

exist for almost $\mu_\pm$-every $\mathbf{y} \in \Gamma_\pm$ [1, Proposition 3.16, Definition 3.17]. Notice that the *trace* $\mathsf{B}^\pm f$ of a given $f \in \mathscr{D}(\mathcal{T}_{\max})$ does not necessarily belong to $L^1_\pm$. Nevertheless one can prove the following [2, Theorem 3.1, Proposition 3.2, Corollary 2.1] :

Theorem 2.2. *Define the following measures over* $\Gamma_\pm$:

$$\mathrm{d}\xi_\pm(\mathbf{y}) = \min\left(\tau_\mp(\mathbf{y}), 1\right)\mathrm{d}\mu_\pm(\mathbf{y}), \qquad \mathbf{y} \in \Gamma_\pm.$$

Then, for any $f \in \mathscr{D}(\mathcal{T}_{\max})$, *the trace* $\mathsf{B}^\pm f$ *belongs to* $Y_\pm := L^1(\Gamma_\pm, \mathrm{d}\xi_\pm)$ *with*

$$\|\mathsf{B}^\pm f\|_{Y_\pm} \leqslant \|f\|_X + \|\mathcal{T}_{\max} f\|_X, \qquad f \in \mathscr{D}(\mathcal{T}_{\max}).$$

Moreover

$$\mathscr{W} := \left\{f \in \mathscr{D}(\mathcal{T}_{\max})\,;\,\mathsf{B}^- f \in L^1_-\right\} = \left\{f \in \mathscr{D}(\mathcal{T}_{\max})\,;\,\mathsf{B}^+ f \in L^1_+\right\}$$

and the Green formula

$$\int_{\boldsymbol{\Omega}} \mathcal{T}_{\max} f \mathrm{d}\mu = \int_{\Gamma_-} \mathsf{B}^- f \mathrm{d}\mu_- - \int_{\Gamma_+} \mathsf{B}^+ f\,\mathrm{d}\mu_+. \qquad (2.6)$$

holds for any $f \in \mathscr{W}$.

We are then in position to define the transport operator associated to a bounded boundary operator as follows:

Definition 2.3. For any bounded linear operator

$$H \in \mathscr{B}(L^1_+, L^1_-)$$

we define *the transport operator $(\mathcal{T}_H, \mathscr{D}(\mathcal{T}_H))$ associated to the boundary condition H* as:

$$\mathscr{D}(\mathcal{T}_H) = \{f \in \mathscr{D}(\mathcal{T}_{\max})\,;\,\mathsf{B}^+ f \in L^1_+ \text{ and } \mathsf{B}^- f = H\mathsf{B}^+ f\},$$
$$\mathcal{T}_H f = \mathcal{T}_{\max} f \quad \forall f \in \mathscr{D}(\mathcal{T}_H). \qquad (2.7)$$

2.2. **Construction of the semigroup associated to boundary operator with unit norm.** We begin by introducing several notations, taken from [2]. For any $\lambda > 0$ one defines the following operators

$$\begin{cases} M_\lambda\,: & Y_- \longrightarrow Y_+ \\ & u \longmapsto [M_\lambda u]\,(\mathbf{y}) = u(\Phi(\mathbf{y}, -\tau_-(\mathbf{y}))) \exp\left(-\lambda\tau_-(\mathbf{y})\right) \chi_{\{\tau_-(\mathbf{y})<\infty\}}, \quad \mathbf{y} \in \Gamma_+ \,; \end{cases}$$

$$\begin{cases} \Xi_\lambda\,: & Y_- \longrightarrow X \\ & u \longmapsto [\Xi_\lambda u]\,(\mathbf{x}) = u(\Phi(\mathbf{x}, -\tau_-(\mathbf{x}))) \exp\left(-\lambda\tau_-(\mathbf{x})\right) \chi_{\{\tau_-(\mathbf{x})<\infty\}}, \quad \mathbf{x} \in \boldsymbol{\Omega} \,; \end{cases}$$

$$\begin{cases} G_\lambda\,: & X \longrightarrow L^1_+ \\ & f \longmapsto [G_\lambda f]\,(\mathbf{z}) = \int_0^{\tau_-(\mathbf{z})} f(\Phi(\mathbf{z}, -s)) \exp(-\lambda s)\mathrm{d}s, \quad \mathbf{z} \in \Gamma_+ \,; \end{cases}$$

and

$$\begin{cases} C_\lambda\,: & X \longrightarrow X \\ & f \longmapsto [C_\lambda f]\,(\mathbf{x}) = \int_0^{\tau_-(\mathbf{x})} f(\Phi(\mathbf{x}, -s)) \exp(-\lambda s)\mathrm{d}s, \quad \mathbf{x} \in \boldsymbol{\Omega} \end{cases}$$

where χ_A denotes the characteristic function of a set A. One has the following where $\mathcal{T}_0$ denotes the transport operator associated to the boundary operator $H \equiv 0$:

Lemma 2.4. *For any $\lambda > 0$, the following hold:*

1. *$M_\lambda \in \mathscr{B}(Y_-, Y_+)$. Moreover, given $u \in Y_-$, $M_\lambda u \in L^1_+$ if and only if $u \in L^1_-$.*
2. *$\Xi_\lambda \in \mathscr{B}(Y_-, X)$. Moreover, the range of Ξ_λ is a subset of $\mathscr{D}(\mathcal{T}_{\max})$ with*

$$\mathcal{T}_{\max}\Xi_\lambda u = \lambda\Xi_\lambda u, \qquad \mathsf{B}^-\Xi_\lambda u = u,\ \mathsf{B}^+\Xi_\lambda u = M_\lambda u, \qquad \forall u \in Y_- \qquad (2.8)$$

3. *$G_\lambda \in \mathscr{B}(X, L^1_+)$. Moreover, G_λ is surjective.*
4. *$C_\lambda \in \mathscr{B}(X)$ with range included in $\mathscr{D}(\mathcal{T}_0)$. Moreover, $C_\lambda = (\lambda - \mathcal{T}_0)^{-1}$ and*

$$G_\lambda f = \mathsf{B}^+ C_\lambda f \quad \text{ for any } f \in X.$$

Notice that, if $H \equiv 0$, it is not difficult to check that $(\mathcal{T}_0, \mathscr{D}(\mathcal{T}_0))$ is the generator of a C_0-semigroup $(U_0(t))_{t \geqslant 0}$ given by

$$U_0(t)f(\mathbf{x}) = f(\Phi(\mathbf{x}, -t))\chi_{\{t < \tau_-(\mathbf{x})\}}(\mathbf{x}), \qquad (\mathbf{x} \in \mathbf{\Omega}, \, f \in X). \qquad (2.9)$$

In all the sequel, we shall assume that $H : L^1_+ \to L^1_-$ is a positive boundary operator of unit norm, i.e.

$$H \in \mathscr{B}(L^1_+, L^1_-) \, ; \, Hf \geqslant 0 \quad \forall f \in L^1_+, f \geqslant 0 \, ;$$
$$\|H\|_{\mathscr{B}(L^1_+, L^1_-)} = \sup_{\|f\|_{L^1_+} = 1} \|Hf\|_{L^1_-} = 1. \quad (2.10)$$

Under such an assumption, for any $0 \leqslant r < 1$, the boundary operator $H_r := rH$ is dissipative with

$$\|H_r\|_{\mathscr{B}(L^1_+, L^1_-)} = r < 1;$$

it is then well-known [2, Corollary 4.1] that the transport operator $\mathcal{T}_{H_r}$ generates a positive contraction semigroup $(V_r(t))_{t \geqslant 0}$ for any $0 \leqslant r < 1$. From [2, Theorem 6.2], one has the following:

Theorem 2.5. *Let H satisfy Assumption 2.10. Then, for any $t \geqslant 0$ and any $f \in X$ the limit $V_H(t)f = \lim_{r \nearrow 1} V_r(t)f$ exists in X and defines a substochastic semigroup $(V_H(t))_{t \geqslant 0}$. If $(\mathcal{A}, \mathscr{D}(\mathcal{A}))$ is the generator of $(V_H(t))_{t \geqslant 0}$, then its resolvent is given by*

$$(\lambda - \mathcal{A})^{-1}f = \lim_{r \nearrow 1}(\lambda - \mathcal{T}_{H_r})^{-1}f = C_\lambda f + \sum_{n=0}^{\infty} \Xi_\lambda H(M_\lambda H)^n G_\lambda f$$
$$\text{for any } f \in X, \, \lambda > 0, \quad (2.11)$$

where the series converges in X. Moreover, $\mathcal{A}$ is an extension of $\mathcal{T}_H$; more precisely

$$\mathscr{D}(\mathcal{T}_H) \subset \mathscr{D}(\mathcal{A}) \subset \mathscr{D}(\mathcal{T}_{\max}) \quad \text{with} \quad \mathcal{A}f = \mathcal{T}_{\max}f \qquad \forall f \in \mathscr{D}(\mathcal{A})$$

and

$$\mathscr{D}(\mathcal{T}_H) = \left\{\varphi \in \mathscr{D}(\mathcal{A}) \, ; \, \mathsf{B}^+\varphi \in L^1_+\right\} = \left\{\varphi \in \mathscr{D}(\mathcal{A}) \, ; \, \mathsf{B}^-\varphi \in L^1_-\right\} = \mathscr{D}(\mathcal{A}) \cap \mathscr{W}$$

3. **A new characterization of** $(V_H(t))_{t \geqslant 0}$**.** In this section, we present a new characterization as well as practical expression of the semigroup $(V_H(t))_{t \geqslant 0}$. Indeed in the following Theorem 3.3 we are able to prove that $(V_H(t))_{t \geqslant 0}$ is the smallest substochastic C_0-semigroup generated by an extension of $\mathcal{T}_H$, while in Theorem 3.5 we show that $(V_H(t))_{t \geqslant 0}$ can be written as the sum of a strongly convergent series. We first need to recall the definition of transport operator associated to an unbounded boundary operator. Precisely, let us introduce $\mathscr{E}$ as the space of elements $(\psi_+, \psi_-) \in Y_+ \times Y_-$ such that $\psi_+ - M_\lambda \psi_- \in L^1(\Gamma_+, \mathrm{d}\mu_+)$ for some/all $\lambda > 0$. We equip $\mathscr{E}$ with the norm

$$\|(\psi_+, \psi_-)\|_{\mathscr{E}} := \|\psi_+\|_{Y_+} + \|\psi_-\|_{Y_-} + \|\psi_+ - M_1\psi_-\|_{L^1_+}$$

that makes it a Banach space. Then, one has the following generalization of Definition 2.3:

Definition 3.1. Given a possibly unbounded operator $\mathcal{K}$ from Y_+ to Y_-, we denote by $\mathscr{D}(\mathcal{K})$ its domain and $\mathscr{G}(\mathcal{K})$ its graph. If $\mathscr{G}(\mathcal{K}) \subset \mathscr{E}$ we can define the transport

operator $\mathcal{T_K}$ associated to the boundary operator $\mathcal{K}$ by $\mathcal{T_K}f = \mathcal{T}_{\max}f$ for any $f \in \mathscr{D}(\mathcal{T_K})$, where

$$\mathscr{D}(\mathcal{T_K}) = \left\{ f \in \mathscr{D}(\mathcal{T}_{\max})\,;\, (\mathsf{B}^+f, \mathsf{B}^-f) \in \mathscr{G}(\mathcal{K}) \right\}.$$

We then have the following

Lemma 3.2. *Let $\mathcal{K}$ be an unbounded operator as in Definition 3.1. For any $\lambda > 0$, the following are equivalent*

1. *$(I - M_\lambda\mathcal{K}) : \mathscr{D}(\mathcal{K}) \mapsto L^1_+$ is bijective;*
2. *$(\lambda I - \mathcal{T_K}) : \mathscr{D}(\mathcal{T_K}) \mapsto X$ is bijective.*

Proof. According to [2, Lemma 4.2], for any $\lambda > 0$ one has $[I - M_\lambda\mathcal{K}]\mathscr{D}(\mathcal{K}) = L^1_+$ if and only if $[\lambda I - \mathcal{T_K}]\mathscr{D}(T_\mathcal{K}) = X$. Therefore we have only to prove that, given $\lambda > 0$, $(I - M_\lambda\mathcal{K}) : \mathscr{D}(\mathcal{K}) \mapsto L^1_+$ is injective if and only if $(\lambda I - \mathcal{T_K}) : \mathscr{D}(\mathcal{T_K}) \mapsto X$ is injective.

Assume now that $(\lambda I - \mathcal{T_K}) : \mathscr{D}(\mathcal{T_K}) \mapsto X$ is injective and let $\psi \in \mathscr{D}(\mathcal{K})$ be a solution to $(I - M_\lambda\mathcal{K})\psi = 0$. Set $f = \Xi_\lambda\mathcal{K}\psi$. One deduces from Lemma 2.4 *(2)* (with $u = \mathcal{K}\psi$) that $f \in \mathscr{D}(\mathcal{T}_{\max})$ with $\mathcal{T}_{\max}f = \lambda f$, $\mathsf{B}^-f = \mathcal{K}\psi$ and $\mathsf{B}^+f = M_\lambda\mathcal{K}\psi = \psi$. In other words, $f \in \mathscr{D}(\mathcal{T_K})$ is a solution to the equation $(\lambda - \mathcal{T_K})f = 0$ and therefore $f = 0$. Since $\psi = \mathsf{B}^+f$, one gets $\psi = 0$ and $(I - M_\lambda\mathcal{K}) : \mathscr{D}(\mathcal{K}) \mapsto L^1_+$ is injective. Conversely, assume $(I - M_\lambda\mathcal{K}) : \mathscr{D}(\mathcal{K}) \mapsto L^1_+$ to be injective and let $f \in \mathscr{D}(\mathcal{T_K})$ be a solution to $(\lambda - \mathcal{T_K})f = 0$. According to [2, Theorem 3.2] (see also Lemma 2.4 *(2)*), $f \in \mathscr{D}(\mathcal{T}_{\max})$ with $\mathsf{B}^+f \in \mathscr{D}(\mathcal{K})$, and $f = \Xi_\lambda\mathcal{K}\mathsf{B}^+f$. Setting then $\psi = \mathsf{B}^+f$, one has $\psi \in \mathscr{D}(\mathcal{K})$ and $(I - M_\lambda\mathcal{K})\psi = 0$. By assumption, $\psi = 0$ and, since $f = \Xi_\lambda\mathcal{K}\psi$, $f = 0$ and $(\lambda I - \mathcal{T_K}) : \mathscr{D}(\mathcal{T_K}) \mapsto X$ is injective. This proves the desired equivalence. □

With this in hands, one can prove the following which somehow characterizes the class of operators sharing the properties of the generator $\mathcal{A}$ (recall that, according to Theorem 2.5, $\mathcal{A}$ satisfies the following properties *(a)–(c)*):

Proposition 1. *Let $\mathcal{A}_0$ be the generator of a strongly continuous substochastic semigroup $(V(t))_{t\geqslant 0}$ in X. Assume further that*

(a) $\mathscr{D}(\mathcal{T}_H) \subseteq \mathscr{D}(\mathcal{A}_0) \subseteq \mathscr{D}(\mathcal{T}_{\max})$
(b) $\mathcal{A}_0 f = \mathcal{T}_{\max}f$ for any $f \in \mathscr{D}(\mathcal{A}_0)$
(c) $\mathscr{D}(\mathcal{T}_H) = \{f \in \mathscr{D}(\mathcal{A}_0) : \mathsf{B}^+f \in L^1_+\} = \{f \in \mathscr{D}(\mathcal{A}_0) : \mathsf{B}^-f \in L^1_-\}$.

Then there exists a boundary linear operator $\mathcal{H}_0$ from Y_+ to Y_- with the following properties:

(i) $L^1_+ = \mathscr{D}(H) \subseteq \mathscr{D}(\mathcal{H}_0)$ *with* $\mathcal{H}_0\psi = H\psi$ *for any* $\psi \in L^1_+$
(ii) $\mathcal{A}_0 = \mathcal{T}_{\mathcal{H}_0}$
(iii) *for any $\lambda > 0$ the mapping $(I - M_\lambda\mathcal{H}_0) : \mathscr{D}(\mathcal{H}_0) \mapsto L^1_+$ is bijective, and*

$$(\lambda - \mathcal{A}_0)^{-1}f = C_\lambda f + \Xi_\lambda\mathcal{H}_0(I - M_\lambda\mathcal{H}_0)^{-1}G_\lambda f. \tag{3.1}$$

(iv) *for any $\lambda > 0$, $u \in L^1_+$, $u \geqslant 0$ one has*

$$(I - M_\lambda\mathcal{H}_0)^{-1}u \geqslant 0 \qquad \mathcal{H}_0(I - M_\lambda\mathcal{H}_0)^{-1}u \geqslant 0 \tag{3.2}$$

Proof. First of all observe that the trace mapping $\mathsf{B}^+ : \mathscr{D}(\mathcal{A}_0) \to Y_+$ is injective. Indeed let $f \in \mathscr{D}(\mathcal{A}_0)$ be such that $\mathsf{B}^+f = 0$. Then assumption *(c)* ensures that

$f \in \mathscr{D}(\mathcal{T}_H)$, so that $\mathsf{B}^- f = H\mathsf{B}^+ f = 0$. In particular, $\|(\mathsf{B}^+ f, \mathsf{B}^- f)\|_{\mathscr{E}} = 0$ and one deduces from [2, Corollary 3.1] that $f = 0$. Let us now introduce the set

$$E_0 := \mathrm{Range}(\mathsf{B}^+|_{\mathscr{D}(\mathcal{A}_0)}) = \{\psi \in Y_+ : \exists g \in \mathscr{D}(\mathcal{A}_0) \quad \text{such that} \quad \psi = \mathsf{B}^+ g\}$$

so that $\mathsf{B}^+ : \mathscr{D}(\mathcal{A}_0) \to E_0 \subseteq Y_+$ is bijective. This allows to define an unbounded linear boundary operator $\mathcal{H}_0 : \mathscr{D}(\mathcal{H}_0) \to Y_-$ as follows:

$$\mathscr{D}(\mathcal{H}_0) = E_0 \qquad \text{and} \qquad \mathcal{H}_0 \psi = \mathsf{B}^- g \qquad \forall \psi \in E_0$$

where g is the unique element of $\mathscr{D}(\mathcal{A}_0)$ such that $\mathsf{B}^+ g = \psi$. Let us prove that $\mathcal{H}_0$ satisfies points $(i) - (iv)$.

(i) Let $h \in L^1_+$ and $\lambda > 0$ be given. Setting $u = (I - M_\lambda H)h \in L^1_+$, by Lemma 2.4 (3), there exists $g \in X$ such that $G_\lambda g = u$. Setting then $f = C_\lambda g + \Xi_\lambda Hh$ one clearly has $f \in \mathscr{D}(\mathcal{T}_{\max})$. Moreover $\mathsf{B}^+ f = G_\lambda g + M_\lambda Hh = u + M_\lambda Hh = h$ and $\mathsf{B}^- f = Hh = H\mathsf{B}^+ f$. In other words, $f \in \mathscr{D}(T_H) \subseteq \mathscr{D}(\mathcal{A}_0)$. Consequently, $h \in E_0$ with $\mathcal{H}_0 h = Hh$ and (i) is proved.

(ii) To prove point (ii), it is enough to show that $\mathscr{D}(\mathcal{A}_0) = \mathscr{D}(\mathcal{T}_{\mathcal{H}_0})$. From the definition of $\mathcal{H}_0$ and the assumption $\mathscr{D}(\mathcal{A}_0) \subseteq \mathscr{D}(\mathcal{T}_{\max})$, one sees that

$$\mathscr{D}(\mathcal{A}_0) \subseteq \{f \in \mathscr{D}(\mathcal{T}_{\max}) \, ; \, \mathsf{B}^+ f \in \mathscr{D}(\mathcal{H}_0) \, , \, \mathsf{B}^- f = \mathcal{H}_0 \mathsf{B}^+ f\} = \mathscr{D}(\mathcal{T}_{\mathcal{H}_0}).$$

Conversely, let $f \in \mathscr{D}(\mathcal{T}_{\max})$ with $\mathsf{B}^+ f \in \mathscr{D}(\mathcal{H}_0)$ and $\mathsf{B}^- f = \mathcal{H}_0 \mathsf{B}^+ f$. By definition of $\mathcal{H}_0$ and since $\mathscr{D}(\mathcal{H}_0) = E_0$, there exists $g \in \mathscr{D}(\mathcal{A}_0)$ such that $\mathsf{B}^+ g = \mathsf{B}^+ f$ and $\mathsf{B}^- g = \mathcal{H}_0 \mathsf{B}^+ f = \mathsf{B}^- f$. Set $h = f - g$. One has $h \in \mathscr{D}(\mathcal{T}_{\max})$ with $\mathsf{B}^+ h = \mathsf{B}^- h = 0$ and again, we can invoke [2, Corollary 3.1] to state that $h = 0$, i.e. $f = g \in \mathscr{D}(\mathcal{A}_0)$, proving the second inclusion.

(iii) Since $\mathcal{A}_0$ is the generator of a substochastic semigroup we can state that for any $\lambda > 0$ and $f \in X$ there exists a unique $g \in \mathscr{D}(\mathcal{A}_0)$ such that $(\lambda - \mathcal{A}_0)g = f$, with moreover $g \geqslant 0$ if $f \geqslant 0$. This means that for any $\lambda > 0$ and $f \in X$ there exists a unique $g \in \mathscr{D}(\mathcal{T}_{\max})$, such that $\mathsf{B}^+ g \in \mathscr{D}(\mathcal{H}_0)$ with g solution to the boundary value problem:

$$(\lambda - \mathcal{T}_{\max})g = f \qquad \mathsf{B}^- g = \mathcal{H}_0 \mathsf{B}^+ g. \tag{3.3}$$

From [2, Theorem 3.2], such a solution g is given by

$$g = C_\lambda f + \Xi_\lambda \mathsf{B}^- g = C_\lambda f + \Xi_\lambda \mathcal{H}_0 \mathsf{B}^+ g, \tag{3.4}$$

and, in particular, $u := \mathsf{B}^+ g \in \mathscr{D}(\mathcal{H}_0)$ satisfies $(I - M_\lambda \mathcal{H}_0)u = G_\lambda f$. Since $(\lambda I - \mathcal{T}_{\mathcal{H}_0}) : \mathscr{D}(\mathcal{T}_{\mathcal{H}_0}) \to X$ is bijective, one deduces from Lemma 3.2 that $(I - M_\lambda \mathcal{H}_0) : \mathscr{D}(\mathcal{H}_0) \to L^1_+$ is bijective. Then, $u = \mathsf{B}^+ g = (I - M_\lambda \mathcal{H}_0)^{-1} G_\lambda f$ which, from (3.4), shows that the solution to (3.3) becomes

$$g = C_\lambda f + \Xi_\lambda \mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} G_\lambda f$$

which is nothing but (3.1).

(iv) Let now $\lambda > 0$ and $u \in L^1_+$ with $u \geqslant 0$ be given. Consider then the function g_λ defined as follows:

$$g_\lambda(\mathbf{x}) = \begin{cases} \dfrac{(1+\lambda)\tau_-(\mathbf{x}) + 1}{\tau_-(\mathbf{x}) + \tau_+(\mathbf{x})} \exp(-\tau_+(\mathbf{x})) \, u(\Phi(\mathbf{x}, \tau_+(\mathbf{x}))) & \text{if} \quad \tau_-(\mathbf{x}) + \tau_+(\mathbf{x}) < \infty, \\ (1+\lambda)\exp(-\tau_+(\mathbf{x}))u(\Phi(\mathbf{x}, \tau_+(\mathbf{x}))) & \text{if} \quad \tau_-(\mathbf{x}) = \infty \\ & \text{and } \tau_+(\mathbf{x}) < \infty, \\ 0 & \text{if} \quad \tau_+(\mathbf{x}) = \infty. \end{cases}$$

One can check easily that $g_\lambda \in X$, $g_\lambda \geqslant 0$ with $G_\lambda g_\lambda = u$. Setting now

$$f_\lambda = (\lambda - \mathcal{A}_0)^{-1} g_\lambda = C_\lambda g_\lambda + \Xi_\lambda \mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} G_\lambda g_\lambda$$

one sees that f_λ is nonnegative, with

$$\mathsf{B}^+ f_\lambda = (I - M_\lambda \mathcal{H}_0)^{-1} G_\lambda g_\lambda = (I - M_\lambda \mathcal{H}_0)^{-1} u \geqslant 0$$

and $\mathsf{B}^- f_\lambda = \mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} G_\lambda g_\lambda = \mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} u \geqslant 0$ which proves the result. □

The above Proposition allows to prove that $(V_H(t))_{t \geqslant 0}$ is the *smallest* substochastic semigroup generated by an extension of $\mathcal{T}_H$. More precisely we have

Theorem 3.3. *Let $(V(t))_{t \geqslant 0}$ be a strongly continuous substochastic semigroup in X with generator $\mathcal{A}_0$ which satisfies the conditions (a)-(c) of Proposition 1. Then, for any $t \geqslant 0$ one has $V(t) \geqslant V_H(t)$, i.e. $V(t)f \geqslant V_H(t)f$ for any nonnegative $f \in X$. In other words, $(V_H(t))_{t \geqslant 0}$ is the smallest substochastic semigroup generated by an extension of $\mathcal{T}_H$.*

Proof. According to the previous Proposition 1, there exists an extension $\mathcal{H}_0$ of H so that the generator $\mathcal{A}_0$ of the semigroup $(V(t))_{t \geqslant 0}$ coincides with the transport operator $T_{\mathcal{H}_0}$, and formula (3.1) holds. Now, since $\mathcal{H}_0 h = H h$ for any $h \in L^1_+$, we have, for $0 < r < 1$ and $H_r = rH$:

$$(I - M_\lambda \mathcal{H}_0)^{-1} - (I - M_\lambda H_r)^{-1} = \left[(I - M_\lambda \mathcal{H}_0)^{-1} (I - M_\lambda H_r) - I \right] (I - M_\lambda H_r)^{-1}$$

$$= (I - M_\lambda \mathcal{H}_0)^{-1} (I - M_\lambda H_r - I + M_\lambda \mathcal{H}_0)(I - M_\lambda H_r)^{-1}$$

$$= (1 - r)(I - M_\lambda \mathcal{H}_0)^{-1} M_\lambda H (I - M_\lambda H_r)^{-1}$$

where we used that the range of $(I - M_\lambda H_r)^{-1}$ is L^1_+. One deduces easily from this that

$$\mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} - H_r (I - M_\lambda H_r)^{-1}$$
$$= (1 - r) \left(\mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} M_\lambda H (I - M_\lambda H_r)^{-1} + H (I - M_\lambda H_r)^{-1} \right).$$

Recalling that $(\lambda - \mathcal{T}_{H_r})^{-1} f = C_\lambda f + \Xi_\lambda H_r (I - M_\lambda H_r)^{-1} G_\lambda f$ (see [2, Eq. (4.6)]), by virtue of (3.1) one has then, for any $f \in X$,

$$(\lambda - \mathcal{A}_0)^{-1} f - (\lambda - \mathcal{T}_{H_r})^{-1} f = \Xi_\lambda \mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} G_\lambda f - \Xi_\lambda H_r (I - M_\lambda H_r)^{-1} G_\lambda f$$

$$= (1 - r)\Xi_\lambda \left(\mathcal{H}_0 (I - M_\lambda \mathcal{H}_0)^{-1} M_\lambda H (I - M_\lambda H_r)^{-1} + H (I - M_\lambda H_r)^{-1} \right) G_\lambda f.$$

If $f \geqslant 0$, according to Proposition 1 *(iv)*, we get $(\lambda - \mathcal{A}_0)^{-1} f \geqslant (\lambda - \mathcal{T}_{H_r})^{-1} f$ for any $0 < r < 1$. This inequality together with (2.11) allow to state that $(\lambda - \mathcal{A}_0)^{-1} \geqslant (\lambda - \mathcal{A})^{-1}$ which gives the result according to the exponential formula. □

We recall now the recent result of the first author [3] about the construction of a suitable strongly continuous family of bounded linear operators in X. First, let

$$\mathcal{D}_0 = \{ f \in \mathscr{D}(\mathcal{T}_{\max}) \ : \ \mathsf{B}^\pm f = 0 \}.$$

The subset $\mathcal{D}_0$ is dense in X (see [3, Proposition 1]). Remember that the semigroup $(U_0(t))_{t \geqslant 0}$ is defined through (2.9). Now, one introduces the following

Definition 3.4. For any $t \geqslant 0$, we define the family $(U_k(t))_{k \in \mathbb{N}}$ by induction as follows: if $f \in \mathcal{D}_0$, $t > 0$ and $k \geqslant 1$, one sets

$$U_k(t)f(\mathbf{x}) = \begin{cases} H(\mathsf{B}^+ U_{k-1}(t - \tau_-(\mathbf{x}))f)(\Phi(\mathbf{x}, -\tau_-(\mathbf{x}))) & \forall \mathbf{x} \in \mathbf{\Omega}, \quad \text{with} \quad \tau_-(\mathbf{x}) \leqslant t, \\ 0 & \forall \mathbf{x} \in \mathbf{\Omega} \quad \text{with} \quad \tau_-(\mathbf{x}) > t. \end{cases}$$
(3.5)

Moreover, for $t = 0$, we set $U_k(0)f = 0$ for any $k \geqslant 1$ and any $f \in X$.

Remark 2. In other words, if we put $\boldsymbol{\Omega}_t := \{\mathbf{x} \in \boldsymbol{\Omega} : \mathbf{x} = \Phi(\mathbf{y}, s), \mathbf{y} \in \Gamma_-, 0 < s < t \wedge \tau_+(\mathbf{y})\}$, then $[U_k(t)f](\mathbf{x})$ may be different from zero only for $\mathbf{x} \in \boldsymbol{\Omega}_t$, being $U_k(t)f(\Phi(\mathbf{y}, s)) = H(\mathsf{B}^+U_{k-1}(t-s)f)(\mathbf{y})$.

Remark 3. Notice that, given $f \in \mathcal{D}_0$ and $t > 0$, one has $\left(\int_0^t U_k(s)f\mathrm{d}s\right)(\mathbf{x}) = 0$ for any $\mathbf{x} \in \boldsymbol{\Omega}$ with $\tau_-(\mathbf{x}) > t$. In particular,

$$\mathsf{B}^+\left(\int_0^t U_k(s)f\mathrm{d}s\right)(\mathbf{z}) = 0 \qquad \forall \mathbf{z} \in \Gamma_+ \, ; \; \tau_-(\mathbf{z}) > t \, , \; k \geqslant 1. \tag{3.6}$$

The properties of the family $(U_k(t))_{t\geqslant 0}$, for given $k \geqslant 1$, have been established in [3]. In particular, for any $f \in \mathcal{D}_0$ and any $t > 0$, one has $U_k(t)f \in X$ with

$$\|U_k(t)f\|_X \leqslant \|H\|_{\mathscr{B}(L_+^1, L_-^1)}^k \|f\|_X = \|f\|_X \qquad \forall k \geqslant 1.$$

Since $\mathcal{D}_0$ is dense in X, one can extend $U_k(t)$ in a bounded linear operator in X, still denoted $U_k(t)$ such that

$$\|U_k(t)\|_{\mathscr{B}(X)} \leqslant 1.$$

Moreover, one has the following

Proposition 2. *For any $k \geqslant 1$, the family $(U_k(t))_{t\geqslant 0}$ enjoys the following properties:*

1. *$(U_k(t))_{t\geqslant 0}$ is a strongly continuous family of operators in X.*
2. *For all $f \in \mathcal{D}_0$ and $t \geqslant 0$ one has $U_k(t)f \in \mathscr{D}(\mathcal{T}_{\max})$ with $\mathcal{T}_{\max}U_k(t)f = U_k(t)\mathcal{T}_{\max}f$.*
3. *For all $f \in \mathcal{D}_0$ and $t \geqslant 0$ the traces $\mathsf{B}^\pm U_k(t)f \in L_\pm^1$ and the mappings $t \mapsto \mathsf{B}^\pm U_k(t)f \in L_\pm^1$ are continuous.*
4. *For any $f \in X$, $t \geqslant 0$ and $s \geqslant 0$ we have $U_k(t+s)f = \sum_{j=0}^k U_j(t)U_{k-j}(s)f$.*
5. *For all $f \in X$ and $t > 0$ one has $\int_0^t U_k(s)f\mathrm{d}s \in \mathscr{D}(\mathcal{T}_{\max})$ with*

$$\mathcal{T}_{\max}\int_0^t U_k(s)f\mathrm{d}s = U_k(t)f.$$

Moreover, $\mathsf{B}^\pm\left(\int_0^t U_k(s)f\mathrm{d}s\right) \in L_\pm^1$ and

$$H\mathsf{B}^+\left(\int_0^t U_{k-1}(s)f\mathrm{d}s\right) = \mathsf{B}^-\left(\int_0^t U_k(s)f\mathrm{d}s\right). \tag{3.7}$$

6. *For any $f \in X$ and $\lambda > 0$, setting $g_k := \int_0^\infty \exp(-\lambda t)U_k(t)f\mathrm{d}t$, one has $g_k \in \mathscr{D}(\mathcal{T}_{\max})$ with*

$$\mathcal{T}_{\max}g_k = \lambda g_k \quad for \quad k \geqslant 1, \qquad while \quad \mathcal{T}_{\max}g_0 = \lambda g_0 - f;$$

and $\mathsf{B}^+g_k = (M_\lambda H)^k G_\lambda f \in L_+^1$ for any $k \geqslant 0$ while $\mathsf{B}^-g_0 = 0$ and $\mathsf{B}^-g_k = H\mathsf{B}^+g_{k-1}$ if $k \geqslant 1$.
7. *For any nonnegative $f \in X$ and any $t \geqslant 0$ and $n \geqslant 1$ one has*

$$\sum_{k=0}^n \|U_k(t)f\|_X = \|f\|_X - \left\|\mathsf{B}^+\int_0^t U_n(s)f\mathrm{d}s\right\|_{L_+^1}$$

$$+ \sum_{k=0}^{n-1}\left[\left\|H\mathsf{B}^+\int_0^t U_k(s)f\mathrm{d}s\right\|_{L_-^1} - \left\|\mathsf{B}^+\int_0^t U_k(s)f\mathrm{d}s\right\|_{L_+^1}\right]. \tag{3.8}$$

In particular,

$$\sum_{k=0}^{n} \|U_k(t)f\|_X \leqslant \|f\|_X - \left\| \mathsf{B}^+ \int_0^t U_n(s)f\mathrm{d}s \right\|_{L_+^1} \leqslant \|f\|_X. \tag{3.9}$$

The above listed properties allow to give a characterization of the semigroup $(V_H(t))_{t \geqslant 0}$ in terms of a strongly convergent expansion series, reminiscent to classical Dyson-Phillips expansion series for additive perturbation:

Theorem 3.5. *For any $f \in X$ and any $t \geqslant 0$, one has*

$$V_H(t)f = \sum_{k=0}^{\infty} U_k(t)f. \tag{3.10}$$

Proof. For any $f \in X$ and any $t \geqslant 0$, set $V(t)f = \sum_{k=0}^{\infty} U_k(t)f$. Notice that the series is convergent in X and the family $(V(t))_{t \geqslant 0}$ defines a substochastic C_0-semigroup in X (see [3, Theorem 4.3] for details). Let us prove that $V(t) = V_H(t)$ for all $t \geqslant 0$. Let $f \in X$ and $\lambda > 0$ be fixed. Set, for any $k \geqslant 1$,

$$g_k = \int_0^{\infty} \exp(-\lambda t)U_k(t)f\mathrm{d}t.$$

Proposition 2 asserts that $g_k \in \mathscr{D}(\mathcal{T}_{\max})$ and satisfies $\mathcal{T}_{\max} g_k = \lambda g_k$ for any $k \geqslant 1$. According to [2, Theorem 2.1] we deduce that, for $k \geqslant 1$, $g_k = \Xi_\lambda H \mathsf{B}^+ g_{k-1} = \Xi_\lambda H (M_\lambda H)^{k-1} G_\lambda f$. Summing this identity, we get that

$$\int_0^{\infty} \exp(-\lambda t)V(t)f\mathrm{d}t = \sum_{k=0}^{\infty} g_k = C_\lambda f + \sum_{k=0}^{\infty} \Xi_\lambda H(M_\lambda H)^k G_\lambda f.$$

Since this last expression coincides with $(\lambda - \mathcal{A})^{-1}f$, one deduces from the injectivity of Laplace transform that $V(t)f = V_H(t)f$ for any $t \geqslant 0$. $\square$

An immediate consequence of the above Theorem 3.5 is given in the following

Corollary 1. *For any $f \in X$ and $\lambda > 0$, as $n \to \infty$, the sum*

$$\sum_{k=0}^{n} \int_0^{\infty} \exp(-\lambda t)U_k(t)f\mathrm{d}t$$

converges to $(\lambda - \mathcal{A})^{-1}f$ in the graph norm of $\mathcal{A}$.

We end this section with a technical result that complements Proposition 2 and shall be useful in the sequel

Lemma 3.6. *Let $f \in X$ be nonnegative and $t > 0$ be given. For any $\mathbf{z} \in \Gamma_+$ and any $k \geqslant 1$ it holds*

$$\left[\mathsf{B}^+ \int_0^t U_k(s)f\mathrm{d}s \right](\mathbf{z}) \leqslant \left[H\mathsf{B}^+ \int_0^t U_{k-1}(s)f\mathrm{d}s \right](\Phi(\mathbf{z}, -\tau_-(\mathbf{z}))).$$

Proof. Let $k \geqslant 1$ and $\mathbf{z} \in \Gamma_+$ be given. If $\tau_-(\mathbf{z}) > t$, one gets from (3.6) that

$$\mathsf{B}^+ \left(\int_0^t U_k(s)f\mathrm{d}s \right)(\mathbf{z}) = 0$$

from which the conclusion clearly holds. Now, if $\tau_-(\mathbf{z}) \leqslant t$, set $\mathbf{y} = \Phi(\mathbf{z}, -\tau_-(\mathbf{z})) \in \Gamma_-$. Since $\int_0^t U_k(s)f\mathrm{d}s \in \mathscr{D}(\mathcal{T}_{\max})$ with $\mathcal{T}_{\max}\int_0^t U_k(s)f\mathrm{d}s = U_k(t)f$, one deduces from Definition 2.1 (see also [1, Theorem 3.6]) that

$$\int_{t_1}^{t_2}[U_k(t)f](\Phi(\mathbf{y},s))\mathrm{d}s = \left[\int_0^t U_k(s)f\mathrm{d}s\right](\Phi(\mathbf{y},t_1)) - \left[\int_0^t U_k(s)f\mathrm{d}s\right](\Phi(\mathbf{y},t_2))$$

for any $0 < t_1 < t_2 < \tau_+(\mathbf{y}) = \tau_-(\mathbf{z}) \leqslant t$. In particular, for nonnegative f we get

$$\left[\int_0^t U_k(s)f\mathrm{d}s\right](\Phi(\mathbf{y},t_2)) \leqslant \left[\int_0^t U_k(s)f\mathrm{d}s\right](\Phi(\mathbf{y},t_1))$$
$$\forall 0 < t_1 < t_2 < \tau_+(\mathbf{y}) = \tau_-(\mathbf{z}) \leqslant t.$$

Letting $t_1 \to 0^+$ and $t_2 \to \tau_+(\mathbf{y})$ and since $\mathbf{z} = \Phi(\mathbf{y}, \tau_+(\mathbf{y}))$, we get

$$\left[\mathsf{B}^+\int_0^t U_k(s)f\mathrm{d}s\right](\mathbf{z}) \leqslant \left[\mathsf{B}^-\left(\int_0^t U_k(s)f\mathrm{d}s\right)\right](\mathbf{y}).$$

Using now (3.7) and the fact that $\mathbf{y} = \Phi(\mathbf{z}, -\tau_-(\mathbf{z}))$ we get the conclusion. $\qquad\square$

4. **Honesty theory.**

4.1. **On some functionals.** For any $f \in \mathscr{D}(\mathcal{T}_{\max})$ we define

$$\mathfrak{a}(f) = -\int_{\mathbf{\Omega}} \mathcal{T}_{\max}f d\mu.$$

while, for any $f \in \mathscr{W}$, we set

$$\mathfrak{a}_0(f) = \int_{\Gamma_+} \mathsf{B}^+ f d\mu_+ - \int_{\Gamma_-} H\mathsf{B}^+ f d\mu_-.$$

Clearly $\mathfrak{a} : \mathscr{D}(\mathcal{T}_{\max}) \to \mathbb{R}$ is a linear functional with $|\mathfrak{a}(f)| \leqslant \|\mathcal{T}_{\max}f\|_X$ for any $f \in \mathscr{D}(\mathcal{T}_{\max})$. Here we are interested in the restriction of $\mathfrak{a}$ to $\mathscr{D}(\mathcal{A})$, that we still denote by $\mathfrak{a}$. Since $\mathcal{A}$ generates a positive contraction semigroup $(V_H(t))_{t\geqslant 0}$ we have

$$\mathfrak{a}(f) = -\int_{\mathbf{\Omega}} \mathcal{A}f\mathrm{d}\mu = \lim_{t\to 0^+} t^{-1}\int_{\mathbf{\Omega}}(f - V_H(t)f)\,\mathrm{d}\mu \geqslant 0$$
$$\forall f \in \mathscr{D}(\mathcal{A})_+ := \mathscr{D}(\mathcal{A}) \cap X_+.$$

Hence $\mathfrak{a} : \mathscr{D}(\mathcal{A}) \to \mathbb{R}$ is a positive linear functional. Furthermore $\mathfrak{a}$ is continuous in the graph norm of $\mathcal{A}$ and its restriction to $\mathscr{D}(\mathcal{T}_H)$ is equal to the restriction of $\mathfrak{a}_0$ to $\mathscr{D}(\mathcal{T}_H)$. Indeed, according to Green's formula (2.6) for all $f \in \mathscr{D}(\mathcal{T}_H)$ we have

$$\mathfrak{a}(f) = \int_{\Gamma_+} \mathsf{B}^+ f d\mu_+ - \int_{\Gamma_-} \mathsf{B}^- f d\mu_- = \int_{\Gamma_+} \mathsf{B}^+ f d\mu_+ - \int_{\Gamma_-} H\mathsf{B}^+ f d\mu_- = \mathfrak{a}_0(f).$$

This basic observation allows to formulate an equivalent to [6, Proposition 4.5] in this boundary perturbation context. Precisely, one has

Proposition 3. *For all $f \in \mathscr{D}(\mathcal{A})$ there exists*

$$\lim_{t\to 0+} \frac{1}{t}\sum_{k=0}^{\infty} \mathfrak{a}_0\left(\int_0^t U_k(s)f\mathrm{d}s\right) =: \widehat{\mathfrak{a}}(f) \tag{4.1}$$

with $|\widehat{\mathfrak{a}}(f)| \leqslant 2(\|f\|_X + \|\mathcal{A}f\|_X)$. Furthermore, if $f \in \mathscr{D}(\mathcal{A})_+$, then

$$0 \leqslant \widehat{\mathfrak{a}}(f) \leqslant \mathfrak{a}(f) \leqslant \|\mathcal{T}_{\max}f\|. \tag{4.2}$$

The proof of Proposition 3 is based upon the following

Lemma 4.1. *For any $f \in X$ and $t > 0$ one has*

$$\left| \sum_{k=0}^{\infty} \mathfrak{a}_0 \left(\int_0^t U_k(s) f \mathrm{d}s \right) \right| \leqslant \|f\|_X. \tag{4.3}$$

If $f \in \mathscr{D}(\mathcal{A})$ then one also has

$$\left| \sum_{k=0}^{\infty} \mathfrak{a}_0 \left(\int_0^t U_k(s) f \mathrm{d}s \right) \right| \leqslant 2t \left(\|f\|_X + \|\mathcal{T}_{\max} f\|_X \right). \tag{4.4}$$

Proof. For simplicity, for any *fixed* $t > 0$, we set

$$G_k(f) = \int_0^t U_k(s) f \mathrm{d}s \qquad \forall k \geqslant 1.$$

According to Proposition 2 (5), $G_k(f) \in \mathscr{D}(\mathcal{T}_{\max})$ for any $f \in X$, $k \geqslant 1$ with moreover $\mathsf{B}^+ G_k(f) \in L^1_+$, i.e. $G_k(f) \in \mathscr{W}$. We begin with assuming $f \in X_+$ and $t > 0$. One can reformulate (3.8) as

$$\sum_{k=0}^{n-1} \mathfrak{a}_0 \left(G_k(f) \right) = \|f\|_X - \sum_{k=0}^{n} \|U_k(t) f\|_X - \|\mathsf{B}^+ G_n(f)\|_{L^1+} \leqslant \|f\|_X \tag{4.5}$$

Therefore, we can see that $\left(\sum_{k=0}^{n} \mathfrak{a}_0 \left(G_k(f) \right) \right)_n$ is an increasing bounded sequence whose limit satisfies

$$\sum_{k=0}^{\infty} \mathfrak{a}_0 \left(G_k(f) \right) \leqslant \|f\|_X - \sum_{k=0}^{\infty} \|U_k(t) f\|_X. \tag{4.6}$$

Now, for general $f \in X$, since $G_k(f) \in \mathscr{W}$, we deduce from [2, Proposition 2.2] that $|G_k(f)| \in \mathscr{W}$ and, since $U_k(s)$ $(0 < s < t, k \geqslant 0)$ is a positive operator, the inequalities

$$|\mathfrak{a}_0 \left(G_k(f) \right)| \leqslant \mathfrak{a}_0 \left(|G_k(f)| \right) \leqslant \mathfrak{a}_0 \left(G_k(|f|) \right) \qquad \forall k \geqslant 1$$

hold. This, together with (4.6) yields (4.3). Before proving (4.4), one notices that the right-hand side of (4.6) for $f \geqslant 0$ is

$$\int_{\mathbf{\Omega}} \left(f - \sum_{k=0}^{\infty} U_k(t) f \right) \mathrm{d}\mu = \int_{\mathbf{\Omega}} (f - V_H(t) f) \, \mathrm{d}\mu = - \int_{\mathbf{\Omega}} \mathcal{A} \left(\int_0^t V_H(s) f \mathrm{d}s \right) \mathrm{d}\mu$$

where we used Theorem 3.5 and the well-know fact (see [16, Lemma 1.3, p. 50]) that, for any C_0-semigroup $(V_H(t)_{t \geqslant 0}$ with generator $\mathcal{A}$, one has $\int_0^t V_H(s) f \mathrm{d}s \in \mathscr{D}(\mathcal{A})$ with $\mathcal{A} \left(\int_0^t V_H(s) f \mathrm{d}s \right) = V_H(t) f - f$ for any $t \geqslant 0$ and any $f \in X$. Since moreover $V_H(t) f - f = \int_0^t V_H(s) \mathcal{A} f \mathrm{d}s$ if $f \in \mathscr{D}(\mathcal{A})$, one gets

$$\sum_{k=0}^{\infty} \mathfrak{a}_0 \left(\int_0^t U_k(s) f \mathrm{d}s \right) \leqslant \mathfrak{a} \left(\int_0^t V_H(s) f \mathrm{d}s \right)$$

$$= - \int_{\mathbf{\Omega}} \left(\int_0^t V_H(s) \mathcal{A} f \mathrm{d}s \right) \mathrm{d}\mu \qquad \forall f \in \mathscr{D}(\mathcal{A}) \cap X_+. \tag{4.7}$$

Let us now fix $f \in \mathscr{D}(\mathcal{A})$ and set $g := (I - \mathcal{A}) f = g^+ - g^-$, where g_+ and g_- denote respectively the positive and negative parts of g. Put also $f_1^{\pm} = (I - \mathcal{A})^{-1} g_{\pm}$ so that

$f = f_1^+ - f_1^-$, where $f_1^\pm$ are belonging to $\mathscr{D}(\mathcal{A})_+$ (notice that $f_1^\pm$ do not necessarily coincide with the positive and negative parts $f^\pm$ of f). One has

$$\|\mathcal{A}f_1^\pm\|_X \leqslant \|f_1^\pm\|_X + \|g^\pm\|_X \leqslant 2\|g^\pm\|_X.$$

Recalling that $\mathcal{A}f_1^\pm = \mathcal{T}_{\max}f_1^\pm$ and using formula (4.7) we get

$$\sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f_1^\pm \mathrm{d}s \right) \leqslant - \int_\Omega \left(\int_0^t V_H(s)\mathcal{A}f_1^\pm \mathrm{d}s \right) \mathrm{d}\mu$$

$$\leqslant \int_0^t \|V_H(s)\mathcal{T}_{\max}f_1^\pm\| \mathrm{d}s \leqslant t\|\mathcal{T}_{\max}f_1^\pm\|_X \leqslant 2t\|g_\pm\|_X$$

where we used that the semigroup $(V_H(t))_{t\geqslant 0}$ is substochastic. Finally, noticing that

$$\left| \sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f\mathrm{d}s \right) \right| \leqslant \sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f_1^+\mathrm{d}s \right) + \sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f_1^-\mathrm{d}s \right)$$

we obtain (4.4) since $\|g^+\|_X + \|g^-\|_X = \|g\|_X \leqslant \|f\|_X + \|\mathcal{T}_{\max}f\|_X$. $\qquad\square$

Proof of Proposition 3. Using Lemma 4.1 together with a repeated use of Proposition 2 *(4)*, it is not difficult to resume the proof of [6, Proposition 4.5] to get the result. Details are omitted. $\qquad\square$

As an immediate consequence of Proposition 3 we deduce the following

Corollary 2. *For any $f \in X$, $t > 0$ and $\lambda > 0$ one has*

$$\widehat{\mathfrak{a}} \left(\int_0^t V_H(s)f\mathrm{d}s \right) = \sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f\mathrm{d}s \right), \tag{4.8}$$

and

$$\widehat{\mathfrak{a}} \left((\lambda - \mathcal{A})^{-1}f \right) = \sum_{k=0}^\infty \left(\int_{\Gamma_+} (M_\lambda H)^k G_\lambda f\mathrm{d}\mu_+ - \int_{\Gamma_-} H(M_\lambda H)^k G_\lambda f\mathrm{d}\mu_- \right). \tag{4.9}$$

Proof. Identity (4.8) is simply deduced from definition (4.1) and from the identity

$$\sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f\mathrm{d}s \right) = \lim_{\tau\to 0+} \frac{1}{\tau} \sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^\tau U_k(r) \left(\int_0^t V_H(s)f\mathrm{d}s \right) \mathrm{d}r \right).$$

which is the equivalent of [6, Eq. (4.14)] in our context. Regarding (4.9), observe that for any $f \in X$, and $\lambda > 0$ one has

$$(\lambda - \mathcal{A})^{-1}f = \int_0^\infty \exp(-\lambda t)V_H(t)f\mathrm{d}t = \lambda \int_0^\infty \exp(-\lambda t) \left(\int_0^t V_H(s)f\mathrm{d}s \right) \mathrm{d}t.$$

Therefore, from (4.8),

$$\widehat{\mathfrak{a}} \left((\lambda - \mathcal{A})^{-1}f \right) = \lambda \int_0^\infty \exp(-\lambda t)\widehat{\mathfrak{a}} \left(\int_0^t V_H(s)f\mathrm{d}s \right) \mathrm{d}t$$

$$= \lambda \int_0^\infty \exp(-\lambda t) \sum_{k=0}^\infty \mathfrak{a}_0 \left(\int_0^t U_k(s)f\mathrm{d}s \right) \mathrm{d}t.$$

Setting, $g_k = \int_0^\infty \exp(-\lambda t) U_k(t) f \mathrm{d}t$ and $\phi_k(t) = \int_0^t U_k(s) f \mathrm{d}s$, one deduces from Proposition 2 *(6)* that, for any $k \geqslant 1$

$$\lambda \int_0^\infty \exp(-\lambda t) \mathsf{B}^+ \phi_k(t) \mathrm{d}t = \mathsf{B}^+ \, g_k = (M_\lambda H)^k G_\lambda f,$$

and, recalling that $\mathfrak{a}_0(\phi_k(t)) = \int_{\Gamma_+} \mathsf{B}^+ \phi_k(t) \mathrm{d}\mu_+ - \int_{\Gamma_-} H \mathsf{B}^+ \phi_k(t) \mathrm{d}\mu_-$ we get (4.9).

$\square$

Remark 4. In the free-streaming context, the identity (4.9) shows that the functional $\widehat{\mathfrak{a}}$ coincides with the functional c_λ defined in [18]. In particular, this shows that the functional c_λ of [18] *does not depend on* λ, answering the question left open in [18, Remark 17]. Moreover, by Proposition 3, we see that the functionals c_λ and $\widehat{c}$ of [18] (corresponding respectively to our $\widehat{\mathfrak{a}}$ and $\mathfrak{a}$) are positive functionals such that $c_\lambda(\varphi) \leqslant \widehat{c}(\varphi)$ for all $\varphi \in \mathscr{D}(\mathcal{A})_+$ which extends the result of [18, Remark 17] valid only for $\varphi \in (\lambda - \mathcal{A})^{-1} X_+$.

Proposition 3 allows to define a third linear positive functional $\Theta : \mathscr{D}(\mathcal{A}) \to \mathbb{R}$ by setting

$$\Theta(f) = \mathfrak{a}(f) - \widehat{\mathfrak{a}}(f) \qquad \text{for any } f \in \mathscr{D}(\mathcal{A}).$$

Clearly, the functional Θ is continuous in the graph norm of $\mathscr{D}(\mathcal{A})$. Other properties of Θ are stated here below.

Corollary 3. *For any $f \in X$, $t > 0$ and $\lambda > 0$ one has*

$$\Theta\left(\int_0^t V_H(s) f \mathrm{d}s \right) = \lim_{n \to \infty} \int_{\Gamma_+} \mathsf{B}^+ \left(\int_0^t U_n(s) f \mathrm{d}s \right) \mathrm{d}\mu_+, \qquad (4.10)$$

and

$$\Theta\left((\lambda - \mathcal{A})^{-1} f \right) = \lim_{n \to \infty} \int_{\Gamma_+} (M_\lambda H)^n G_\lambda f \mathrm{d}\mu_+. \qquad (4.11)$$

In particular, both the limits appearing in (4.10) and (4.11) exist and are finite for any $f \in X$.

Proof. As in the proof of Corollary 2, for fixed $f \in X$, $\lambda > 0$ and $t \geqslant 0$ set

$$g_k = \int_0^\infty \exp(-\lambda t) U_k(t) f \mathrm{d}t \qquad \text{and} \qquad \phi_k(t) = \int_0^t U_k(s) f \mathrm{d}s, \qquad k \geqslant 0.$$

Notice that $\phi_k(t) \in \mathscr{W}$ for any $t > 0$ and any $k \geqslant 1$. One checks then easily thanks to Proposition 2 *(7)* that, for any $n \geqslant 1$

$$\sum_{k=0}^n \mathfrak{a}\left(\phi_k(t) \right) = \sum_{k=0}^{n-1} \mathfrak{a}_0\left(\phi_k(t) \right) + \int_{\Gamma_+} \mathsf{B}^+ \phi_n(t) \mathrm{d}\mu_+.$$

One deduces easily (4.10) from this last identity combined with (4.8) and the fact that $\left(\sum_{k=0}^n \phi_k(t) \right)_n$ converges to $\int_0^t V_H(s) f \mathrm{d}s$ in the graph norm of $\mathcal{A}$. In the same way, noticing that for any $n \in \mathbb{N}$ one has

$$\sum_{k=0}^n \mathfrak{a}\left(g_k \right) = \sum_{k=0}^{n-1} \mathfrak{a}_0\left(g_k \right) + \int_{\Gamma_+} \mathsf{B}^+ g_n \mathrm{d}\mu_+,$$

one readily gets (4.11) using now (4.9) together with the fact that $\left(\sum_{k=0}^n g_k \right)_n$ converges to $(\lambda - \mathcal{A})^{-1} f$ in the graph norm of $\mathcal{A}$ as $n \to \infty$ (see Corollary 1). $\square$

The above results yield the following

Proposition 4. *For any $f \in \mathscr{D}(\mathcal{T}_H)$ one has $\widehat{\mathfrak{a}}(f) = \mathfrak{a}(f) = \mathfrak{a}_0(f)$. Consequently,*

$$\Theta(f) = 0 \qquad \forall f \in \mathscr{D}\left(\overline{\mathcal{T}_H}\right). \tag{4.12}$$

Proof. By definition and since Θ is continuous over $\mathscr{D}(\mathcal{A})$ endowed with the graph norm, it is enough to prove that $\widehat{\mathfrak{a}}(f) = \mathfrak{a}(f)$ for any $f \in \mathscr{D}(\mathcal{T}_H)$. For any $\lambda > 0$, since the operator $G_\lambda : X \to L^1_+$ is surjective, one deduces from (4.11) that the limit $\lim_{n\to\infty} \int_{\Gamma_+} (M_\lambda H)^n h \mathrm{d}\mu_+$ exists and is finite for any $h \in L^1_+$. Now, given $f \in \mathscr{D}(\mathcal{T}_H)$, set $g = (\lambda - \mathcal{A})f$. Since $\mathsf{B}^+ f \in L^1_+$ the limit $\lim_{n\to\infty} \int_{\Gamma_+} (M_\lambda H)^n \mathsf{B}^+ f \mathrm{d}\mu_+$ exists and is finite while, from $f = (\lambda - \mathcal{A})^{-1} g$ one deduces that

$$\sum_{k=0}^{n-1} (M_\lambda H)^k G_\lambda g = \mathsf{B}^+ f - (M_\lambda H)^n \mathsf{B}_+ f.$$

Therefore, the sequence $\left(\sum_{k=0}^{n-1} \int_{\Gamma_+} (M_\lambda H)^k G_\lambda g \mathrm{d}\mu_+\right)_n$ is converging. In particular,

$$\lim_{n\to\infty} \int_{\Gamma_+} (M_\lambda H)^n G_\lambda g \mathrm{d}\mu_+ = 0.$$

From (4.11), this limit coincides with $\Theta\left((\lambda - \mathcal{A})^{-1} g\right) = \Theta(f)$ which shows the result. $\qquad\square$

Now one proves that, somehow, (4.12) is a characterization of $\mathscr{D}(\overline{\mathcal{T}_H})$, at least for nonnegative f:

Proposition 5. *If $f \in \mathscr{D}(\mathcal{A})_+$ is such that $\Theta(f) = 0$, then, $f \in \mathscr{D}(\overline{\mathcal{T}_H})$.*

Proof. The proof is inspired by the analogous result for additive perturbation [19, Proposition 1.6]. Let $f \in \mathscr{D}(\mathcal{A})_+$ be given such that $\Theta(f) = 0$, i.e. $\widehat{\mathfrak{a}}(f) = \mathfrak{a}(f)$. Since $\lambda(\lambda - \mathcal{A})^{-1} f \to f$ in the graph norm of $\mathscr{D}(\mathcal{A})$ as $\lambda \to \infty$, we get that

$$\lim_{\lambda\to\infty} \Theta\left(\lambda(\lambda - \mathcal{A})^{-1} f\right) = \Theta(f) = 0.$$

Now, since

$$\Theta\left(\lambda(\lambda - \mathcal{A})^{-1} f\right) = \lim_{n\to\infty} \lambda \int_{\Gamma_+} (M_\lambda H)^n G_\lambda f \mathrm{d}\mu_+$$

we see that, for any $\varepsilon > 0$ we can find $\lambda > 1$ and $N \geqslant 1$ such that

$$\|\lambda(\lambda - \mathcal{A})^{-1} \mathcal{T}_{\max} f - \mathcal{T}_{\max} f\|_X < \varepsilon; \quad \|\lambda(\lambda - \mathcal{A})^{-1} f - f\|_X < \varepsilon$$

$$\text{and} \int_{\Gamma_+} (M_\lambda H)^n G_\lambda f \mathrm{d}\mu_+ < \frac{\varepsilon}{\lambda} \qquad \forall n \geqslant N. \tag{4.13}$$

For such $\lambda > 1$ and $N \geqslant 1$, we construct a sequence (φ_n) in $\mathscr{W}$ with the following properties

$$\mathsf{B}^- \varphi_n = 0; \quad \mathsf{B}^+ \varphi_n = (M_\lambda H)^n G_\lambda f; \quad \|\varphi_n\|_X \leqslant \|(M_\lambda H)^n G_\lambda f\|_{L^1_+}$$

$$\text{and} \qquad \|\mathcal{T}_{\max} \varphi_n\|_X \leqslant \|(M_\lambda H)^n G_\lambda f\|_{L^1_+}.$$

The existence of such a sequence is ensured by [2, Proposition 2.3]. Then, for any $n \geqslant 1$, we set

$$u_n = C_\lambda f + \sum_{k=0}^{n-1} \Xi_\lambda H (M_\lambda H)^k G_\lambda f - \varphi_n.$$

Clearly, $u_n \in \mathscr{D}(\mathcal{T}_{\max})$ for any $n \geqslant 1$ with

$$\mathcal{T}_{\max} u_n = \lambda \left(C_\lambda f + \sum_{n=0}^{n-1} \Xi_\lambda H (M_\lambda H)^n G_\lambda f \right) - f - \mathcal{T}_{\max} \varphi_n;$$

$$\mathsf{B}^+ u_n = \sum_{k=0}^{n-1} (M_\lambda H)^k G_\lambda f \quad \text{and} \quad \mathsf{B}^- u_n = \sum_{k=0}^{n-1} H (M_\lambda H)^k G_\lambda f = H \mathsf{B}^+ u_n$$

i.e. $u_n \in \mathscr{D}(\mathcal{T}_H)$ for all $n \geqslant 1$. Considering that

$$(\lambda - \mathcal{A})^{-1} f = C_\lambda f + \sum_{k=0}^{\infty} \Xi_\lambda H (M_\lambda H)^k G_\lambda f,$$

we can choose $n \geqslant N$ such that

$$\left\| (\lambda - \mathcal{A})^{-1} f - C_\lambda f - \sum_{k=0}^{n-1} \Xi_\lambda H (M_\lambda H)^k G_\lambda f \right\|_X < \frac{\varepsilon}{\lambda^2}.$$

With such choice, since $\lambda(\lambda - \mathcal{A})^{-1} \mathcal{T}_{\max} f = \lambda^2 (\lambda - \mathcal{A})^{-1} f - \lambda f$, we check that

$$\|\lambda u_n - f\|_X \leqslant \|\lambda(u_n - (\lambda - \mathcal{A})^{-1} f)\|_X + \|\lambda(\lambda - \mathcal{A})^{-1} f - f\|_X < 4\varepsilon$$
$$\text{and} \qquad \|\lambda \mathcal{T}_{\max} u_n - \mathcal{T}_{\max} f\|_X < 4\varepsilon.$$

Since $\lambda u_n \in \mathscr{D}(\mathcal{T}_H)$, this shows that $f \in \mathscr{D}(\overline{\mathcal{T}_H})$. □

4.2. **Honesty criteria.** Here we want to improve the honesty theory developed in [18]. First of all we adapt the definition of honesty, established in the additive perturbation framework in [19, 6].

Definition 4.2. Let $f \in X_+$ be given. Let $J \subseteq [0, \infty)$ be an interval. Then, the trajectory $(V_H(t)f)_{t \geqslant 0}$ is said to be honest on J if

$$\|V_H(t)f\|_X = \|V_H(s)f\|_X - \widehat{\mathfrak{a}} \left(\int_s^t V_H(r)f \mathrm{d}r \right), \qquad \forall \quad s, t \in J, s \leqslant t.$$

The trajectory is said to be honest if it is honest on $[0, \infty)$. The whole C_0-semigroup $(V_H(t))_{t \geqslant 0}$ will be said honest if all the trajectories are honest.

In the following, we establish thanks to the representation series (3.10) an approach to honesty on subinterval $J \subseteq [0, \infty)$ which is completely new in the context of boundary perturbation. The proof is inspired by the recent similar results obtained in the additive perturbation framework thanks to Dyson-Phillips series (see the concept of so-called 'mild honesty' in [6, Section 4]). More precisely, we have the following honesty criteria, analogous to [6, Theorem 4.8]:

Theorem 4.3. *Given $f \in X_+$ and $J \subseteq [0, \infty)$, the following statements are equivalent*

1) the trajectory $(V_H(t)f)_{t \geqslant 0}$ is honest on J;

2) $\displaystyle \lim_{n \to \infty} \left\| \mathsf{B}^+ \int_s^t U_n(r)f \mathrm{d}r \right\|_{L_+^1} = 0$ for any $s, t \in J, s \leqslant t$;

3) $\displaystyle \int_s^t V_H(r)f \mathrm{d}r \in \mathscr{D}(\overline{\mathcal{T}_H})$ for any $s, t \in J, s \leqslant t$;

4) the set $\left(\mathsf{B}^+ \int_s^t U_n(r)f\,\mathrm{d}r\right)_n$ is relatively weakly compact in L^1_+ for any $s,t \in J, s \leqslant t$.

Proof. Let $f \in X_+$, $J \subseteq [0,\infty)$ and $s,t \in J, s \leqslant t$ be given. Recall that

$$\mathfrak{a}\left(\int_s^t V_H(r)f\,\mathrm{d}r\right) = \|V_H(s)f\|_X - \|V_H(t)f\|_X\,.$$

so that, according to Definition 4.2, the trajectory $(V_H(t)f)_{t\geqslant 0}$ is honest on J if and only if

$$\Theta\left(\int_s^t V_H(r)f\,\mathrm{d}r\right) = 0 \quad \forall s,t \in J, s \leqslant t.$$

According to (4.10), this is equivalent to *2)*, i.e. *1)* $\Leftrightarrow$ *2)*. Since moreover

$$\int_s^t V_H(r)f\,\mathrm{d}r \in \mathscr{D}(\mathcal{A})_+,$$

statements *1)* and *3)* are equivalent by virtue of Corollary 4.12 and Proposition 5. Clearly *2)* implies *4)*. Assume now that the set $\left(\mathsf{B}^+ \int_s^t U_n(r)f\,\mathrm{d}r\right)_n$ is relatively weakly compact in L^1_+. Let us show that $\lim_{n\to\infty} \|\mathsf{B}^+ \int_s^t U_n(r)f\,\mathrm{d}r\|_{L^1_+} = 0$. According to (4.10), the limit

$$\lim_{n\to\infty} \left\|\mathsf{B}^+ \int_s^t U_n(r)f\,\mathrm{d}r\right\|_{L^1_+} := \ell(s,t)$$

exists. By Theorem 2.2, we also have

$$\left\|\mathsf{B}^+ \int_s^t U_n(r)f\,\mathrm{d}r\right\|_{Y_+} \leqslant \left\|\int_s^t U_n(r)f\,\mathrm{d}r\right\|_X + \left\|\mathcal{T}_{\max} \int_s^t U_n(r)f\,\mathrm{d}r\right\|_X$$

$$= \left\|\int_s^t U_n(r)f\,\mathrm{d}r\right\|_X + \|U_n(t)f\|_X - \|U_n(s)f\|_X\,,$$

and, since the series $\sum_n \int_s^t U_n(r)f\,\mathrm{d}r$, $\sum_n U_n(s)f$ and $\sum_n U_n(t)f$ are converging (towards $\int_s^t V_H(r)f\,\mathrm{d}r$, $V_H(s)f$ and $V_H(t)f$ respectively), one deduces that the right-hand-side is converging to 0 as $n \to \infty$ and

$$\lim_{n\to\infty} \left\|\mathsf{B}^+ \int_s^t U_n(r)f\,\mathrm{d}r\right\|_{Y_+} = 0. \tag{4.14}$$

Now, by assumption *4)*, there exists a subsequence $\left(\mathsf{B}^+ \int_s^t U_{n_k}(r)f\,\mathrm{d}r\right)_k$ which converges weakly to, say, $g_{s,t} \in L^1_+$. For any $i \in \mathbb{N}$ we set $\Gamma_{i,+} = \{z \in \Gamma_+ : \tau_-(\mathbf{z}) > \frac{1}{i}\}$ and denote by χ_i the characteristic function of the set $\Gamma_{i,+}$. Then for any $i \in \mathbb{N}$, the limit

$$\lim_{k\to\infty} \int_{\Gamma_{i,+}} \mathsf{B}^+ \int_s^t U_{n_k}(r)f\,\mathrm{d}r\,\mathrm{d}\mu_+ = \lim_{k\to\infty} \int_{\Gamma_+} \chi_i(\mathbf{z})\left(\mathsf{B}^+ \int_s^t U_{n_k}(r)f\,\mathrm{d}r\right)(\mathbf{z})\,\mathrm{d}\mu_+(\mathbf{z})$$

$$= \int_{\Gamma_{i,+}} g_{s,t}\,\mathrm{d}\mu_+.$$

Thus, from (4.14),

$$\lim_{k\to\infty} \int_{\Gamma_{i,+}} \mathsf{B}^+ \int_s^t U_{n_k}(r)f\,\mathrm{d}r\,\mathrm{d}\xi_+ = 0 \qquad \forall i \in \mathbb{N}.$$

Since, for any fixed $i \in N$ and any $\mathbf{z} \in \Gamma_{i,+}$ one has $\mathrm{d}\xi_+(\mathbf{z}) \geqslant \frac{1}{i}\mathrm{d}\mu_+(\mathbf{z})$ so that

$$\int_{\Gamma_{i,+}} g_{s,t}\mathrm{d}\mu_+ = 0.$$

Since $g_{s,t}$ is nonnegative on $\Gamma_+ = \bigcup_{i=1}^{\infty} \Gamma_{i,+}$, we deduce that $g_{s,t}(\mathbf{z}) = 0$ for μ_+-almost every $\mathbf{z} \in \Gamma_+$. In other words, the unique possible weak limit is $g_{s,t} = 0$ and therefore $\ell(s,t) = 0$, i.e. *2)* holds. □

Remark 5. We deduce directly from the above, with $J = [0,\infty)$ that the C_0-semigroup $(V_H(t))_{t\geqslant 0}$ is honest if and only if $\lim_{n\to\infty} \|\mathsf{B}^+ \int_0^t U_n(s)f\mathrm{d}s\|_{L_+^1} = 0$ for any $f \in X$ and $t > 0$.

Remark 6. Recall that in [19], in the free-streaming case, the defect function $[0,\infty) \ni t \to \eta_f(t)$ has been defined, for each fixed $f \in (\lambda - \mathcal{A})^{-1}X_+$, by $\eta_f(t) := \|V_H(t)f\| - \|f\| + c_\lambda(\int_0^t V_H(s)f\mathrm{d}s)$. We have already observed (see Remark 4) that c_λ of [19] corresponds to our functional $\widehat{\mathfrak{a}}$. Hence the defect function can be defined for each fixed $f \in X_+$, as

$$\eta_f(t) = -\Theta\left(\int_0^t V_H(s)f\mathrm{d}s\right) = -\lim_{n\to\infty}\left\|\mathsf{B}^+\int_0^t U_n(s)f\mathrm{d}s\right\|_{L_+^1}.$$

Such a representation of η_f allows to deduce immediately that the mapping $t \to \eta_f(t)$ is nonpositive and nonincreasing. Moreover, if the trajectory $(V_H(t)f)_{t\geqslant 0}$ is not honest then there exists $t_0 \geqslant 0$ such that $\eta_f(t) = 0$ for $0 \leqslant t \leqslant t_0$ and $\eta_f(t) < 0$ for all $t > t_0$. Setting $g = V_H(t_0)f \in X_+$. Then for any $t > 0$ one has

$$\eta_f(t + t_0) = -\Theta\left(\int_0^{t+t_0} V_H(s)f\mathrm{d}s\right)$$

$$= -\Theta\left(\int_0^{t_0} V_H(s)f\mathrm{d}s\right) - \Theta\left(\int_{t_0}^{t+t_0} V_H(s)f\mathrm{d}s\right)$$

$$= -\Theta\left(\int_0^t V_H(s)g\mathrm{d}s\right) = \eta_g(t) < 0,$$

i.e., with the terminology of [19], the trajectory $(V_H(t)g)_{t\geqslant 0}$ is immediately dishonest.

For any subinterval $J \subseteq [0,\infty)$ we denote by

$$X_J := \{f \in X_+ \; ; \; (V_H(t)f)_{t\geqslant 0} \text{ is honest on } J\}$$

and, whenever $J = [0,\infty)$, we simply denote $X_h = X_{[0,\infty)}$ the set of initial positive data giving rise to honest trajectories. Moreover, arguing exactly as in [6, Proposition 3.13], one sees that X_h is invariant under $(V_H(t))_{t\geqslant 0}$. Moreover, arguing exactly as in [19, Proposition 2.4], one has

Proposition 6. *For any subinterval $J \subseteq [0,\infty)$, one has $\widehat{X}_J := \mathrm{span}(X_J) = X_J - X_J$ is a closed lattice ideal of X whose positive cone is X_J. In particular, $\widehat{X}_h = \mathrm{span}(X_h)$ is a closed lattice ideal in X which is invariant under $(V_H(t))_{t\geqslant 0}$ and $(\widehat{X}_h)_+ = X_h$.*

We recall now that a positive semigroup $(T(t))_{t\geqslant 0}$ in X is said to be irreducible if there is no trivial closed ideal of X (i.e. different from X and $\{0\}$) which is invariant under $T(t)$ for all $t \geqslant 0$. We have then the following to be compared to [18, Theorem 19 & Remark 20]:

Proposition 7. *Let $g \in X_+$, $g \neq 0$ such that the trajectory $(V_H(t)g)_{t \geqslant 0}$ is honest.*

    1. If $(V_H(t))_{t \geqslant 0}$ is irreducible then the whole *semigroup $(V_H(t))_{t \geqslant 0}$ is honest.*

    2. If g is quasi-interior *then the* whole *semigroup $(V_H(t))_{t \geqslant 0}$ is honest.*

Proof. Let $g \neq 0$ such that $(V_H(t)g)_{t \geqslant 0}$ is honest be given.

(1) One has then $\widehat{X_h} \neq \{0\}$. Since $\widehat{X_h}$ is an ideal invariant under $(V_H(t))_{t \geqslant 0}$, if $(V_H(t))_{t \geqslant 0}$ is irreducible, this shows that necessarily $\widehat{X_h} = X$ and, in particular, $X_+ = X_h$.

(2) If g is quasi-interior, since $g \in \widehat{X_h}$ one has $\widehat{X_h} = X$ and the conclusion follows.
$\square$

    We have the following practical criterion extending [18, Theorem 21 & Corollary 23]

Proposition 8. *Assume that there exists some quasi-interior $h \in L^1_+$ such that*

$$H \, h(\Phi(\mathbf{z}, -\tau_-(\mathbf{z}))) \chi_{\{\tau_-(\mathbf{z}) < \infty\}} \leqslant h(\mathbf{z}) \qquad \text{for almost every } \mathbf{z} \in \Gamma_+. \qquad (4.15)$$

Then, the whole semigroup $(V_H(t))_{t \geqslant 0}$ is honest.

Proof. Let $h \in L^1_+$ satisfying the above assumption be given. Define then

$$f(\mathbf{x}) = \begin{cases} \dfrac{\tau_-(\mathbf{x})}{\tau_-(\mathbf{x}) + \tau_+(\mathbf{x})} \exp(-\tau_+(\mathbf{x})) \, h(\Phi(\mathbf{x}, \tau_+(\mathbf{x}))) & \text{if} \quad \tau_-(\mathbf{x}) + \tau_+(\mathbf{x}) < \infty, \\ \exp(-\tau_+(\mathbf{x})) \, h(\Phi(\mathbf{x}, \tau_+(\mathbf{x}))) & \text{if} \quad \tau_-(\mathbf{x}) = \infty \quad \text{and } \tau_+(\mathbf{x}) < \infty, \end{cases}$$

and f chosen freely on $\mathbf{\Omega}_{+\infty}$ in such a way that $f \in X$ is quasi-interior. One sees easily (see [2, Proposition 2.3] for details) that $\mathsf{B}^+ f = h$. Moreover, since $\tau_\pm(\mathbf{x}_t) = \tau_\pm(\mathbf{x}) \pm t$ and $\Phi(\mathbf{x}_t, \tau_+(\mathbf{x}_t)) = \Phi(\mathbf{x}, \tau_+(\mathbf{x}))$ for any $\mathbf{x} \in \mathbf{\Omega}$, $t > 0$, $\mathbf{x}_t = \Phi(\mathbf{x}, -t)$, one checks easily that, for any $\mathbf{x} \in \mathbf{\Omega}_+$, it holds

$$U_0(t)f(\mathbf{x}) = \frac{\tau_-(\mathbf{x}) - t}{\tau_-(\mathbf{x}) + \tau_+(\mathbf{x})} \exp(-t - \tau_+(\mathbf{x})) \, h(\Phi(\mathbf{x}, \tau_+(\mathbf{x})) \chi_{\{t < \tau_-(\mathbf{x})\}}$$

if $\mathbf{x} \in \mathbf{\Omega}_+ \cap \mathbf{\Omega}_-$, while

$$U_0(t)f(\mathbf{x}) = \exp(-t - \tau_+(\mathbf{x})) \, h(\Phi(\mathbf{x}, \tau_+(\mathbf{x})))$$

whenever $\mathbf{x} \in \mathbf{\Omega}_+ \cap \mathbf{\Omega}_{-\infty}$. Therefore, one sees that for any $t > 0$, $U_0(t)f(\mathbf{x}) \leqslant f(\mathbf{x})$ for almost every $\mathbf{x} \in \mathbf{\Omega}_+$. Let $t > 0$ be fixed. According to Lemma 3.6, one has

$$\left[\mathsf{B}^+ \int_0^t U_1(s)f\mathrm{d}s \right](\mathbf{z}) \leqslant \left[H\mathsf{B}^+ \int_0^t U_0(s)f\mathrm{d}s \right](\Phi(\mathbf{z}, -\tau_-(\mathbf{z})) \qquad \forall \mathbf{z} \in \Gamma_+.$$

Since $U_0(s)f \leqslant f$ on $\mathbf{\Omega}_+$ we get

$$\left[\mathsf{B}^+ \int_0^t U_1(s)f\mathrm{d}s \right](\mathbf{z}) \leqslant t \left[H \, \mathsf{B}^+ f \right](\Phi(\mathbf{z}, -\tau_-(\mathbf{z})) = tH \, h(\Phi(\mathbf{z}, -\tau_-(\mathbf{z}))) \quad \forall \mathbf{z} \in \Gamma_+.$$

From (4.15), one gets therefore

$$\left[\mathsf{B}^+ \int_0^t U_1(s)f\mathrm{d}s \right](\mathbf{z}) \chi_{\{\tau_-(\mathbf{z}) < \infty\}} \leqslant th(\mathbf{z}) \qquad \text{for a. e. } \mathbf{z} \in \Gamma_+.$$

Recalling that $\left[\mathsf{B}^+ \int_0^t U_1(s)f\mathrm{d}s \right](\mathbf{z}) = 0$ if $\tau_-(\mathbf{z}) > t$, we get therefore that

$$\left[\mathsf{B}^+ \int_0^t U_1(s)f\mathrm{d}s \right](\mathbf{z}) \leqslant th(\mathbf{z}) \qquad \text{for a. e. } \mathbf{z} \in \Gamma_+.$$

Using repeatedly Lemma 3.6 together with (4.15) we get that

$$\mathsf{B}^+ \int_0^t U_n(s) f \mathrm{d}s \leqslant th \qquad \forall t > 0, n \geqslant 1.$$

This shows that, for any $t > 0$, the family $\left(\mathsf{B}^+ \int_s^t U_n(r) f \mathrm{d}r \right)_n$ is equi-integrable and thus relatively weakly compact in L_+^1 and, according to Theorem 4.3, the trajectory $(V_H(t)f)_{t \geqslant 0}$ is honest. Since f is quasi-interior, Proposition 7 yields the conclusion. □

Besides the semigroup approach that we developed in the previous lines, it is also possible to develop a resolvent approach to honesty, as the one developed in [18] for the free-streaming case and in [5] for conservative boundary conditions. Such an approach provides necessary and sufficient conditions for a trajectory to be honest which are different from the one listed above. They can be seen as the analogue of [6, Theorem 3.5 & Theorem 3.11] which are established in the additive perturbation framework. Since we decided to mainly focus on the semigroup approach, we only state the result for the sake of completeness but omit the details of the proof which can be adapted without major difficulty from [18] and [6]:

Theorem 4.4. *Given $f \in X_+$, the following statements are equivalent*

1) the trajectory $(V_H(t)f)_{t \geqslant 0}$ is honest;

2) $\Theta \left((\lambda - \mathcal{A})^{-1} f \right) = 0$ for all $\lambda > 0$;

3) $\lim_{n \to \infty} \|(M_\lambda H)^n G_\lambda f\|_{L_+^1} = 0$ for some $\lambda > 0$;

4) $(\lambda - \mathcal{A})^{-1} f \in \mathscr{D}(\overline{T_H})$ for some $\lambda > 0$;

5) the set $((M_\lambda H)^n G_\lambda f))_n$ is relatively weakly compact in L_+^1 for some $\lambda > 0$.

In particular, the whole C_0-semigroup $(V_H(t))_{t \geqslant 0}$ is honest if and only if $\mathcal{A} = \overline{T_H}$.

Remark 7. Notice that, if one of the equivalent conditions *2)–5)* holds true of some $\lambda > 0$ then it can be proved that it holds true for *any* $\lambda > 0$ (see [9] for details).

Remark 8. It is possible to provide sufficient conditions for a trajectory to be honest which are reminiscent to those given in [19, Proposition 2.6]. Namely,

1. given $f \in X_+$, if there exists $\lambda > 0$ such that $(M_\lambda H) G_\lambda f \leqslant G_\lambda f$, then the trajectory $(V_H(t)f)_{t \geqslant 0}$ is honest;
2. if $g \in \mathscr{D}(\overline{\mathcal{T}_H})$ is such that $\mathcal{T}_H g \leqslant \lambda g$ for some $\lambda > 0$, then $g \in X_+$ and the trajectory $(V_H(t)g)_{t \geqslant 0}$ is honest.

5. **Some examples.** We illustrate here our approach by two examples. These two examples are dealing with the free-streaming equation conservative boundary and, as so, have already been dealt with in our previous contribution [5]. The scope here is to show that our new approach, based upon the semigroup representation (3.10), allows not only to recover, by different means, the results of [5], but also to characterize, in both examples, new interesting properties.

5.1. **An instructive one dimensional example revisited.** We revisit here a one-dimensional example introduced in [20, Example 4.12, p. 76]. This example has been revisited recently in both [5, 18]. Given two real nondecreasing real sequences $(a_k)_{k \geqslant 0}$ and $(b_k)_{k \geqslant 0}$ with

$$a_k < b_k < a_{k+1} \qquad \forall k \geqslant 0, \qquad \lim_{k \to \infty} a_k = \infty$$

set

$$\boldsymbol{\Omega} = \bigcup_{k=0}^{\infty} (a_k, b_k) =: \bigcup_{k=0}^{\infty} I_k.$$

We assume then μ to be the Lebesgue measure on $\mathbb{R}$ and consider the constant field $\mathscr{F} : \mathbb{R} \to \mathbb{R}$ given by $\mathscr{F}(x) = 1$ for all $x \in \mathbb{R}$. In such a case, the flow $\Phi(x, t)$ is given by

$$\Phi(x, t) = x + t \qquad \text{for any } x, t \in \mathbb{R},$$

and

$$\Gamma_- = \{a_k, k \in \mathbb{N}\}, \quad \Gamma_+ = \{b_k, k \in \mathbb{N}\}, \quad \tau_-(x) = x - a_k \quad \forall a_k < x < b_k, \quad k \in \mathbb{N}.$$

The measures $\mathrm{d}\mu_\pm$ are then the counting measures over $\Gamma_\pm$. We define then $H \in \mathscr{B}(L_+^1, L_-^1)$ by

$$H\psi(a_k) = \begin{cases} 0 & \text{if } k = 0, \\ b_{k-1} & \text{if } k > 0 \end{cases} \tag{5.1}$$

for any $\psi \in L_+^1$. It is clear that H is a positive boundary operator with unit norm. We then explicit the strongly family of operators $\{(U_k(t))_{t \geqslant 0} \, ; \, k \in \mathbb{N}\}$ as defined in Definition 3.4. To this aim for any $k \in \mathbb{N}$, set $\Delta_k = b_k - a_k$. For $f \in \mathcal{D}_0$ and $t > 0$ one easily sees that

$$U_0(t)f(x) = \begin{cases} f(x - t) & \text{if } 0 < t < x - a_k, \\ 0 & \text{otherwise}, \end{cases} \tag{5.2}$$

which yields

$$\mathsf{B}^+ U_0(t)f(b_k) = \begin{cases} f(b_k - t) & \text{if } 0 < t < \Delta_k, \\ 0 & \text{otherwise}. \end{cases} \tag{5.3}$$

By induction one can easily show that for $n \geqslant 1$, $k \geqslant 0$, $a_k < x < b_k$ one has

$$U_n(t)f(x) = \begin{cases} f(b_{k-n} - a_k + x + \sum_{j=k-n+1}^{k-1} \Delta_j - t) \\ \text{if } k \geqslant n \text{ and } \quad x - a_k + \sum_{j=k-n+1}^{k-1} \Delta_j < t < x - a_k + \sum_{j=k-n}^{k-1} \Delta_j, \\ 0 \qquad \text{otherwise} \end{cases} \tag{5.4}$$

so that

$$\mathsf{B}^+ U_n(t)f(b_k) = \begin{cases} f(b_{k-n} + \sum_{j=k-n+1}^{k} \Delta_j - t) \\ \text{if } \quad k \geqslant n \quad \text{and} \quad \sum_{j=k-n+1}^{k} \Delta_j < t < \sum_{j=k-n}^{k} \Delta_j, \\ 0 \qquad \text{otherwise}. \end{cases} \tag{5.5}$$

Because of this we have for all $f \in X$

$$\mathsf{B}^+ \left(\int_0^t U_n(s)f \, \mathrm{d}s \right)(b_k) = \begin{cases} \int_{a_{k-n} \vee (b_{k-n} + \sum_{j=k-n+1}^{k} \Delta_j - t)}^{b_{k-n}} f(s) \, \mathrm{d}s \\ \text{if} \quad k \geqslant n \quad \text{and} \quad t > \sum_{j=k-n+1}^{k} \Delta_j, \\ 0 \qquad \text{otherwise}. \end{cases} \tag{5.6}$$

Now we are able to prove the following where $(V_H(t))_{t \geqslant 0}$ is the C_0-semigroup constructed through Theorem 2.5 and given by $V_H(t) = \sum_{n=0}^{\infty} U_n(t)$ $(t \geqslant 0)$:

Proposition 9. *The C_0-semigroup $(V_H(t))_{t \geq 0}$ is honest if and only if*

$$\Delta := \sum_{k=0}^{\infty} (b_k - a_k) = \infty \tag{5.7}$$

If $\Delta < \infty$, define

$$J_k := \left[\sum_{j=k+1}^{\infty} \Delta_j, \sum_{j=k}^{\infty} \Delta_j \right] \subset [0, \infty) \qquad \text{for any} \quad k \in \mathbb{N}.$$

Then, given $f \in X_+$, the trajectory $(V_H(t)f)_{t \geq 0}$ is honest on J_k if and only if $\int_{a_k}^{b_k} f(s)\mathrm{d}s = 0$ which is equivalent to $f(s) = 0$ for almost every $s \in I_k$.

Remark 9. The first part of the Proposition is a well-known fact, first proven in [20] and revisited recently in [5, 18]. The second part, on the contrary, is new to our knowledge and provides a criterion for 'local' honesty.

Proof. Thanks to formula (5.6) we can state that for all $f \in X_+$, $n \geq 1$ and $t > 0$

$$\left\| \mathsf{B}^+ \int_0^t U_n(s)f\mathrm{d}s \right\|_{L_+^1} = \sum_{k=0}^{\infty} \left[\int_{a_k \vee (b_k + \sum_{j=k+1}^{k+n} \Delta_j - t)}^{b_k} f(s)\mathrm{d}s \right] \chi_{\{t > \sum_{j=k+1}^{k+n} \Delta_j\}}$$
$$\leq \sum_{k=0}^{\infty} \int_{a_k}^{b_k} f(s)\mathrm{d}s = \|f\|_X. \tag{5.8}$$

Assume first that $\Delta = +\infty$. In such a case for any $f \in X_+$, $k \in \mathbb{N}$ we have:

$$\lim_{n \to \infty} \left[\int_{a_k \vee (b_k + \sum_{j=k+1}^{k+n} \Delta_j - t)}^{b_k} f(s)\mathrm{d}s \right] \chi_{\{t > \sum_{j=k+1}^{k+n} \Delta_j\}} = 0$$

which, thanks to the dominated convergence theorem yields

$$\lim_{n \to \infty} \left\| \mathsf{B}^+ \int_0^t U_n(s)f\mathrm{d}s \right\|_{L_+^1} = 0.$$

By virtue of Corollary 5 the trajectory $(V_H(t)f)_{t \geq 0}$ is then honest. Since $f \in X_+$ is arbitrary, the semigroup itself is honest.

Now consider the case $\Delta < +\infty$. In such a case, given $k_0 \in \mathbb{N}$, let $t \in J_{k_0} \subseteq [0, \Delta] \subseteq [0, \infty)$. It is easy to see that

$$\lim_{n \to \infty} \left[\int_{a_k \vee (b_k + \sum_{j=k+1}^{k+n} \Delta_j - t)}^{b_k} f(s)\mathrm{d}s \right] \chi_{\{t > \sum_{j=k+1}^{k+n} \Delta_j\}}$$
$$= \begin{cases} 0 & \text{if } k = 0, \ldots, k_0 - 1 \geq 0, \\ \int_{a_k}^{b_k} f(s)\mathrm{d}s & \text{if } k \geq k_0 + 1, \end{cases}$$

while

$$\lim_{n \to \infty} \left[\int_{a_{k_0} \vee (b_{k_0} + \sum_{j=k_0+1}^{k_0+n} \Delta_j - t)}^{b_{k_0}} f(s)\mathrm{d}s \right] \chi_{\{t > \sum_{j=k_0+1}^{k_0+n} \Delta_j\}}$$
$$= \int_{(b_{k_0} + \sum_{j=k_0+1}^{\infty} \Delta_j - t)}^{b_{k_0}} f(s)\mathrm{d}s.$$

Then, from the dominated convergence theorem and (5.8)) we get

$$\lim_{n\to\infty} \left\| \mathsf{B}^+ \int_0^t U_n(s)f\mathrm{d}s \right\|_{L^1_+} = \int_{(b_{k_0}+\sum_{j=k_0+1}^\infty \Delta_j - t)}^{b_{k_0}} f(s)\mathrm{d}s + \sum_{k=k_0+1}^\infty \int_{a_k}^{b_k} f(s)\mathrm{d}s.$$

Therefore, for any $t_1, t_2 \in J_{k_0}$ with $t_1 < t_2$ one has

$$\lim_{n\to\infty} \left\| \mathsf{B}^+ \int_{t_1}^{t_2} U_n(s)f\mathrm{d}s \right\|_{L^1_+} = \int_{(b_{k_0}+\sum_{j=k_0+1}^\infty \Delta_j - t_2)}^{(b_{k_0}+\sum_{j=k_0+1}^\infty \Delta_j - t_1)} f(s)\mathrm{d}s \leqslant \int_{a_{k_0}}^{b_{k_0}} f(s)\mathrm{d}s$$

where the last inequality is an identity if $t_1 = \sum_{j=k_0+1}^\infty \Delta_j$ and $t_2 = \sum_{j=k_0}^\infty \Delta_j$. Using Corollary 5 again, this shows that the trajectory $(V_H(t)f)_{t\geqslant 0}$ is honest on J_{k_0} if and only if $\int_{a_{k_0}}^{b_{k_0}} f(s)\mathrm{d}s = 0$ and, being f nonnegative, this is equivalent to $f(s) = 0$ for almost every $s \in I_{k_0}$. □

Remark 10. As an immediate consequence of the obtained result we can state the following: if $\Delta < \infty$, then no trajectory $(V_H(t)f)_{t\geqslant 0}$ $(f \in X_+)$ is honest on an interval $J \supseteq [0, \Delta]$.

Remark 11. Finally observe that, in case $\Delta < +\infty$, for all $f \in X$ and $t \geqslant \Delta$ one has $V_H(t)f \equiv 0$. Indeed for $t \geqslant \Delta$ one has $U_0(t)f \equiv 0$. Furthermore for any $n \geqslant 1$, $k \geqslant 0$, one has $t > \sum_{j=k}^{k+n} \Delta_j$. This implies $U_n(t)f(x) = 0$ for all $n \geqslant 1$, $f \in \mathcal{D}_0$ and $x \in \mathbf{\Omega}$ which gives the result.

5.2. **Kinetic equation with specular reflections.** In this second example, we consider the physically relevant case of free-streaming semigroup associated to specular reflections. Such a model, as well-known [3], is strongly related to the so-called billiard flow which is a well-known dynamical system studies in ergodic theory [14]. We do not provide here any new honesty criterion but show how the result we obtained before yields possibly new property of the billiard flow. More precisely, we consider now a transport equation in $\mathbb{R}^N$ with $N = 2d$, $d \in \mathbb{N}$ and consider then

$$\mathbf{\Omega} = D \times \mathbb{R}^d \tag{5.9}$$

where D is a smooth open bounded and convex subset of $\mathbb{R}^d$. Any $\mathbf{x} \in \mathbf{\Omega}$ can be written $\mathbf{x} = (x, v)$, with $x \in D$, $v \in \mathbb{R}^d$ and consider the measure $\mathrm{d}\mu(\mathbf{x}) = \mathrm{d}x \otimes \mathrm{d}\varrho(v)$, where $\mathrm{d}\varrho$ is a positive Radon measure on $\mathbb{R}^d$ with support V. We assume for simplicity that V and $\mathrm{d}\varrho$ to be orthogonally invariant. Assume the field $\mathscr{F} : \mathbf{x} \in \mathbb{R}^N \to \mathbb{R}^N$ to be given by $\mathscr{F}(\mathbf{x}) = (v, 0)$ for all $\mathbf{x} = (x, v) \in R^d \times R^d$. Classically, the associated flow is given by $\Phi(\mathbf{x}, t) = (x + vt, v)$ for all $\mathbf{x} = (x, v)$ and $t \in \mathbb{R}$. In this case,

$$\Gamma_\pm = \{\mathbf{y} = (x, v) \in \partial D \times V \,;\, \pm v \cdot n(x) > 0\}, \qquad \mathrm{d}\mu_\pm(\mathbf{y}) = |v \cdot n(x)|\mathrm{d}\varrho(v)\mathrm{d}\sigma(x) \tag{5.10}$$

where $n(x)$ denotes the outward normal unit vector at $x \in \partial D$ and $\mathrm{d}\sigma(\cdot)$ is the Lebesgue surface measure on ∂D. We consider here the boundary operator H to be associated by the specular reflection, i.e.

$$H\psi(\mathbf{y}) = \psi(x, v - 2(v \cdot n(x))n(x)), \qquad \psi \in L^1_+, \ \mathbf{y} = (x, v) \in \Gamma_-.$$

It is known that H is a positive and conservative operator, i.e. $\|H\psi\|_{L^1_-} = \|\psi\|_{L^1_+}$ for any nonnegative $\psi \in L^1_+$. In particular, H has unit norm. As in the previous

example, let us characterize the families $(U_k(t))_{t\geqslant 0}$, $k \in \mathbb{N}$. Observe that for $\mathbf{x} = (x, v) \in \overline{D} \times V$ we can define the sequence of *rebound times*:

$$t_1(\mathbf{x}) = \tau_-(\mathbf{x}), \qquad t_2(\mathbf{x}) = t_1(\mathbf{x}) + \tau_-(\mathbf{x}_1) \quad \dots$$

$$t_k(\mathbf{x}) = t_{k-1}(\mathbf{x}) + \tau_-(\mathbf{x}_{k-1}) = \sum_{j=1}^{k-1} \tau_-(\mathbf{x}_j)$$

where , setting $\mathbf{x}_0 = \mathbf{x}$, $x_0 = x$, $v_0 = v$, one has, for any $j = 1, \dots, k$:

$$\mathbf{x}_j = (x_j, v_j) \in \Gamma_+, \qquad \text{with} \quad x_j = x_{j-1} - \tau_-(\mathbf{x}_{j-1})v_{j-1}, \in \partial D,$$
$$v_j = v_{j-1} - 2(v_{j-1} \cdot n(x_{j-1}))n(x_{j-1}).$$

With this notations, setting also $t_0 = 0$, we have for any $f \in \mathcal{D}_0$, $k \geqslant 0$, $\mathbf{x} = (x, v) \in \Omega$

$$U_k(t)f(\mathbf{x}) = f(x_k - (t - t_k(\mathbf{x}))v_k, v_k)\,\chi_{\{t_k(\mathbf{x})\leqslant t < t_{k+1}(\mathbf{x})\}}.$$

Recall that $V_H(t) = \sum_k U_k(t)$ and, since for a given $t > 0$ and a given $\mathbf{x} \in \Omega$, there exists a unique $k \in \mathbb{N}$ such that $t \in [t_k(\mathbf{x}), t_{k+1}(\mathbf{x})]$, one has $V_H(t)f(\mathbf{x}) = U_k(t)f(\mathbf{x})$ (of course, such k depends on $\mathbf{x}$). This implies that, for any $t \geqslant 0$,

$$V_H(t)f = f \circ \vartheta_t$$

where $\{\vartheta_t; t \in \mathbb{R}\}$ is the one-parameter group of transformations on $\overline{D} \times V$ corresponding to the so-called *billiard flow* [14] (see also [4]). The following is taken from [20, 5] and is proven through the resolvent approach:

Proposition 10. *Assume that* $d\sigma(\partial D) < \infty$. *If there exists some nonnegative* $\psi \in L^1(V, d\varrho(v))$ *with* $\psi(v) = \psi(|v|)$ *for any* $v \in V$ *and*

$$\int_V (1 + |v|)\,\psi(|v|)d\varrho(v) < \infty$$

then the semigroup $(V_H(t))_{t\geqslant 0}$ *is honest.*

Proof. The proof follows from Proposition 8 since, as in [18, Corollary 2.3], the mapping $h(x, v) = \psi(|v|)$ provides a quasi-interior element of L^1_+ satisfying (4.15). $\square$

Remark 12. From Remark 5 we deduce that for any $t > 0$ and $f \in X_+$

$$\lim_{k\to\infty} \left\| \mathsf{B}^+ \int_0^t U_k(s)f ds \right\|_{L^1_+} = 0,$$

i.e.

$$\lim_{k\to\infty} \int_{\Gamma_+} \left(\int_{t_k(\mathbf{x})}^{t\wedge t_{k+1}(\mathbf{x})} f(x_k - v_k(s - t_k(\mathbf{x}))v_k, v_k)ds \right) \chi_{\{t_k(\mathbf{x})<t\}} d\mu_+(\mathbf{x}) = 0.$$

We have not been able to find such a property of the billiard flow in the literature.

5.3. **Kinetic equations with Maxwell-like boundary conditions.** We consider here, as a final example, the case of the free-streaming semigroup associated to Maxwell-like boundary conditions which is particularly relevant in the kinetic theory of gases [13]. For such a model, the mathematical framework is the one described in the previous paragraph, in particular $\boldsymbol{\Omega}$ is given by (5.9) and $\Gamma_{\pm}$ are described by (5.10). The boundary operator is given by

$$H\psi(x,v) = \alpha(x)\psi(x, v - 2(v \cdot n(x))n(x)) +$$

$$(1 - \alpha(x))M(v)\int_{\{w \in V \,;\, w \cdot n(x) > 0\}} \psi(x, w)|w \cdot n(x)|\mathrm{d}w, \qquad (x, v) \in \Gamma_{-} \quad (5.11)$$

where, for simplicity $V = \mathbb{R}^d$ endowed with the Lebesgue measure, the so-called accommodation coefficient $\alpha \,:\, \partial D \to \mathbb{R}$ is measurable and $0 \leqslant \alpha(x) \leqslant 1$ for $\mathrm{d}\sigma$-almost every $x \in \partial D$ (where $\mathrm{d}\sigma$ is the Lebesgue measure over ∂D). The distribution M is the Maxwell distribution

$$M(v) = C_d \exp\left(-\frac{|v|^2}{\theta}\right) \qquad v \in V$$

where $\theta > 0$ is the temperature of the surface ∂D (assumed here to be constant for simplicity) while $C_d > 0$ is a normalisation constant such that

$$\int_{\{v = (v_1, \ldots, v_d) \in V \,, v_1 > 0\}} M(v)v_1 \mathrm{d}v = 1.$$

This boundary operator models the fact that, when hitting the boundary ∂D a fraction $\alpha(\cdot)$ of particles undergoes a specular reflection while the remaining fraction $1 - \alpha(x)$ is diffused with the Maxwellian distribution M of the wall ∂D (see [13] for more details). Then, one has the following:

Proposition 11. *Let H be given by* (5.11). *If $\mathrm{d}\sigma(\partial D) < \infty$ then the semigroup $(V_H(t))_{t \geqslant 0}$ is honest.*

Proof. Again, the proof follows from Proposition 8. Indeed, one checks as in [5, Proposition 5.8] that the mapping $h(x,v) = M(v)$ provides a quasi-interior element of L^1_+ satisfying (4.15). □

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