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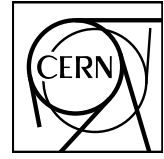
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## Precision measurement of the mass difference between light nuclei and anti-nuclei

ALICE Collaboration\*

### Abstract

The measurement of the mass differences for systems bound by the strong force has reached a very high precision with protons and anti-protons. The extension of such measurement from (anti-)baryons to (anti-)nuclei allows one to probe any difference in the interactions between nucleons and anti-nucleons encoded in the (anti-)nuclei masses. This force is a remnant of the underlying strong interaction among quarks and gluons and can be described by effective theories, but cannot yet be directly derived from quantum chromodynamics. Here we report a measurement of the difference between the ratios of the mass and charge of deuterons ( $d$ ) and anti-deuterons ( $\bar{d}$ ), and  ${}^3\text{He}$  and  ${}^3\bar{\text{He}}$  nuclei carried out with the ALICE (A Large Ion Collider Experiment) detector in Pb-Pb collisions at a centre-of-mass energy per nucleon pair of 2.76 TeV. Our direct measurement of the mass-over-charge differences confirm CPT invariance to an unprecedented precision in the sector of light nuclei. This fundamental symmetry of nature, which exchanges particles with anti-particles, implies that all physics laws are the same under the simultaneous reversal of charge(s) (charge conjugation C), reflection of spatial coordinates (parity transformation P) and time inversion (T).

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\*See Appendix A for the list of collaboration members

The measurement of the mass differences for systems bound by the strong force has reached a very high precision with protons and anti-protons[1, 2]. The extension of such measurement from (anti-)baryons to (anti-)nuclei allows one to probe any difference in the interactions between nucleons and anti-nucleons encoded in the (anti-)nuclei masses. This force is a remnant of the underlying strong interaction among quarks and gluons and can be described by effective theories[3], but cannot yet be directly derived from quantum chromodynamics. Here we report a measurement of the difference between the ratios of the mass and charge of deuterons ( $d$ ) and anti-deuterons ( $\bar{d}$ ), and  ${}^3\text{He}$  and  ${}^3\bar{\text{He}}$  nuclei carried out with the ALICE (A Large Ion Collider Experiment)[4] detector in Pb-Pb collisions at a centre-of-mass energy per nucleon pair of 2.76 TeV. Our direct measurement of the mass-over-charge differences confirm CPT invariance to an unprecedented precision in the sector of light nuclei[5, 6]. This fundamental symmetry of nature, which exchanges particles with anti-particles, implies that all physics laws are the same under the simultaneous reversal of charge(s) (charge conjugation C), reflection of spatial coordinates (parity transformation P) and time inversion (T). Heavy ions are collided at very high energies at the CERN Large Hadron Collider (LHC) to study matter at extremely high temperatures and densities. Under these conditions heavy-ion collisions are a copious source of matter and anti-matter particles and thus are suitable for an experimental investigation of their properties such as mass and electric charge. In relativistic heavy-ion collisions, nuclei and corresponding anti-nuclei are produced with nearly equal rates[7]. Their yields have been measured at the Relativistic Heavy Ion Collider (RHIC) by the STAR[8] and PHENIX[9] experiments and at the LHC by the ALICE[4] experiment. To date, the heaviest anti-nucleus which has been observed[7] is  ${}^4\bar{\text{He}}$  (anti- $\alpha$ ); meanwhile, for lighter nuclei and anti-nuclei, which are more copiously produced, a detailed comparison of their properties is possible. This comparison represents an interesting test of CPT symmetry in an analogous way as done for elementary fermions[10, 11] and bosons[12], and for QED[13, 14] and QCD systems[1, 2, 15–17] (a particular example for the latter being the measurements carried out on neutral kaon decays[18]), with different levels of precision which span over several orders of magnitude. All these measurements can be used to constrain, for different interactions, the parameters of effective field theories that add explicit CPT violating terms to the Standard Model Lagrangian, such as the Standard Model Extension[19] (SME).

The measurements reported in this paper are based on the high-precision tracking and identification capabilities of the ALICE experiment[20]. The main detectors employed in this analysis are the ITS[21] (Inner Tracking System) for the determination of the interaction vertex, the TPC[22] (Time Projection Chamber) for tracking and specific energy loss ( $dE/dx$ ) measurements, and the TOF[23] (Time Of Flight) detector to measure the time  $t_{\text{TOF}}$  needed by each track to traverse the detector. The combined ITS and TPC information is used to determine the track length ( $L$ ) and the rigidity ( $p/z$ , where  $p$  is the momentum and  $z$  the electric charge in units of the elementary charge  $e$ ) of the charged particles in the solenoidal 0.5 T magnetic field of the ALICE central barrel (pseudorapidity  $|\eta| < 0.8$ ). Based on these measurements, we can extract the squared mass-over-charge ratio  $\mu_{\text{TOF}}^2 \equiv (m/z)_{\text{TOF}}^2 = (p/z)^2 [(t_{\text{TOF}}/L)^2 - 1/c^2]$ . The choice of this variable is motivated by the fact that  $\mu^2$  is directly proportional to the square of the time of flight, allowing to better preserve its Gaussian behaviour.

The high precision of the TOF detector, which determines the arrival time of the particle with a resolution of 80 ps[20], allows us to measure a clear signal for (anti-)protons, (anti-)deuterons and (anti-) ${}^3\text{He}$  nuclei over a wide rigidity range ( $1 < p/|z| < 4 \text{ GeV}/c$ ). The main source of background, which is potentially of the same order of the signal, arises from tracks erroneously associated to a TOF hit. To reduce this contamination, a  $2\sigma$  cut (where  $\sigma$  is the standard deviation) around the expected TPC  $dE/dx$  signal is applied. Such a requirement strongly suppresses (to below 4%) this background for rigidities below  $p/|z| < 2.0 \text{ GeV}/c$  for (anti-)deuterons and for all rigidities for (anti-) ${}^3\text{He}$  (to below 1%). For each of the species under study, the mass is extracted by fitting the mass-squared distributions in narrow  $p/|z|$  and  $\eta$  intervals, using a Gaussian with a small exponential tail that reflects the time signal distribution of the TOF detector. Examples of the mass-squared distributions for (anti-)deuterons and (anti-) ${}^3\text{He}$  candidates are reported in Fig. 1 in selected rigidity intervals.

Using mass differences, rather than absolute masses, allows us to reduce the systematic uncertainties related to tracking, spatial alignment (affecting the measurement of the track momentum and length) and time calibration. Despite that, residual effects are still present, due to imperfections in the detector alignment and the description of the magnetic field, which can lead to position-dependent systematic uncertainties. In terms of relative uncertainties, the ones affecting the measurement of the momentum are the largest and independent of the mass, and are the same for all positive (negative) particles in a given momentum interval. It is therefore possible to correct the (anti-)deuteron and the (anti-) $^3\text{He}$  masses by scaling them with the ratio between the (anti-)proton masses recommended by PDG[24] ( $\mu_{p(\bar{p})}^{\text{PDG}}$ ) and the ones measured in the analysis presented here ( $\mu_{p(\bar{p})}^{\text{TOF}}$ ), i.e.  $\mu_{A(\bar{A})} = \mu_{A(\bar{A})}^{\text{TOF}} \times (\mu_{p(\bar{p})}^{\text{PDG}} / \mu_{p(\bar{p})}^{\text{TOF}})$ . These correction factors, which depend on the rigidity, deviate from unity by at most 1%. Conversely, systematic effects connected to the track-length measurement are mass dependent and cannot be completely accounted for using the above correction. However, they are expected to be symmetric for positive and negative particles when inverting the magnetic field. Any residual asymmetry is therefore indicative of remaining systematic uncertainties related to the detector conditions. In order to estimate them and keep these effects under control, both nuclei and anti-nuclei measurements are performed for two opposite magnetic field configurations and then averaged. Their half difference is taken as the estimate of this systematic uncertainty. Other sources of systematic uncertainties are evaluated by varying energy loss corrections applied to the reconstructed momentum, the range and the shape of the background function assumed in the fit of the mass-squared distributions and the track selection criteria. In particular, TPC  $dE/dx$  cuts are varied between one and four standard deviations to probe the sensitivity of the fit results on the residual background, and a tracking quality cut on the distance of closest approach of the track to the vertex is varied to evaluate the influence of secondary particles on the measurement. The sources of systematic uncertainties are found to be fully correlated among all the rigidity intervals, except for those due to the fit procedure and the TPC selection criteria where the uncertainties are uncorrelated. For deuterons and anti-deuterons, the largest relative systematic uncertainties on  $\Delta\mu/\mu$  come from the detector alignment ( $\sim 0.7 \times 10^{-4}$ ), the TPC selection criteria ( $\sim 0.7 \times 10^{-4}$ ) and the secondaries ( $\sim 1.0 \times 10^{-4}$ ). For  $^3\text{He}$  and  $^3\bar{\text{He}}$ , they come from the energy loss corrections ( $\sim 0.7 \times 10^{-3}$ ), the fit procedure ( $\sim 0.5 \times 10^{-3}$ ) and the TPC selection criteria ( $\sim 0.4 \times 10^{-3}$ ).

The (anti-)deuteron and (anti-) $^3\text{He}$  masses are measured as the peak position of the fitting curves of the mass-squared distribution. The mass-over-charge ratio differences between the deuteron and  $^3\text{He}$  and their respective anti-particle are then evaluated as a function of the rigidity of the track, as shown in Fig. 2. The measurements in the individual rigidity intervals are combined, taking into account statistical and systematic uncertainties (correlated and uncorrelated), and the final result is shown in the same figure with one and two standard deviation uncertainty bands. The measured mass-over-charge ratio differences are

$$\Delta\mu_{d\bar{d}} = [1.7 \pm 0.9(\text{stat.}) \pm 2.6(\text{syst.})] \times 10^{-4} \text{ GeV}/c^2, \quad (1)$$

$$\Delta\mu_{^3\text{He}^3\bar{\text{He}}} = [-1.7 \pm 1.2(\text{stat.}) \pm 1.4(\text{syst.})] \times 10^{-3} \text{ GeV}/c^2, \quad (2)$$

corresponding to

$$\frac{\Delta\mu_{d\bar{d}}}{\mu_d} = [0.9 \pm 0.5(\text{stat.}) \pm 1.4(\text{syst.})] \times 10^{-4}, \quad (3)$$

$$\frac{\Delta\mu_{^3\text{He}^3\bar{\text{He}}}}{\mu_{^3\text{He}}} = [-1.2 \pm 0.9(\text{stat.}) \pm 1.0(\text{syst.})] \times 10^{-3}, \quad (4)$$

where  $\mu_d$  and  $\mu_{^3\text{He}}$  are the values recommended by CODATA[25]. The mass-over-charge differences are compatible with zero within the estimated uncertainties, in agreement with CPT invariance expectations.

Given that  $z_{\bar{d}} = -z_d$  and  $z_{^3\bar{\text{He}}} = -z_{^3\text{He}}$  as for the proton and anti-proton[1, 2], the mass-over-charge differences in Eq. 1 and Eq. 2 and the measurement of the mass differences between proton and anti-proton[1, 2] and between neutron and anti-neutron[15, 16] can be used to derive the relative binding

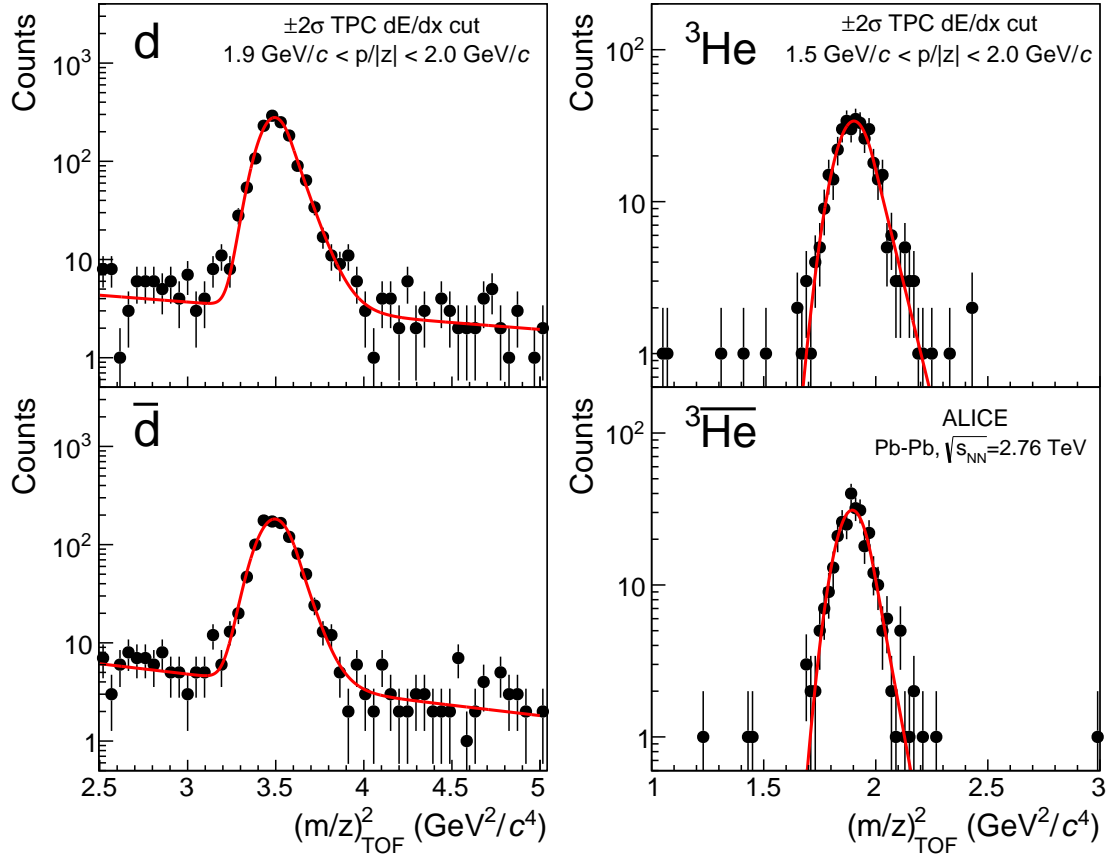
energy differences between the two studied particle species. We obtain

$$\frac{\Delta\epsilon_{d\bar{d}}}{\epsilon_d} = -0.04 \pm 0.05 \text{ (stat.)} \pm 0.12 \text{ (syst.)}, \quad (5)$$

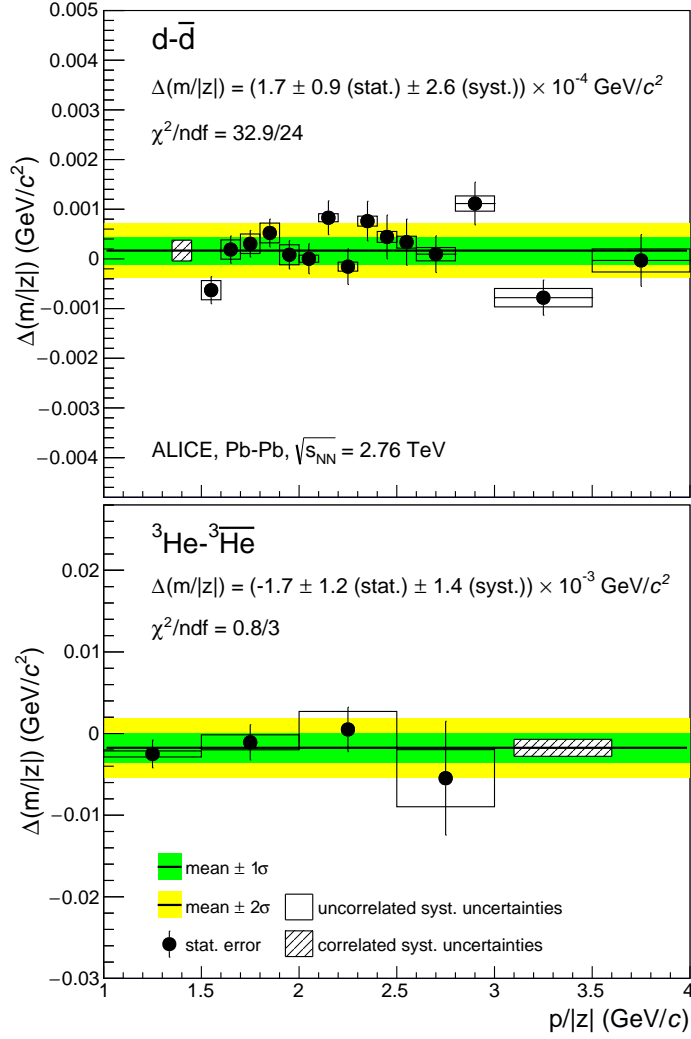
$$\frac{\Delta\epsilon_{{}^3\text{He}{}^3\bar{\text{He}}}}{\epsilon_{{}^3\text{He}}} = 0.24 \pm 0.16 \text{ (stat.)} \pm 0.18 \text{ (syst.)}, \quad (6)$$

where  $\epsilon_A = Zm_p + (A - Z)m_n - m_A$ , being  $m_p$  and  $m_n$  the proton and the neutron mass values recommended by PDG[24] and  $m_A$  the mass value of the nucleus with atomic number  $Z$  and mass number  $A$ , recommended by CODATA[25]. This quantity allows one to explicitly isolate possible violations of the CPT symmetry in the (anti-)nucleon interaction from the one connected to the (anti-)nucleon masses, the latter being constrained with a precision of  $7 \times 10^{-10}$  for the proton/anti-proton system[1, 2]. Our results and the comparisons with previous mass difference measurements for (d- $\bar{d}$ )[26, 27] and ( ${}^3\text{He}$ - ${}^3\bar{\text{He}}$ )[28], as well as binding energy measurements for (d- $\bar{d}$ )[29, 30] are reported in Fig. 3.

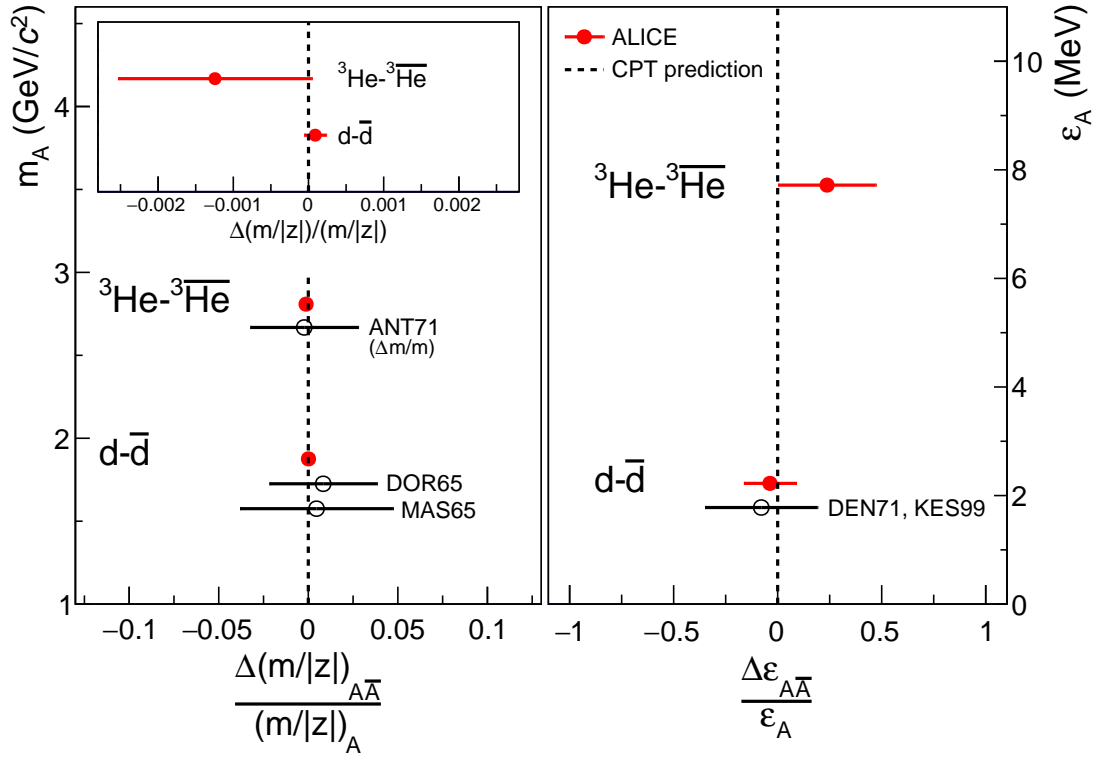
We have shown that the copious production of (anti-)nuclei in relativistic heavy-ion collisions at the LHC represents a unique opportunity to test the CPT invariance of nucleon-nucleon interaction using light nuclei. In particular, we have measured the mass-over-charge ratio differences for deuteron and  ${}^3\text{He}$ . The values are compatible, within uncertainties, with zero and represent a CPT invariance test in systems bound by nuclear forces. The results reported here (Fig. 3, left) represent the highest precision direct measurements of mass differences in the sector of nuclei and they improve by one to two orders of magnitude analogous results originally obtained more than 40 years ago[26–28], and precisely 50 years ago for the anti-deuteron[26, 27]. Remarkably such an improvement is reached in an experiment which is not specifically dedicated to test the CPT invariance in nuclear systems. In the forthcoming years the increase in luminosity and center-of-mass energy at the LHC will allow to push forward the sensitivity of these measurements, and possibly to extend the study to (anti-) ${}^4\text{He}$ . Given the equivalence between mass and binding energy differences, our results also improve (Fig. 3, right) by a factor two the constraints on CPT invariance inferred by existing measurements[29, 30] in the (anti-)deuteron system. The binding energy difference has been determined for the first time in the case of (anti-) ${}^3\text{He}$ , with a relative precision comparable to the one obtained in the (anti-)deuteron system.



**Fig. 1:** Examples of squared mass-over-charge ratio distributions for deuterons (left) and  ${}^3\text{He}$  (right) in selected rigidity intervals. Particle and anti-particle spectra are in the top and bottom plots, respectively. The fit function (red curve) also includes, for the (anti-)deuteron case, an exponential term to describe the background. In the rigidity intervals shown here the background is about 4% for (anti-)deuterons, while it is 0.7% for  ${}^3\text{He}$  and  ${}^3\bar{\text{He}}$ .



**Fig. 2:** The  $d-\bar{d}$  (top) and  ${}^3\text{He}-{}^3\overline{\text{He}}$  (bottom) mass-over-charge ratio difference measurements as a function of the particle rigidity. Vertical bars and open boxes show the statistical and the uncorrelated systematic uncertainties (standard deviations), respectively. Both are taken into account to extract the combined result in the full rigidity range, together with the correlated systematic uncertainty, which is shown as a box with tilted lines. Also shown are the  $1\sigma$  and  $2\sigma$  bands around the central value, where  $\sigma$  is the sum in quadrature of the statistical and systematic uncertainties.



**Fig. 3:** The ALICE measurements for  $d\bar{d}$  and  ${}^3\text{He}\text{-}{}^3\bar{\text{He}}$  mass-over-charge ratio differences compared with CPT invariance expectation (dotted lines) and existing mass measurements MAS65[26], DOR65[27] and ANT71[28] (left panel). The inset shows the ALICE results on a finer  $\Delta(m/z)/(m/z)$  scale. The right panel shows our determination of the binding energy differences compared with direct measurements from DEN71[29] and KES99[30]. Error bars represent the sum in quadrature of the statistical and systematic uncertainties (standard deviations).



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## A The ALICE Collaboration

J. Adam<sup>39</sup>, D. Adamová<sup>82</sup>, M.M. Aggarwal<sup>86</sup>, G. Aglieri Rinella<sup>36</sup>, M. Agnello<sup>110</sup>, N. Agrawal<sup>47</sup>, Z. Ahammed<sup>130</sup>, I. Ahmed<sup>16</sup>, S.U. Ahn<sup>67</sup>, I. Aimo<sup>93,110</sup>, S. Aiola<sup>135</sup>, M. Ajaz<sup>16</sup>, A. Akindinov<sup>57</sup>, S.N. Alam<sup>130</sup>, D. Aleksandrov<sup>99,99</sup>, B. Alessandro<sup>110</sup>, D. Alexandre<sup>101</sup>, R. Alfaro Molina<sup>63</sup>, A. Alici<sup>104,12</sup>, A. Alkin<sup>3</sup>, J. Alme<sup>37</sup>, T. Alt<sup>42</sup>, S. Altinpinar<sup>18,18</sup>, I. Altsybeev<sup>129</sup>, C. Alves Garcia Prado<sup>118</sup>, C. Andrei<sup>77</sup>, A. Andronic<sup>96</sup>, V. Anguelov<sup>92</sup>, J. Anielski<sup>53</sup>, T. Antičić<sup>97</sup>, F. Antinori<sup>107</sup>, P. Antonioli<sup>104</sup>, L. Aphecetche<sup>112</sup>, H. Appelshäuser<sup>52</sup>, S. Arcelli<sup>28</sup>, N. Armesto<sup>17</sup>, R. Arnaldi<sup>110</sup>, T. Aronsson<sup>135</sup>, I.C. Arsene<sup>22</sup>, M. Arslanodk<sup>52</sup>, A. Augustinus<sup>36</sup>, R. Averbeck<sup>96</sup>, M.D. Azmi<sup>19,19</sup>, M. Bach<sup>42</sup>, A. Badalà<sup>106</sup>, Y.W. Baek<sup>43</sup>, S. Bagnasco<sup>110</sup>, R. Bailhache<sup>52</sup>, R. Bala<sup>89</sup>, A. Baldisseri<sup>15</sup>, M. Ball<sup>91</sup>, F. Baltasar Dos Santos Pedrosa<sup>36</sup>, R.C. Baral<sup>60</sup>, A.M. Barbano<sup>110</sup>, R. Barbera<sup>29</sup>, F. Barile<sup>33</sup>, G.G. Barnaföldi<sup>134</sup>, L.S. Barnby<sup>101</sup>, V. Barret<sup>69</sup>, P. Bartalini<sup>7</sup>, J. Bartke<sup>115</sup>, E. Bartsch<sup>52</sup>, M. Basile<sup>28</sup>, N. Bastid<sup>69</sup>, S. Basu<sup>130</sup>, B. Bathen<sup>53</sup>, G. Batigne<sup>112</sup>, A. Batista Camejo<sup>69</sup>, B. Batyunya<sup>65</sup>, P.C. Batzing<sup>22</sup>, I.G. Bearden<sup>79</sup>, H. Beck<sup>52</sup>, C. Bedda<sup>110</sup>, N.K. Behera<sup>48,47</sup>, I. Belikov<sup>54</sup>, F. 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Paticchio<sup>103</sup>, B. Paul<sup>100</sup>, T. Pawlak<sup>132</sup>, T. Peitzmann<sup>56</sup>, H. Pereira Da Costa<sup>15</sup>, E. Pereira De Oliveira Filho<sup>118</sup>, D. Peresunko<sup>75,99</sup>, C.E. Pérez Lara<sup>80</sup>, V. Peskov<sup>52,52</sup>, Y. Pestov<sup>5</sup>, V. Petráček<sup>39</sup>, V. Petrov<sup>111</sup>, M. Petrovici<sup>77</sup>, C. Petta<sup>29</sup>, S. Piano<sup>109</sup>, M. Pikna<sup>38</sup>, P. Pillot<sup>112</sup>, O. Pinazza<sup>104,36</sup>, L. Pinsky<sup>120</sup>, D.B. Piyathana<sup>120</sup>, M. Płoskoń<sup>73</sup>, M. Planinic<sup>127</sup>, J. Pluta<sup>132</sup>, S. Pochybova<sup>134</sup>, P.L.M. Podesta-Lerma<sup>117,117</sup>, M.G. Poghosyan<sup>85</sup>, B. Polichtchouk<sup>111</sup>, N. Poljak<sup>127</sup>, W. Poonsawat<sup>113</sup>, A. Pop<sup>77</sup>, S. Porteboeuf-Houssais<sup>69</sup>, J. Porter<sup>73</sup>, J. Pospisil<sup>82</sup>, S.K. Prasad<sup>4</sup>, R. Preghenella<sup>104,104,36</sup>, F. Prino<sup>110</sup>, C.A. Pruneau<sup>133</sup>, I. Pshenichnov<sup>55</sup>, M. Puccio<sup>110</sup>, G. Puddu<sup>25</sup>, P. Pujahari<sup>133</sup>, V. Punin<sup>98</sup>, J. Putschke<sup>133</sup>, H. Qvigstad<sup>22</sup>, A. Rachevski<sup>109</sup>, S. Raha<sup>4</sup>, S. Rajput<sup>89</sup>, J. Rak<sup>121</sup>, A. Rakotozafindrabe<sup>15</sup>, L. Ramello<sup>32</sup>, R. Raniwala<sup>90</sup>, S. Raniwala<sup>90</sup>, S.S. Räsänen<sup>45</sup>, B.T. Rascanu<sup>52</sup>, D. Rathee<sup>86</sup>, V. Razazi<sup>25</sup>, K.F. Read<sup>123</sup>, J.S. Real<sup>70</sup>, K. Redlich<sup>76</sup>, R.J. Reed<sup>133</sup>, A. Rehman<sup>18</sup>, P. Reichelt<sup>52</sup>, M. Reicher<sup>56</sup>, F. Reidt<sup>92,36</sup>, X. Ren<sup>7</sup>, R. Renfordt<sup>52</sup>, A.R. Reolon<sup>71</sup>, A. Reshetin<sup>55</sup>, F. Rettig<sup>42</sup>, J.-P. Revol<sup>12</sup>, K. Reygers<sup>92</sup>, V. Riabov<sup>84</sup>, R.A. Ricci<sup>72</sup>, T. Richert<sup>34</sup>, M. Richter<sup>22,22</sup>, P. Riedler<sup>36</sup>, W. Riegler<sup>36</sup>, F. Riggi<sup>29</sup>, C. Ristea<sup>61</sup>, A. Rivetti<sup>110</sup>, E. Rocco<sup>56</sup>, M. Rodríguez Cahuantzi<sup>11,2,11</sup>, A. Rodríguez Manso<sup>80</sup>, K. Røed<sup>22</sup>, E. Rogochaya<sup>65</sup>, D. Rohr<sup>42</sup>, D. Röhrich<sup>18</sup>, R. Romita<sup>122</sup>, F. Ronchetti<sup>71</sup>, L. Ronflette<sup>112</sup>, P. Rosnet<sup>69</sup>, A. Rossi<sup>36</sup>, F. Roukoutakis<sup>87</sup>, A. Roy<sup>48</sup>, C. Roy<sup>54</sup>, P. Roy<sup>100</sup>, A.J. Rubio Montero<sup>10</sup>, R. Rui<sup>26</sup>, R. Russo<sup>27</sup>, E. Ryabinkin<sup>99</sup>, Y. Ryabov<sup>84</sup>, A. Rybicki<sup>115</sup>, S. Sadovsky<sup>111</sup>, K. Šafařík<sup>36</sup>, B. Sahlmüller<sup>52</sup>, P. Sahoo<sup>48</sup>, R. Sahoo<sup>48</sup>, S. Sahoo<sup>60</sup>, P.K. Sahu<sup>60</sup>, J. Saini<sup>130</sup>, S. Sakai<sup>71</sup>, M.A. Saleh<sup>133</sup>, C.A. Salgado<sup>17</sup>, J. Salzwedel<sup>20</sup>,

S. Sambyal<sup>89</sup>, V. Samsonov<sup>84</sup>, X. Sanchez Castro<sup>54</sup>, L. Šándor<sup>58</sup>, A. Sandoval<sup>63</sup>, M. Sano<sup>126</sup>, G. Santagati<sup>29</sup>, D. Sarkar<sup>130</sup>, E. Scapparone<sup>104</sup>, F. Scarlassara<sup>30</sup>, R.P. Scharenberg<sup>94</sup>, C. Schiaua<sup>77</sup>, R. Schicker<sup>92</sup>, C. Schmidt<sup>96</sup>, H.R. Schmidt<sup>35</sup>, S. Schuchmann<sup>52</sup>, J. Schukraft<sup>36</sup>, M. Schulc<sup>39</sup>, T. Schuster<sup>135</sup>, Y. Schutz<sup>112,36</sup>, K. Schwarz<sup>96</sup>, K. Schweda<sup>96</sup>, G. Scioli<sup>28</sup>, E. Scomparin<sup>110</sup>, R. Scott<sup>123</sup>, K.S. Seeder<sup>118</sup>, J.E. Seger<sup>85</sup>, Y. Sekiguchi<sup>125</sup>, I. Selyuzhenkov<sup>96,96</sup>, K. Senosi<sup>64</sup>, J. Seo<sup>66,95</sup>, E. Serradilla<sup>10,63</sup>, A. Sevcenco<sup>61</sup>, A. Shabanov<sup>55</sup>, A. Shabetai<sup>112</sup>, O. Shadura<sup>3</sup>, R. Shahoyan<sup>36</sup>, A. Shangaraev<sup>111</sup>, A. Sharma<sup>89</sup>, N. Sharma<sup>60,123</sup>, K. Shigaki<sup>46</sup>, K. Shtejer<sup>9,27</sup>, Y. Sibiriyak<sup>99</sup>, S. Siddhanta<sup>105</sup>, K.M. Sielewicz<sup>36</sup>, T. Siemiarczuk<sup>76</sup>, D. Silvermyr<sup>83,34</sup>, C. Silvestre<sup>70</sup>, G. Simatovic<sup>127</sup>, G. Simonetti<sup>36</sup>, R. Singaraju<sup>130</sup>, R. Singh<sup>89,78</sup>, S. Singha<sup>78,130</sup>, V. Singhal<sup>130</sup>, B.C. Sinha<sup>130</sup>, T. Sinha<sup>100</sup>, B. Sitar<sup>38</sup>, M. Sitta<sup>32</sup>, T.B. Skaali<sup>22</sup>, M. Slupecki<sup>121</sup>, N. Smirnov<sup>135</sup>, R.J.M. Snellings<sup>56</sup>, T.W. Snellman<sup>121</sup>, C. Sogaard<sup>34</sup>, R. Soltz<sup>74</sup>, J. Song<sup>95</sup>, M. Song<sup>136</sup>, Z. Song<sup>7</sup>, F. Soramel<sup>30</sup>, S. Sorensen<sup>123</sup>, M. Spacek<sup>39</sup>, E. Spiriti<sup>71</sup>, I. Sputowska<sup>115</sup>, M. Spyropoulou-Stassinaki<sup>87</sup>, B.K. Srivastava<sup>94</sup>, J. Stachel<sup>92</sup>, I. Stan<sup>61</sup>, G. Stefanek<sup>76</sup>, M. Steinpreis<sup>20</sup>, E. Stenlund<sup>34</sup>, G. Steyn<sup>64</sup>, J.H. Stiller<sup>92</sup>, D. Stocco<sup>112</sup>, P. Strmen<sup>38</sup>, A.A.P. Suaide<sup>118</sup>, T. Sugitate<sup>46</sup>, C. Suire<sup>50</sup>, M. Suleymanov<sup>16</sup>, R. Sultanov<sup>57</sup>, M. Šumbera<sup>82</sup>, T.J.M. Symons<sup>73</sup>, A. Szabo<sup>38</sup>, A. Szanto de Toledo<sup>118,i</sup>, I. Szarka<sup>38</sup>, A. Szczepankiewicz<sup>36</sup>, M. Szymanski<sup>132</sup>, J. Takahashi<sup>119</sup>, N. Tanaka<sup>126</sup>, M.A. Tangaro<sup>33</sup>, J.D. Tapia Takaki<sup>ii,50</sup>, A. Tarantola Peloni<sup>52</sup>, M. Tariq<sup>19</sup>, M.G. Tarzila<sup>77</sup>, A. Tauro<sup>36</sup>, G. Tejeda Muñoz<sup>2</sup>, A. Telesca<sup>36</sup>, K. Terasaki<sup>125</sup>, C. Terrevoli<sup>30,25</sup>, B. Teyssier<sup>128</sup>, J. Thäder<sup>96,73</sup>, D. Thomas<sup>56,116</sup>, R. Tieulent<sup>128</sup>, A.R. Timmins<sup>120</sup>, A. Toia<sup>52</sup>, S. Trogolo<sup>110</sup>, V. Trubnikov<sup>3</sup>, W.H. Trzaska<sup>121</sup>, T. Tsuji<sup>125</sup>, A. Tumkin<sup>98</sup>, R. Turrisi<sup>107</sup>, T.S. Tveter<sup>22</sup>, K. Ullaland<sup>18</sup>, A. Uras<sup>128,128</sup>, G.L. Usai<sup>25</sup>, A. Utrobicic<sup>127</sup>, M. Vajzer<sup>82</sup>, M. Vala<sup>58</sup>, L. Valencia Palomo<sup>69</sup>, S. Vallero<sup>27</sup>, J. Van Der Maarel<sup>56</sup>, J.W. Van Hoorne<sup>36</sup>, M. van Leeuwen<sup>56</sup>, T. Vanat<sup>82</sup>, P. Vande Vyvre<sup>36</sup>, D. Varga<sup>134</sup>, A. Vargas<sup>2</sup>, M. Vargyas<sup>121</sup>, R. Varma<sup>47</sup>, M. Vasileiou<sup>87</sup>, A. Vasiliev<sup>99</sup>, A. Vauthier<sup>70</sup>, V. Vechernin<sup>129</sup>, A.M. Veen<sup>56</sup>, M. Veldhoen<sup>56</sup>, A. Velure<sup>18</sup>, M. Venaruzzo<sup>72</sup>, E. Vercellin<sup>27</sup>, S. Vergara Limón<sup>2</sup>, R. Vernet<sup>8</sup>, M. Verweij<sup>133</sup>, L. Vickovic<sup>114</sup>, G. Viesti<sup>30,i</sup>, J. Viinikainen<sup>121</sup>, Z. Vilakazi<sup>124</sup>, O. Villalobos Baillie<sup>101</sup>, A. Vinogradov<sup>99</sup>, L. Vinogradov<sup>129</sup>, Y. Vinogradov<sup>98</sup>, T. Virgili<sup>31</sup>, V. Vislavicius<sup>34</sup>, Y.P. Viyogi<sup>130</sup>, A. Vodopyanov<sup>65</sup>, M.A. Völkl<sup>92</sup>, K. Voloshin<sup>57</sup>, S.A. Voloshin<sup>133</sup>, G. Volpe<sup>36,134</sup>, B. von Haller<sup>36</sup>, I. Vorobyev<sup>91</sup>, D. Vranic<sup>96,36</sup>, J. Vrláková<sup>40</sup>, B. Vulpescu<sup>69</sup>, A. Vyushin<sup>98</sup>, B. Wagner<sup>18</sup>, J. Wagner<sup>96</sup>, H. Wang<sup>56</sup>, M. Wang<sup>7,112</sup>, Y. Wang<sup>92</sup>, D. Watanabe<sup>126</sup>, M. Weber<sup>36,120</sup>, S.G. Weber<sup>96</sup>, J.P. Wessels<sup>53</sup>, U. Westerhoff<sup>53</sup>, J. Wiechula<sup>35</sup>, J. Wikne<sup>22</sup>, M. Wilde<sup>53</sup>, G. Wilk<sup>76</sup>, J. Wilkinson<sup>92</sup>, M.C.S. Williams<sup>104</sup>, B. Windelband<sup>92</sup>, M. Winn<sup>92</sup>, C.G. Yaldo<sup>133</sup>, Y. Yamaguchi<sup>125</sup>, H. Yang<sup>56,56</sup>, P. Yang<sup>7</sup>, S. Yano<sup>46</sup>, S. Yasnopolskiy<sup>99</sup>, Z. Yin<sup>7</sup>, H. Yokoyama<sup>126</sup>, I.-K. Yoo<sup>95</sup>, V. Yurchenko<sup>3</sup>, I. Yushmanov<sup>99</sup>, A. Zaborowska<sup>132</sup>, V. Zaccolo<sup>79</sup>, A. Zaman<sup>16</sup>, C. Zampolli<sup>104</sup>, H.J.C. Zanolli<sup>118</sup>, S. Zaporozhets<sup>65</sup>, A. Zarochentsev<sup>129</sup>, P. Závada<sup>59</sup>, N. Zaviyalov<sup>98</sup>, H. Zbroszczyk<sup>132</sup>, I.S. Zgura<sup>61</sup>, M. Zhalov<sup>84</sup>, H. Zhang<sup>7</sup>, X. Zhang<sup>73</sup>, Y. Zhang<sup>7</sup>, C. Zhao<sup>22</sup>, N. Zhigareva<sup>57</sup>, D. Zhou<sup>7</sup>, Y. Zhou<sup>56</sup>, Z. Zhou<sup>18</sup>, H. Zhu<sup>7</sup>, J. Zhu<sup>7,112</sup>, X. Zhu<sup>7</sup>, A. Zichichi<sup>12,28</sup>, A. Zimmermann<sup>92</sup>, M.B. Zimmermann<sup>53,36</sup>, G. Zinovjev<sup>3</sup>, M. Zyzak<sup>42</sup>

## Affiliation notes

<sup>i</sup> Deceased

<sup>ii</sup> Also at: University of Kansas, Lawrence, Kansas, United States

## Collaboration Institutes

- <sup>1</sup> A.I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation, Yerevan, Armenia
- <sup>2</sup> Benemérita Universidad Autónoma de Puebla, Puebla, Mexico
- <sup>3</sup> Bogolyubov Institute for Theoretical Physics, Kiev, Ukraine
- <sup>4</sup> Bose Institute, Department of Physics and Centre for Astroparticle Physics and Space Science (CAPSS), Kolkata, India
- <sup>5</sup> Budker Institute for Nuclear Physics, Novosibirsk, Russia
- <sup>6</sup> California Polytechnic State University, San Luis Obispo, California, United States
- <sup>7</sup> Central China Normal University, Wuhan, China
- <sup>8</sup> Centre de Calcul de l'IN2P3, Villeurbanne, France
- <sup>9</sup> Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba
- <sup>10</sup> Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain
- <sup>11</sup> Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico
- <sup>12</sup> Centro Fermi - Museo Storico della Fisica e Centro Studi e Ricerche "Enrico Fermi", Rome, Italy
- <sup>13</sup> Chicago State University, Chicago, Illinois, USA

- 14 China Institute of Atomic Energy, Beijing, China
- 15 Commissariat à l’Energie Atomique, IRFU, Saclay, France
- 16 COMSATS Institute of Information Technology (CIIT), Islamabad, Pakistan
- 17 Departamento de Física de Partículas and IGFAE, Universidad de Santiago de Compostela, Santiago de Compostela, Spain
- 18 Department of Physics and Technology, University of Bergen, Bergen, Norway
- 19 Department of Physics, Aligarh Muslim University, Aligarh, India
- 20 Department of Physics, Ohio State University, Columbus, Ohio, United States
- 21 Department of Physics, Sejong University, Seoul, South Korea
- 22 Department of Physics, University of Oslo, Oslo, Norway
- 23 Dipartimento di Elettrotecnica ed Elettronica del Politecnico, Bari, Italy
- 24 Dipartimento di Fisica dell’Università ‘La Sapienza’ and Sezione INFN Rome, Italy
- 25 Dipartimento di Fisica dell’Università and Sezione INFN, Cagliari, Italy
- 26 Dipartimento di Fisica dell’Università and Sezione INFN, Trieste, Italy
- 27 Dipartimento di Fisica dell’Università and Sezione INFN, Turin, Italy
- 28 Dipartimento di Fisica e Astronomia dell’Università and Sezione INFN, Bologna, Italy
- 29 Dipartimento di Fisica e Astronomia dell’Università and Sezione INFN, Catania, Italy
- 30 Dipartimento di Fisica e Astronomia dell’Università and Sezione INFN, Padova, Italy
- 31 Dipartimento di Fisica ‘E.R. Caianiello’ dell’Università and Gruppo Collegato INFN, Salerno, Italy
- 32 Dipartimento di Scienze e Innovazione Tecnologica dell’Università del Piemonte Orientale and Gruppo Collegato INFN, Alessandria, Italy
- 33 Dipartimento Interateneo di Fisica ‘M. Merlin’ and Sezione INFN, Bari, Italy
- 34 Division of Experimental High Energy Physics, University of Lund, Lund, Sweden
- 35 Eberhard Karls Universität Tübingen, Tübingen, Germany
- 36 European Organization for Nuclear Research (CERN), Geneva, Switzerland
- 37 Faculty of Engineering, Bergen University College, Bergen, Norway
- 38 Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovakia
- 39 Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Prague, Czech Republic
- 40 Faculty of Science, P.J. Šafárik University, Košice, Slovakia
- 41 Faculty of Technology, Buskerud and Vestfold University College, Vestfold, Norway
- 42 Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 43 Gangneung-Wonju National University, Gangneung, South Korea
- 44 Gauhati University, Department of Physics, Guwahati, India
- 45 Helsinki Institute of Physics (HIP), Helsinki, Finland
- 46 Hiroshima University, Hiroshima, Japan
- 47 Indian Institute of Technology Bombay (IIT), Mumbai, India
- 48 Indian Institute of Technology Indore, Indore (IITI), India
- 49 Inha University, Incheon, South Korea
- 50 Institut de Physique Nucléaire d’Orsay (IPNO), Université Paris-Sud, CNRS-IN2P3, Orsay, France
- 51 Institut für Informatik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 52 Institut für Kernphysik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 53 Institut für Kernphysik, Westfälische Wilhelms-Universität Münster, Münster, Germany
- 54 Institut Pluridisciplinaire Hubert Curien (IPHC), Université de Strasbourg, CNRS-IN2P3, Strasbourg, France
- 55 Institute for Nuclear Research, Academy of Sciences, Moscow, Russia
- 56 Institute for Subatomic Physics of Utrecht University, Utrecht, Netherlands
- 57 Institute for Theoretical and Experimental Physics, Moscow, Russia
- 58 Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia
- 59 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- 60 Institute of Physics, Bhubaneswar, India
- 61 Institute of Space Science (ISS), Bucharest, Romania
- 62 Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 63 Instituto de Física, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 64 iThemba LABS, National Research Foundation, Somerset West, South Africa

- 65 Joint Institute for Nuclear Research (JINR), Dubna, Russia
- 66 Konkuk University, Seoul, South Korea
- 67 Korea Institute of Science and Technology Information, Daejeon, South Korea
- 68 KTO Karatay University, Konya, Turkey
- 69 Laboratoire de Physique Corpusculaire (LPC), Clermont Université, Université Blaise Pascal, CNRS-IN2P3, Clermont-Ferrand, France
- 70 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS-IN2P3, Grenoble, France
- 71 Laboratori Nazionali di Frascati, INFN, Frascati, Italy
- 72 Laboratori Nazionali di Legnaro, INFN, Legnaro, Italy
- 73 Lawrence Berkeley National Laboratory, Berkeley, California, United States
- 74 Lawrence Livermore National Laboratory, Livermore, California, United States
- 75 Moscow Engineering Physics Institute, Moscow, Russia
- 76 National Centre for Nuclear Studies, Warsaw, Poland
- 77 National Institute for Physics and Nuclear Engineering, Bucharest, Romania
- 78 National Institute of Science Education and Research, Bhubaneswar, India
- 79 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- 80 Nikhef, National Institute for Subatomic Physics, Amsterdam, Netherlands
- 81 Nuclear Physics Group, STFC Daresbury Laboratory, Daresbury, United Kingdom
- 82 Nuclear Physics Institute, Academy of Sciences of the Czech Republic, Řež u Prahy, Czech Republic
- 83 Oak Ridge National Laboratory, Oak Ridge, Tennessee, United States
- 84 Petersburg Nuclear Physics Institute, Gatchina, Russia
- 85 Physics Department, Creighton University, Omaha, Nebraska, United States
- 86 Physics Department, Panjab University, Chandigarh, India
- 87 Physics Department, University of Athens, Athens, Greece
- 88 Physics Department, University of Cape Town, Cape Town, South Africa
- 89 Physics Department, University of Jammu, Jammu, India
- 90 Physics Department, University of Rajasthan, Jaipur, India
- 91 Physik Department, Technische Universität München, Munich, Germany
- 92 Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- 93 Politecnico di Torino, Turin, Italy
- 94 Purdue University, West Lafayette, Indiana, United States
- 95 Pusan National University, Pusan, South Korea
- 96 Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt, Germany
- 97 Rudjer Bošković Institute, Zagreb, Croatia
- 98 Russian Federal Nuclear Center (VNIIEF), Sarov, Russia
- 99 Russian Research Centre Kurchatov Institute, Moscow, Russia
- 100 Saha Institute of Nuclear Physics, Kolkata, India
- 101 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- 102 Sección Física, Departamento de Ciencias, Pontificia Universidad Católica del Perú, Lima, Peru
- 103 Sezione INFN, Bari, Italy
- 104 Sezione INFN, Bologna, Italy
- 105 Sezione INFN, Cagliari, Italy
- 106 Sezione INFN, Catania, Italy
- 107 Sezione INFN, Padova, Italy
- 108 Sezione INFN, Rome, Italy
- 109 Sezione INFN, Trieste, Italy
- 110 Sezione INFN, Turin, Italy
- 111 SSC IHEP of NRC Kurchatov institute, Protvino, Russia
- 112 SUBATECH, Ecole des Mines de Nantes, Université de Nantes, CNRS-IN2P3, Nantes, France
- 113 Suranaree University of Technology, Nakhon Ratchasima, Thailand
- 114 Technical University of Split FESB, Split, Croatia
- 115 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Cracow, Poland
- 116 The University of Texas at Austin, Physics Department, Austin, Texas, USA
- 117 Universidad Autónoma de Sinaloa, Culiacán, Mexico



- 118 Universidade de São Paulo (USP), São Paulo, Brazil
- 119 Universidade Estadual de Campinas (UNICAMP), Campinas, Brazil
- 120 University of Houston, Houston, Texas, United States
- 121 University of Jyväskylä, Jyväskylä, Finland
- 122 University of Liverpool, Liverpool, United Kingdom
- 123 University of Tennessee, Knoxville, Tennessee, United States
- 124 University of the Witwatersrand, Johannesburg, South Africa
- 125 University of Tokyo, Tokyo, Japan
- 126 University of Tsukuba, Tsukuba, Japan
- 127 University of Zagreb, Zagreb, Croatia
- 128 Université de Lyon, Université Lyon 1, CNRS/IN2P3, IPN-Lyon, Villeurbanne, France
- 129 V. Fock Institute for Physics, St. Petersburg State University, St. Petersburg, Russia
- 130 Variable Energy Cyclotron Centre, Kolkata, India
- 131 Vinča Institute of Nuclear Sciences, Belgrade, Serbia
- 132 Warsaw University of Technology, Warsaw, Poland
- 133 Wayne State University, Detroit, Michigan, United States
- 134 Wigner Research Centre for Physics, Hungarian Academy of Sciences, Budapest, Hungary
- 135 Yale University, New Haven, Connecticut, United States
- 136 Yonsei University, Seoul, South Korea
- 137 Zentrum für Technologietransfer und Telekommunikation (ZTT), Fachhochschule Worms, Worms, Germany