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## Characterization of a neutron imaging setup at the INES facility

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# UNIVERSITÀ DEGLI STUDI DI TORINO

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# 1 CHARACTERIZATION OF A NEUTRON

## 2 IMAGING SETUP AT THE INES FACILITY

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## 22 ABSTRACT

23 The Italian Neutron Experimental Station (INES) located at the ISIS pulsed neutron source  
24 (Didcot, United Kingdom) provides a thermal neutron beam mainly used for diffraction analysis.  
25 A neutron transmission imaging system was also developed for beam monitoring and for aligning  
26 the sample under investigation. Although the time-of-flight neutron diffraction is a consolidated  
27 technique, the neutron imaging setup is not yet completely characterized and optimized. In this  
28 paper the performance for neutron radiography and tomography at INES of two scintillator screens  
29 read out by two different commercial CCD cameras is compared in terms of linearity, signal-to-  
30 noise ratio, effective dynamic range and spatial resolution. In addition, the results of neutron  
31 radiographies and a tomography of metal alloy test structures are presented to better characterize  
32 the INES imaging capabilities of metal artifacts in the cultural heritage field.

33 *Keywords: Neutron imaging; Metal alloy; Cultural heritage*

## 34 1. Introduction

35 Analysis techniques based on thermal and epithermal neutrons represent a  
36 powerful tool to characterize metal alloy artifacts in the cultural heritage field.  
37 Neutrons are a non-invasive diagnostic tool and can penetrate thick layers of  
38 samples. Neutron transmission imaging, obtained with radiographic or

tomographic techniques, allows to investigate the inner structure of artworks, determining, as an example, the materials thickness and the conservation condition [1, 2]. All this information can be helpful to assist restoration activities, to increase the knowledge of the work of art and to help dating or attributing an artifact through the understanding of the manufacturing techniques used.

The Italian Neutron Experimental Station (INES) is a neutron diffractometer located at the ISIS pulsed neutron source [3]. A low-cost imaging device is currently installed and used for beam monitoring and samples positioning. Preliminary measurements were performed by Bartoli et al. [4] to test the imaging setup. Salvato et al. realized an upgrade of the system providing the first neutron tomography investigation at INES [5].

The aim of this paper is to fully characterize the present imaging setup in terms of linearity, signal-to-noise ratio, effective dynamic range and spatial resolution. A comparison with a different converter screen and a different CCD camera will be shown. Moreover, the analysis of neutron radiography and tomography of metal alloy test structures will be presented in order to assess the diagnostic potential of the INES imaging apparatus on metal artifacts in the cultural heritage field.

56

## 57 **2. Experimental set up**

The Italian Neutron Experimental Station, INES, provides a stable pulsed neutron beam with an approximately uniform intensity distribution over a cross section, at the sample position, of 35x35 mm<sup>2</sup> (40x40 mm<sup>2</sup> including also non-uniform borders). The divergence of the beam, in terms of L/D ratio, is about 90.

The neutron imaging setup currently installed consists of a scintillator converter screen, a mirror and a commercial CCD camera [6]. The scintillator is made of ZnS/<sup>6</sup>LiF, 225 μm thick, layered on an aluminum substrate; in the following we will refer to it as scintillator 1. It is positioned orthogonal to the incident neutron beam, at a distance from the sample ranging from 1 cm to 10 cm and it emits visible light centered at a wavelength of about 520 nm, in order to best match the CCD sensitivity curve. A mirror reflects the emitted light toward a black and white camera (The Imaging Source DMK21BF04) equipped with an optical system of f=16mm and F/1.4. The camera works at room temperature without any image intensifier and is based on a Sony ICX098BL CCD sensor with 640x480

72 pixels of 5.6  $\mu\text{m}$  size, one pixel corresponding to 100  $\mu\text{m}$  in the image field, read  
73 out by a 10 bit ADC whose most significant 8 bits encode the CCD output signal.

74 In order to test how the performance depends on the components of the existing  
75 system, a second converter screen, named scintillator 2 in the following, and a  
76 different CCD camera were also tested. The converter has the same composition  
77 as scintillator 1, embedded in a stable plastic matrix, and has a total thickness of  
78 450  $\mu\text{m}$ . The CCD camera is a Manta G-032B (Allied Vision Technology)  
79 equipped with the same optical system of  $f=16$  mm and  $F/1.4$ . The CCD sensor is  
80 a Sony ICX424 with 656x492 pixel of 7.4  $\mu\text{m}$  size, one pixel corresponding to  
81 100  $\mu\text{m}$  in the image field, read out by a 12 bit ADC.

82 An x-y translator, positioned inside the sample tank, is available for sample  
83 alignment [6]. A precision rotary stage (Newport URS150-BPP) was added to the  
84 existing setup to perform a neutron tomography.

85 In order to test the diagnostic potential of the apparatus on metal artifacts, test  
86 samples were prepared by the INFN (National Institute of Nuclear Physics)  
87 mechanical workshop in Torino. The densities and material compositions of these  
88 samples are reported in Table 1.

89 In particular, to probe the imaging capabilities of the apparatus as a function of  
90 material composition and thickness, three step wedges of steel, brass and  
91 quaternary bronze alloy were prepared (figure 1 a). These three materials were  
92 chosen as they constitute by far the most common metal alloys used in the past for  
93 cultural artifacts. Each step is a square of 10x10  $\text{mm}^2$  with a thickness  
94 incrementing in each step by 2 mm, from 2 mm up to a maximum of 30 mm. In  
95 addition, to investigate the capability of neutron tomography of metal alloy  
96 objects, a 25 mm edge bronze cube (figure 1 b) was built to fit in the spot size of  
97 the INES neutron beam. The cube faces are held together by means of screws and  
98 welds and have a thickness of 4 mm, reduced to 2 mm for two opposite side faces,  
99 to mimic the typical bronze thicknesses of statues. Two circular holes of 1 mm  
100 diameter were added on one of the two thinner faces to simulate air bubbles and  
101 three square base rods, 3x3x15  $\text{mm}^3$  size, two made of steel and one of aluminum,  
102 were inserted in the central part to simulate the presence of inner structures.

### 103    **3. Results**

104    The two scintillator screens coupled with both the CCD cameras were tested in  
105    order to determine the linearity, the signal-to-noise ratio, the effective dynamic  
106    range and the spatial resolution separately for all the combinations. Later, the best  
107    combinations were used to analyze the metallic test samples.

108

#### 109    **3.1 Linearity**

110    For this, a set of images of the light emitted by the converter screen illuminated by  
111    the open neutron beam was collected for increasing exposure times to simulate the  
112    behavior with increasing signal intensities. A square region of interest (ROI) of  
113    10x10 pixels, corresponding to an area of approximately 1 mm<sup>2</sup> of the image  
114    field, was selected in the central area of the images, where the neutron beam is  
115    uniform, and was used to calculate the average ADC value (or gray level) and its  
116    standard deviation (or noise). These quantities are derived with the assumption  
117    that the pixels gray level fluctuations in the ROI of a single image are equivalent  
118    to the gray level fluctuations of a single pixel in a sequence of identical images.

119    The measurements with the DMK camera, with both the scintillator screens, were  
120    performed by setting the electronic gain to 24.6 dB (available range is 0 - 36 dB)  
121    and increasing the exposure times up to the maximum allowed of 30 s. Similar  
122    measurements with the Manta camera were performed setting the electronic gain  
123    to 15 dB (available range is 0 dB- 36 dB) and increasing the exposure time up to  
124    the maximum available of 60 s. The results are shown in figure 2 and 3 where the  
125    gray levels are displayed as a function of the exposure time. It should be noticed  
126    that, at the maximum exposure time, the signal range of the DMK camera covers  
127    only about one half of the available 256 gray levels. On the contrary, at the  
128    maximum exposure time, the Manta camera reaches with the scintillator 2 the  
129    saturation level of 4095; this value is therefore not included in the analysis.

130    The data indicate a good linearity of the imaging systems as resulting from the fits  
131    to a straight line shown in the figures 2-3, the maximum deviation from linearity  
132    being 1 gray level for DMK with both scintillators and 25 and 12 gray levels for  
133    Manta with scintillator 1 and scintillator 2 respectively. In addition, the slope of  
134    the fit for the thicker converter, scintillator 2, is always larger, as expected as a  
135    consequence of the larger conversion efficiency. For a fixed exposure time this

136 leads to a slightly larger effective dynamic range, as will be shown in the  
137 following section.

138

### 139 **3.2 Effective dynamic range**

140 First the signal-to-noise ratio was studied, where the signal and the noise were  
141 determined for increasing exposure times as described in the previous section. It is  
142 found that, for all the different setup configurations, the signal-to-noise ratio as a  
143 function of the gray level signal can be approximated by a power curve with a  
144 fractional exponent; an example it is shown in figure 4 for the Manta camera  
145 coupled with scintillator 2. This behavior indicates that the dominant contribution  
146 to the noise is originated from the statistical fluctuations in the number of photons  
147 hitting the CCD cells and thus increases with the signal level; in fact, assuming a  
148 pure Poisson statistical process, the signal-to-noise ratio should behave like a  
149 power function with exponent equal to 0.5.

150 In order to compare the effective dynamic range of the different combinations of  
151 camera and converter, the method described in [8] was used. The effective  
152 dynamic range can be expressed as the number of effective distinguishable gray  
153 levels  $L$  of the system, depending on both the signal  $s$  and the noise  $N(s)$ , through  
154 the formula

$$155 \quad L = \int \frac{ds}{N(s)} \quad (1)$$

156 where the integration is performed over the useful signal range. Since this range  
157 increases with the exposure time and given the poissonian nature the noise, the  
158 comparison should be performed for equal exposure times, i.e. for equal number  
159 of neutrons hitting the converter, and the signal integration ranges should be  
160 chosen accordingly.

161 As an example, figure 5 shows  $1/\text{noise}$  as a function of the gray level for the same  
162 camera and scintillator of figure 4. For each scintillator-camera combination the  
163 data were fitted with a power function and, for each exposure time, the result was  
164 integrated from dark signal (8.33 and 9.13 gray levels for DMK and Manta  
165 respectively) up to the measured signal level shown in the figures 2 and 3.

166 The results are shown in figure 6 where the effective number of gray levels  $L$  is  
167 represented as a function of the exposure time. For each exposure time the  
168 scintillator 2 is providing a slightly better dynamic range than the scintillator 1

and so does the Manta camera compared to the DMK. The best combination is found to be the Manta CCD camera coupled with scintillator 2 which reaches about 80 effective gray levels before saturation. On the contrary the present INES imaging system comprising the DMK camera is limited in dynamic range mainly because of the shorter maximum exposure time of the camera. Table 2 summarizes the maximum dynamic ranges, corresponding to the maximum exposure time before saturation, expressed for all the combinations both in terms of effective number of gray levels  $L$  and in dB through the formula

$$D(\text{dB}) = 20 \log L \quad (2)$$

It should be observed that even if these results were obtained in a limited central area of the image of  $1 \text{ mm}^2$ , their validity can be extended to the entire beam active area as the maximum non-uniformity was measured to be below 1% (5%) sampling a region of  $1 \text{ mm}^2$  in a central area of  $24 \times 24 \text{ mm}^2$  ( $35 \times 35 \text{ mm}^2$ ).

### 3.3 Spatial resolution

The spatial resolution of the radiographic system was evaluated by analyzing the image of a high contrast sharp edge, realized with a Gadolinium slab, in terms of the Edge Spread Function (ESF) and the Modular Transfer Function (MTF) [9]. The edge was placed at around 10 cm far from the scintillator in order to reproduce the measurement conditions required in the tomography set-up. Since the Gd slab was not perfectly aligned with the CCD, the sharp edge image was fitted with a modified Cauchy functional form:

$$\phi(x, y)_{a,b,\beta,\lambda,x_0} = a \left\{ \frac{1}{2} + \frac{1}{\pi} \arctan[\lambda(x - \beta y - x_0)] \right\} + b \quad (3)$$

where  $a$  and  $b$  are constants,  $x$  and  $y$  are the coordinates in the image plane,  $x_0$  is the edge position,  $\beta$  is the edge slope and  $\lambda$  is a parameter related to the full width half maximum (FWHM) of the ESF derivative. In fact it can be shown that

$$\text{FWHM} = \frac{2}{\lambda} \cos(\tan^{-1} \beta) \approx \frac{2}{\lambda} \left( 1 - \frac{\beta^2}{2} \right) \quad (4)$$

where the last approximation is valid for small values of  $\beta$ .

The spatial resolution FWHM obtained with this method for all combinations of CCD cameras and scintillators is summarized in Table 3.

The corresponding MTF was obtained from the ESF derivative by Fourier transform and the spatial frequency corresponding to the MTF value of 0.1 was



201 considered as the spatial resolution of the imaging system in lp/mm. The results  
202 are summarized in Table 4.

203 As expected, the scintillator 1, being thinner than scintillator 2, shows a better  
204 spatial resolution. It is also found that the spatial resolution of the Manta camera  
205 appears worse than with the DMK; this can be partially explained by considering  
206 that the optical system used for the experiment was specifically optimized for the  
207 use with the DMK camera.

208

### 209 **3.4 Study of the test samples**

210 In order to test the imaging capabilities of the apparatus as a function of metal  
211 type and thickness, neutron radiographies of the three step wedges made of steel,  
212 brass and quaternary bronze alloy were acquired. Additional dark and open beam  
213 images were acquired for later correction. Only the Manta CCD camera coupled  
214 with scintillator 2 was used in order to exploit the highest effective dynamic  
215 range. Images were acquired setting the integration time to 45 s to avoid  
216 saturation effects, and the electronic gain to 15 dB.

217 To reduce the randomly distributed noise due to scattered gamma rays, the raw  
218 images were processed with a radius 3 median filter [10]. The final positive  
219 radiographic image ( $I_F$ ) was obtained by subtracting from the raw image ( $I$ ) the  
220 dark image ( $D$ ) and normalizing to the beam image ( $B$ ) to correct for the  
221 inhomogeneities of the beam, according to the formula:

$$222 \quad I_F = \frac{(I) - (D)}{(B) - (D)} \quad (5)$$

223 Figure 7 shows the radiographic image of the three step wedges and the  
224 corresponding gray level profiles, calculated averaging over 20 rows of pixels in  
225 the centre of the wedge. Here, zero level corresponds to complete absorption  
226 (black) and one to no absorption (white). Given the good linearity of the system,  
227 the profiles can also be interpreted as neutron beam intensity profiles. As  
228 expected, for all the three wedges, the gray level of each step decreases with the  
229 increase of the thickness (right to left in the figure), up to the largest thickness  
230 where the noise contribution starts to dominate. It can be deduced from this figure  
231 that only those metallic art objects having a metal thickness not exceeding 30 mm  
232 can be conveniently investigated with this system. The separation among the three  
233 profiles originates from the differences of the neutron cross sections for the

different wedge materials. For reference, figure 8 displays the total cross sections of the three constituent materials as a function of the neutron energy calculated with the MCNP4C2 simulation code [11] in the INES neutron energy range (7.8 meV – 5 eV). Such a difference in the cross section, combined with a 3D tomographic reconstruction, could be exploited to differentiate the constituting metal of hidden parts of the object.

Figure 9 shows, for each of the three wedges, the attenuation of the neutron beam as a function of the thickness of the wedge step. Since the intensity  $I$  of a neutron beam traversing a material thickness  $x$  is attenuated according to the exponential law  $I = I_0 e^{-\Sigma x}$ , it is possible by performing a fit to extract the macroscopic cross section  $\Sigma$  separately for the three materials. Table 5 shows the results in terms of cross section over density ( $\Sigma/\rho$ ) compared to the predictions obtained by processing the ENDF/B VI microscopic cross sections with the MCNP4C2 code and extracting the expected  $\Sigma$  at the neutron beam energy of 1 eV; this energy corresponds to the average of the INES beam energy spectrum at the sample position [3]. The agreement is found to be satisfactory within the statistical uncertainty.

Finally, in order to investigate the potential of neutron tomography on metal alloy objects, a tomography of the bronze cube was performed. Here the DMK camera and the scintillator 1 were used as they showed to provide the best spatial resolution. The cube, placed 10 cm far from the scintillator screen, was rotated by angular steps of  $0.7^\circ$  up to a total of  $179.9^\circ$  and a projection was acquired at each angular position using an exposure time of 16 s.

To perform the 3D volume reconstruction, the algorithm of Filtered Back-Projection (FBP), developed for X-ray tomography and implemented in the Imgreco software [12], was used. Parallel beam geometry was assumed which allows the reconstruction using projections over  $180^\circ$  [13].

The tomography of the cube proved the possibility to resolve details at a millimetric scale inside metal structures. Figure 10 shows a horizontal section and a 3D rendering of the cube. Welds and holes are clearly visible in the horizontal section. Note that in the figure the aluminum column is invisible as the aluminum cross section is about one order of magnitude lower than steel cross section. However, changing the gray level threshold, the 3D rendering of the cube worsens but the aluminum column becomes visible.

## 268 **4. Conclusions**

269 Series of measurements were performed to characterize the imaging capabilities  
270 of the INES setup for objects of interest in the cultural heritage field. Different  
271 combinations using a thicker scintillator screen and a different CCD camera with  
272 a larger maximum bit depth were also tried to find possible improvements to the  
273 existing system.

274 For all the combinations a good linearity of the imaging system was obtained.  
275 The best dynamic range of about 80 effective gray levels was obtained with a  
276 Manta 12 bit CCD camera coupled with the thicker scintillator. On the contrary  
277 the present INES imaging system, comprising a DMK 8 bit CCD camera and the  
278 thinner scintillator, was limited to about 43 effective gray level mainly because of  
279 the short maximum exposure time of this camera. However this combination  
280 showed the best spatial resolution of 182  $\mu\text{m}$  (FWHM), corresponding to 4.0  
281 lp/mm.

282 The combination featuring the best dynamic range was used to investigate the  
283 imaging capabilities for metal alloy objects as a function of the metal thickness  
284 and alloy composition. For this purpose, three step wedges made of steel, brass  
285 and bronze were analyzed, showing that the investigation of metallic objects up to  
286 a thickness of about 30 mm can be achieved at this facility. The neutron  
287 attenuation coefficients for the three metal alloys were also measured from the  
288 attenuation curves and compared with the predictions of the ENDF/BVI library,  
289 showing a good agreement.

290 Finally, the tomography of a hollow bronze cube with wall thicknesses in the  
291 range 2-4 mm, typical of the bronze statues, and with inner insertions of steel and  
292 aluminum parts, was carried on using the present INES setup in order to exploit  
293 the best spatial resolution. The results showed that the presence of metallic  
294 structures inside the bronze shell can be easily revealed with this technique, and  
295 that holes and welding material at a millimetric scale can be revealed.

## 296 **Acknowledgements**

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## 300    **References**

- 301    1 Lehmann E H et al. Nuclear Instruments and Methods in Physics Research A “Non-invasive  
302    studies of objects from cultural heritage” (2005), 542 68–75.
- 303    2 Lehmann E H et al. Archaeometry “Investigation of the content of ancient Tibetan metallic  
304    Buddha statues by means of neutron imaging methods” (2010) 52, 3:416-428.
- 305    3 Grazzi F et al. Nuovo Cimento “Preliminary results of the Italian neutron experimental station  
306    INES at ISIS: Archaeometric applications” (2007) 30 59-65.
- 307    4 Bartoli L et al. Nuclear Instruments and Methods in Physics Research A “Test measurements  
308    with a new imaging alignment camera at ISIS” (2008) 595, 643-646.
- 309    5 Salvato G et al. Nuovo Cimento “Neutron Tomography at INES: first experimental results”  
310    (2008) 31 C, N. 4.
- 311    6 Grazzi F et al. Review of Scientific Instruments “A neutron imaging device for sample  
312    alignment in a pulsed neutron scattering instrument” (2009) 80 1-4.
- 313    7 Data provided by the material supplier.
- 314    8 Bettuzzi M et al. Proceedings of SPIE “Effective dynamic range measurement for a CCD in full-  
315    field industrial X-ray imaging applications” (2007) 6616 66161L.
- 316    9 Hall E L “Computer image processing and recognition” (1979) Academic Press, New York.
- 317    10 <http://rsb.info.nih.gov/ij/>
- 318    11 Briesmeister J F E “MCNP—a general Monte Carlo N-particle transport code” Version 4C,  
319    LA-13709 (2000) Los Alamos National Laboratory Report.
- 320    12 Developed by D Schneberk at Lawrence Livermore National Laboratory, see also Martz H E et  
321    al. “X-ray Imaging” (2009) CRC Press Inc, Bosa Roca.
- 322    13 Kak A C and Slaney M “Principles of Computerized Tomographic Imaging” (1988), IEEE  
323    Press.

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## 337 **Figures Captions**

338 **Fig. 1** The three step wedges (a) and the bronze cube (b).

339

340 **Fig. 2** DMK camera: linearity of the signal as a function of the exposure time.

341

342 **Fig. 3** Manta camera: linearity of the signal as a function of the exposure time.

343

344 **Fig. 4** Manta camera: signal to noise ratio as a function of the gray level signal.

345

346 **Fig. 5** Manta camera:  $1/\text{noise}$  as a function of the gray level.

347

348 **Fig. 6** Effective number of gray levels as a function of the exposure time for the four different  
349 combinations of camera and converter.

350

351 **Fig. 7** Neutron radiography of the three step wedges (above) and corresponding neutron absorption  
352 profiles (below). The upper scale of the plot indicates the step thickness.

353

354 **Fig. 8** Total neutron cross sections of the constituent materials of the step wedges calculated with  
355 MCNP4C2 simulation code.

356

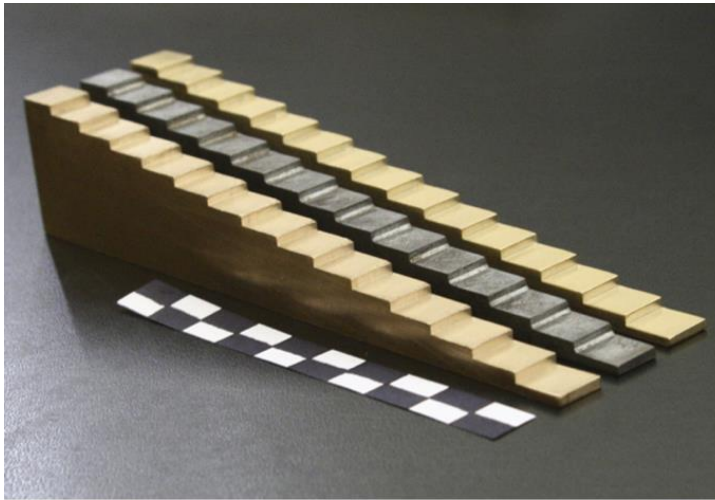
357 **Fig. 9** Neutron attenuation curves for the three wedges as a function of the step thickness.

358

359 **Fig. 10** Cube phantom tomography reconstruction: section (a) and 3D rendering (b).

360

361 **Figure 1**



( a )

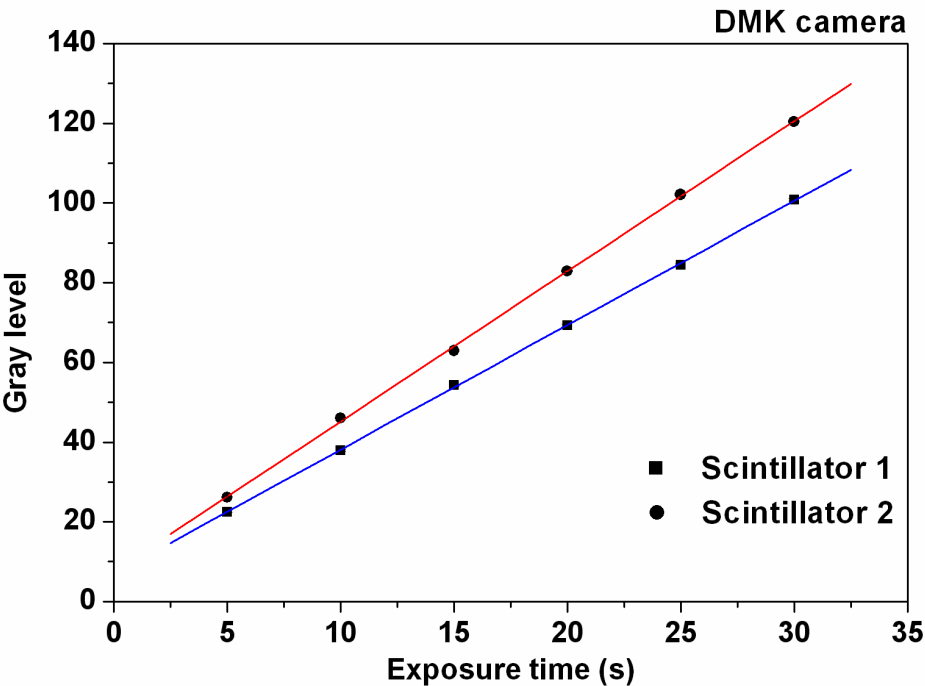


( b )

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363

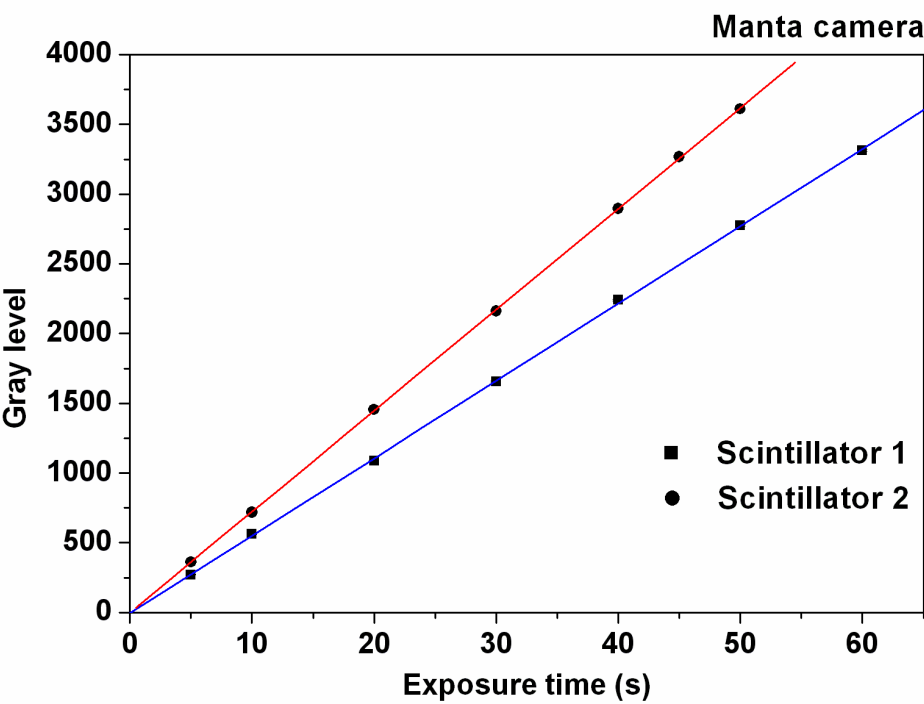
364 **Figure 2**



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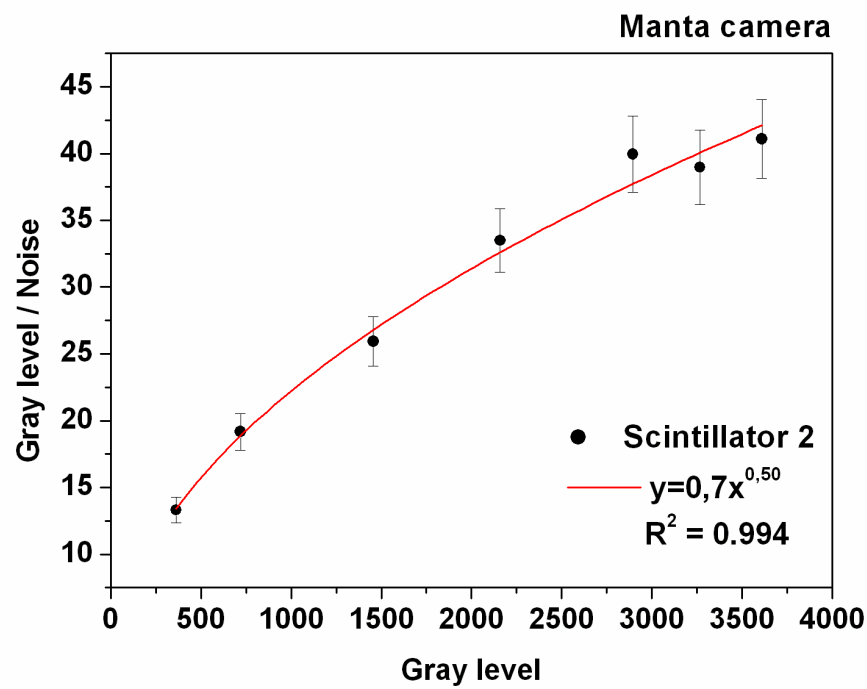
367 **Figure 3**



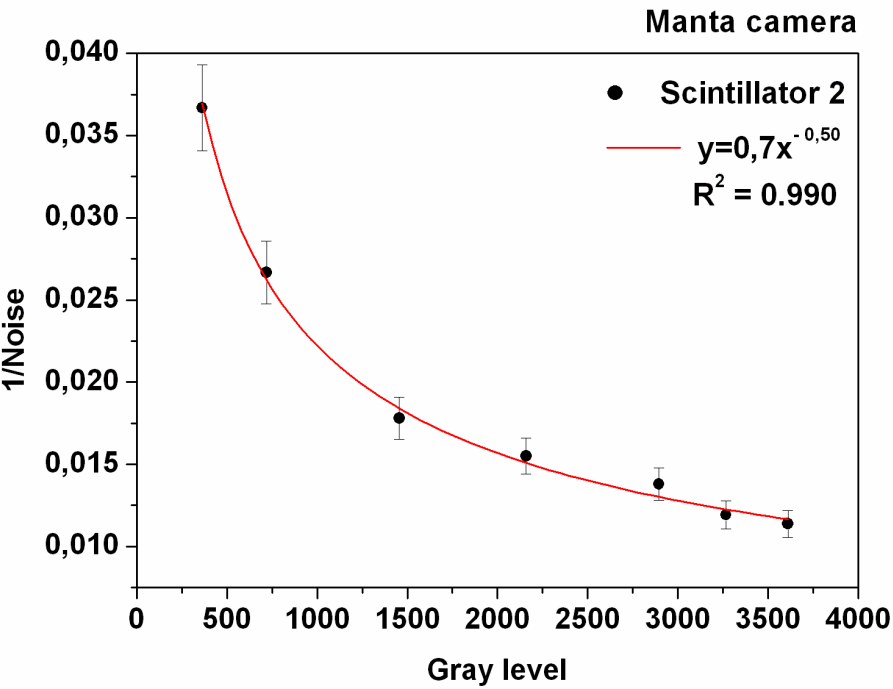
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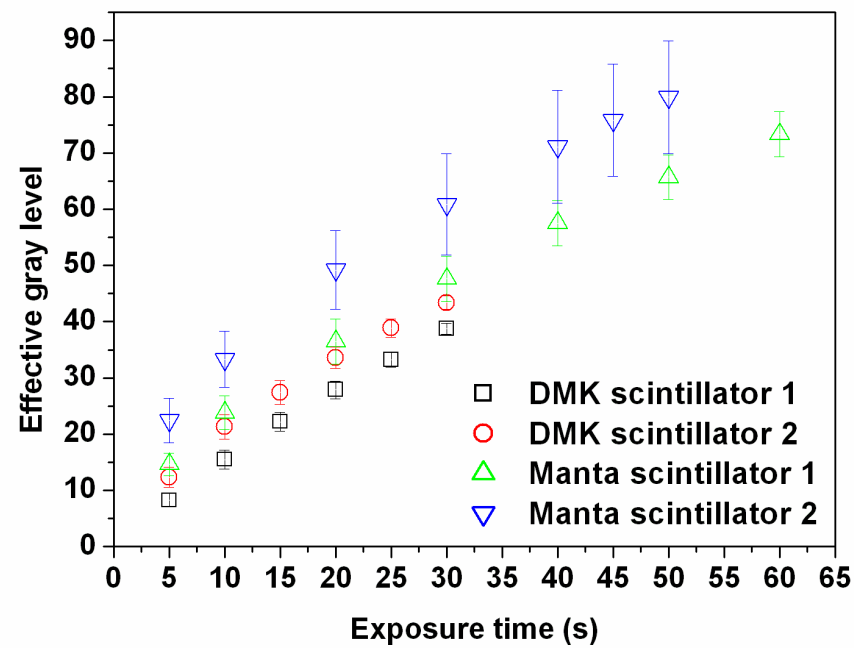
372 **Figure 5**



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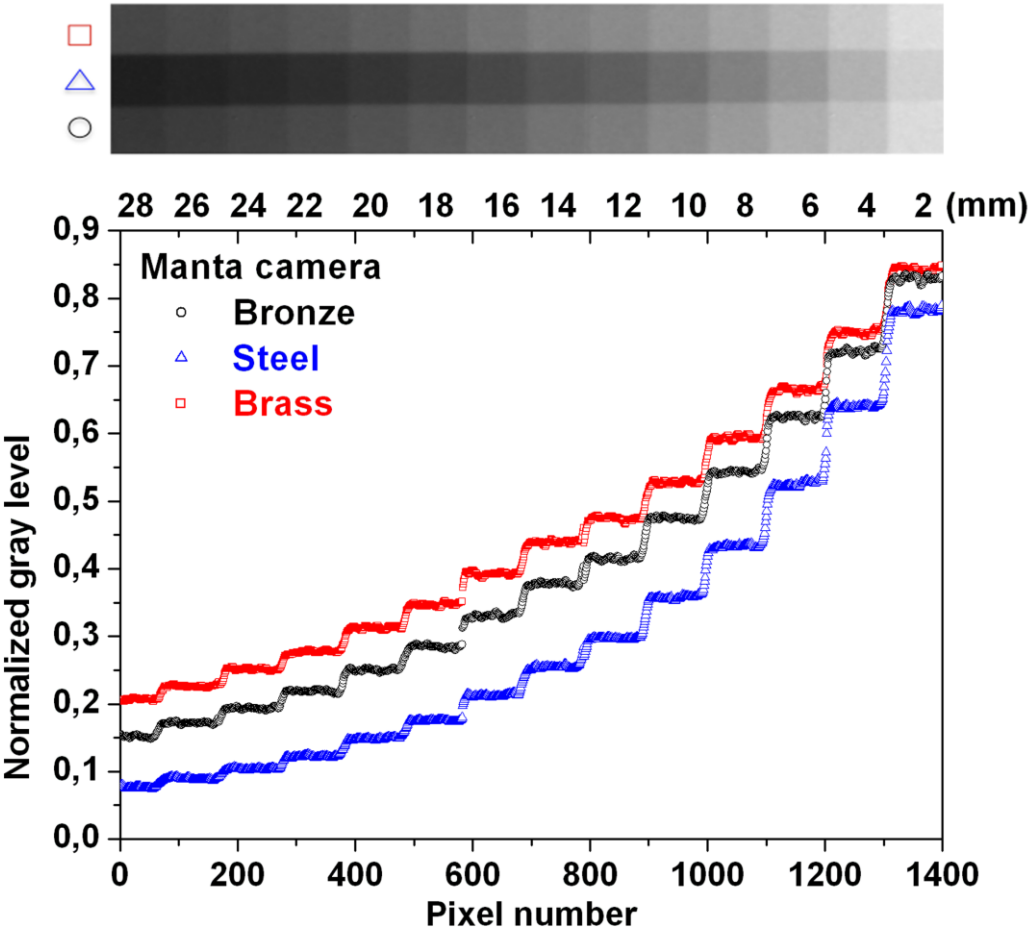
375 **Figure 6**



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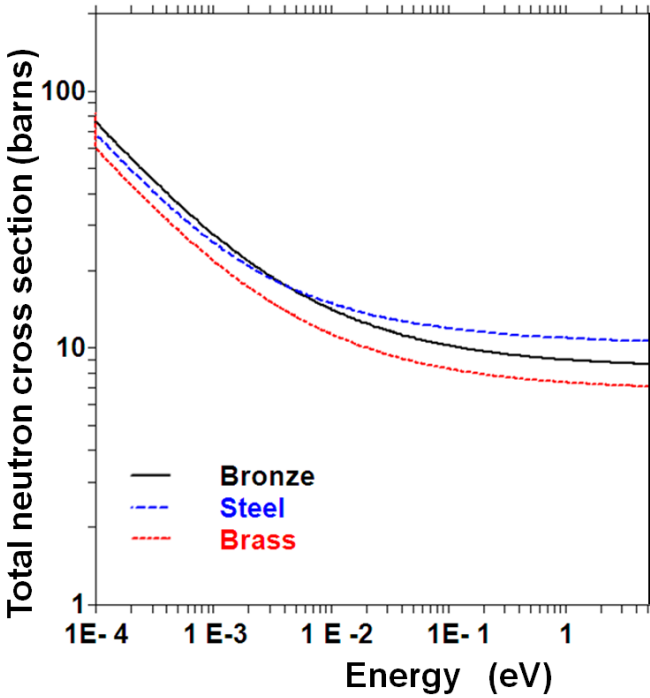
378 **Figure 7**



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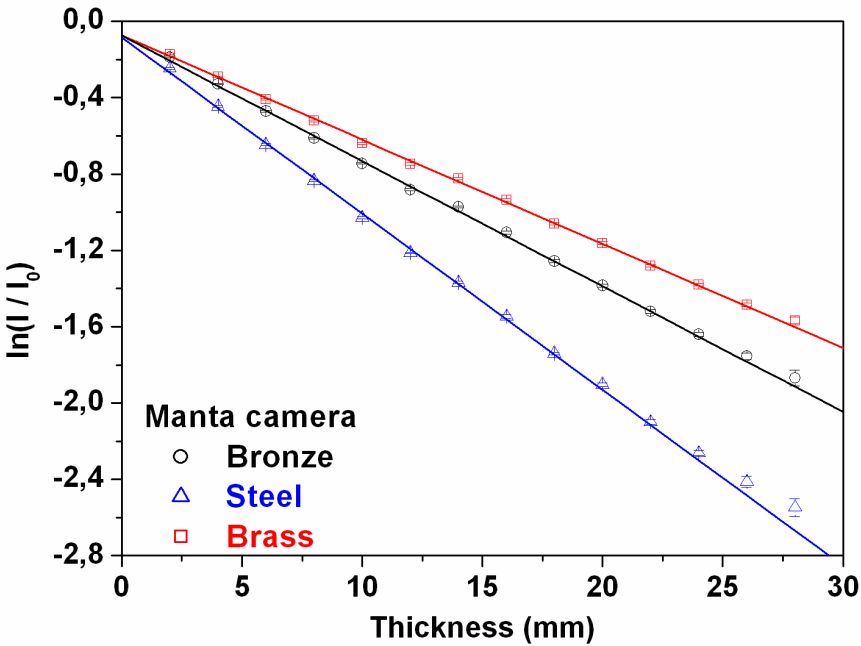
381 **Figure 8**



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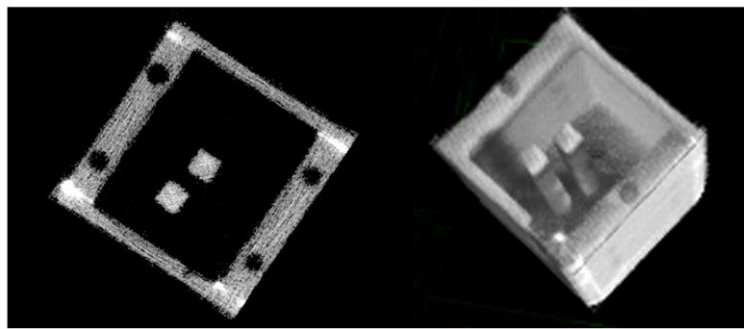
384 **Figure 9**



385

386

387 **Figure 10**



(a)

(b)

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391 Table1: Samples density and elemental composition [7].

392

Material	Density [g cm <sup>-3</sup> ]	Elemental composition
C40 carbon steel	7.85	Fe; C (0.38-0.43%); Si (0.15-0.35%); Mn (0.6-0.8%); P (max 0.02%); S (max 0.008%); Cr (0.15-0.35%)
Turning brass	8.29	Cu; Zn (39%); Pb (3%)
Bronze (85-5-5-5)	8.88	Cu; Pb (5.81%); Sn (5.05%); Zn (5.46 %); Fe (0.11%); Ni (1.21%); Al (trace)

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396 Table 2: Maximum effective dynamic range in terms of effective number of gray  
397 levels and in dB.

	DMK		Manta	
Scintillator 1	39±2	(31.8±0.4) dB	73±4	(37.3±0.5) dB
Scintillator 2	43±1	(32.7±0.2) dB	80±10	(38.1±1.1) dB

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400

401 Table 3: Spatial resolution, defined as FWHM of the ESF derivative, measured for  
402 both cameras coupled with both scintillators.

FWHM (μm)	DMK	Manta
Scintillator 1	182±11	261±9
Scintillator 2	306±14	389±13

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404

405 Table 4: Spatial resolution defined as MTF value at 10% calculated for both  
406 cameras coupled with both scintillators.

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MTF 10% (lp/mm)	DMK	Manta
Scintillator 1	4.0±0.2	2.74±0.10
Scintillator 2	2.37±0.11	1.84±0.06



408

409 Table 5:  $\Sigma/\rho$  values obtained with Manta CCD camera compared to the

410 predictions.

411

Material	$\rho$ (g cm <sup>-3</sup> )	Measured $\Sigma/\rho$ (cm <sup>2</sup> g <sup>-1</sup> )	ENDF/B-VI $\Sigma/\rho$ @ 1 eV (cm <sup>2</sup> g <sup>-1</sup> )
Steel	7.85	(11.9±0.4) 10 <sup>-2</sup>	12.02 10 <sup>-2</sup>
Brass	8.29	(6.6±0.2) 10 <sup>-2</sup>	6.64 10 <sup>-2</sup>
Quaternary bronze alloy	8.88	(7.4±0.2) 10 <sup>-2</sup>	7.24 10 <sup>-2</sup>

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