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UNIVERSITÀ DEGLI STUDI DI TORINO

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Heat capacity, configurational heat capacity and fragility of hydrous magmas

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Abstract

The glassy and liquid heat capacities of four series of dry and hydrous natural glasses and magma as a function of temperature and water content (up to 19.9 mol%) were investigated using differential scanning calorimetry (DSC). The analyzed compositions are basalt, latite, trachyte and pantellerite. The results of this study indicate that the measured heat capacity of glasses (C_{pg}) is a linear function of composition and is well reproduced by the empirical model of Richet (1987). For the investigated glasses, the partial molar heat capacity of water can be considered as independent of composition, in agreement with Bouhifd et al. (2006). For hydrous liquids, the heat capacity (C_{pliq}) decreases nonlinearly with increasing water content. Previously published models, combined with the partial molar heat capacity of water from the literature, are not able to reproduce our experimental data in a satisfactory way. We estimated the partial molar heat capacity of water (C_{pH_2O}) in hydrous magma over a broad compositional range. The proposed value is $41 \pm 3 \text{ J mol}^{-1} \text{ K}^{-1}$. Water strongly affects the configurational heat capacity at the glass transition temperature [$C_{pconf}(T_g)$]. An increase of $C_{pconf}(T_g)$ with water content was measured for the polymerized liquids (trachyte and pantellerite), while the opposite behavior was observed for the most depolymerized liquids (basalt and latite). Structural and rheological implications of this behavior are discussed in light of the presented results.

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1. INTRODUCTION

Physical and thermodynamic properties of silicate melts are key factors governing petrological and volcanic processes such as melt generation, transport, emplacement, crystallization, degassing and eruptive style (e.g., Lange and Carmichael, 1990; Dingwell et al., 1996; Ochs and Lange, 1999; Papale, 1999; Snyder, 2000; Spera, 2000; Dingwell, 2006). Several studies performed on silicate glasses and liquids have demonstrated that anhydrous chemical composition and water, the most abundant volatile in magmatic systems, strongly affect the structure (e.g., Mysen et al., 1982; Stolper, 1982a,b; Mysen, 1997; Mercier et al., 2009; Xue, 2009), the physical properties (viscosity, diffusivity, density, heat capacities; Persikov et al., 1990; Ochs and Lange, 1999; Whittington et al., 2000, 2009; Romano et al., 2001; Bouhifd et al., 2006; Vetere et al., 2007; Behrens and Zhang, 2009; Giordano et al., 2009; Fanara et al., 2012; Di Genova et al., 2013, 2014), the phase relationships and crystallization behavior of magmas (Fenn, 1977; Muncill and Lasaga, 1988; Davis et al., 1997; Vona and Romano, 2013). The isobaric heat capacity (C_p) is a thermodynamic property related to variations in the internal energy in a system and to structural changes of the silicate framework in response to temperature variations. Understanding how C_p varies with composition and water content is crucial in order to construct thermal models of petrological and volcanic processes, such as Ghiorso and Sack, 1995; Papale, 1997; Moretti and Papale, 2004; Papale et al., 2006; Richet et al., 2006; James et al., 2009; Whittington et al., 2009; Iacono-Marziano et al., 2012). In addition, knowledge of the heat capacity of liquids helps to quantify the configurational entropy of magmas through the Adam and Gibbs (1965) theory of relaxation processes (Richet, 1984). It has been observed that the partial molar heat capacity of a single oxide ($C_{p,i}$) in a glass is independent of composition (Stebbins et al., 1984; Richet, 1987). The three-term Maier Kelley equation, based on the additive models of calculation from oxide components, well reproduces the variation of heat capacity of glasses (C_{pg}) with temperature (Richet, 1987; Bouhifd et al., 2006). Heat capacity measurements on

liquids performed on a simple Al-free silicate system have shown that C_p can be considered a linear function of oxide concentration, resulting in predictive models that assume ideal mixing of oxide components, each of which associated with a partial molar heat capacity (e.g., Stebbins et al., 1984; Richet and Bottinga, 1985; Lange and Navrotsky, 1992). However, simple additive models fail to reproduce the heat capacities of aluminosilicate liquids, for which non-linear variations of heat capacity with compositions have been observed. This behavior has been generally attributed to the non-ideal mixing of oxide components (Stebbins et al., 1984; Richet and Bottinga, 1985; Richet and Neuville, 1992; Bouhifd et al., 1998, 2012). In addition, despite the recent advances and large amount of data produced in the literature over the last few decades, there is no unanimous consensus concerning the variation of heat capacity for hydrous natural glasses and liquids.

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as phase equilibria calculations, theoretical investigation of physical properties of liquids and water exsolution during ascent-driven degassing (e.g., Burnham and Davis, 1974; Burnham and Davis (1974) and Clemens and Navrotsky (1987) determined a partial molar heat capacity for water in hydrous melt of $80 \text{ J mol}^{-1} \text{ K}^{-1}$ by calorimetric measurements and thermodynamic calculations on the albite-water system. Furthermore, Clemens and Navrotsky (1987) showed that this value is close to that predicted for pure water at P and T, suggesting that the partial molar heat capacity of water has little dependence on composition. About twenty years later, Bouhifd et al. (2006) measured the heat capacity of water on hydrated iron-free polymerized aluminosilicate glasses and liquids of albitic, phonolitic, trachytic and leucogranitic compositions. They argued, based on the Richet (1987) model, that water has a small effect on the heat capacities of silicate glasses, showing that the partial molar heat capacity of water could be considered independent of glass composition varying as a function of temperature as follows: $C_{pH_2O} = -122.319 + 341.631 \cdot 10^{-3} T + 63.4426 \cdot 10^5 \cdot T^{-2} \text{ (J mol}^{-1} \text{ K}^{-1})$. For liquids, Bouhifd et al. (2006) reported a value of partial molar heat capacity of water ($C_{pH_2O} = 85 \text{ J mol}^{-1} \text{ K}^{-1}$) independent of liquid composition and water content for the studied compositions. Recently, Bouhifd et al. (2012) presented new data for more depolymerized iron-free analogs (liquid and glasses) of tephritic and foiditic composition. The partial molar heat capacity of water calculated by Bouhifd et al. (2006) for polymerized silicate glasses is similar to the values observed for these depolymerized compositions, suggesting, that the partial molar heat capacity of water in glasses is independent of composition. However, for depolymerized liquids, the authors reported a partial molar heat capacity of water of $237 \text{ J mol}^{-1} \text{ K}^{-1}$, suggesting a dependency of the heat capacity of water upon the degree of polymerization of the melts. In order to address the effect of water on the heat capacity of silicate glasses and melts, we have investigated the C_{pg} and C_{pliq} of dry and hydrous remelted volcanic materials. We selected four natural compositions: basalt from Mt. Etna (ETN, Giordano and Dingwell, 2003a; Vona et al., 2011); latite from Fondo Riccio – Phlegrean fields (FR, Giordano et al., 2006; Misiti et al., 2011; Di Genova et al., 2014); trachyte from Agnano Monte Spina – Phlegrean fields (AMS-B1, Romano et al., 2003) and pantellerite from Pantelleria Island (PS, Di Genova et al., 2013). The compositions selected allow to explore the wide chemical variation characterizing natural magmas and to address in detail the effect of specific chemical parameters on the internal energy of the liquids. At the same time, as the rheological behavior of those liquids is well characterized in the literature, a comparison between viscosity and configurational heat capacity allow us to calculate and discuss the configuration entropy and fragility of these liquids and their variation as a function of water content in the magma.

2. EXPERIMENTAL METHODS

2.1. Synthesis of hydrous samples

The starting glassy dry materials were obtained by direct fusion of natural pantellerite (PS), trachyte (AMS-B1), latite (FR) and basalt (ETN) magmas. The natural materials were melted in a large thin-walled Pt crucible in a Nabertherm MoSi₂ box furnace at $T = 1670\text{--}1870 \text{ K}$ for a few hours. The melts were then quenched in air, and samples were removed from the crucibles. Chips of glass were loaded into a Pt80Rh20 cylindrical crucible and remelted in a box furnace heated by MoSi₂ elements (Theta) and stirred in air using a PaarRheolab QC rheometer (EVPLab Rome Tre University). Stirring allowed complete homogenization and removal of bubbles from the samples. Subsequently, the melt was cooled

in air and the glasses obtained were used for the chemical characterization of dry material. The major element composition of the glasses was determined at the CNR – Istituto di Geologia Ambientale e Geoingegneria in Rome, using a Cameca SX50 electron microprobe, with 5 WDS spectrometers and EDS Link eXL system, under analytical conditions of 15 kV accelerating voltage, 15 nA beam current and 10 μ m beam diameter to reduce K and Na volatilization. Chemical analyses are reported in [Table 1](#).

In order to synthesize hydrous samples, the glassy dry materials were powdered and loaded together with known amounts of distilled water into Au–Pd capsules with 2.9 mm outer diameter, 2.5 mm inner diameter and 20 mm height. A rapid-quench-internally-heated pressure vessel (IHPV, CNRS/INSU-Universit  d’Orl ans) was used to synthesize water-bearing samples, containing up to 19.9 mol% H₂O, at $P = 4$ kbar and $T = 1470$ – 1670 K ([Table 2](#)). To check that water did not leak from the capsules, the capsules were weighed before and after being placed in an oven for at least one hour and no weight difference was observed. Homogeneity and absolute water content of the hydrated samples were determined by Karl-Fischer Titration (Hannover University), following the method described by [Behrens et al. \(1996\)](#) and [Nowak and Behrens \(1997\)](#). To correct for unextracted water after KFT, a quantity of 0.10 wt% was added to the measured values ([Behrens and Stuke, 2003](#)) for AMS-B1, PS and FR samples. In order to check water homogeneity distribution and possible loss of water during the calorimetric and viscosity measurements, Fourier Transform Infrared (FTIR) Spectroscopy was performed before and after each experiment. 4500 cm^{-1} (OH groups) and 5200 cm^{-1} (H₂O molecules) peak heights of the near-infrared (NIR) absorption bands were used to analyze the water content of the samples using the Lambert–Beer law ([Stolper, 1982a](#)). Infrared absorption spectra were collected on doubly polished glass chips with thickness ranging from 260 to 350 μ m. The water content of the nominally dry starting glasses was determined by measuring the peak height of the mid-infrared (MIR) absorption band at 3550 cm^{-1} . FTIR spectra were collected at the Laboratory of Infrared Spectroscopy, Department of Science (University Roma Tre), using a NicPlan IR-microscope equipped with a liquid nitrogen-cooled MCT detector. The nominal resolution was 4 cm^{-1} and 128 scans were averaged for each sample and for the background. Simple linear baselines were fitted to both NIR peaks (TT baseline according to [Ohlhorst et al., 2001](#)). Molar absorption coefficients used for each composition are reported in [Table 2](#). Uncertainty of the results is about 2% based on the reproducibility of measurements and on the error associated with the background subtraction procedure. Archimedean buoyancy was performed to measure densities of dry and hydrous glasses. Thickness was measured with a digital micrometer with an uncertainty of ± 5 μ m. Results of the water determinations are reported in [Table 2](#), together with synthesis conditions. Finally, in order to check possible loss of iron during synthesis, we performed microprobe analysis on hydrous samples. The results show that, for hydrated silicate glasses, the FeO content is identical to that of the starting anhydrous glass within the error of the electron microprobe measurements (see [Appendix, Table A1](#)).

2.2. Calorimetry

Calorimetric measurements were performed by Differential Scanning Calorimeter technique (Netzsch SC 404 Pegasus, EVPLab Rome Tre University). Prior to analysis of the samples, temperature and sensitivity calibration was performed using melting temperatures of standard materials (inorganic salts of Rb, K, Cs and Ba) up to 1270 K. Subsequently, a baseline measurement was taken where two empty Pt/Rh crucibles were loaded into the DSC and then the DSC was calibrated against the C_p of a single sapphire crystal ([Robie et al., 1979](#)). Doubly polished glass samples 30–60 mg in weight were placed in Pt crucibles under a constant Argon flow rate (20 lt/min). Each sample was heated at 1 K/min heating rate from room temperature to 323 K and kept at this temperature for approximately one hour in order to achieve DSC signal equilibrium. Subsequently, in order to structurally relax the high pressure liquids, temperature was raised with a heating step of 20 K/min from $T = 323$ K up to a maximum of 50 K above the estimated glass transition temperature (T_g) of the sample. The cooling rate after the experiment matched the heating rate (20 K/min).

Table 1. Dry composition (mol%) of glasses. Fe²⁺/Fe^{tot} ratio of PS was measured by redox titration as reported in [Di Genova et al. \(2013, 2014\)](#).

a Gram formula weight on the basis of one mole of oxides.

b Number of atoms per gfw.

Measurements of C_p were performed in 3 subsequent thermal treatments where the heating rates always matched the previous cooling rates (20–10–5 K/min). The heat capacity versus temperature paths did not show any evidence of water exsolution, crystallization or weight loss during the experiments. Post run

optical analysis performed after each experiment confirmed the absence of thermal alteration of the samples.

3. RESULTS

3.1. Heat capacity and glass transition temperatures (T_g)

The heat capacities and the glass transition temperatures (T_g) of the silicate liquids investigated in this study are reported in Table 3 and shown in Fig. 1(A–D) as a function of composition, temperature and water content. The heat capacity (C_p) measurements were carried out on nominally anhydrous and hydrous samples containing up to 19.9 mol% H₂O in the temperature range of 574–1002 K.

Except for liquids of some specific composition such as titanosilicates, borosilicates and highly polymerized aluminosilicates (e.g., Richet and Bottinga, 1984, 1986; Courtial and Richet, 1993; Lange and Navrotsky, 1993; Richet et al., 1997), it is generally observed that the heat capacity of silicate liquids (C_{pliq}) is constant with temperature, and this allows for extrapolation over a large T interval (e.g., Richet and Bottinga, 1986; Lange and Navrotsky, 1993; Bouhifd et al., 2006). In the investigated temperature range above T_g , no variation of C_{pliq} was observed, in agreement with literature studies. Basalt liquid from Etna (ETN) shows the highest dry liquid heat capacity $C_{\text{pliq}} = 96.6 \text{ J mol}^{-1} \text{ K}^{-1}$, followed by the latitic sample from Fondo Riccio (FR, Phlegrean Fields) with a $C_{\text{pliq}} = 94.8 \text{ J mol}^{-1} \text{ K}^{-1}$, pantellerite from Pantelleria (PS) with a $C_{\text{pliq}} = 91.6 \text{ J mol}^{-1} \text{ K}^{-1}$ and trachytic liquid from Agnano Monte Spina, Phlegrean Fields (AMS-B1) with a $C_{\text{pliq}} = 90.5 \text{ J mol}^{-1} \text{ K}^{-1}$. As can be seen from the C_p curves in Fig. 1, the liquid heat capacity of the silicate melts (C_{pliq}) decreases with increasing water content for each composition. The variations of C_{pliq} as a function of water content are shown in Fig. 2. Upon initial dissolution of water (up to 5 mol% H₂O), a small decrease of heat capacity can be observed for all liquids. As additional amounts of H₂O are introduced in the melts, the C_{pliq} seems to decrease more strongly converging to similar values for all compositions. From the heat capacity curves we derived glass transition temperatures, defined as the kinetic transition from solidlike behavior to liquid-like behavior (Dingwell et al., 1996). The glass-transition interval was constrained using the peak (T_{gpeak}), onset (T_{gonset}) and stable liquid (T_{gliquid}) position of the heat capacity curves (Fig. 3). The values are listed in Table 3. The values of the measured glass transition temperatures, as collected at the peak and at the onset are reported in Fig. 4 for each specific composition, H₂O content and thermal treatment. Liquid T_g 's are affected by larger error compared to the onset and peak measurements and for this reason they have not been considered in the following discussion. Peak and onset T_g s correlate very well showing similar trends of variation as a function of anhydrous composition, H₂O content and thermal treatment. In general, peak T_g 's show consistently higher values (approximately 50 K) compared to onset T_g s. For each composition, higher T_g values are attained for higher cooling/heating rates.

Table 2. Measured water contents (FTIR and KFT), synthesis condition (pressure and temperature) of the hydrous samples and room temperature density of hydrated samples.

Molar absorption coefficients for the combination band of OH₂ at 4500 cm⁻¹ and H₂O_m at 5200 cm⁻¹: PS (Behrens and Zhang, 2009), AMSB1 (Misiti et al., 2006) and ETN (Lesne et al., 2011).

The duration of each experiment is 24 h.

a Water content (mol%) measured by Infrared spectroscopy.

b Water content (mol%) measured by Karl-Fischer titration (KFT).

The effect of water on the T_g s can be better evaluated by direct comparison of all liquids in Fig. 5. In this figure, the value of T_{gpeak} is chosen for convenience, as the measurement of T_{gpeak} is associated with the lowest standard error, and cooling/heating rates of 20 K/min were considered. The pantellerite samples (PS) show the lowest measured value of T_g both for anhydrous and for hydrous conditions.

For instance, the dry T_{gpeak} exhibited by the PS pantellerite ($T_g = 875 \text{ K}$), measured at 20 K/min, is about 110 K lower than the T_{gpeak} shown by the AMS-B1 trachyte (987 K) and FR latite sample (985 K), and is approximately 80 K lower than the T_g of the ETN basalt ($T_g = 954 \text{ K}$) (Table 3). Glass transition temperatures also appear to be strongly dependent on water content, decreasing as water content increases in the melt (Figs. 4 and 5). The decrease appears to be more marked at lower concentrations of H₂O. Among the liquids shown in Fig. 5, the PS pantellerite shows the lowest T_g over the entire range of water content investigated ($T_{\text{gpeak}} = 636 \text{ K}$ at H₂O = 11.95 mol%). The decrease in glass transition temperature due to the addition of water in the melt is most pronounced for the latitic sample (FR) that shows a decrease in T_{gpeak} of 294 K when 12.55 mol% H₂O is added. AMS-B1 trachyte shows a decrease in T_{gpeak} of 275 K when 12.46 mol% H₂O is added and ETN basalt exhibits a decrease in T_{gpeak} of 224 K with the introduction of 11.99 mol% H₂O.

3.2. Water speciation

Water dissolved in the melt was measured both via Karl Fisher Titration and Fourier-Transform Infrared Spectroscopy. Spectroscopic measurements of water in the glasses Latite

Fig. 1. Measured heat capacity for (A) PS pantellerite, (B) AMS-B1 trachyte, FR latite and ETN basalt.

Fig. 2. Measured heat capacity for the selected samples as a function of water content. quenched at 20 K/min cooling rate after calorimetric measurements allowed for the determination of water species in the liquid frozen at T_g.

Table 2 illustrates the concentration of molecular water, hydroxyl groups and total water content expressed both in weight percent and in molar concentration (two oxygen basis), as determined by KFT analyses and FTIR. The molar absorption coefficients used for each investigated composition are also reported in Table 2. The concentrations of molecular water and hydroxyl groups (expressed as mol% H₂O) as a function of increasing water content are also shown in Fig. 6. The concentrations of both species vary with composition with the molecular water increasing and the hydroxyl groups leveling off as the total water content in the melt increases. The concentration of molecular and hydroxyl groups is the same for water content of approximately 7 mol%. Similar trends of OH₋ and molecular H₂O species are largely shown in literature for a wide range of silicate composition (e.g., Stolper, 1982b; Zhang et al., 1997; Behrens and Nowak, 2003; Liu et al., 2004; Ni et al., 2009; Lesneet al., 2011).

Fig. 3. Measured heat capacity (C_p) for dry pantellerite sample (PS DRY) after cooling and heating at 20 K/min. In figure are shown the selected glass transition temperatures: T_{gonset}, T_{gpeak} and T_{gliq}.

Table 3. Measured glass transition temperatures (onset, peak, liquid) and liquid, glass and configuration heat capacities of the studied samples.

a Data in J mol⁻¹ K⁻¹.

b Data in J g⁻¹ atom⁻¹ K⁻¹.

4. DISCUSSION

4.1. Heat capacity of glasses

In order to parameterize and compare our results with the model of Richet (1987), the measured heat capacity of silicates glasses (C_{pg}, Table 3) have been described by the Maier Kelley equation as follow:

$$C_{pg} = a + bT + c/T^2 \quad (1)$$

where the a, b and c coefficients are reported in Table 4 together with the standard error. For the anhydrous glasses, calculated C_p using Eq. (1) is predicted by the Richet (1987) model within 1.8%. Using the measured heat capacity of hydrous glasses, together with the Richet (1987) model for the anhydrous composition and considering a temperature dependence of C_{pgH₂O} as reported in Eq. (1), we derived the partial molar heat capacity of water in silicate glasses (C_{pgH₂O}, R² = 0.92) as following:

$$C_{pgH_2O} = 112.52 - 59.10 \cdot 10^{-3} T - 77.20 \cdot 10^5/T^2 \quad (2)$$

Fig. 4. Measured T_{gpeak} as a function of different water content. The numbers reported in figures (20–10–5) indicate the different coolingheating thermal treatment adopted in C_p measurements.

Fig. 5. Effect of water on measured T_{gpeak} after cooling and heating at 20 K/min. The glass transition temperatures for pantelleritic melts are significantly lower (~100 K) than the measured T_{gpeak} for AMS-B1, FR and ETN samples.

Fig. 6. Measured water speciation in quenched hydrous glasses as determined from infrared spectroscopy measurements.

In Fig. 7A and B we report the measured and calculated C_{pg} as a function of water content at 500 and 450 K. The C_p of all hydrous glasses are reproduced within 4%, confirming that the partial molar heat capacity of water is independent of composition in silicate glasses over a broad range of polymerization (0.6 NBO/T 6.1.5), in agreement with Bouhifd et al. (2006, 2012). Finally, we investigated the proximity

to the Dulong and Petit harmonic limit ($\sim 25 \text{ J g}^{-1} \text{ atom}^{-1} \text{ K}^{-1}$) of our measured glassy heat capacities [$C_{pg}(\text{Tg})$]. In fact, the glass transition of silicates usually takes place when the heat capacity becomes close to the Dulong and Petit harmonic limit (Haggerty et al., 1968; Richet and Bottinga, 1986), which is a theoretical upper limit to the heat capacity at constant volume (C_v). Given the small thermal expansivities of silicate glasses, this harmonic limit should also apply to C_p (Toplis et al., 2001). This has been confirmed by Toplis et al. (2001) and Bouhifd et al. (2006, 2012) for anhydrous multicomponent silicate glasses. However, Bouhifd and coauthors show that the hydrous glasses exhibit glass transition temperatures which take place before achieving the Dulong and Petit harmonic limit. Both aspects are confirmed by our results. In fact, in Table 3 we report the $C_{pg}(\text{Tg})$ expressed in $\text{J g}^{-1} \text{ atom}^{-1} \text{ K}^{-1}$ and as can be seen, for the anhydrous samples, the measured glassy heat capacities are close to the Dulong and Petit harmonic limit. On the other hand hydrous glasses containing around 3.5 wt% H_2O ($\sim 12 \text{ mol}\%$ H_2O) show that the glass transition takes place when the heat capacity is $\sim 93\%$ of the Dulong–Petit limit.

4.2. Heat capacity of anhydrous liquids and comparison with literature data

The measured heat capacities of anhydrous silicate liquids (C_{liq} , Table 3) have been compared with predicted values using the models of Stebbins et al. (1984), Lange and Navrotsky (1992) and Richet and Bottinga (1985) with the mean C_p of Al_2O_3 reported by Courtial and Richet (1993). In Table 5, the measured C_{liq} of the present study are compared with those predicted by these three models. Fetot is partitioned in equal amounts as FeO and Fe_2O_3 (in wt%; Giordano et al., 2006). We observed significant differences, expressed as $D\% = [100 * (\text{Measured } C_p - \text{Calculated } C_p) / \text{Measured } C_p]$, between the measured and calculated values. The measured anhydrous C_{liq} for the PS sample is reproduced within $\pm 1.6\%$ by all models. However, the Lange and Navrotsky (1992) and Stebbins et al. (1984) models strongly overestimate (by 7.7% and 4.8%, respectively) the measured C_{liq} for AMS-B1, while the Richet and Bottinga (1985) model overestimates the measured C_{liq} for AMS-B1 by 3.5%. As far as the more depolymerized liquids are concerned, the Richet and Bottinga (1985) and Stebbins et al. (1984) models excellently reproduce (within 0.9%) the measured data for the ETN basalt and for the FR latite (within 0.5% and 0.4%, respectively). The Lange and Navrotsky (1992) model overestimates our data to a greater extent, with an error of 2.1% and 2.7% for ETN and FR, respectively. In summary, the Richet and Bottinga (1985) model seems to reproduce more closely our anhydrous dataset, showing the largest error (3.5%) for the AMS-B1. The Lange and Navrotsky (1992) and Stebbins et al. (1984) models largely fail to reproduce the measured C_p for AMS-B1 and FR. A possible explanation for this resides in the high Al_2O_3 content of these liquids, much higher (12.29 and 11.68 mol% for AMS-B1 and FR respectively) compared to those for pantelleritic and basaltic compositions (5.96 and 10.76 mol%, respectively). Both Stebbins et al. (1984) and Lange and Navrotsky (1992) models are in fact ideal solution models in which the partial molar heat capacities of Al_2O_3 content are temperature independent. On the other hand, the model of Richet and Bottinga (1985), implemented by Courtial and Richet (1993), considers specific interactions between aluminum and alkali or other alkaline-earth elements which result in a temperature-dependent term for the heat capacity of Al_2O_3 , more realistic for Al-rich compositions such as AMS-B1 and FR. The temperature dependence of the heat capacity of Al_2O_3 also correlates with analogous complex volumetric and rheological behavior for this oxide. It is beyond the scope of this paper to address the issue of the structural role of the aluminum in melts, but certainly the reason of its behavior has to reside in the complex structural role of this element [$\text{Al}/(\text{Al} + \text{Si})$ order–disorder, Al triclusters, distribution in Q species], in its capability to strongly interact and form complexes with alkalis and to assume different coordination states in different chemical environments (Mysen and Toplis, 2007).

Table 4. Coefficients for the Maier Kelley equation (Eq. (1)) for the glasses studied in this work.

Fig. 7. Measured and calculated (up to 100 wt% of water) glassy heat capacity at 500 K (A) and 450 K (B).

4.3. Heat capacity of water in silicate liquids

Water generally decreases the heat capacity of silicate liquids (Fig. 1). In Fig. 2 the variations of C_{liq} as a function of mole percent of water are shown. The decrease of C_p as a function of water content for all compositions investigated here is not linear, being slight upon initial dissolution of water and much stronger as additional H_2O is introduced in the melts. We compared our results for hydrous liquids with literature data. For the anhydrous silicate melt we considered the partial molar heat capacity of single oxides from the model of Richet and Bottinga (1985) which best reproduces our data. The partial molar heat capacity of water is poorly constrained in the literature. Bouhifd et al. (2006) published a value of 85

J mol⁻¹ K⁻¹ for the heat capacity of water for polymerized melts (0 < NBO/T < 0.2), in agreement with the studies of [Burnham and Davis \(1974\)](#) and [Clemens and Navrotsky \(1987\)](#). Recently, a much greater value for the partial molar heat capacity of water of 237 J mol⁻¹ K⁻¹ was determined by the same authors ([Bouhifd et al., 2012](#)) for depolymerized melts of synthetic tephritic and foiditic composition (0.86 < NBO/T < 1.51).

In [Table 5](#) we report calculations performed considering the [Bouhifd et al. \(2006\)](#) value of 85 J mol⁻¹ K⁻¹ for partial molar heat capacity of water for polymerized compositions and the [Richet and Bottinga model \(1985\)](#) for partial molar heat capacity of the other oxides. The parameterization is not able to reproduce our experimental data in a satisfactory way and the error between the measured and calculated values increases with increasing water content. For liquids with water contents of 12 mol% (3.5 wt% H₂O), the difference between measured and calculated values ranges between 3.7% and 8.5%, with a maximum error up to 13.3% for the FR 6.3 sample. The calculated liquid heat capacities for our depolymerized melts considering the partial molar heat capacity of water of [Bouhifd et al. \(2012\)](#) in depolymerized liquids (CpH₂O = 237 J mol⁻¹ K⁻¹), and the [Richet and Bottinga model \(1985\)](#) for partial molar heat capacity of the other oxides, are reported in [Table 6](#). As can be seen from [Table 6](#), the calculated Cps severely overestimate (up to 50%) our experimental data. Recent work by [Robert et al. \(2014\)](#) also reports a value for the partial molar heat capacity of H₂O for basaltic compositions of 86 J mol⁻¹ K⁻¹, similar to [Bouhifd et al. \(2006\)](#). Comparison between our experimental results and literature data suggest that estimated values of partial molar heat capacities of water from literature are not able to reproduce our experimental results in any appreciable way. The causes of this lack of agreement could be diverse. They may be related to the paucity of data at high water content (higher than 7 mol% of H₂O) or to the approximation resulting from linearly fitting the heat capacity versus water curves. Moreover, the discrepancy could arise from compositional dependence terms for the heat capacity of water, as suggested by [Bouhifd et al. \(2012\)](#), or deviations for the ideal behavior of the solution assumed in the analytical approach, or the error deriving from the comparison of natural versus synthetic compositions. In order to check if the presence of iron and, in particular, its redox state could explain the offsets between the measured hydrous heat capacities and calculated with literature models, we performed further calculations varying the Fe²⁺/Fe^{tot}. We took into account the measured oxidation state (Fe²⁺/Fe^{tot}) of a representative set of our investigated glasses (iron rich, FR and PS) provided by [Di Genova et al. \(2013, 2014\)](#). The results point out that the offsets between measurements and literature models do not result from uncertainties in the ferric–ferrous ratio. For example, we considered FR 6.3 sample (measured C_{pliq} = 81 J mol⁻¹ K⁻¹), (showing the highest error in predicting the C_{pliq} (up to 17%) with literature models), and calculated C_{pliq} using the measured Fe²⁺/Fe^{tot}. No improvements between these and the previous calculations (performed considering Fe^{tot} equally partitioned between Fe²⁺ and Fe³⁺) were observed. In fact, using the measured Fe²⁺/Fe^{tot} ratio of 0.7 ([Di Genova et al., 2014](#)), the [Lange and Navrotsky \(1992\)](#) model provides a C_{pliq} of 95.4 J mol⁻¹ K⁻¹, higher with respect to the value of C_{pliq} obtained with the same model (C_{pliq} = 95 J mol⁻¹ K⁻¹) and considering Fe^{tot} equally partitioned as FeO and Fe₂O₃. Similar results are obtained when employing [Stebbins et al. \(1984\)](#) and [Richet and Bottinga \(1985\)](#) models. Given the poor correlation of our data with existing literature, we decided to derive the heat capacity of water independently for each specific composition investigated in this study. Treating the hydrous melt as a pseudo-binary mixture of an anhydrous silicate melt end-member and a pure water end-member, one can write ([Bouhifd et al., 2006](#)):

$$C_{p\text{liq}} = X_{\text{H}_2\text{O}} C_{p\text{H}_2\text{O}} + (1 - X_{\text{H}_2\text{O}}) C_{p\text{liq,Anhydrous}} \quad (3)$$

where X_{H₂O} is the mole fraction of water, CpH₂O is the partial molar heat capacity of water and C_{pliq,Anhydrous} is the heat capacity of anhydrous liquid.

Table 5 Measured and calculated liquid heat capacity for the studied samples.

a Data from [Di Genova et al. \(2013, 2014\)](#).

b Calculated values using the [Lange and Navrotsky \(1992\)](#) model.

c Calculated values using the [Stebbins et al. \(1984\)](#) model.

d Calculated values using the [Richet and Bottinga \(1985\)](#) model with Cp of Al₂O₃ reported by [Courtial and Richet \(1993\)](#).

e Calculated values using Eqs. (2) and (3).

* 100 * (Measured Cp – Calculated Cp)/Measured Cp.

Table 6

Measured and calculated liquid heat capacity for the studied samples using a partial molar heat capacity for water in depolymerized melts of $237 \text{ J mol}^{-1} \text{ K}^{-1}$ reported in [Bouhifd et al. \(2012\)](#).

a Calculated values using the [Lange and Navrotsky \(1992\)](#) model.

b Calculated values using the [Stebbins et al. \(1984\)](#) model.

c Calculated values using the [Richet and Bottinga \(1985\)](#) model with C_p of Al_2O_3 reported by [Courtial and Richet \(1993\)](#).

* $100 * (\text{Measured } C_p - \text{Calculated } C_p) / \text{Measured } C_p$.

Using Eq. (3) and our experimental data, we derived partial molar $C_{p\text{H}_2\text{O}}$ from the different liquids varying from $26.05 \text{ J mol}^{-1} \text{ K}^{-1}$ (FR latite) to $56.64 \text{ J mol}^{-1} \text{ K}^{-1}$ (PS pantellerite) with intermediate values for ETN ($44.58 \text{ J mol}^{-1} \text{ K}^{-1}$) and AMS-B1 liquids ($42.12 \text{ J mol}^{-1} \text{ K}^{-1}$). Although it is possible to discern a trend of decreasing partial molar $C_{p\text{H}_2\text{O}}$ with NBO/T, the data are not robust enough to express this relationship in a quantitative way and additional studies on natural liquids are required to further investigate this issue. For this, we decided to neglect any compositional dependence and fit all data together. Using the Eq. (3) and our experimental data, we therefore derived $41 \pm 3 \text{ J mol}^{-1} \text{ K}^{-1}$ as the partial molar heat capacity for water in natural silicate melts. In Table 5 we reported the calculated liquid heat capacity using Eq. (3) together with the calculated partial molar heat capacity for water. The relative average error for the investigated data is on the order of 1.1%. Therefore, using the anhydrous heat capacity reported in Table 5 for the investigated liquids with Eq. (3) and the derived partial molar heat capacity for water, it is possible to calculate liquid heat capacity variations with H_2O content. Moreover, we also decided to derive the $C_{p\text{H}_2\text{O}}$ including the heat capacity data provided by [Bouhifd et al. \(2006, 2012\)](#). Combining Eq. (3) with experimental data, we derived a $C_{p\text{H}_2\text{O}}$ of $60 \text{ J mol}^{-1} \text{ K}^{-1}$. With this value we can reproduce the 39 experimental data with relative average error of 2.7%. This value is slightly higher than the previously reported ($41 \text{ J mol}^{-1} \text{ K}^{-1}$), calculated taking into account only the data presented in this work, and is lower than the values of $85 \text{ J mol}^{-1} \text{ K}^{-1}$ and $237 \text{ J mol}^{-1} \text{ K}^{-1}$ previously reported by [Bouhifd et al. \(2006, 2013\)](#) for synthetic polymerized and depolymerized composition respectively and the values of $79 \text{ J mol}^{-1} \text{ K}^{-1}$ previously published by [Clemens and Navrotsky \(1987\)](#). The origin of these differences is unclear but it is likely related to the different nature of the liquids investigated. In our case, we studied natural compositions and the presence of multicomponent melts certainly introduces more complexity in the energetics of the system and possibly more compositional dependent terms, negligible or not present when considering simplified synthetic compositions such as those reported in literature. For instance, the presence of iron and its partitioning in Fe^{2+} and Fe^{3+} can affect the energy of the chemical interactions and thus the results.

For these reasons, and given the paucity of heat capacity data for natural compositions, we consider reasonable to consider the value of $41 \text{ J mol}^{-1} \text{ K}^{-1}$ as the partial molar heat capacity of water for natural silicate melts. New C_p data from natural compositions are required in order to both enrich the $C_{p\text{H}_2\text{O}}$ database and to obtain a clearer picture of the partial molar heat capacity of water in natural silicate melts. Even though the approximation of a linear fitting of the data leads to excellent reproducibility of our results within the range of water content investigated in our study, contribution to the heat capacity of molecular ($\text{H}_2\text{O}_{\text{mol}}$) water versus hydroxyl species (OH_-) can also be extracted. In fact, when looking in detail at Fig. 2, the relationship between the measured C_{pliq} and H_2O content is nonlinear. The heat capacity of liquids tends to remain constant from anhydrous conditions to up to approximately 6 mol% H_2O , and then decreases as further water is dissolved into the melt. Similarly, it can also be observed (Fig. 6) that the relative concentration of water species in the melt varies in a nonlinear fashion. The relative abundance of $\text{H}_2\text{O}_{\text{mol}}$ and OH_- molar species tends to be similar when 7 mol% H_2O is added to the melt, whereas a clear increase of $\text{H}_2\text{O}_{\text{mol}}$ content, with respect to hydroxyl groups, is observed as an increasing amount of water is dissolved (from 7 to 15 mol% H_2O). It should be noted that the measurements of the water speciation on glasses quenched from high temperature experiments do not represent isothermal values (e.g. [Romano et al., 1995](#); [Shen and Keppler, 1995](#); [Behrens and Nowak, 2003](#)), as the water species keep reacting during quenching and the reaction is actually frozen only at the glass transition temperature, which differs as a function of water content itself. However, as the calorimetric measurements are performed in the proximity of the glass transition temperature, and assuming that temperature does not significantly affect the anhydrous liquid heat capacities, comparison between the speciation curve and the calorimetric curve, although not isothermal, is still possible. By assuming this, the nonlinear relationship between the C_{pliq} and H_2O content (Fig. 2) can then be related to the nonlinear dependence of speciation on total water concentration, as also reported in [Bouhifd et al. \(2006, 2012\)](#). Taking into account the measured C_{pliq} shown in Table 3 and the measured water

speciation data reported in Table 2, we have calculated the partial molar heat capacity for molecular water and OH⁻ ions, for each composition, modifying Eq. (3) as follow:

$$C_{p,liq} = X_{H_2O\ mol} \cdot C_{p,H_2O\ mol} + X_{OH^-} \cdot C_{p,OH^-} + (1 - X_{H_2O}) \cdot C_{p,liq,Anhydrous} \quad (4)$$

where $X_{H_2O\ mol}$ and X_{OH^-} are the mole fraction of molecular water and OH⁻ ions respectively, $C_{p,H_2O\ mol}$ is the partial molar heat capacity of molecular water, C_{p,OH^-} is the partial molar heat capacity of OH⁻ ions and $C_{p,liq,Anhydrous}$ is the heat capacity of anhydrous liquid and X_{H_2O} is as in Eq. (3) the mole fraction of total water content. The calculated values of $C_{p,H_2O\ mol}$ and C_{p,OH^-} are reported in Table 4 as a function of the SM parameter from Giordano and Dingwell (2003a), which is calculated as the sum of the network modifier oxides on a molar basis ($Na_2O + K_2O + CaO + MgO + MnO + FeO_{tot}/2$ mol%), and is related to the degree of polymerization of silicate melts. It should be noted that Fondo Riccio samples (FR) are not included in this parameterization because no molar absorptivity coefficients for molecular water and hydroxyl groups (e5200 and e4500) are available in the literature. The partial molar heat capacity of H₂O mol and OH⁻ species appears to be dependent on chemical composition. The partial molar heat capacity of molecular water is very small for all compositions investigated, with the highest $C_{p,H_2O\ mol}$ (21 J mol⁻¹ K⁻¹) showed by the most depolymerized sample (ETN). This could reflect the basically passive role of water molecules dispersed in the silicate network silicate structure. The PS sample, characterized by the lowest network modifier content (16 mol%), shows the highest C_{p,OH^-} (142 J mol⁻¹ K⁻¹) with respect to AMS-B1 (SM = 21 mol%; C_{p,OH^-} = 80 J mol⁻¹ K⁻¹) and ETN (SM = 32 mol%; C_{p,OH^-} = 68 J mol⁻¹ K⁻¹). This trend could be due to the different energy of interaction between hydroxyl groups and silicate network. C_{p,OH^-} is higher for polymerized liquids as the dissolution of water into these melts strongly disrupts the structure creating a large number of configurational states. The energy needed to achieve these new configurational states is therefore higher compared to the case of the basalt, where an already depolymerized liquid interacts with water with weaker modification of its structure. We have to stress that the associated uncertainties to the calculated partial molar heat capacities (especially for H₂O mol) are quite high and for this, the interpretation of the obtained results, although very intriguing, may remain qualitative. More heat capacity measurements, combined with water speciation analysis, are required in order to better constrain the possible dependence of partial molar heat capacity of H₂O mol and OH⁻ species from chemical composition.

4.4. Configurational heat capacity of hydrous magmas

The heat capacity of a liquid is considered to be the sum of a vibrational and a configurational contribution. The configurational contribution on heat capacity [$C_{p,conf}(T)$] is a measure of the structural changes that occur on heating and represents the energy used to change the structure of a liquid in response to temperature variations. Assuming that the $C_{p,g}(T_g)$, which represents the vibrational contribution to the heat capacity, do not vary significantly above the glass transition temperature (Richet and Bottinga, 1986), $C_{p,conf}(T)$ can be taken as the difference between the heat capacity of the melt ($C_{p,liq}$) and the heat capacity of the glass at glass transition temperature [$C_{p,g}(T_g)$]:

$$C_{p,conf}^{(T)} = C_{p,liq}(T) - C_{p,g}(T_g) \quad (5)$$

Fig. 8 shows the configurational heat capacities as a function of mole% H₂O for the studied liquids. As previously observed (Richet and Bottinga, 1985; Bouhifd et al., 1998; Toplis et al., 2001; Webb, 2008), the $C_{p,conf}(T_g)$ increases with decreasing SiO₂ content and increases with Al content and cation field strength of the atoms in the melt. $C_{p,conf}(T)$ has contributions from both chemical (related to all possible configurational states between oxygens and cations) and topological (distribution of TO distances and TOT angles) and the exact identification of the compositional dependence of these parameters is somewhat difficult. In general, the $C_{p,conf}(T)$ increase with decreasing SiO₂ content (and decreasing of the SiO₂/Al₂O₃ ratio) is ascribed to a general decrease of the overall strength of TO bonds, and to a correlated increase of the TO and TOT distribution (topological contribution to the configurational entropy). The role of alkali versus alkaline earth cations and Al/Si order disorder is more complex. According to Webb (2008, and references therein), it is well known that the configurational heat capacities of Na₂O–Al₂O₃–SiO₂ melts are lower than those of the CaO–Al₂O₃–SiO₂ melts. Explanations for this behavior are various and are related to the differences between the field strength of these cations (Bouhifd et al., 1998). To understand how the $C_{p,conf}(T)$ varies with composition, and hence silicate structure, we can refer to the fragility of silicates melts because it is well known that the configurational entropy affects the fragility of melts (e.g., Toplis et al., 1997; Neuville, 2006; Webb et al.,

2007; Webb, 2008; Di Genova et al., 2013). The fragility describes the departure of the liquid viscosity from an Arrhenian temperature dependence (Angell, 1991), and within the framework of Adam and Gibbs (1965) theory, is related to the temperature dependence of configurational entropy (S_c), which is in turn governed by the magnitude of the configurational heat capacity (Richet, 1984):

$$\log_{10} \dot{\gamma} = \frac{Ae}{B(T - T_g)} + \frac{Cp_{conf}}{R(T - T_g)}$$

where Ae and Be are temperature independent constants (Richet, 1984). Fragility can be expressed via the steepness index m (Plazek and Ngai, 1991; Bohmer and Angell, 1992) which describes the gradient of the viscosity curve at the glass transition temperature (T_g) on a reduced temperature scale:

$$m = \left. \frac{d(\log_{10} \tau)}{d\left(\frac{T_g}{T}\right)} \right|_{T=T_g} \quad (6)$$

where τ is the average relaxation time and T_g the glass transition temperature (at $s = 100$ s). The calculated values of fragility using viscosity measurements and Eq. (6) for our liquids are $m = 46.1$, 33.6, 24.8 and 21.4 for ETN, FR, PS and AMS-B1 melts respectively (Giordano and Dingwell, 2003a; Di Genova et al., 2013, 2014). In Fig. 9 we report the calculated fragility for the melts studied in this work as a function of $Cp_{conf}(T_g)$. Moreover, for the purposes of comparison, we report the calculated fragility from Whittington et al. (2001) for the trachytic (Tr_W) and phonolitic (Ph_W) samples which have been subjected to calorimetric studies in Bouhifd et al. (2006). This figure shows that fragile melts have a higher configurational heat capacity compared to strong melts as predicted by the Adam and Gibbs (1965) theory. Furthermore, Toplis et al. (1997) provided an equation that shows that the fragility is proportional to the ratio between the $Cp_{conf}(T_g)$ and configurational entropy at T_g . In general, many structural studies as well as molecular dynamic simulations, have suggested that alkaline-earth bearing liquids are more fragile than alkali-bearing liquids (Romano et al., 2001; Whittington et al., 2001; Neuville, 2006; Webb et al., 2007; Webb, 2008; Cheng et al., 2012; Di Genova et al., 2013). The high electronegativity of the alkaline earth cations causes a general lengthening of the T–O average bond strength and a decrease in the mean TOT bond angle compared to the alkali-bearing case. Moreover, being divalent, the alkaline earth cations require two Al cations to be electrically balanced and the formation of Al–O–Al linkages further destabilizes the structure. The presence of divalent cations and the coupling of Al tetrahedral therefore imposes a strain in the network which facilitates its breakdown at high temperature, causing fragile behavior (Scamehorn and Angell, 1991). An increase of Al with respect to Si has a similar effect of lengthening TO bonds and decreasing TOT bond angle, thus destabilizing the network (Geisinger et al., 1985; Navrotsky et al., 1985). Moreover, it has been shown that both an increase of Al/Si ratio and Ca + Mg/Na + K ratio can increase the chemical contribution to the configurational heat capacity via an increase in Q species disorder (Mysen, 1997; Neuville, 2006; Cheng et al., 2012).

In summary, low silica content, high Al/Si ratio and Ca + Mg/Na + K ratio can increase the fragility via an increase in both chemical (Q species disorder, Al/Si order–disorder) and topological (TO bond distances, TOT bond angles and TO and TOT distribution) disorder as registered by an equivalent increase in S_c and Cp_{conf} . The fragility of our samples can be easily explained in terms of different silica content, and relative proportions of alkali and alkaline earth cations. We observe a steady increase in fragility from AMS-B1 ($m = 21.4$), PS ($m = 24.8$), Ph_W ($m = 30.5$), Tr_W ($m = 32.9$), FR ($m = 33.6$) to ETN ($m = 46$), accompanied to a general increase of $Cp_{conf}(T_g)$ (AMS-B1 = 9.8; PS = 11.2; Ph_W = 11.7; Tr_W = 12.8; FR = 16.4; ETN = 22.4 J mol⁻¹ K⁻¹). A first-order parameter expressing the increase in fragility in our liquids is the NBO/T, non-bridging oxygen per tetrahedra on a molar basis (Mysen, 1988) (0.11; 0.10; 0.19; 0.21; 0.20; 0.45) and the silica content, gradually decreasing from PS to ETN. An increase in Al/Si ratio and in Ca + Mg/Na + K ratio from PS to ETN parallel the decrease in SiO₂ and can further explain the increase in fragility and Cp in the same direction. Although the trends observed can, in general, be easily interpreted in terms of chemistry and structural relationships, the extent of such variations are in some cases more difficult to rationalize. For instance, PS pantelleritic liquid, with much higher SiO₂ content than AMS-B1, very low Ca + Mg/Na + K ratio and very low Al/Si would have been expected to be stronger than AMS-B1 liquid. In contrast, there is a large increase in m and $Cp_{conf}(T_g)$ from AMS-B1 to FR, despite the relatively small decrease in SiO₂, increase in Ca + Mg/Na + K and the same values for Al/Si.

Obviously more structural and energetic parameters should be taken into account in order to completely define the behavior of these complex liquids. In Fig. 8 we illustrate the variation of configurational heat capacity as a function of total water content. Also in this case a compositional dependence is apparent. For the most polymerized alkaline melts (AMS-B1 and PS) a slight increase of $Cp_{conf}(T_g)$ can be observed as water is added to the melt. The increase in configurational heat capacity is consistent with the

increases in fragility determined by viscosity measurements (Romano et al., 2001; Di Genova et al., 2013). In Fig. 8 the variation of the Cpconf (Tg) as a function of H₂O for the most depolymerized alkaline earth melts (ETN and FR) is also presented. For these depolymerized compositions, after an initial increase up to ~6 mol% of water, the Cpconf (Tg) strongly decreases as further H₂O is added to the melt, suggesting, contrary to the case for polymerized compositions, that water would decrease the fragility. In the literature, the effect of H₂O on melt fragility is still debated and both a decrease and an increase in fragility as a function of water content have been reported (Richet et al., 1996; Whittington et al., 2000; Romano et al., 2001; Whittington et al., 2001; Bouhifd et al., 2012; Robert et al., 2012; Di Genova et al., 2013). We derived the fragility variations as a function of water content for FR (Misiti et al., 2011; Di Genova et al., 2014) and AMS-B1 (Romano et al., 2001; Misiti et al., 2006) only. For these two compositions, in fact, a complete data set of high-T, low-T, anhydrous and hydrous viscosity measurements exist and this allows to minimize errors included in the determination of fragility and hence possible artifacts in the results. In fact, as reported in Angell (2002), the value of the fragility (here defined as steepness index m, Eq. (6)) obtained by a relatively small viscosity range near the glass transition temperature depends on the range of data fitted. For this reason, we chose not to include PS and ETN sample in the calculation, as for these two compositions hydrous viscosity data at high T are not available. Using the well-known three parameter Vogel–Fulcher–Tammann (VFT) equation to describe the temperature dependence of viscosity [$\log \eta = A + B/(T - C)$], the steepness index m can be expressed as follows (Plazek and Ngai, 1991; Bohmer and Angell, 1992; Giordano and Dingwell, 2003b):

$$m = \frac{B}{T_{12} \left(1 - \frac{C}{T_{12}}\right)^2} \quad (7)$$

Table 7 Calculated partial molar heat capacity of molecular water (CpH₂O_{mol}) and OH₋ (CpOH₋) using Eq. (4).

a NBO/T = Mysen (1988).

b Average relative error.

c Sum of structure modifying oxides (Giordano and Dingwell, 2003).

Fig. 8. Configurational heat capacities at glass transition temperature of the melts investigated in this work as a function of water content.

Fig. 9. Configurational heat capacities at glass transition temperature of the melts investigated in this work as a function of fragility (m).

where B and C are VFT parameters and T₁₂ represents the temperature at which the viscosity (η) is equal to 10¹² Pa sec. In order to evaluate the effect of water content on the B and C parameter, and hence on the fragility (m) of the melts, we adopted a modified VFT equation reported in Di Genova et al. (2013):

$$\log \eta = A + \frac{[b_3 + b_4 \log(1 + H_2O)]}{T - [c_3 + c_4 \log(1 + H_2O)]} \quad (8)$$

Using Eq. (8) we parameterized the viscosity data by Misiti et al. (2011) and Di Genova et al. (2014) for FR samples and by Romano et al. (2001) and Misiti et al. (2006) for AMS-B1 samples, and finally calculated the fragility using Eq. (7). In Table 8 fitting parameters for Eq. (8) are reported. In Table 9 and Fig. 10 the calculated fragility for the latite (FR) and trachyte (AMS-B1) liquids are reported as a function of water content. It is evident from Fig. 10 that water acts in different ways decreasing the fragility in the case of the FR latitic liquid, and increasing m for the AMS-B1 trachytic liquid. This different behavior might be related to the different dissolution mechanisms of water in silicate melts characterized by different chemical compositions (Giordano et al., 2009; Di Genova et al., 2013). Water dissolves in melts as both hydroxyl groups and molecular water reacting with the silicate network in different ways (Stolper, 1982a,b; Silver and Stolper, 1985, 1989; Silver et al. 1990). According to the classical model of water dissolution (Burnham, 1975; Stolper, 1982a,b; Sykes and Kubicki, 1993, 1994; McMillan, 1994), water dissolves into polymerized compositions such as AMS-B1 and PS melts by creating NBOs in terms of SiOH and AlOH species and disrupting the silicate network. The depolymerization reaction not only decreases the overall T–O strength of the melt (increasing in topological contribution to Cpconf and Sc)

but also creates a large number of configurational states available for the activation of viscous flow (increasing the chemical contribution to C_{pconf} and Sc). This is in good

Table 8 Fit parameters of FR (Misiti et al., 2011) and AMS-B1 (Romano et al., 2001 and Misiti et al., 2006) according to Eq. (8). Numbers in parenthesis are standard deviation.

Table 9 Calculated steepness index m (e.g., Plazek and Ngai, 1991) for AMS-B1 and FR sample as a function of water content.

Fig. 10. Configurational heat capacities at glass transition temperature [$C_{pconf}(T_g)$] and fragility (m defined by Eq. (7)) as a function of water content. In agreement with the previously calculated values of C_{pOH} for AMS-B1 and PS melts (80 and 145 J mol⁻¹ K⁻¹) which suggest that, for these polymerized liquids, the dissolution of water into the melt creates a large number of configurational states. This emerges also from our heat capacity measurements, which show that the configurational heat capacity (and the configurational entropy) increases resulting in an increase in fragility, as derived for the AMS-B1 liquid. For depolymerized melts, different reactions involving a polymerization of the silicate network and creation of free OH⁻ upon dissolution of water have been proposed (Fraser, 1977; Xue and Kanzaki, 2004, 2008; Moretti, 2005; Mysen and Cody, 2005). In principle, all the reactions, leading to either the polymerization or depolymerization of the silicate network, can operate simultaneously, and the net effect of water would depend on the predominance of one reaction over the others. In general, the polymerization reaction is expected to be more abundant in the presence of NBOs, and therefore in depolymerized melts, whereas the depolymerization reaction should be dominant in the opposite case of polymerized melts. The effect on fragility would also be the opposite. In the case of depolymerized compositions, the polymerization induced by the dissolution of water would create a stronger melt with fewer configurational states available for viscous flow and stronger T–BO bonds, decreasing C_{pconf} and Sc and decreasing fragility, as observed for the FR liquid. Further to this behavior, we note that lower C_{pOH} values are calculated for our depolymerized liquid. In conclusion the increase in $C_{pconf}(T_g)$ caused by water dissolution in trachyte is consistent with the fragility increase determined by viscosity data provided by Romano et al. (2001) and Misiti et al. (2006) and discussed in Giordano et al. (2009), while the opposite behavior is observed for the more depolymerized liquid (FR).

5. VOLCANOLOGICAL IMPLICATIONS: ERUPTIVE THERMAL BUDGET

Heat capacity represents a fundamental parameter to understand the evolution of multicomponent and multiphase equilibria energetics and therefore the kinetics and dynamics of igneous and volcanic processes. Transport properties of silicate melts, such as thermal diffusivity (D =mm²/s) and thermal conductivity (k = W/m K) represent the main properties controlling the kinetics of crystal growth and the conductive heat flow of a magmatic system, respectively. Thermal diffusivity (D) is commonly expressed by the following equation (Snyder et al., 1994; Whittington et al., 2009):

$$D = k / \rho C_p \quad (9)$$

where k is thermal conductivity, ρ is density (kg/m³) and C_p is heat capacity (J kg⁻¹ K⁻¹). As an example to demonstrate the need and importance of accurate C_p measurements, we calculated the thermal diffusivity of our Ca–Mg rich liquid composition (ETN and FR) using both the liquid heat capacity data presented in this paper and C_p values calculated by Lange and Navrotsky (1992) model for dry and hydrous conditions. We use liquid thermal conductivities provided by Snyder et al. (1994) on molten CaMgSi₂O₆ at T between 1673 and 1873 K. The calculated D using data presented in this work are always higher compared to calculations performed using the Lange and Navrotsky (1992) model for the liquid heat capacity, with differences up to 8.4% for the highest water content considered (3.8 wt%), corresponding to the calculated error reported in Table 5 for ETN samples. Those differences will influence significantly any model involving the flow of heat, such as cooling time-scale of a magma body, or convection and crystallization models. In general, experimental data (e.g. Romine et al., 2012) demonstrate that the effect of T is much stronger on C_p than on ρ and D , and that therefore C_p variations, especially at the glass to liquid transition, will influence thermal conductivity the most, making accurate measurements of C_p essential for proper calculation of thermal conductivity. Heat capacity data are also essential for quantifying magma degassing and heat flux during steady-state volcanic activity, which in turn control the style and energy of the eruption. During transient activity, for instance, the energy released by the expansion of the gas and by the cooling of magma juveniles, is converted into energy used

to accelerate the gas and magma in the explosion, to raise the magmatic material within the gravity field, and to overcome the air drag acting on individual pieces of solid material. [Harris and Stevenson \(1997\)](#) presented a model for magma degassing in order to constrain the steady-state magma system of Stromboli (Italy). In order to calculate the total thermal flux, the heat carried by expelled ejecta (Q_{ejecta}) during eruptions must be taken into account. [Harris and Stevenson \(1997\)](#) calculated Q_{ejecta} as follow:

$$Q_{\text{ejecta}} = (C_p \Delta T_e + C_L) M_e \quad (10)$$

where C_p is the heat capacity of ejecta (1150 J kg⁻¹ K⁻¹), ΔT_e is the ejecta cooling from eruption (1000 K) to room temperature, C_L is the latent heat of crystallization (3 * 10⁵ J/kg) and M_e is the mass flux of erupted material (6 kg/s). Using these data with Eq. (10) the authors calculated $Q_{\text{ejecta}} = 8$ MW. Using the measured heat capacity in this work for anhydrous basalt we obtain a dramatic increase of 38% in heat flux, that is $Q_{\text{ejecta}} = 11$ MW which results in a much stronger energy release and eruption energy. Proper assessment of heat capacity values, for glasses and liquids, for molten material and gases, are therefore pivotal to the modeling of the energy budget in an eruption, therefore defining eruptive style and hazard assessment issues.

6. CONCLUSIONS

Our study highlights a very complex energetic behavior for multicomponent hydrous silicate melts. In general, the anhydrous and hydrous heat capacities of glasses are in good agreement with previous studies confirming that the partial molar heat capacity of water is independent of glass composition, in a wide range of silicate glasses (0.6 NBO/ T 6.5). Anhydrous liquid heat capacities are also well reproduced by literature models, considering a temperature dependence for $C_p\text{Al}_2\text{O}_3$ as reported by [Courtial and Richet \(1993\)](#). The behavior of water in silicate liquids is however more complex. We derived partial molar $C_p\text{H}_2\text{O}$ of 41 ± 3 J mol⁻¹ K⁻¹ independent of composition for the natural liquids investigated. Partial molar heat capacities of molecular water and hydroxyl groups were also derived for the compositions investigated in this study. The partial molar heat capacity of molecular water is very small for all compositions investigated, possibly reflecting the passive role of water molecules dispersed in the silicate network. The partial molar heat capacity of OH⁻ species appears to be dependent on chemical composition varying as a function of polymerization degree of the melt. PS sample, characterized by the lowest network modifier content (SM = 16 mol%), shows the highest $C_p\text{OH}^-$ (142 J mol⁻¹ K⁻¹) with respect to AMS-B1 (SM = 21 mol%; $C_p\text{OH}^-$ 80 J mol⁻¹ K⁻¹) and ETN (SM = 32 mol%; $C_p\text{OH}^-$ 68 J mol⁻¹ K⁻¹). This trend could be due to the different energy of interaction between hydroxyl groups and silicate network. Anhydrous configurational heat capacity $C_p\text{conf}$ (T_g) and fragility (m) increase with decreasing SiO₂ and increasing Al content. We observed that alkaline-earth bearing (ETN and FR) liquids are more fragile than alkali-bearing liquids (AMS-B1 and PS). The effect of water on $C_p\text{conf}$ (T_g) and therefore on fragility was also parameterized in terms of composition. For the most polymerized melts (AMS-B1 and PS) a slightly increase of $C_p\text{conf}$ (T_g), and hence in fragility, can be observed as water is added to the melt, while for depolymerized compositions (ETN and FR) after a small initial increase, the $C_p\text{conf}$ (T_g) and fragility strongly decrease as further water is added to the melt. We argue that this behavior might be related to the different dissolution mechanisms of the water in silicate melts characterized by different chemical compositions. The effect of water on melt fragility would depend on the predominance of the different polymerization and/or depolymerization reactions. The effect of water on magma fragility should also be taken into account when modeling degassing processes and eruption energy. In general, during degassing, magma viscosity increases and this leads to an exponential interrelated increase of exsolution and viscosity up to fragmentation conditions. It is generally considered ([Dingwell et al., 1996](#); [Whittington et al., 2001](#); [Vetere et al., 2007](#); [Giordano et al., 2008](#); [Misiti et al., 2011](#); [Di Genova et al., 2013](#)) that the increase in viscosity due to exsolution is stronger for polymerized compositions (rhyolites) than for depolymerized liquids (basalts), which translates into higher energy of degassing and fragmentation for the former. However, the trends reported in previous parameterizations could at least be in part affected by a lack of measurements in the hydrous high- T conditions and therefore by an incorrect estimate and modeling of fragility. When corrected for the increase in fragility at decreasing water content, as observed in our measurements of C_p for depolymerized liquids, the viscosity increase of basalt during degassing may be more rapid than previously expected. Moreover, degassing induced crystallization is also known to be more pronounced for basaltic compositions ([Applegarth et al., 2013](#)), leading to a further increase in viscosity. The combined effects of degassing-induced crystallization and degassing induced increase in fragility, may be so strong to induce rapid rheological changes with temperature prone to affect the eruption dynamics.

Table A1. Chemical composition (wt%) of glasses.

a Water content (wt%) measured by Karl-Fischer titration (KFT).

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APPENDIX A

See [Table A1](#).

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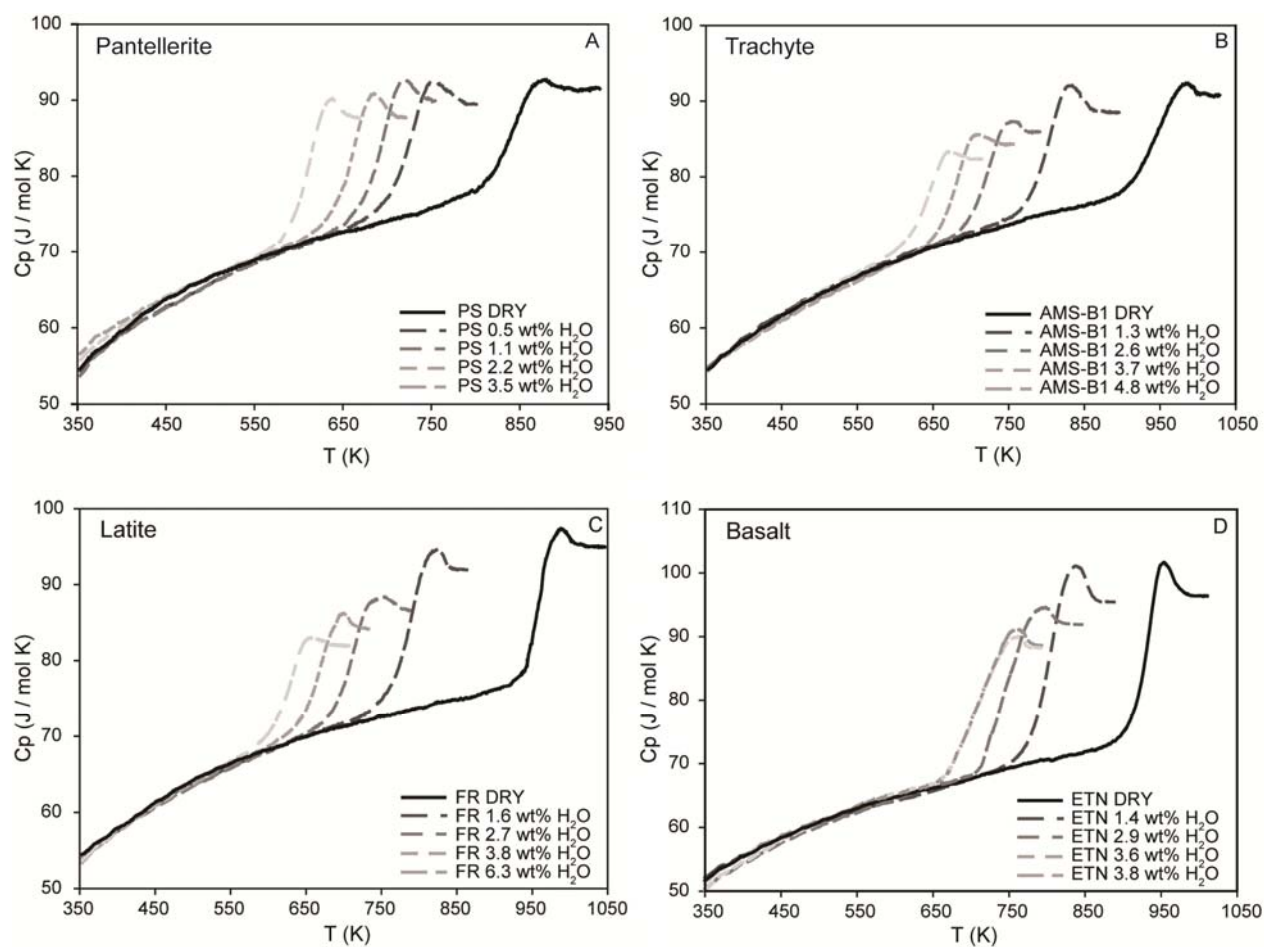


Fig. 1: Measured heat capacity for (A) PS pantellerite, (B) AMS-B1 trachyte, FR latite and ETN basalt.

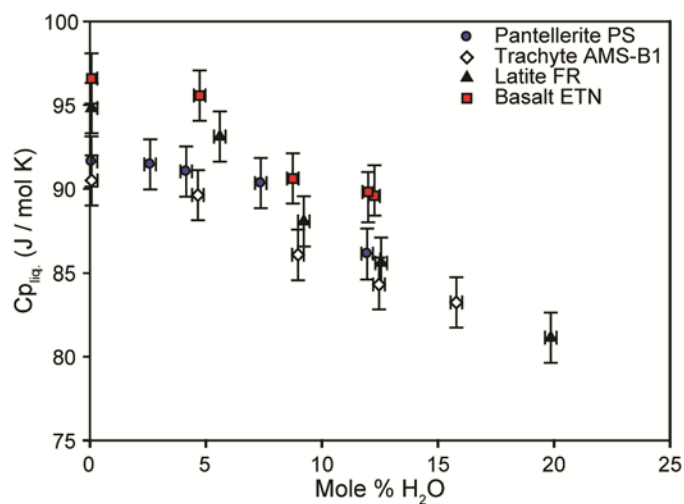


Fig. 2: Measured heat capacity for the selected samples as a function of water content.

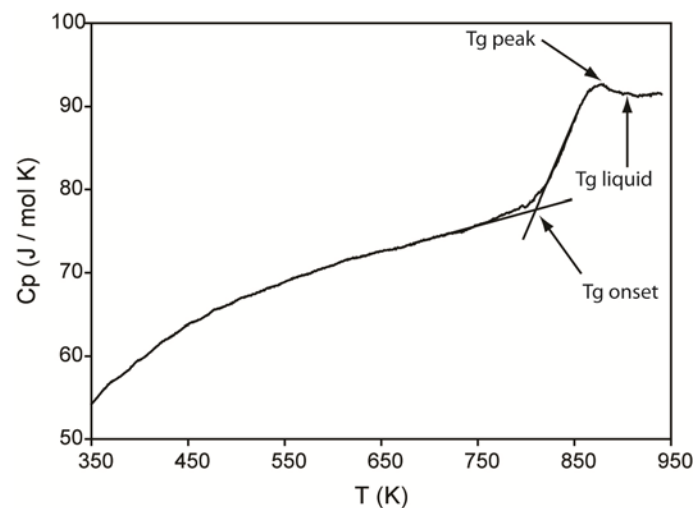


Fig.3: Measured heat capacity (C_p) for dry pantellerite sample (PS DRY) after cooling and heating at 20 K/min. In figure are shown the selected glass transition temperatures: $T_{g\text{onset}}$, $T_{g\text{peak}}$ and $T_{g\text{liq}}$.

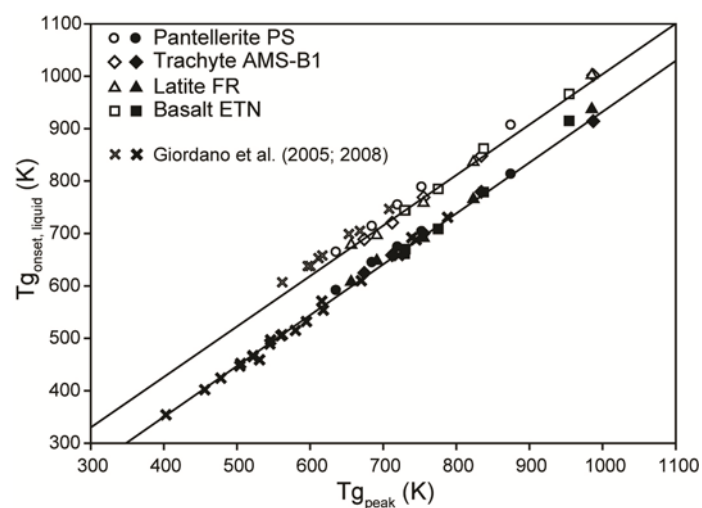


Fig.4: Comparison between measured glass transition temperatures in this study and reported in Giordano et al. (2005, 2008) as taken at onset, peak and stable liquid position of the C_p curves. Dark symbols represent the $T_{g\text{onset}}$ while the empty symbols represent $T_{g\text{liq}}$.

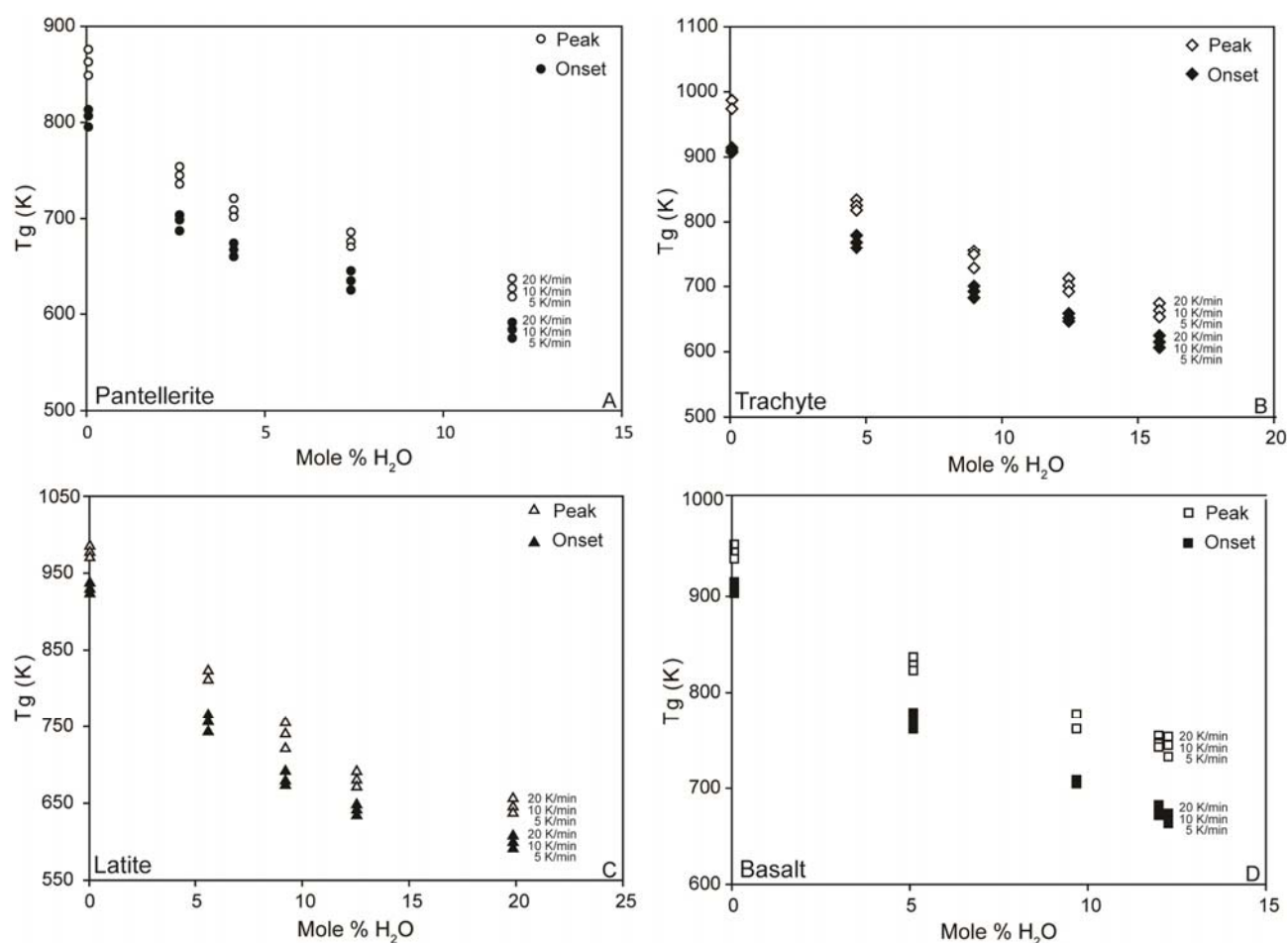


Fig.5: Measured $T_{g_{peak}}$ as a function of different water content. The numbers reported in figures (20-10-5) indicate the different cooling-heating thermal treatment adopted in Cp measurements.

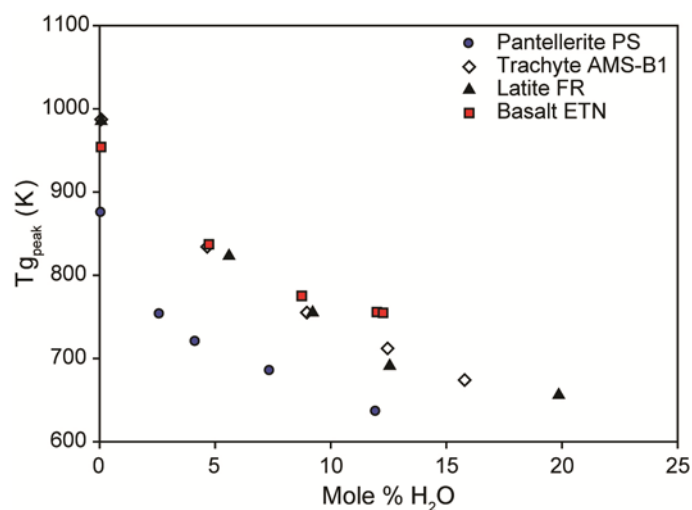


Fig.6: Effect of water on measured $T_{g_{peak}}$ after cooling and heating at 20 K/min. The glass transition temperatures for pantelleritic melts are significantly lower (~ 100 K) than the measured $T_{g_{peak}}$ for AMS-B1, FR and ETN samples.

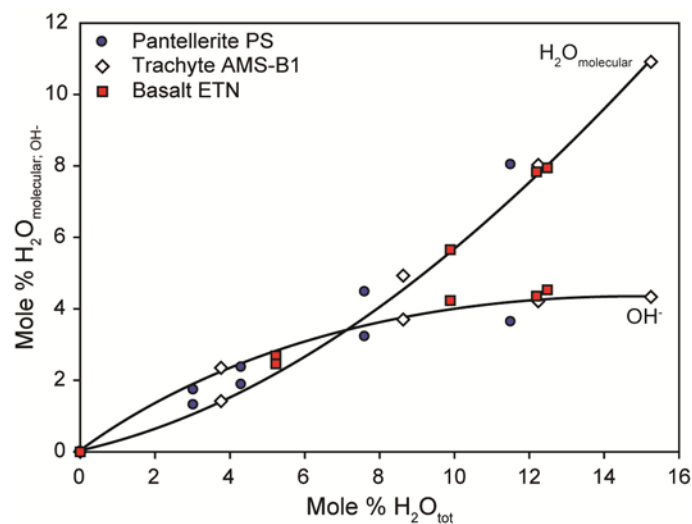


Fig.7: Measured water speciation in quenched hydrous glasses as determined from infrared spectroscopy measurements.

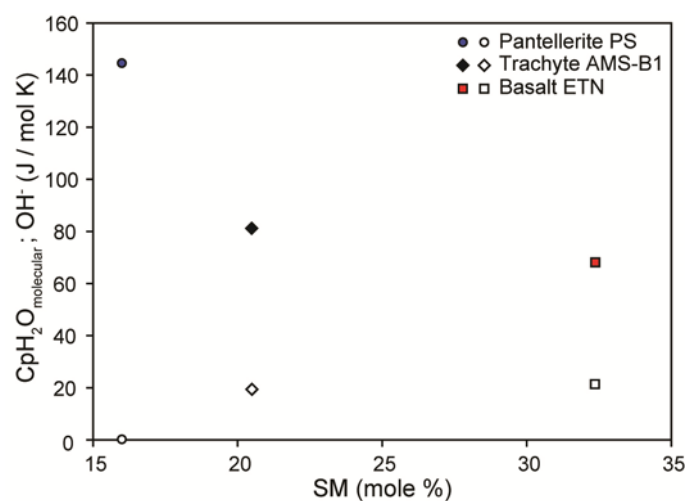


Fig.8: Calculated partial molar heat capacity of $\text{H}_2\text{O}_{\text{mol}}$ and OH^- species using Eq. 4 as a function of SM parameter. Filled symbols refer to Cp_{OH^-} , while empty symbols represent the $\text{Cp}_{\text{H}_2\text{O}_{\text{mol}}}$.

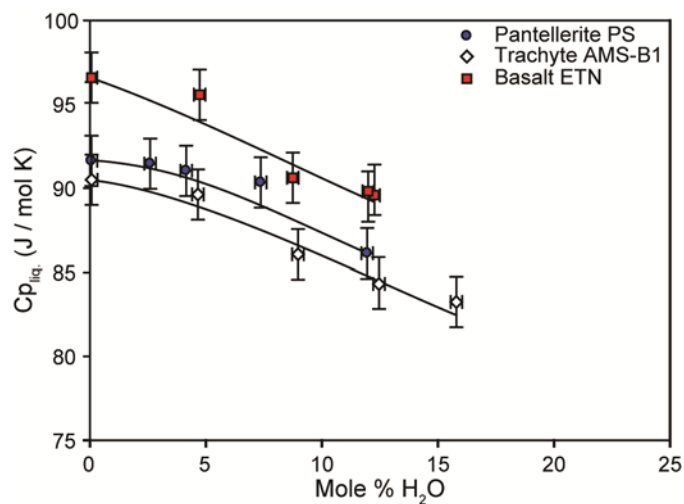


Fig. 9: Heat capacities of hydrous liquids ($C_{p,liq.}$) as a function of water content. Lines show the calculated heat capacities for liquids using Eq. 4 and measured partial molar heat capacities of OH^- and $H_2O_{mol.}$

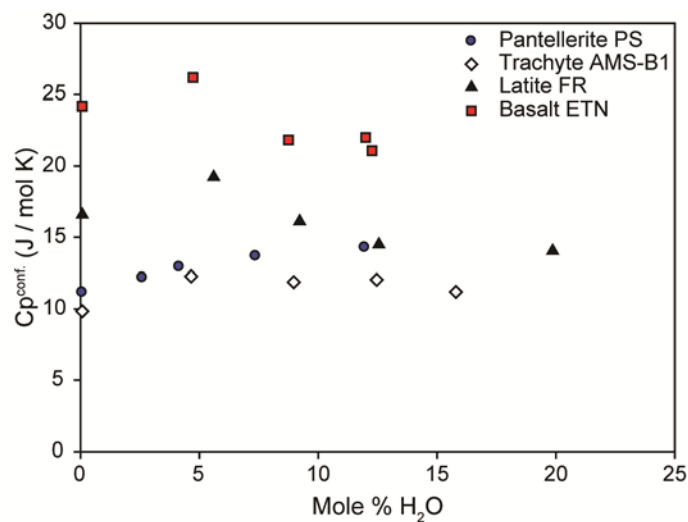


Fig.10: Configurational heat capacities at glass transition temperature of the melts investigated in this work as a function of water content.

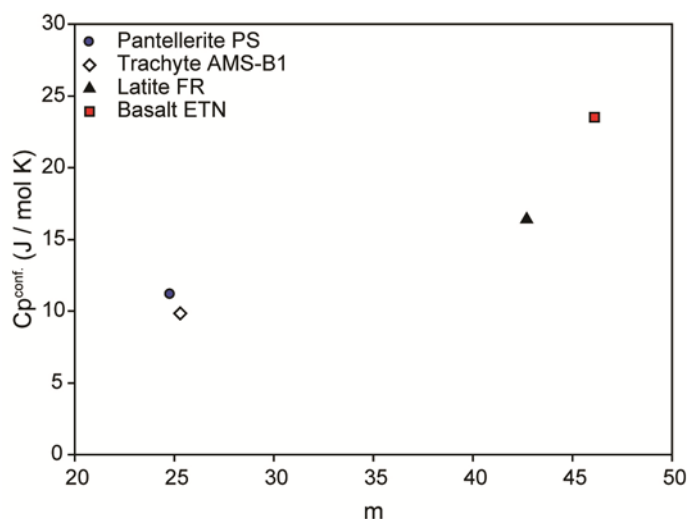


Fig.11: Configurational heat capacities at glass transition temperature of the melts investigated in this work as a function of fragility (m).

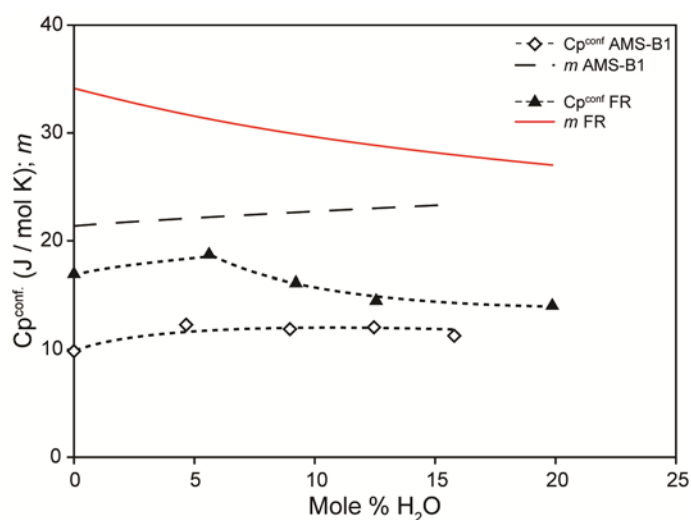


Fig. 12: Configurational heat capacities at glass transition temperature [$C_p^{\text{conf}}(\text{Tg})$] and fragility (m defined by Eq. 7) as a function of water content.

Table 1

Dry composition (mol %) of glasses

	PS	AMS-B1	FR	ETN
SiO ₂	77.76 (±0.27)	65.43 (±0.21)	62.33 (±0.24)	52.48 (±0.16)
TiO ₂	0.41 (±0.4)	0.33 (±0.02)	0.68 (±0.05)	1.37 (±0.06)
Al ₂ O ₃	5.96 (±0.05)	12.29 (±0.07)	11.68 (±0.06)	10.76 (±0.12)
FeO _{TOT}	5.21 (±0.15)	4.28 (±0.05)	6.16 (±0.09)	9.06 (±0.17)
Fe ₂ O ₃	2.30 (±0.15)	-	-	-
FeO	3.14 (±0.11)	-	-	-
MnO	0.36 (±0.06)	0.10 (±0.02)	0.16 (±0.04)	0.23 (±0.03)
MgO	0.13 (±0.02)	2.47 (±0.03)	3.95 (±0.06)	8.89 (±0.09)
CaO	0.68 (±0.06)	5.14 (±0.06)	6.60 (±0.11)	11.73 (±0.11)
Na ₂ O	6.29 (±0.09)	4.09 (±0.06)	4.92 (±0.08)	3.90 (±0.09)
K ₂ O	2.89 (±0.09)	5.71 (±0.05)	3.21 (±0.09)	1.28 (±0.05)
P ₂ O ₅	-	0 (±0.00)	0.22 (±0.04)	0.26 (±0.06)
gfw (g) ^a	66.40	67.20	66.15	64.39
N ^b	3.12	3.13	3.07	2.93

^a Gram formula weight on the basis of one mole of oxides.

^b Number of atoms per gfw.

Fe²⁺/Fe^{tot} ratio of PS was measured by redox titration as reported in Di Genova et al. (2013).

Table 2

Measured water contents (FTIR and KFT), synthesis condition (pressure and temperature) of the hydrous samples and room temperature density of hydrated samples.

Sample	FTIR ^a		KFT ^b	P (kbar)	T (K)	ρ (g l ⁻¹)
	H ₂ O _{molec.}	OH ⁻				
PS 0.5	1.38 (±0.04)	1.65 (±0.06)	2.6 (±0.15)	4	1473	2522 (±1)
PS 1.1	1.95 (±0.06)	2.36 (±0.09)	4.15 (±0.26)	4	1473	2507 (±1)
PS 2.2	4.47 (±0.13)	3.14 (±0.11)	7.36 (±0.21)	4	1473	2473 (±4)
PS 3.5	8.03 (±0.19)	3.48 (±0.14)	11.95 (±0.27)	4	1473	2420 (±2)
AMS-B1 1.3	1.42 (±0.06)	2.35 (±0.07)	4.66 (±0.26)	2.5	1373	2515 (±2)
AMS-B1 2.6	4.93 (±0.09)	3.70 (±0.06)	8.97 (±0.30)	2.5	1373	2472 (±1)
AMS-B1 3.7	8.03 (±0.17)	4.21 (±0.12)	12.46 (±0.27)	2.5	1373	2464 (±3)
AMS-B1 4.8	10.92 (±0.21)	4.33 (±0.13)	15.79 (±0.34)	2.5	1373	2418 (±2)
FR 1.6	-	-	5.6 (±0.19)	2.5	1473	-
FR 2.7	-	-	9.22 (±0.21)	2.5	1473	-
FR 3.8	-	-	12.55 (±0.16)	2.5	1473	-
FR 6.3	-	-	19.86 (±0.25)	2.5	1473	-
ETN 1.4	2.44 (±0.04)	2.59 (±0.06)	4.73 (±0.16)	2.5	1473	2760 (±1)
ETN 2.9	5.46 (±0.19)	4.23 (±0.13)	8.75 (±0.13)	2.5	1473	2741 (±2)
ETN 3.6	7.83 (±0.16)	4.36 (±0.11)	11.99 (±0.17)	2.5	1473	2719 (±3)
ETN 3.8	7.95 (±0.13)	4.53 (±0.12)	12.26 (±0.16)	2.5	1473	2718 (±2)

^a Water content (mol%) measured by Infrared spectroscopy.

^b Water content (mol%) measured by Karl-Fischer titration (KFT).

Molar absorption coefficients for the combination band of OH⁻ at 4500 cm⁻¹ and H₂O_m at 5200 cm⁻¹: PS (Behrens and Zhang, 2009), AMS-B1 (Misiti et al., 2006) and ETN (Lesne et al., 2011).

The duration of each experiment is 24 h.

Table 3

Measured glass transition temperatures (onset, peak, liquid) and liquid, glass and configuration heat capacities of the studied samples.

Sample	T _g ^{onset} (K)			T _g ^{peak} (K)			T _g ^{liquid} (K)			C _{pliq.} ^a	C _{pg} (T _g ^{onset}) ^a	C _{pconf.} ^a
	20 (K/min)	10 (K/min)	5 (K/min)	20 (K/min)	10 (K/min)	5 (K/min)	20 (K/min)	10 (K/min)	5 (K/min)			
PS DRY	813	806	794	875	862	848	906	892	876	91.6	80.5	11.2
PS 0.5	703	698	686	753	744	735	788	766	753	91.5	79.3	12.2
PS 1.1	674	666	659	720	708	701	754	733	723	91.0	78.1	12.9
PS 2.2	644	634	624	685	675	669	713	691	-	90.3	76.7	13.7
PS 3.5	591	583	574	636	626	617	664	647	-	86.1	71.8	14.3
AMS-B1 DRY	914	910	907	987	974	974	1002	987	978	89.4	79.5	9.8
AMS-B1 1.3	779	768	759	834	825	818	848	848	834	89.6	77.3	12.3
AMS-B1 2.6	700	692	682	755	749	728	769	755	740	86.1	74.2	11.9
AMS-B1 3.7	659	651	647	712	701	692	721	708	709	84.3	72.3	12.0
AMS-B1 4.8	625	615	606	674	663	653	689	676	666	83.2	72.2	11.0
FR DRY	937	929	923	985	977	970	1002	993	984	94.8	78.4	16.4
FR 1.6	765	756	743	823	811	811	836	832	831	93.1	73.9	19.2
FR 2.7	691	679	673	755	740	721	758	742	726	88.1	71.9	16.1
FR 3.8	648	641	634	691	680	671	696	695	677	85.6	71.1	14.5
FR 6.3	608	599	591	656	645	637	678	649	642	81.1	67.1	14.0
ETN DRY	915	909	903	954	946	939	966	-	-	96.6	72.2	24.4
ETN 1.4	779	769	762	837	829	823	862	858	-	95.6	68.6	26.9
ETN 2.9	709	707	705	775	760	753	785	-	-	91.1	68.9	22.2
ETN 3.6	663	656	654	760	753	748	778	-	-	90.4	68.1	22.3
ETN 3.8	661	655	651	759	750	745	775	-	-	89.5	67.6	21.9

^a Data in J mol⁻¹ K⁻¹

Table 4

Coefficients for the Maier Kelley equation (Eq. 1) for the glasses studied in this work.

Sample	a	10^3 b	10^{-5} c	st. err ^a
PS DRY	70.0	13.1	-27.9	0.11
PS 0.5	58.9	28.3	-17.1	0.10
PS 1.1	56.5	31.6	-16.3	0.13
PS 2.2	64.1	20.1	-21.9	0.22
PS 3.5	63.5	22.1	-23.0	0.10
AMS DRY	61.4	20.6	-18.1	0.13
AMS-B1 1.3	59.4	23.8	-16.5	0.09
AMS-B1 2.6	49.5	36.1	-9.2	0.07
AMS-B1 3.7	46.1	41.0	-7.3	0.07
AMS-B1 4.8	46.9	42.6	-9.2	0.08
FR DRY	64.6	15.5	-20.5	0.17
FR 1.6	57.7	24.8	-16.0	0.07
FR 2.7	51.6	32.5	-11.3	0.07
FR 3.8	52.6	32.8	-13.3	0.08
FR 6.3	52.1	34.6	-13.3	0.06
ETN DRY	56.5	20.4	-4.7	0.12
ETN 1.4	53.4	23.4	-11.4	0.15
ETN 2.9	51.8	27.3	-13.4	0.08
ETN 3.6	54.8	22.9	-13.0	0.10
ETN 3.8	61.4	13.7	-18.1	0.10

^a Standar error of estimation

Table 5

Measured and calculated liquid heat capacity for the studied samples.

Sample	Measured Cp (J mol ⁻¹ K ⁻¹)	Calculated ^a Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]	Calculated ^b Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]	Calculated ^c Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]	Calculated ^d Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]
PS DRY	91.6	93.1	-1.6	90.4	1.3	90.1	1.6	91.6	0.0
PS 0.5	91.5	92.8	-1.5	90.3	1.3	89.8	1.8	90.1	1.5
PS 1.1	91.0	92.7	-1.9	90.2	0.9	89.7	1.4	89.5	1.7
PS 2.2	90.3	92.5	-2.4	90.0	0.4	89.5	0.9	87.8	2.8
PS 3.5	86.1	92.1	-7.0	89.8	-4.2	89.3	-3.7	85.8	0.3
AMS DRY	90.5	97.5	-7.7	94.8	-4.8	93.7	-3.5	90.5	0.0
AMS-B1 1.3	89.6	96.8	-8.0	94.3	-5.2	92.6	-3.3	88.7	1.1
AMS-B1 2.6	86.1	96.3	-11.9	93.9	-9.1	91.9	-6.8	86.3	-0.2
AMS-B1 3.7	84.3	95.9	-13.8	93.6	-11.0	91.5	-8.5	84.5	-0.2
AMS-B1 4.8	83.2	95.5	-14.8	93.3	-12.1	91.2	-9.5	83.0	0.3
FR DRY	94.8	97.4	-2.7	95.2	-0.4	94.4	0.5	94.8	0.0
FR 1.6	93.1	96.7	-3.8	94.6	-1.6	93.3	-0.2	91.8	1.4
FR 2.7	88.1	96.3	-9.3	94.2	-7.0	92.9	-5.5	89.9	-2.1
FR 3.8	85.6	95.8	-12.0	93.9	-9.7	92.5	-8.1	88.1	-3.0
FR 6.3	81.1	95.0	-17.1	93.2	-14.9	91.9	-13.3	84.2	-3.8
ETN DRY	96.6	98.6	-2.1	97.5	-0.9	95.7	0.9	96.6	0.0
ETN 1.4	95.6	97.9	-2.5	96.9	-1.3	95.1	0.5	93.7	2.0
ETN 2.9	91.1	97.3	-6.8	96.3	-5.6	94.6	-3.8	91.1	0.0
ETN 3.6	90.4	97.0	-7.3	96.0	-6.2	94.5	-4.5	89.9	0.6
ETN 3.8	89.5	96.9	-8.4	95.9	-7.3	94.4	-5.5	89.7	-0.3

^a Calculated values using the Lange and Navrotsky (1992) model^b Calculated values using the Stebbins et al. (1984) model^c Calculated values using the Richet and Bottinga (1985) model with Cp of Al₂O₃ reported by Courtial and Richet (1993)^d Calculated values using Eq. 2 and Eq. 3^{*} 100* (Measured Cp – Calculated Cp)/Measured Cp

Table 6

Measured and calculated liquid heat capacity for the studied samples using a partial molar heat capacity for water in depolymerized melts of 237 J/mol K reported in Bouhifd et al., 2012.

Sample	Measured Cp (J mol ⁻¹ K ⁻¹)	Calculated ^a Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]	Calculated ^b Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]	Calculated ^c Cp (J mol ⁻¹ K ⁻¹)	Δ% [*]
FR 1.6	93.1	105.3	-13.1	103.2	-10.8	101.9	-9.4
FR 2.7	88.1	110.2	-25.2	108.2	-22.9	106.8	-21.3
FR 3.8	85.6	114.8	-34.1	112.8	-31.8	111.4	-30.2
FR 6.3	81.1	124.0	-52.9	122.2	-50.7	120.9	-49.0
ETN 1.4	95.6	105.4	-10.3	104.3	-9.2	102.6	-7.3
ETN 2.9	91.1	112.0	-22.9	110.9	-21.7	109.3	-19.9
ETN 3.6	90.4	114.8	-27.0	113.8	-25.9	112.3	-24.2
ETN 3.8	89.5	115.6	-29.2	114.6	-28.1	113.1	-26.4

^a Calculated values using the Lange and Navrotsky (1992) model^b Calculated values using the Stebbins et al. (1984) model^c Calculated values using the Richet and Bottinga (1985) model with Cp of Al₂O₃ reported by Courtial and Richet (1993)^d 100* (Measured Cp – Calculated Cp)/Measured Cp

Table 7

Calculated partial molar heat capacity of molecular water ($C_{p_{H_2O_{mol}}}$) and OH^- ($C_{p_{OH^-}}$) using Eq. 4.

	NBO/T ^a	$C_{p_{H_2O_{mol}}}$	$C_{p_{OH^-}}$	Err. ^b	SM ^c
AMS-B1	0.10	18.87	80.33	0.56	20.51
PS	0.11	0.00	144.73	0.39	16.02
ETN	0.45	21.29	68.22	0.74	32.39

^aNBO/T = (Mysen, 1988).

^bAverage relative error.

^c Sum of structure modifying oxides (Giordano and Dingwell, 2003).

Table 8

Fit parameters of FR (Misiti et al., 2011) and AMS-B1 (Romano et al., 2001 and Misiti et al., 2006) according to Eq. (8). Numbers in parenthesis are standard

Parameters	FR		AMS-B1	
A_{VFT}	-4.55		-4.55	
b_1	7631.61	(±236.88)	11495.04	(±277.84)
b_2	-2041.55	(±766.30)	-5528.46	(±549.27)
c_1	488.45	(±18.90)	202.47	(±21.20)
c_2	-305.85	(±81.89)	-27.57	(±4.18)
St. Err.	0.51		0.29	

Table 9

Calculated steepness index m (e.g., Plazek and Ngai, 1991) for AMS-B1 and FR sample as a function of water content.

Sample	H ₂ O (mol%)	m
AMS-B1 DRY	-	21.4
AMS-B1 1.3	4.66	22.1
AMS-B1 2.6	8.97	22.6
AMS-B1 3.7	12.46	23.0
AMS-B1 4.8	15.79	23.4
FR DRY	-	34.1
FR 1.6	5.61	31.2
FR 2.7	9.22	29.9
FR 3.8	12.55	28.9
FR 6.3	19.86	27.0