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**Simulation and experiment reveal a complex scenario for the adsorption of an antifungal drug in ordered mesoporous silica**

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# Simulation and experiment reveal a complex scenario for the adsorption of an antifungal drug in ordered mesoporous silica

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## ABSTRACT

Ordered mesoporous silicas have been widely investigated as drug carriers in several fields, from tissue engineering to cancer therapy. The knowledge of the specific interactions between the surface of mesoporous silicas and drugs is necessary to guide development of new and improved drug delivery systems. However, such knowledge is still scarce, due to the arduous interpretation of experimental results. In this work, we characterize the incorporation of clotrimazole, a common antifungal drug, inside ordered mesoporous silica by means of a joint computational and experimental approach. Experimentally the drug was loaded through supercritical CO<sub>2</sub> and its adsorption investigated through infrared spectroscopy, N<sub>2</sub> adsorption isotherms and thermogravimetric analysis. Modelling involved static and dynamic Density Functional Theory simulations of clotrimazole adsorbed on realistic models of amorphous silica surfaces. A good agreement between the computational and experimental results was obtained, concerning the energies of adsorption, the infrared spectra and the distribution of drug inside the mesopores. However, a complete interpretation of the experimental results was possible only when simultaneously considering all the complex aspects of the drug-silica interaction. Indeed, the combination of both approaches allowed us to describe the drug-silica interface as a mix of multiple interaction configurations, based on a subtle balance of hydrogen bonding and dispersion interactions. Furthermore, at high drug loading, clotrimazole molecules are statistically distributed on the pore walls forming an adsorbed molecular layer. Finally, notwithstanding the

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stable interactions, the drug still exhibits a significant mobility at room temperature, moving on a complex potential energy surface, as revealed by molecular dynamics simulations.

## INTRODUCTION

Over the past three decades, a rapid growth has affected the research area of drug delivery, aiming at optimizing drug efficiency while simultaneously reducing adverse collateral effects.<sup>1</sup> Several studies have reported that pharmacokinetics, drug efficiency and suppression of undesired side effects in different pathological conditions can be improved by correct timing of drug administration and controlled kinetics of drug release.<sup>2</sup>

Recently, the interest has concerned the use of mesoporous materials as controlled drug delivery matrices thanks to their uniform mesoporous structures, high surface areas, tunable pore sizes and well-defined surface properties.<sup>3-6</sup>

In this context, Ordered Mesoporous Silicas (OMSs) have been widely investigated as drug carriers in several fields, from tissue engineering to cancer therapy.<sup>7</sup> Initially, the research on OMSs for drug delivery was focused on the achievement of controlled release formulations. The release kinetics of drugs by OMSs depends on several carrier properties, including pore size, pore connectivity and the chemical composition of the surface.<sup>8</sup> A recent emerging feature of OMS carriers is the enhanced oral bioavailability of molecules poorly soluble in water.<sup>9</sup> It has been shown that both small and large molecular drugs can be entrapped within the mesopores by an impregnation process and liberated via a diffusion controlled mechanism.<sup>10</sup>

In the first pioneering work by Vallet-Regi et al.,<sup>4</sup> drug incorporation in MCM-41 was carried out by adsorption from a solution using hexane as solvent. Since the long-term toxicity of n-hexane in humans is well known,<sup>11,12</sup> many other solvents have been studied for the incorporation. The incorporation by supercritical carbon dioxide (scCO<sub>2</sub>) is an alternative to adsorption or impregnation from a liquid solution.<sup>13-15</sup> Carbon dioxide is one of the most commonly used fluids in supercritical fluid technology. Its main advantages are a critical temperature close to ambient temperature (304.25 K)

1 and a not too high critical pressure (7.39 MPa).<sup>16</sup> In addition, it is non-flammable, has low cost and  
2 low toxicity. Being a supercritical fluid, its physical properties are halfway between a gas and a liq-  
3 uid; in particular, it has a solvent power like a liquid and a high diffusivity like a gas. At pressures  
4 and temperatures not too far from its critical point, a supercritical fluid has a high compressibility,  
5 therefore its density and hence its solvent power are easily adjustable over a wide range with a min-  
6 imal change in temperature or pressure.  
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13 When a drug is incorporated inside a material, such as mesoporous silica, the interactions occurring  
14 at its surface are of great importance. They can deeply influence stability, absorption and manufac-  
15 turability of the formulation<sup>17</sup> and, for drug delivery systems, are essential for determining the final  
16 performance of the product, i.e. maximum loading, release profile and shelf life. For these reasons,  
17 the development of improved pharmaceutical formulations requires a thorough knowledge of the  
18 physico-chemical features of these interactions. However, despite the scientific and technological  
19 relevance of this topic, the atomistic details of the implied interactions are rarely investigated. Exper-  
20 imentally, most studies deal with ibuprofen as a model drug incorporated in OMSs.<sup>18,19</sup> Nevertheless,  
21 the observed ibuprofen behavior might not be applicable to other drugs and this may become a seri-  
22 ous problem in the translation from basic research to clinical. Indeed, to our knowledge only few pa-  
23 pers deal with the details of other drug-silica interactions.<sup>20,21</sup> Molecular modeling can be an im-  
24 portant tool in addressing the problem of studying drug-silica interactions.<sup>22</sup> Some of us have already  
25 simulated, through ab-initio calculations, ibuprofen interacting with both silica surfaces and a realis-  
26 tic mesoporous silica model, providing a detailed description of the system and some clues to the in-  
27 terpretation of ambiguous experimental results.<sup>23-25</sup>  
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46 One of the characteristics of mesoporous silica materials is the presence of a high concentration of  
47 silanol groups (SiOH) exposed on the pore walls. It has been shown that the adsorption of molecules  
48 on a silica surface (such as the pore walls) is mainly guided by the formation of H-bonds with si-  
49 lanols, but, especially with hydrophobic molecules, also London type interactions (dispersion) can  
50 play an important role.<sup>23</sup>  
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1 MSU-H is a two-dimensional, hexagonally OMS with large uniform mesopores synthesized through  
2 a nonionic supramolecular assembly pathway using sodium silicate as a silica source and the triblock  
3 copolymer Pluronic P123 (EO<sub>20</sub>PO<sub>70</sub>EO<sub>20</sub>) as a structure-directing agent. The framework structure of  
4 MSU-H is analogous to that of SBA-15 and consists of large ordered pores connected by micropores  
5 in the pore walls. Morphology of MSU-H differs from that of SBA-15, in that SBA-15 mesostruc-  
6 tures usually assemble into larger and more monolithic particles.<sup>26</sup>

7 Clotrimazole (CTZ) is an active principle ingredient poorly soluble in water and present in antifun-  
8 gal topical medications for the treatment of fungal infections (of both humans and animals).<sup>27</sup> It is  
9 commonly available as an over-the-counter drug in various forms including creams, vaginal tablets  
10 and as troche or throat lozenge. Topically, CTZ is used for vulvovaginal candidiasis (yeast infection)  
11 or skin yeast infections. Troche or throat lozenge preparations are used for oropharyngeal candidiasis  
12 (oral thrush) or prophylaxis against oral thrush in neutropenic patients. CTZ is usually used 5 times  
13 daily for 14 days for oral thrush, twice daily for 2 to 8 weeks for skin infections, and once daily for 3  
14 or 7 days for vaginal infections. For these reasons, finding controlled release system can give a good  
15 improvement to pharmaceutical formulations reducing the number of needed applications.<sup>28</sup>

16 The aim of this work is the investigation of clotrimazole (1-[(2-chlorophenyl)(diphenyl)methyl]-1H-  
17 imidazole) incorporated inside MSU-H ordered mesoporous material, by means of a joint computa-  
18 tional and experimental approach.

19 In the present work, CTZ is incorporated into MSU-H with scCO<sub>2</sub> following a previous work<sup>15</sup> and  
20 characterized through Fourier Transform InfraRed Spectroscopy (FTIR), N<sub>2</sub> adsorption isotherms  
21 and ThermoGravimetric analysis (TG). These results are compared and interfaced with static and  
22 dynamic quantum mechanical simulations, focused on the interactions between silanols groups and  
23 CTZ, evaluating also the contribution of dispersive forces being CTZ a highly hydrophobic mole-  
24 cule. Quantum mechanical modeling is used to help the interpretation of the experimental data, and  
25 to obtain a complete description of the MSU-H-CTZ system. Finally, our broader aim is to exploit  
26 the knowledge gained on this system for explaining the amorphization, stabilization and interaction  
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phenomena of drugs on silica surfaces.

## COMPUTATIONAL AND EXPERIMENTAL DETAILS

### Materials

CTZ and Ordered Mesoporous Silica (MSU-H type) were purchased from Sigma-Aldrich. Carbon dioxide with a purity of 99.5% was supplied by SIAD.

### Drug loading

The drug was loaded by the scCO<sub>2</sub> incorporation process at 373 K, 50 MPa for 12 hours. More details on the scCO<sub>2</sub> process are reported in a previous work.<sup>15</sup>

### Computational details

All the calculations were performed within the Density Functional Theory (DFT). Concerning static calculations, the developmental version of the CRYSTAL14 code<sup>29</sup> in its massively parallel version<sup>30</sup> was adopted and the computational approach is the same of Ref.<sup>23</sup>. Briefly, the chosen functional was the Perdew, Burke, and Enzerhof GGA (Generalized Gradient Approximation) exchange-correlation functional (PBE),<sup>31</sup> including the empirical Grimme's D2 correction,<sup>32</sup> to describe the dispersive interactions (vdW). In the following, the superscript D means that Grimme's correction is included. Split valence double- (for Si atoms) and triple- $\zeta$  (for other atoms) Gaussian type basis sets plus polarization functions were used to describe the systems.<sup>33,34</sup> Chlorine atoms were represented with a 86-311G\* basis set.<sup>35</sup> Only the atomic coordinates of the two more superficial layers of each silica slab in the docking geometries were optimized, to compensate for the reduced thickness of the models. Starting geometries were generated so to maximize the interactions between the drug and the surface. Interaction energies, per unit cell per adsorbate molecule ( $\Delta E$ ), were calculated and corrected for the basis set superposition error (BSSE) according to the counter-poise methodology described in previous papers by Delle Piane et al.<sup>23,25</sup> and reported in Supporting Information.

Harmonic frequencies were calculated with CRYSTAL14 at  $\Gamma$  point and the infrared intensity for each normal mode was obtained by computing the dipole moment variation along the normal mode,

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adopting the Berry phase method.<sup>36</sup> For the simulation of the IR spectra of the different structures, only a fragment consisting of the most interesting chemical groups has been considered for constructing the Hessian matrix and will be defined for each case in Results and Discussion.

Enthalpies ( $\Delta H$ ) of adsorption at standard temperature (298 K) were obtained from the vibrational partition functions, by applying the Zero Point Energies ( $\Delta ZPE$ ) and thermal ( $\Delta E_T$ ) corrections to the BSSE corrected electronic adsorption energies ( $\Delta E^C$ ) as  $\Delta H = \Delta E^C + \Delta ZPE + \Delta E_T$ .

Ab-initio molecular dynamics (AIMD) simulations were performed using the CP2K code.<sup>37</sup> The Quickstep technique<sup>38</sup> with a mixed plane wave and Gaussian basis set methodology (Gaussian and Plane Wave method, GPW) was employed to calculate the electronic structure. We used the PBE functional, with the Goedecker–Teter–Hutter pseudopotentials<sup>39</sup> and a triple- $\zeta$  basis set with polarization functions (TZVP)<sup>40</sup> augmented with the empirical Grimme’s D2 correction.<sup>32</sup> The cutoff for the plane wave basis was set to 400 Ry. AIMD simulations were run at 300 K in the NVT ensemble, using the Canonical Sampling through Velocity Rescaling (CSVR) thermostat.<sup>41</sup> A time step of 0.5 fs was chosen. All simulations were equilibrated at 300 K with a more stringent thermostat (time constant: 10 fs) for about 1 ps and then the production phase was run for at least 10 ps with a more relaxed thermostat (time constant: 50 fs). Since CP2K requires 3D periodic systems, a value of  $c = 35$  Å was chosen to separate the slab replicas with enough vacuum. In all cases, only the superficial layer of the silica slab and the drug molecules were allowed to move.

### Characterization

Samples were characterized by means of Fourier Transform Infrared spectroscopy (FTIR), nitrogen adsorption isotherms and thermogravimetry (TG) analysis. For FTIR measurements, powders were pressed in self-supporting wafers and spectra were recorded at room temperature with a Bruker Tensor 27 spectrometer operating at  $2 \text{ cm}^{-1}$  resolution, after outgassing the sample at 373 K for one hour (residual pressure equal to 0,1 Pa). FTIR spectrum of crystalline CTZ was recorded on the powder dispersed in potassium bromide (KBr). FTIR spectrum of CTZ in solution was recorded on a diluted carbon tetrachloride ( $\text{CCl}_4$ ) solution ( $1 \text{ g}\cdot\text{L}^{-1}$ ). Nitrogen adsorption isotherms were measured using a



Quantachrome AUTOSORB-1 instrument after degassing at 373K for 2 hours. Brauner-Emmet-Teller (BET) specific surface areas (SSA) were calculated in the relative pressure range 0,04-0,1 and the pore size distribution was determined through the NLDFT (Non Linear Density Functional Theory) method, using the NLDFT equilibrium model for cylindrical pores.<sup>42</sup> TG analyses were carried out between 298 K and 1073 K in air (flow rate 100 mL·min<sup>-1</sup> with a heating rate of 10 K·min<sup>-1</sup>) using a SETARAM 92 instrument to evaluate the quantity of incorporated drug. The desorption analysis were performed using a SETARAM 92 instrument following the procedure explained by Verevkin et al.<sup>43</sup> and Price et al.<sup>44</sup> Using the Clausius-Clapeyron relation, equation (1), the enthalpies of vaporization ( $\Delta H_{vap}$ ) at the average temperature of investigation were obtained.

$$\ln \left( \sqrt{T} \frac{dm}{dt} \right) = A + \frac{\Delta H_{vap}}{RT} \quad (1)$$

where  $dm/dt$  is the mass loss rate at the specified temperature, kg·s<sup>-1</sup>;  $R$  is the universal gas constant, J·mol<sup>-1</sup>·K<sup>-1</sup> and  $T$  is the temperature of the isothermal experiment, K. Subsequently, vaporization enthalpies were reported to 298.15 K using a general methods of correction, reported by Chickos et al. (the Sidgwick's rule):<sup>45</sup>

$$\Delta H_{vap}(298.15) = \Delta H_{vap}(T) + 0.0545 (T - 298.15) \quad (2)$$

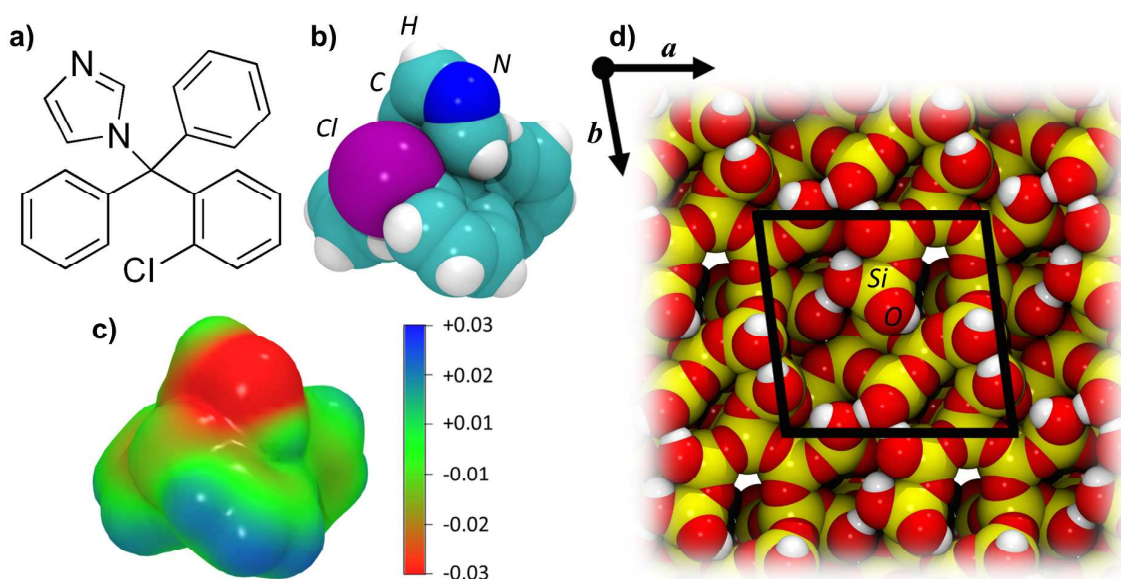
In equation (2),  $T$  is the temperature of measurement or mean temperature of measurement if  $\Delta H_{vap}(T)$  has been obtained from a Clausius-Clapeyron treatment of vapor pressures. The experimental conditions used for the reliable determination of vaporization enthalpies of low volatile molecular compounds are following the references of Verevkin et al.<sup>43</sup> A calibration curve with phenol has been done. The uncertainty of temperature calibration was less than 1 K.

## RESULTS AND DISCUSSION

### Clotrimazole: molecule and crystal

CTZ was modeled both in gas phase and as a crystal, before studying its adsorption on amorphous silica. The starting point was the X-ray experimental structure by Song et al.<sup>46</sup> CTZ has a tetrahedral structure (Figure 1.a), with the central  $sp^3$  carbon linked to two phenyl rings, one chlorophenyl ring

and one imidazole ring. The CTZ molecule, optimized at PBE-D2 level of theory, is reported in Figure 1.b. The electrostatic potential mapped on the PBE electron density (Figure 1.c) clearly shows the nucleophilic character of the external nitrogen in the imidazole ring group, which is expected to behave as H-bond acceptor when interacting with the surface silanols. The rest of the molecule is generally apolar, as reflected by the experimentally measured low solubility in water (less than  $0.01 \text{ g}\cdot\text{L}^{-1}$ ),<sup>27</sup> and is expected to engage in dispersive interactions with the silica surface.



**Figure 1.** **a)** chemical structure of the CTZ molecule ( $\text{H}_{17}\text{N}_2\text{C}_{22}\text{Cl}$ ). **b)** 3D space filling model of the CTZ tetrahedral structure with the three phenyl rings almost perpendicular to the imidazole ring. **c)** CTZ electrostatic potential mapped on the electron density showing the high electronegativity nature of the external nitrogen in the imidazole ring group. Blue, green and red colors correspond to positive, neutral, and negative values of the electrostatic potential (range values: MIN  $-0.03 \text{ au}$ ; MAX  $+0.03 \text{ au}$ ). **d)** top view of the amorphous silica surface model used in this paper, cell borders in black ( $a=12.6 \text{ \AA}$ ,  $b=12.8 \text{ \AA}$  and  $\alpha=83.1^\circ$ ,  $\text{H}_{22}\text{O}_{63}\text{Si}_{26}$ ).

The CTZ crystal has been optimized with and without Grimme's correction and compared to experimental results (Figure S1.a and Table S1, in Supporting Information). The structure is triclinic ( $P\bar{1}$ ,  $a = 8.76 \text{ \AA}$ ,  $b = 10.55 \text{ \AA}$ ,  $c = 10.61 \text{ \AA}$  and  $\alpha = 114.1^\circ$ ,  $\beta = 97.0^\circ$ ,  $\gamma = 97.5^\circ$ ) with two drug molecules per unit cell. If dispersion is not included in the calculation, the cell volume is overestimated by

+25.2 %, while inclusion of dispersive contributions lead to a cell contraction of -8.4% with respect to single crystal X-ray data.<sup>46</sup> The PBE-D2 cohesive energy of crystalline CTZ has been computed as  $-146.5 \text{ kJ}\cdot\text{mol}^{-1}$  (Table 1) and, unsurprisingly, the interactions occurring in the crystal are dominated by dispersion (78%). Indeed, the plain PBE cohesive energy is only  $-31.8 \text{ kJ}\cdot\text{mol}^{-1}$ . The corresponding  $\Delta H^D$  at 298 K, including dispersion (obtained from the full set of harmonic frequencies of the crystal), is  $-131.7 \text{ kJ}\cdot\text{mol}^{-1}$ .

**Table 1. Energetics of the clotrimazole and MSU-H-CTZ systems**

system	$\Delta E^a$	$\Delta E^D^b$	$\Delta H^{298^c}$
clotrimazole bulk <sup>d</sup>	-31.8	-146.5	-131.7
clotrimazole @ (100) <sup>e</sup>	-	-111.3	-104.7
clotrimazole @ (010)	-	-124.5	-117.1
clotrimazole @ (001)	-	-118.9	-111.9
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imidazole (I) <sup>f</sup>	-61.5	-134.1	-126.8
imidazole (II)	-58.4	-102.6	-99.1
phenyls (I)	-19.7	-87.6	-80.0
phenyls (II)	-12.3	-78.4	-73.6
molecular layer (I)	-	-117.1	-101.7
molecular layer (II)	-	-116.8	-99.5
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<i>CTZ bulk EXP. (<math>-\Delta H_V^{298}</math>)<sup>g</sup></i>			<i>-91.6</i>
<i>CTZ incorporated EXP. (<math>-\Delta H_V^{298}</math>)<sup>h</sup></i>			<i>-91.8</i>

<sup>a</sup> Computed (only for selected systems) electronic interaction energy, without accounting for dispersion, corrected for BSSE. <sup>b</sup> As <sup>a</sup> including dispersion. <sup>c</sup> Computed enthalpy of adsorption (dispersion included), at T = 298K. <sup>d</sup> Cohesive energy of crystalline CTZ, computed with respect to a free molecule in gas phase. <sup>e</sup> Interaction of one CTZ molecule adsorbed on the (100), (010) and (001) faces of the CTZ crystal. <sup>f</sup> CTZ/silica systems: for naming refer to Figure 2. <sup>g</sup> Experimentally measured vaporization enthalpy of CTZ bulk, extrapolated at T = 298K. <sup>h</sup> Experimentally measured vaporization enthalpy of CTZ incorporated in MSU-H, extrapolated at T = 298K. All values are in  $\text{kJ}\cdot\text{mol}^{-1}$ .

Experimentally, the lack of crystallization of CTZ inside mesoporous silica is observed. Furthermore, mesoporous materials are known to cause amorphization of crystalline drugs. This phenome-

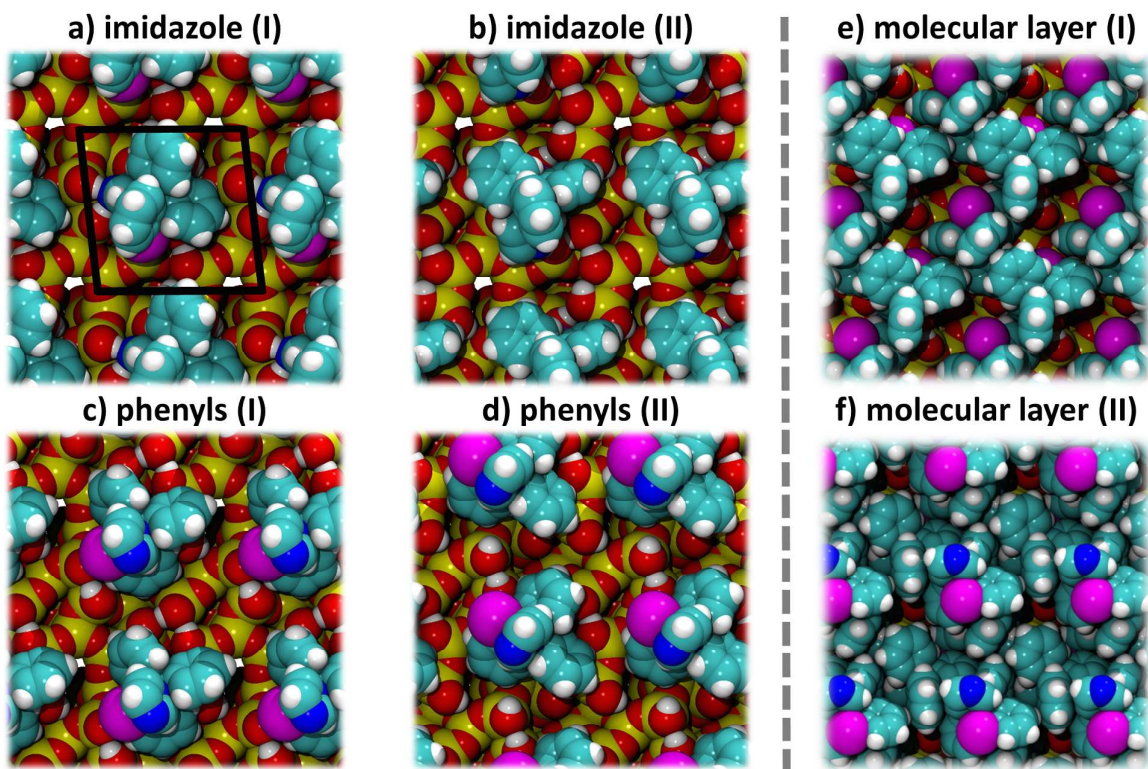
1 non has been described in literature both as a confinement effect<sup>47-50</sup> and as a competition between  
2 crystal cohesion and adsorption on the silica surface.<sup>23</sup> To describe this competition, we have mod-  
3 elled the interaction of CTZ molecules with three crystalline surfaces of the same CTZ crystal. The  
4 three chosen surfaces were the (100), (010) and (001), which have been modeled through a periodic  
5 slab approach with a slab thickness equal to two unit cells. These surfaces represent the three main  
6 defect-free crystal faces and were optimized at the PBE-D2 level. The surface energies were 0.10,  
7 0.11 and 0.09 J m<sup>-2</sup>, respectively. The adsorption of CTZ on the three crystal surfaces has been simu-  
8 lated by adding one drug molecule in interaction with the slab in the same position occupied in the  
9 bulk structure and re-optimizing the structure (Figure S1.b in Supporting Information). The added  
10 CTZ molecule maximizes the contact with the surface molecules mainly through dispersion and  
11 electrostatics interactions. These are in particular  $\pi$ - $\pi$  edge to face interactions of phenyls and chlo-  
12 rophenyls. Other  $\pi$ - $\pi$  interactions (e.g. face to face) are not present due to steric hindrance and rigidi-  
13 ty of the molecule. Table 1 reports the computed interaction energies of CTZ on the three crystal sur-  
14 faces together with the corresponding enthalpies at 298 K. Unsurprisingly, the interaction energies of  
15 CTZ with its crystal surfaces are comparable to each other with an average value of -113 kJ mol<sup>-1</sup>,  
16 lower than the crystal cohesive energy (-132 kJ mol<sup>-1</sup>) due to the reduced number of intermolecular  
17 interactions of the adsorbed drug with the surface compared to those in the bulk.

18 TG desorption analysis of crystalline CTZ (Table 1) shows an experimental enthalpy of vaporization  
19 of about 92 kJ mol<sup>-1</sup>, reasonably close to the theoretical values. The comparison between computed  
20 and experimental cohesive energy for the CTZ crystal shows some overestimation due to the PBE-  
21 D2 method. We proved in the past that the same applies to B3LYP-D2 in the case of molecular crys-  
22 tals and we proposed a slightly different formulation known as B3LYP-D2\* which brings computed  
23 cohesive energy in better agreement with the experiment. Unfortunately this correction is not availa-  
24 ble for PBE-D2 so that we should be aware of a somehow systematic overestimation of the comput-  
25 ed data also for the interaction with the surface.<sup>51</sup>

### Clotrimazole adsorption on the silica pore wall

Experimentally, CTZ was loaded into MSU-H through  $\text{scCO}_2$ , achieving a maximum drug loading of 34% by mass (vide infra).

In the simulations, CTZ was adsorbed on a silica surface model described in a previous study by some of us<sup>52</sup> and already employed to simulate the adsorption of ibuprofen and aspirin.<sup>23</sup> The use of a flat surface in the present work, at variance with an explicit model of MSU-H, is justified by the curvature of the MSU-H pore, whose diameter (8.5 nm) is much larger than the CTZ molecule: indeed, a flat surface model represents a very good approximation of what the drug is “seeing” inside the pores. This surface is represented in Figure 1.d and exhibits a silanol density of  $4.5 \text{ OH}\cdot\text{nm}^{-2}$ , close to the experimentally measured value for fully hydroxylated surfaces ( $4.9 \text{ OH}\cdot\text{nm}^{-2}$ ).<sup>53</sup> The silica surface model contains 111 atoms in the unit cell ( $a=12.6 \text{ \AA}$ ,  $b=12.8 \text{ \AA}$  and  $\alpha=83.1^\circ$ , cell composition:  $\text{H}_{22}\text{O}_{63}\text{Si}_{26}$ ). The surface exposes eight silanols per the unit cell. Of these, only one is free, while the others are interacting through H-bonds. Particularly, three SiOHs cooperates in forming a stable H-bonded chain.



**Figure 2.** View along the z axis of the fully optimized six different geometries of adsorption. **a)** imidazole (I): Imidazole ring pinched by two silanols; this structure started as chlorine interacting with a silanol, and evolved in a shifted position during geometry optimization, due to the high H-bond acceptor character affinity of imidazole's nitrogen. **b)** imidazole (II): imidazole ring interacting with one ending chain silanol. **c)** phenyls (I): phenyl rings relating with silanols and one SiOH-Cl weak H-Bond. **d)** phenyls (II): phenyl rings interacting with silanols. **e)** molecular layer (I): two molecules per silica unit cell. **f)** molecular layer (II): three molecule per unit cell.

CTZ was manually docked on the surface, aiming at maximizing the interactions between exposed silanols and the different drug's functional groups. In order to match the experimental conditions, six main starting geometries have been studied. Four of them are characterized by one molecule per silica unit cell (drug loading ~13% by mass), while the other two simulate a molecular layer as observed in previous results<sup>15</sup> with two and three molecules per unit cell (drug loading ~27% and ~41% by mass, respectively). All models are shown in top view in Figure 2: a) imidazole (I) and b) imidazole (II) are structures with CTZ interacting through its imidazole ring, while in c) phenyls (I) and d) phenyls (II) the molecule is adsorbed through its apolar portion; e) molecular layer (I) and f) molecular layer (II) are the highest loading structures. The 13% loading models (a-d) have been optimized both with and without Grimme's dispersion correction to measure the role of dispersion, while for the high loadings only PBE-D2 has been used.

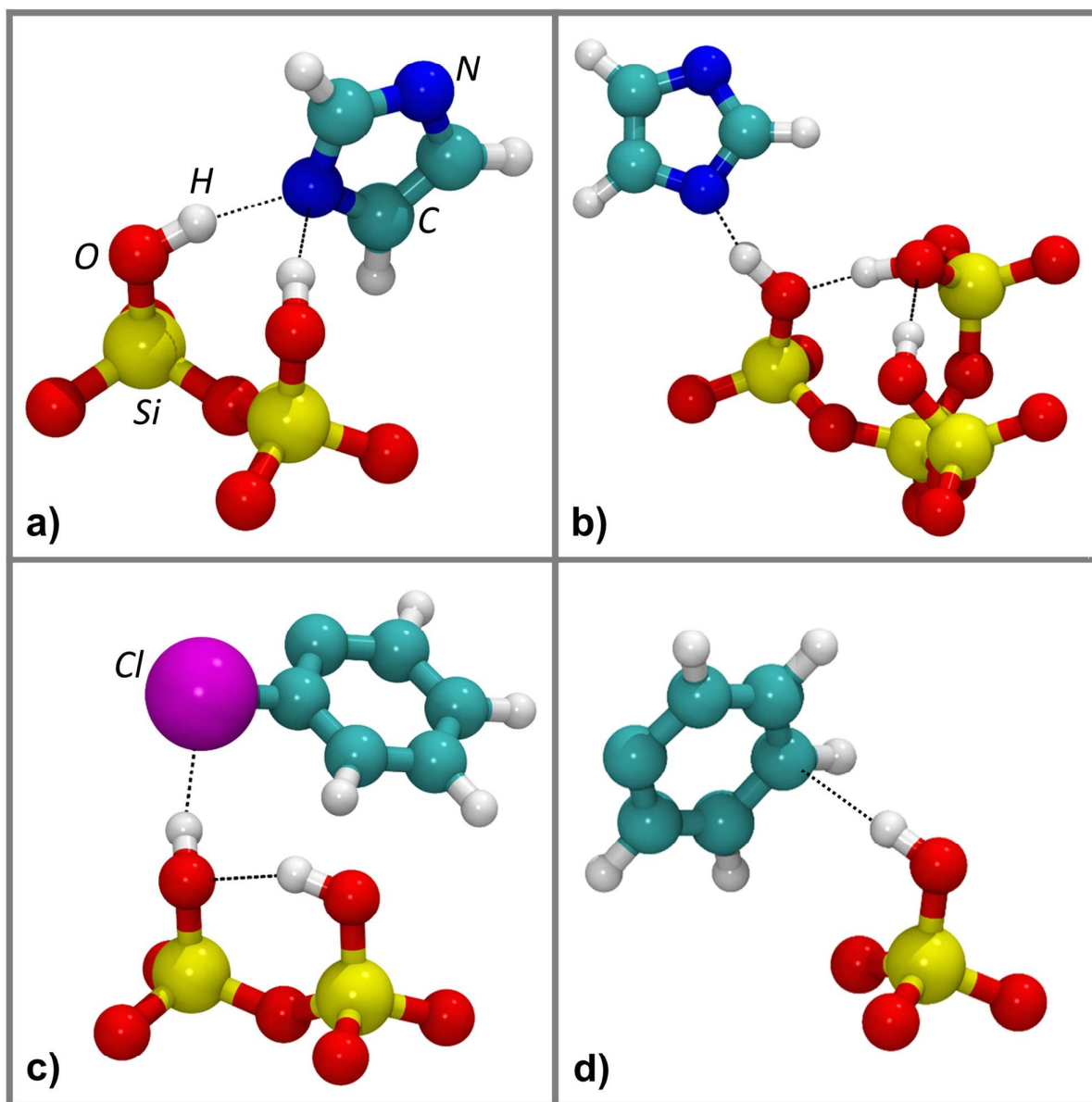
### Interactions between clotrimazole and the silica pore wall

Figure 3 reports the main types of interaction between CTZ and silanols as resulted from the simulations. Both the imidazole ring and the chlorophenyl group of CTZ can be involved in H-bonds with the surface hydroxyls (Figure 3.a-c) but silanol- $\pi$  bonds interactions are also established (Figure 3.d).

The imidazole ring can form both two (Figure 3.a) and one (Figure 3.b) H-bonds with the surface. In the first case (observed only in the imidazole (I) structure) the imidazole's nitrogen is pinched between two silanols and the two H-bonds have the same length (1.84 Å, for the imidazole (I) case),

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suggesting an equivalency of the two interactions. In most cases, however (imidazole (II) and in the molecular layer geometries), CTZ forms only one H-bond through its imidazole ring, usually with the terminal silanol of an H-bonded chain, since these terminal groups are known to be good H-bond donors due to their increased acidity.<sup>22,23,25</sup> Notably, the H-bond length in the imidazole (II) case (1.51 Å) is shorter (and the interaction stronger) than in the imidazole (I) situation (1.84 Å).



**Figure 3.** Types of interactions: **a)** Imidazole-CTZ double H-bonds (extracted from the imidazole (I) structure). **b)** Imidazole-CTZ single H-bond (extracted from the imidazole (II) structure). **c)** Chlorophenyl-CTZ interaction (extracted from the phenyls (I) structure). **d)** SiOH- $\pi$  interaction (extracted from the phenyls (II)

structure).

In agreement with previous results for other more polar drugs<sup>23</sup>, the pattern of mutual H-bonds between surface silanols is restructured in response to the approach of molecules rich of H-bond donors/acceptors functional groups. This feature is present in all the explored dockings. For instance, in the imidazole (I) case (represented in the local view of Figure 3.a) originally the two silanols were interacting with each other. Since an H-bond interaction between vicinal silanols is known to be quite weak,<sup>22</sup> this mutual interaction is easily broken by the incoming CTZ molecule. On the other hand, in cases that include an interaction as reported in Figure 3.b, the effect is more subtle: the H-bond with CTZ is able to enhance the interactions towards the imidazole ring along the whole silanol chain.

H-bonds of the Si-O-H---Cl type (Figure 3.c) are weaker than those with imidazole (2.24 Å in the phenyls (I) case). Interestingly, in phenyls (I) the distance between the two SiOHs is not perturbed by the new interaction.

Several SiOH- $\pi$  (surface-CTZ, Figure 3.d) have been observed in almost all the different structures, while  $\pi$ - $\pi$  edge to face lateral interactions (CTZ-CTZ) characterize the molecular layer models. The mean value of SiOH- $\pi$  bonds is 2.10 Å and accordingly to Tsuzuki et al.<sup>54</sup> the  $\pi$ - $\pi$  bonds have a mean value of 4.60 Å. Comparing bond lengths in the different models, we suggest that SiOH- $\pi$  interactions are half way between SiOH-N and SiOH-Cl H-bonds, as regards interaction strength.

### **Energetics of the clotrimazole-silica interaction**

Table 1 reports the computed interaction energies with ( $\Delta E^D$ ) and without ( $\Delta E$ ) dispersion (when available) and the adsorption enthalpies at standard temperature ( $\Delta H^{298}$ ), including dispersion.

All simulations reveal that the adsorption of CTZ on amorphous silica is a strongly exothermic process. Nevertheless, different geometries and consequently interaction types correspond to different energies of adsorption, ranging from -134.1 kJ·mol<sup>-1</sup> for the imidazole (I) case to -78.4 kJ·mol<sup>-1</sup> for



the phenyls (II) geometry.

Comparing  $\Delta E$  and  $\Delta E^D$ , London forces constitute the predominant contribution to the interaction in all systems, except for the imidazole (I) case. Of particular interest are the phenyls (I) and (II) structures: in these cases the interaction energies without dispersion are small, while including vdW interactions increases these energies by more than +400 and +600%, respectively. Dispersion plays a more important role in CTZ adsorption on hydrophilic silica than it has been reported for more polar drugs such as ibuprofen and aspirin.<sup>23,25</sup> For the molecular layer geometries, the average interaction energies (and enthalpies) per molecule have very similar values despite the fact that the third additional molecule in the molecular layer (II) is not involved in H-bonds with the surface. This is a further evidence that, for an apolar molecule, the lateral self-interaction (dispersion-driven) is similar to that of the molecule with the surface. In all cases the CTZ deformation energies are negligible due to the rigidity of the molecule with a mean value of  $7 \text{ kJ}\cdot\text{mol}^{-1}$ .

Table 1 reveals that all the computed interaction energies and enthalpies are close to each other, enlightening a possible competition between crystalline and adsorbed CTZ. Assuming that the supercritical incorporation is a step route controlled by dissolution and adsorption of single molecules, this process is almost isoenergetic, since for each adsorption on silica an average of  $110 \text{ kJ}\cdot\text{mol}^{-1}$  are gained and for each desorption from the crystal  $115 \text{ kJ}\cdot\text{mol}^{-1}$  are lost. This could be the explanation of the lack of crystallization in pores smaller than 20 times the molecule diameter explained by Sliwiska-Bartkowiak et al.<sup>47</sup> and observed by other authors.<sup>48-50</sup>

TG desorption analysis of CTZ in MSU-H (34% by mass) produces an enthalpy of vaporization of  $91.6 \text{ kJ}\cdot\text{mol}^{-1}$  as reported in Table 1. As abovementioned (and also reported in Table 1), a similar analysis on crystalline CTZ results in a value of  $91.8 \text{ kJ}\cdot\text{mol}^{-1}$ . Thus, also the experimental data suggest that the two processes are almost isoenergetic, supporting the hypothesis on silica-induced drug amorphization.

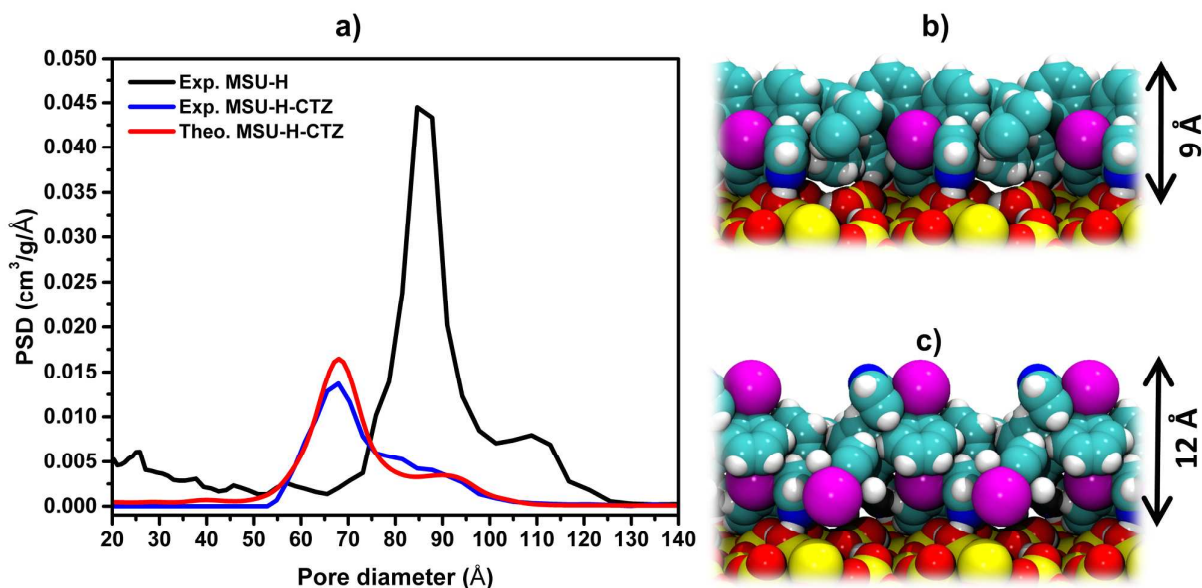
The experimental enthalpy of vaporization for the MSU-H-CTZ case is very close (in absolute value) to the computed adsorption enthalpies for both CTZ molecular layers ( $-101 \text{ kJ}\cdot\text{mol}^{-1}$ ). The com-

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parison of the experimental value ( $91.6 \text{ kJ}\cdot\text{mol}^{-1}$ ) with the data obtained for the different low loading geometries suggests which interaction reported in Figure 3 is more represented in the real sample. The closest value is obtained for the imidazole (II) case, i.e. adsorption driven by one H-bond between the imidazole group and a surface silanol. On the other hand, situations like imidazole (I), although theoretically very stable, seem to be rarer events in the real material due to the lower hydrophilicity of the experimental silica, the steric hindrance of the molecules and the low configurational entropy content of this unique configuration. Interestingly, the computed interaction enthalpy of the phenyls (I) model is also quite close to the experimental value, suggesting that phenyls-driven interactions may be present in a significant portion of the adsorbed CTZ population. In summary, all configurations of Figure 3.a-d are probably represented in the real sample, since their average  $\Delta H^D$  ( $-94.9 \text{ kJ}\cdot\text{mol}^{-1}$ ) is very close to the experimental measure.

#### **Model of adsorption: nitrogen adsorption and TG analysis**

Experimental  $\text{N}_2$  adsorption isotherms (Figure S2.a in Supporting Information), through Non Linear Density Functional Theory (NLDFT), allow the calculation of the experimental pore size distributions (PSDs) before and after the  $\text{scCO}_2$  incorporation of CTZ. The results are reported in Figure 4.a. As expected, drug incorporation reduces the pore diameter from the  $86 \text{ \AA}$  value of bare MSU-H (black line) to a lower value of  $67 \text{ \AA}$  (blue line).



**Figure 4.** Adsorption model. **a)** Comparison of Pore Size Distributions (PSDs) as calculated from the two molecular layer models (red) and as measured before (black) and after (blue) the  $\text{scCO}_2$  incorporation. **b)** and **c)** side views along the  $a$  direction of molecular layers (I) and (II), respectively.

As already described in a previous work by some of us,<sup>15</sup> the drug can be assumed as adsorbed mainly in the form of a molecular layer. Using a bare value of thicknesses of the single molecular layers I and II (Figure 4.b-c), the corresponding new PSDs can be calculated (Figure S2.b). Clearly, with this approach, the calculated PSDs are unable to describe the experimental results.

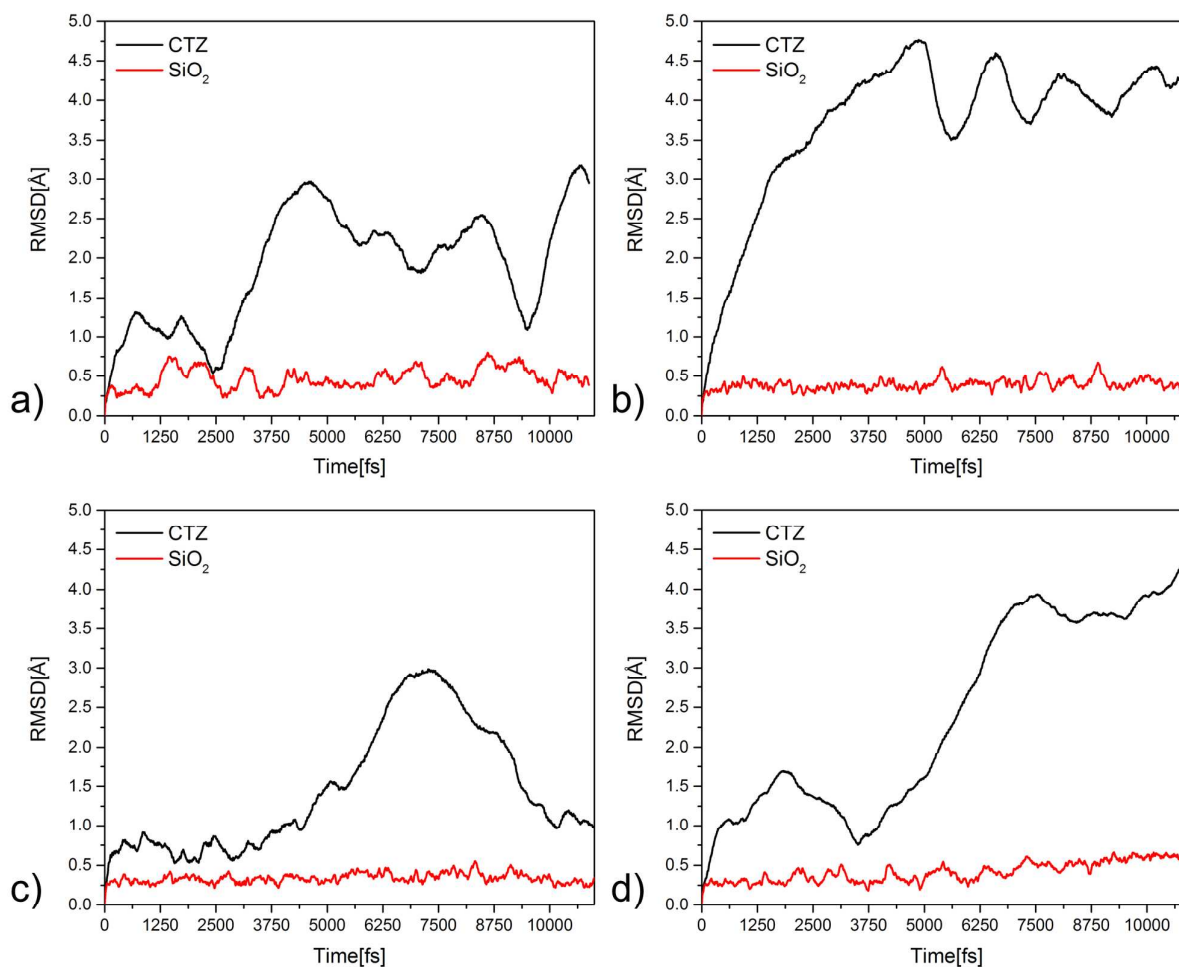
On the other hand, experimental TG analysis reports an incorporated quantity of 34% by mass, which is in between the loading calculated for layer (I) and layer (II), which are  $\sim 27\%$  and  $\sim 41\%$  by mass, respectively. By calculating the experimental planar concentration of CTZ (molecules per  $\text{nm}^2$ ) and comparing it to our theoretical models of molecular layers (2 CTZs/cell, Figure 2.e and 3 CTZs/cell, Figure 2.f), it is shown that the real system can be described by a 50:50 mixture of the two molecular layer geometries. Indeed, the unit cell of the simulated silica surface has an area of  $1.6 \text{ nm}^2$  and the experiment reports 2.5 CTZ molecules per  $1.6 \text{ nm}^2$ . As a consequence, a simulated MSU-H-CTZ PSD has to be calculated, starting from the experimental PSD of bare MSU-H, assuming thicknesses representative of both the molecular layer (I) and (II) models. These thicknesses have been evaluated following the Connolly surfaces<sup>55</sup> before and after CTZ adsorption of the computed

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models, with the purpose to take into account the vdW molecular volume and the roughness generated by the statistical distribution of 2 and 3 CTZ molecules per  $1.6 \text{ nm}^2$ . Therefore, the molecular layer surfaces have been discretized in 677 squares ( $0.25 \text{ \AA}^2$ ) in order to evaluate the thicknesses, for each point, between the starting silica model and the molecular layers. Subsequently, a new PSD has been calculated for each couple of evaluated thicknesses and all curves have been combined together in the final theoretical PSD of Figure 4.a (red). Such procedure results in an impressive agreement between simulation and experiment, validating the data interpretation.

### Mobility of adsorbed clotrimazole

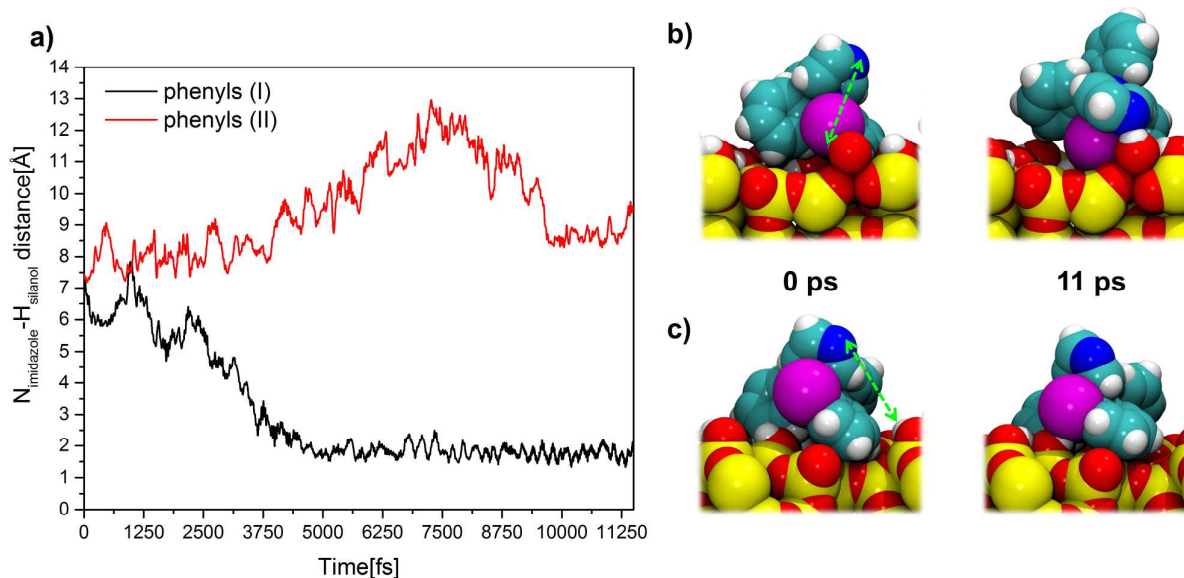
Ab-Initio Molecular Dynamics (AIMD) has been performed on the different statically optimized geometries in order to evaluate the stability of the local minimum structures of Figure 2. In Figure 5 the Root Mean Square Deviations (RMSD) of the atomic positions, during the production, with respect to the first frame, are reported for the four AIMD simulations (imidazole (I), phenyls (I), phenyls (II) and molecular layer (II)). RMSDs are separated in the CTZ and silica contributions. In all cases, the mobility of the adsorbed drugs is higher than that of the silica surface. The latter seems equilibrated during all simulations.



**Figure 5.** Silica ( $\text{SiO}_2$ ) and clotrimazole (CTZ) RMSD (in Å) with respect to the initial structure for all AIMD simulations. For the silica surface, only the atoms free to move in the simulation have been considered. **a)** imidazole (I); **b)** phenyls (I); **c)** phenyls (II); **d)** molecular layer (II). See Figure 3 for the corresponding structural details.

The imidazole (I) structure, which is the most stable configuration of CTZ on silica according to Table 1, shows a general stability of the adsorbed molecule during the 11 ps simulation. Although Figure 5.a reports a significant CTZ mobility, the molecule fluctuates around the initial adsorption geometry, without evolving into a new configuration. As reported in Figure S3 in Supporting Information, while equivalent H-bond interactions were present in the static model (Figure 3.a), during the AIMD one H-bond remains stable while the second one shows instability: in many points during the time evolution the interaction is lost and the silanol is free to rotate as shown by small bumps in the  $\text{SiO}_2$  RMSD. However, this interaction is always recovered after a certain time proving its rele-

vance at room temperature, although the two H-bonds are not really equivalent, as resulting from static simulations.



**Figure 6.** AIMD of phenyls (I) and (II): **a)** exposed nitrogen – SiOH distance over time: phenyls (I) in red, phenyls (II) black. **b)** 3D space filling models of phenyls (I) at 0 and 11 ps. **c)** 3D space filling model of phenyls (II) at 0 and 11 ps. The nitrogen – SiOH distance plotted in the graph is highlighted in green.

On the contrary, the phenyls (I) structure proved to be unstable at room temperature. The RMSD computed along this AIMD simulation (Figure 5.b) shows the very high mobility of the drug (the highest among all considered cases) on an otherwise equilibrated surface. Particularly, CTZ departs from the initial configuration to reach a new equilibrium about 4 ps later. Figure 6.a traces this movement by plotting the distance between the exposed imidazole's nitrogen and a surface SiOH against time (black line), while Figure 6.b shows the starting and ending structures: the initially present chlorine-silanol interaction is immediately lost and the molecule rotates in order to establish a new SiOH-imidazole H-bond. This new interaction (average length: 1.80 Å) is formed within the first 4 ps and kept until the end of the simulation. This complex process, involving a movement of the whole molecule, appears almost barrierless, being observable in a very short time scale, suggesting that the phenyls (I) configuration transforms rapidly in a more stable imidazole (II) geometry, if

1 the surface morphology allows for it. The new geometry, Figure 6.b (right), was re-optimized and the  
2 corresponding interaction energy ( $\Delta E^D$ ) was computed as  $-106 \text{ kJ}\cdot\text{mol}^{-1}$  with respect to the initial  
3 value of  $-88 \text{ kJ}\cdot\text{mol}^{-1}$ , meaning that the system was able to evolve autonomously to a more stable  
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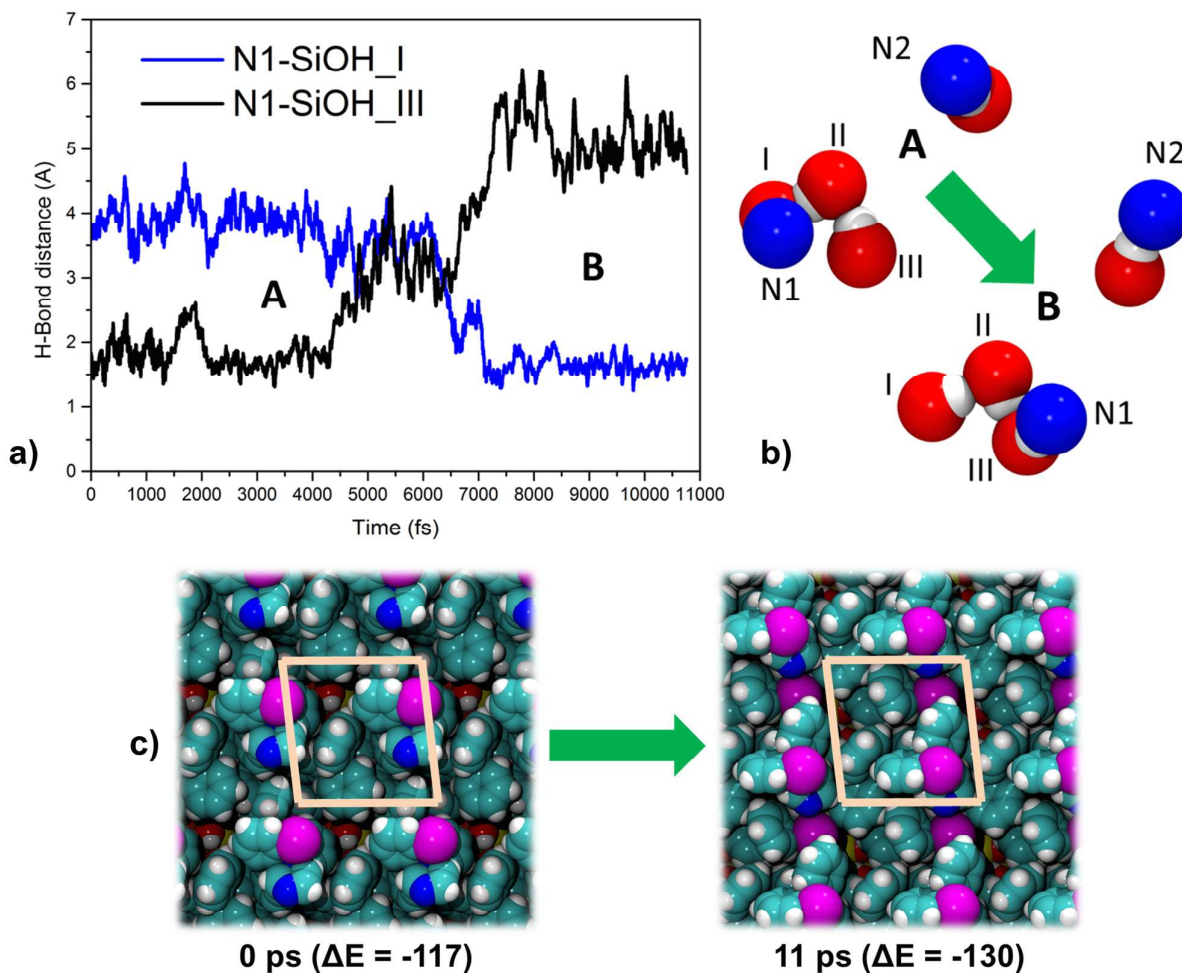
7  
8 To see whether such instability is a common feature of phenyls-driven adsorptions, an AIMD simu-  
9 lation on the phenyls (II) structure has been performed. Along the 11 ps simulation, this model was  
10 unable to rotate and the CTZ's RMSD (Figure 5.c) reveals a much higher stability than the phenyls  
11 (I) case. As can be also noticed from the starting and ending point of the simulation (Figure 6.c), the  
12 molecule simply fluctuates around the initial position. The plot in time of the same imidazole-SiOH  
13 distance adopted for the phenyls (I) simulation (Figure 6.a, red line) shows that no new H-bond in-  
14 teraction is formed and the imidazole ring remains unbound. These results suggest that dispersive  
15 forces and SiOH- $\pi$  interactions are able alone to preserve the phenyls (II) geometry at room tempera-  
16 ture. The AIMD simulations on the phenyls structures show that, at least in part of the cases, if an in-  
17 coming CTZ molecule adsorbs on silica through its phenyl rings it may keep this configuration  
18 without forming imidazole interactions, although such structure is thermodynamically less stable.  
19 This further supports the explanation of TG desorption data (vide supra).  
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### 38 **Molecular layer adsorption**

39 AIMD simulation on the molecular layer (II) structure shows high mobility of the three CTZs (per  
40 unit cell) on the surface. The RMSD of the atomic positions along the AIMD simulation (Figure 5.d)  
41 shows a large movement between 4 and 7 ps that results in a new configuration, with a RMSD value  
42 of  $4.3 \text{ \AA}$  with respect to the starting CTZ conformation after 11 ps. The process is described in Fig-  
43 ure 7, where the exposed nitrogen atoms of the imidazole rings of the three adsorbed molecules in  
44 each unit cell are referred as N1, N2 and N3. In the starting configuration (Figure 7.b, case A, and  
45 Figure 7.c, left), one CTZ molecule is H-bonded through its imidazole's nitrogen (N1) to the termi-  
46 nal silanol of a chain made up of three silanols, here referred as SiOH\_I, SiOH\_II, SiOH\_III. The 11  
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1 ps of AIMD result in a global shift of the adsorbed molecular layer, as reported in Figure 7.c, taking  
2 as reference the unit cell borders. The driving force for this phenomenon is the loss of the N1-  
3 SiOH\_III H-bond, Figure 7.b (case B), with the concurrent N1-SiOH\_I H-bond formation, through a  
4 flip of the whole silanol chain. The graph in Figure 7.a, reporting the N1-SiOH\_III and N1-SiOH\_I  
5 distances, clearly shows a transition state where the chain is lost and the imidazole's nitrogen (N1) is  
6 equidistant from SiOH(I) and (III). Looking at the potential energy fluctuations during this simula-  
7 tion (Figure S4 in Supporting Information), we estimate, in a very approximate way, the electronic  
8 activation energy of this process as  $29.7 \text{ kJ}\cdot\text{mol}^{-1}$  that can be considered an upper limit for the real  
9 value. The new computed interaction energy  $\Delta E^D$ , after optimizing the final point of the AIMD simu-  
10 lation, is  $-130 \text{ kJ}\cdot\text{mol}^{-1}$  and the contribution for the silica deformation, due to the flip of the silanols  
11 chain, is  $+44 \text{ kJ}\cdot\text{mol}^{-1}$ .





**Figure 7.** AIMD of the molecular layer (II) structure: **a)** N1-SiOH(I) and N1-SiOH(III) bond distances in time. **b)** starting (A) and final (B) positions of N1 and N2 with respect to the involved SiOHs. **c)** Top views of the initial (left) and final (right) configurations in the AIMD simulation, with the corresponding interaction energies per CTZ molecule ( $\text{kJ}\cdot\text{mol}^{-1}$ ); cell borders in pink.

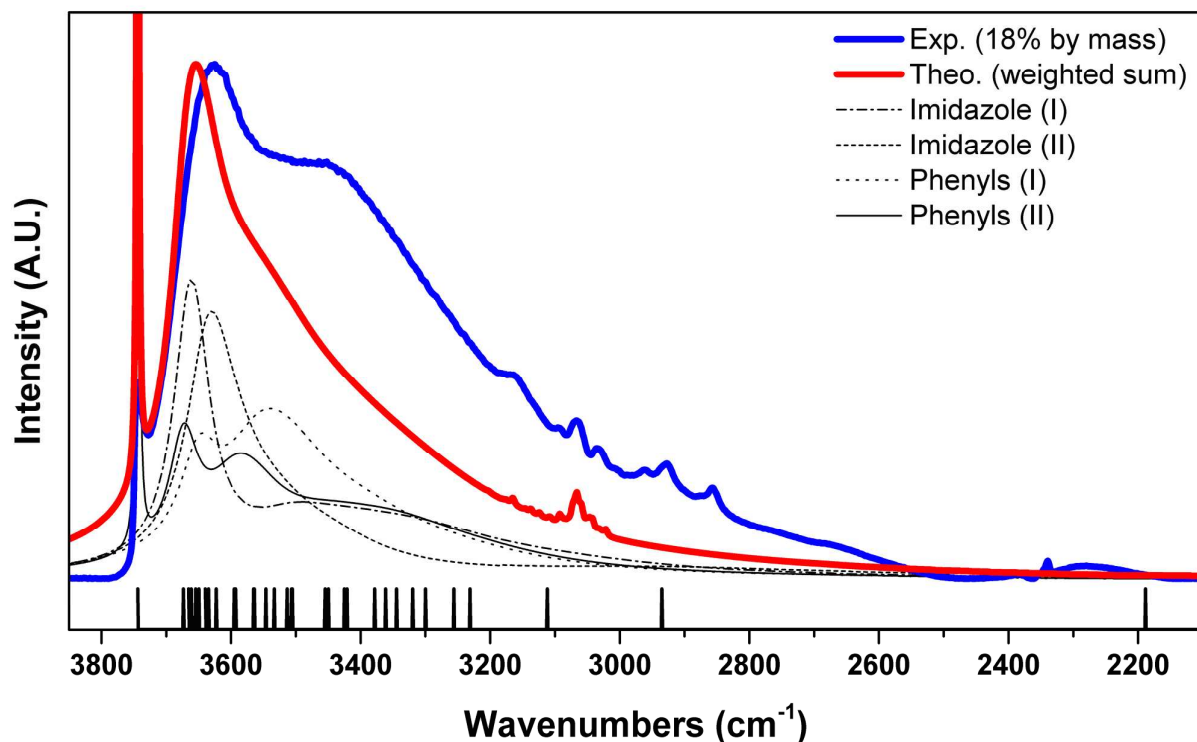
These AIMD results are in accordance with the experimental finding of a “liquid like” behavior of ibuprofen adsorbed in MCM-41.<sup>18,19</sup> Considering that CTZ is much more hydrophobic than ibuprofen it is no surprise that CTZ molecular layers are very mobile, despite the significant underneath interactions. This mobility can be represented as a walking step process guided by local changes in the H-bond interactions, helped by a high flexibility of the silica surface silanols.

### Experimental and theoretical FT-IR interpretation

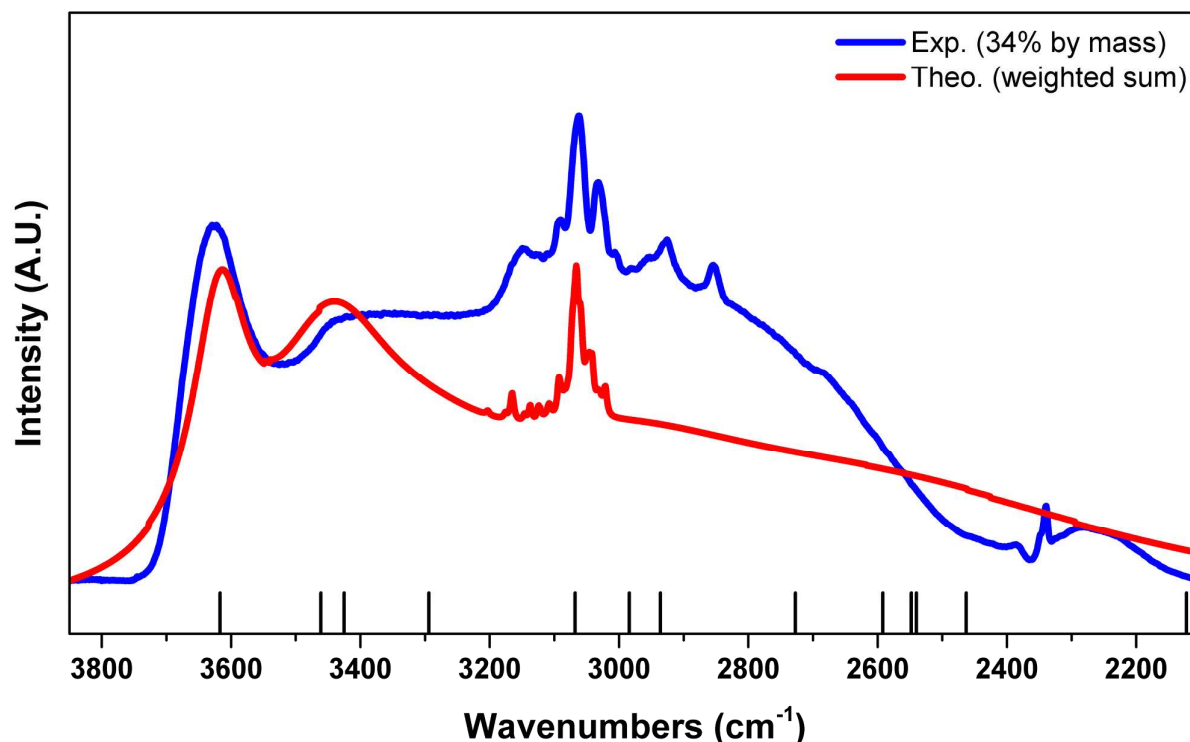
Experimental and theoretical vibrational spectra of CTZ in molecular, crystalline and adsorbed envi-

ronments have been obtained (Figure S5 in Supporting Information). Table S2 in Supporting Information reports a complete assignment of CTZ vibrational frequencies, characterized by three main areas of interest: aromatic (C-H) stretching region ( $3000 - 3250 \text{ cm}^{-1}$ ), aromatic (C=C, C=N) stretching bands ( $1300 - 1600 \text{ cm}^{-1}$ ) and aromatic (C-H) bending peaks ( $800 - 600 \text{ cm}^{-1}$ ).

Figures 8 and 9 report both the experimental (blue line) and the theoretical (red line) IR spectra of CTZ in interaction with silica at low and high coverage. The low coverage experimental spectrum has been acquired on a MSU-H-CTZ sample with a drug loading of 18% by mass, while the high coverage corresponds to a drug loading of 34% by mass. As regards the simulated spectra, the CTZ modes have been obtained in both cases from a vibrational analysis in the harmonic approximation of the drug in the four single molecule configurations (Figure 2.a-d), combined by Boltzmann weighting their contribution to the infrared intensity according to the computed interaction energies. These modes were corrected for anharmonicity, using a 0.960 scaling factor. The theoretical SiO-H stretching mode contributions for the low coverage spectrum (Figure 8) were obtained from the four single molecule geometries (Figure 2.a-d) plus the new configuration of phenyls (I) after AIMD. For the high coverage spectrum (Figure 9), they were obtained from the two molecular layer structures (Figure 2.e-f) plus the new configuration of molecular layer (II) after AIMD. In all cases, the SiO-H stretching modes have been corrected with the Pimentel's empirical rule to account for the broadening effect due to H-bonding.<sup>56</sup> The uncorrected computed SiO-H frequencies have been reported in the lower part of both Figure 8 and 9 as vertical lines. The final theoretical spectra (red lines) of Figure 8 and 9 are then the combination of these SiO-H stretching modes and the CTZ signals. Figure 8 also includes the theoretical spectra computed for the four low loading configurations, limited to the SiO-H stretching modes.



**Figure 8.** Experimental and simulated IR: experimental FTIR spectrum of MSU-H-CTZ at 18% by mass (blue); theoretical IR spectrum of CTZ combined with SiOH vibrational contribution (red); simulated IR spectrum of SiOH vibrational contribution of: imidazole (I) (dot-dashed black line); imidazole (II) (dashed black line); phenyls (I) (dotted black line); phenyls (II) (black line). At the bottom of the graph, sticks represent the theoretical SiO-H stretching frequencies without Pimentel's correction.



**Figure 9.** Experimental and simulated IR: experimental FTIR spectrum of MSU-H-CTZ at 34% by mass, after degas (blue); simulated IR spectrum of CTZ combined with SiOH vibrational contribution (red). At the bottom of the graph, sticks represent the theoretical SiO-H stretching frequencies without Pimentel's correction.

For both the low and high loading cases, the agreement between theory and experiment is remarkable and helps the interpretation of the signals. Considering CTZ modes, the 3175-3100 cm<sup>-1</sup> experimental bands are due to the ensemble of the C-H symmetric stretching of SiOH- $\pi$  interacting aromatic rings. Peaks between 3000-2800 cm<sup>-1</sup> are due to alkyl impurities in the silica samples. As regards SiO-H signals, the sharp peak at 3750 cm<sup>-1</sup> (only in the low loading case) corresponds to isolated silanols. The broad band between 3700 and 3550 cm<sup>-1</sup> corresponds to the silanol signals of SiOH- $\pi$  interactions. All silanols signals at lower wavenumbers (3400-2500 cm<sup>-1</sup>) are due to SiOH-imidazole interactions. As expected, the SiO-H stretching frequency is highly affected by this interaction, with bathochromic shifts up to more than a thousand wavenumbers, as shown by the individual signals reported in both figures (stick lines) and particularly for the high loading case. The higher intensity of this broad band in both experimental spectra of Figure 8 and 9 with respect to the simu-

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lation is due to the Pimental's approximation which cannot account for all the subtle features due to anharmonic coupling and intensity stealing between modes in H-bond interactions.

A first effect of increasing the drug loading is the loss of the isolated silanol peak, suggesting that all silanols become involved in interactions (see spectra of Figure 8 and 9). A second result is an intensification of the SiOH-imidazole signals with respect to the SiOH- $\pi$  ones, suggesting that a decreasing amount of silanols is able to interact with the phenyl groups and an increasing amount of SiOH-imidazole H-bonds is formed. Such behavior seems counterintuitive with respect to the energetics of Table 1: computed  $\Delta H^D$ s suggest that by increasing drug loading more stable configurations (imidazole-driven) should be populated earlier than less stable ones (phenyls-driven), while FT-IR shows that SiOH-phenyls signals are more prominent in the low loading case than the high loading one. This apparent disagreement can be resolved by taking into account the complexity of CTZ adsorption on the silica pore walls. Comparison between computed and experimental enthalpies of adsorption (Table 1) suggests that at least some energetically very stable imidazole-driven interactions (such as that of Figure 3.a) are quite rare in a real sample as they depend on the presence of specific surface sites. Moreover, given the structure of the CTZ molecule, also when it adsorbs through its imidazole ring, more silanols are influenced by its phenyls and chlorophenyls moieties than by the SiOH-N interaction. This can be deduced from the individual simulated spectra of Figure 8, in which signals corresponding to SiOH- $\pi$  are present in all four configurations, independently on what drives the interaction. Furthermore, the strong bathochromic shift caused by the SiOH-N H-bonds corresponds to a significant signal broadening, so that at low loading these modes are smeared out over the whole spectrum. Finally, AIMD simulations proved that some phenyls-driven adsorptions, although based on relatively weak interactions, are indeed stable at room temperature: given the bulky conformation of CTZ (Figure 1.c) such configurations might be highly probable during incorporation, also at low loadings.

## CONCLUSIONS

1 We have reported the concurrent experimental and theoretical (both static and dynamic) characteri-  
2 zation of a common antifungal drug, clotrimazole (CTZ), incorporated in ordered mesoporous silica  
3 of the MSU-H type.

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6 Comparison between experimental and computed interaction energies, pore size distributions and vi-  
7 brational spectra resulted in a very good agreement. This suggests that DFT simulation of drug ad-  
8 sorption on realistic amorphous silica surface models can be used to predict the behavior of drugs in-  
9 side the pores of mesoporous silica materials, provided that a good sampling of the different interac-  
10 tions is achieved.

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13 As regards the MSU-H-CTZ system, results shown here reveal that CTZ is adsorbed, over the incor-  
14 poration with  $\text{scCO}_2$ , following a homogeneous and statistical distribution on the silica surface. This  
15 picture is explained only considering together the molecular layer models. Thanks to the quantum  
16 mechanical simulations, we were able to describe the experimentally measured change in pore size  
17 distribution by considering the thickness of the molecular layer at high loading.

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20 Simulations and their comparison with experimental data suggest that CTZ interacts in multiple ori-  
21 entations with the silica surface: the strongest interaction includes an H-bond between the nitrogen  
22 atom of the imidazole ring with terminal silanols of H-bonded chains already present on the silica  
23 walls, but other H-bonds that involve the phenyl rings and the chlorine atom are also present. Com-  
24 parison with experiment shows that a very stable configuration found by the simulations is probably  
25 quite rare in the real system due to the specificity of the adsorbing sites.

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28 AIMD simulations provided a direct evidence of high mobility of drugs adsorbed into mesoporous  
29 silica, supporting experimental results in literature. We also revealed that large displacements paral-  
30 lel to the surface of adsorbed molecular layers may be triggered by local changes in the interactions.  
31 Energetically, it was confirmed that London forces play a key role in drug adsorption on silica and  
32 that they are indeed the true driving force of the process in the case of apolar molecules such as  
33 CTZ, also assisting the stabilization of molecular layers on the surface. Experimental and theoretical  
34 data also agree in that CTZ adsorption on silica and CTZ interaction on its own crystal surfaces are  
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1 almost isoenergetic processes, thus suggesting a possible explanation for the lack of crystallization  
2 inside the pores of MSU-H.  
3

4 A comprehensive picture of this complex scenario was achieved only through the joint of simulation  
5 and experiment and the present methodology paves the way to improve our understanding of the in-  
6 teractions between drugs and delivery systems.  
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## 10 11 12 13 **ASSOCIATED CONTENT**

### 14 **Supporting Information**

15 Computational details, crystalline structure, surface models, nitrogen adsorption isotherms, AIMD  
16 simulations details, FT-IR spectra and bands attribution, absolute energies, fractional coordinates of  
17 all the optimized structures and video of the significant AIMD simulations. This material is available  
18 free of charge via Internet at <http://pubs.acs.org>.  
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### 28 **AKCNOWLEDGMENTS**

29 The vast majority of the calculations have been carried out on the FERMI supercomputer, in Bolo-  
30 gna, Italy, thanks to the ISCRA (Italian SuperComputing Resource Allocation) class B projects  
31 IBUEXCIP and SIL4SKIN. Models have been visualized and manipulated by MOLDRAW.<sup>57</sup> AIMD  
32 simulations have been analyzed through VMD.<sup>58</sup> Figures have been rendered with the Tachyon ray  
33 tracer,<sup>59</sup> built into VMD.  
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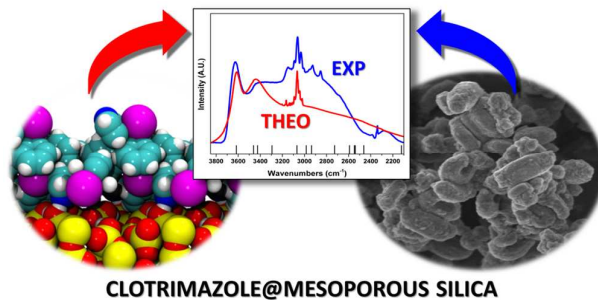
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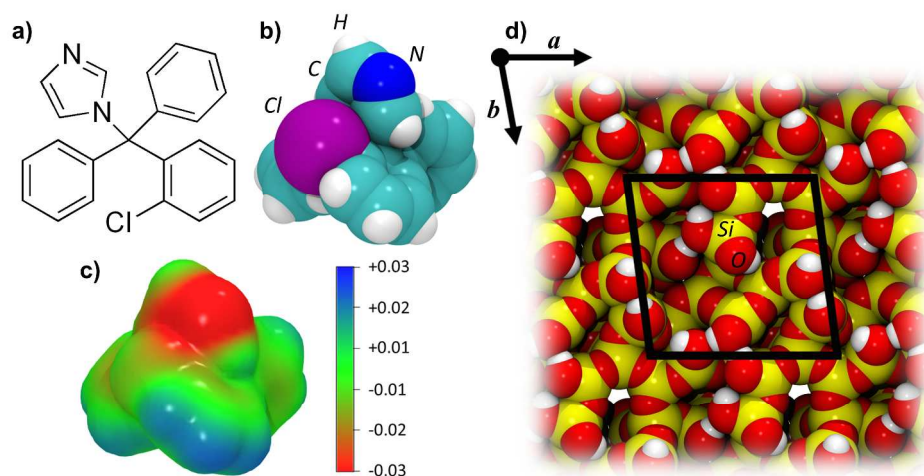
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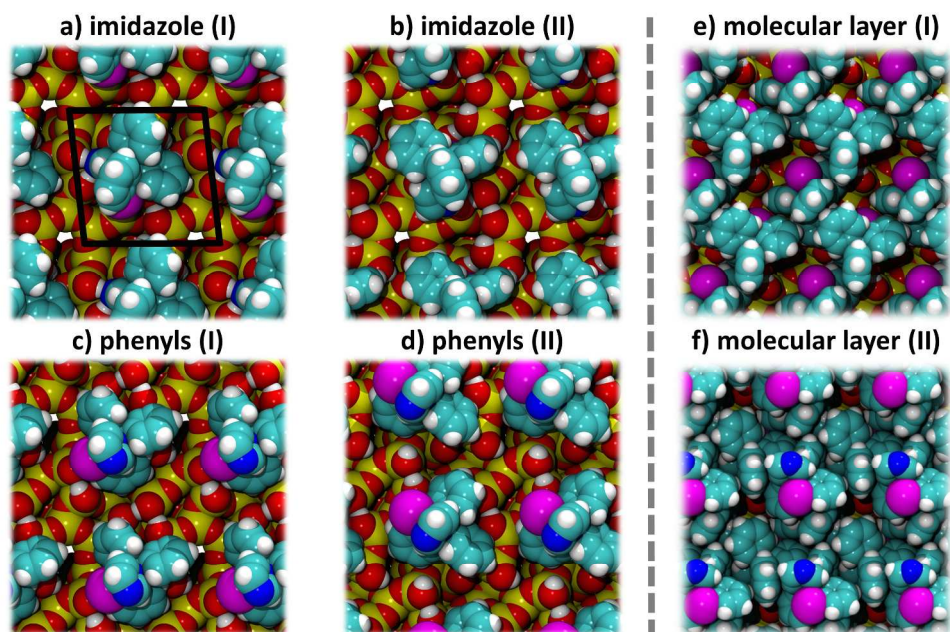
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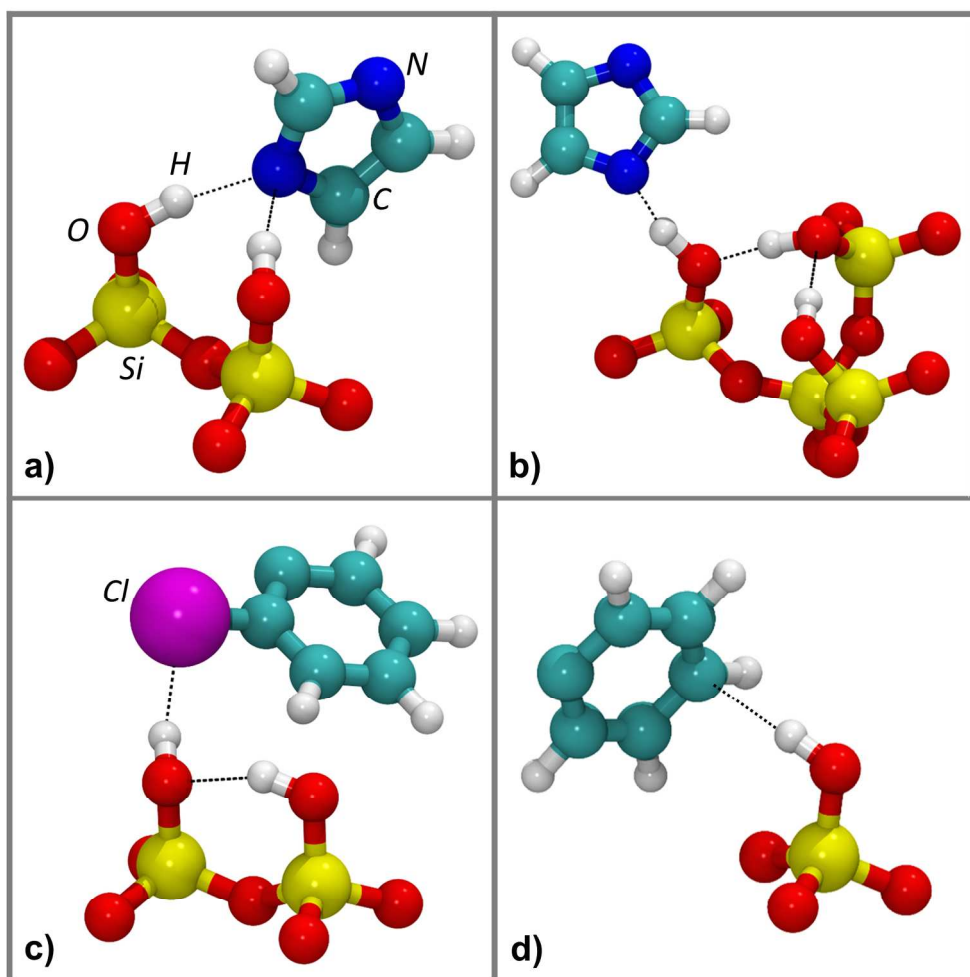
## GRAPHICS TABLE OF CONTENTS





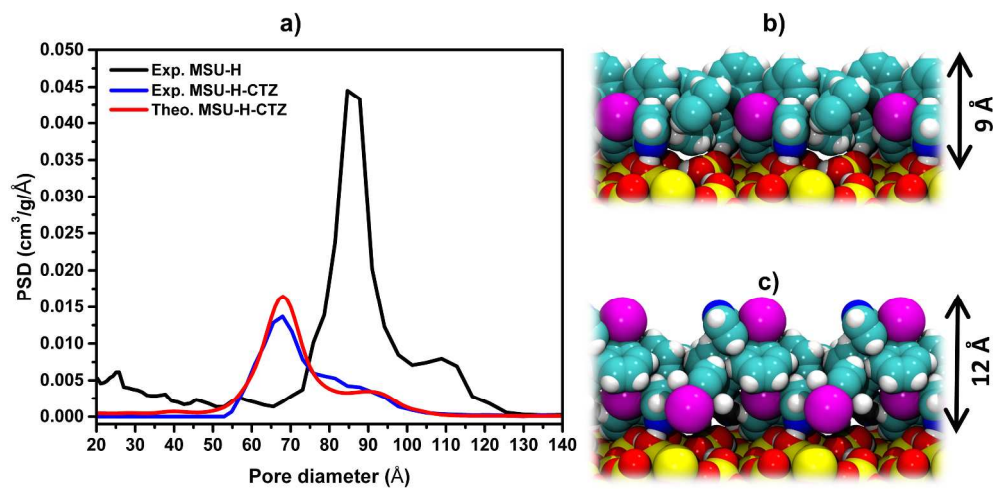
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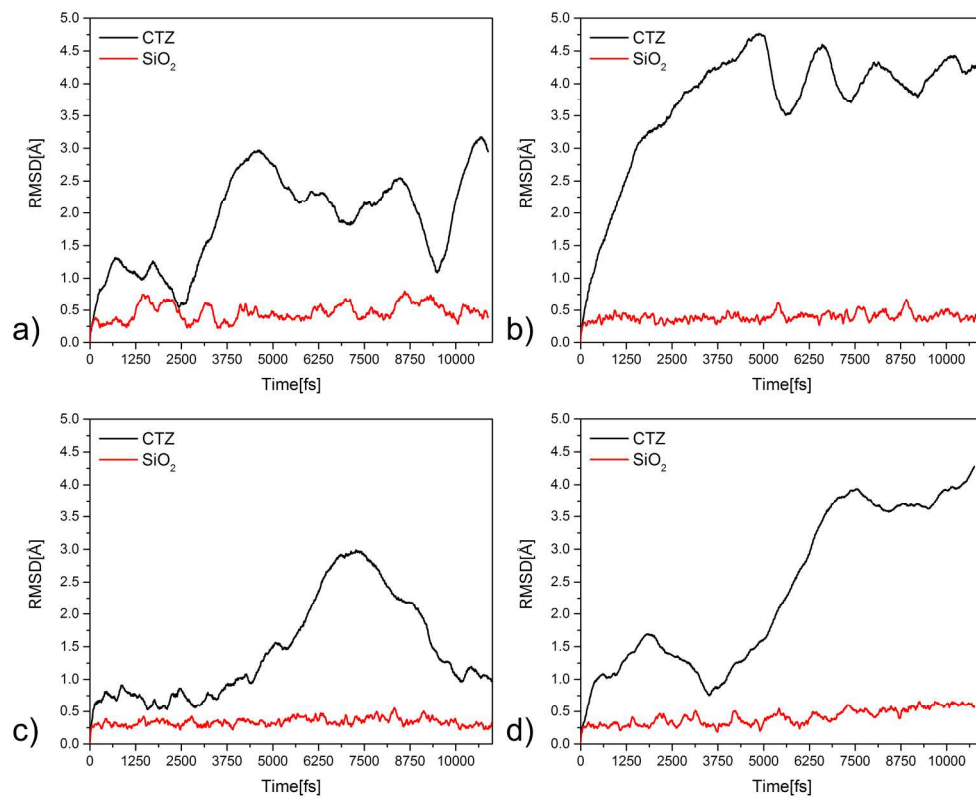


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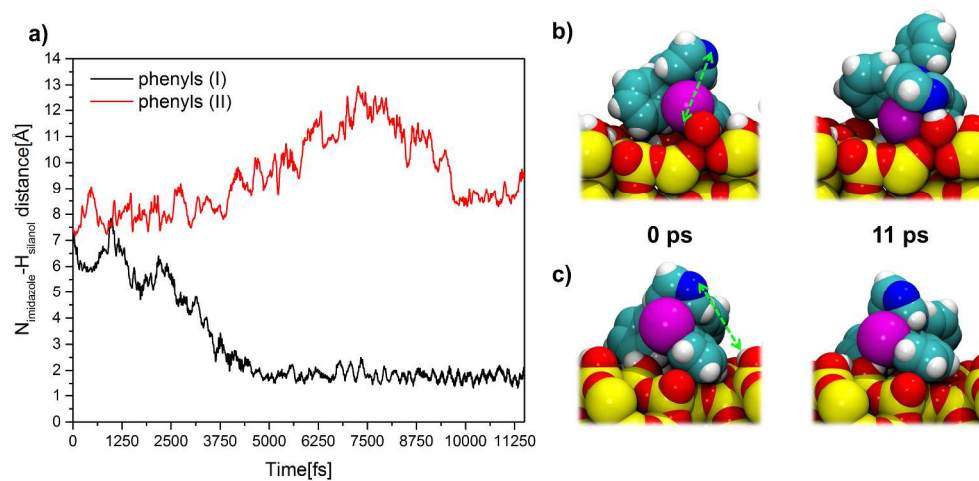




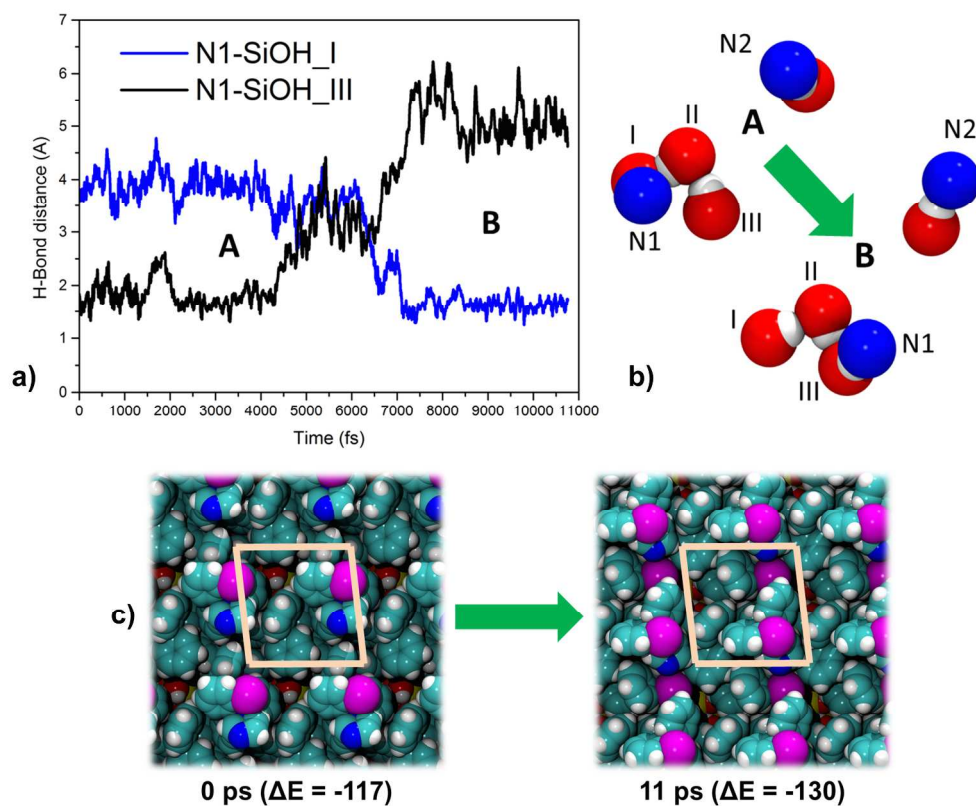
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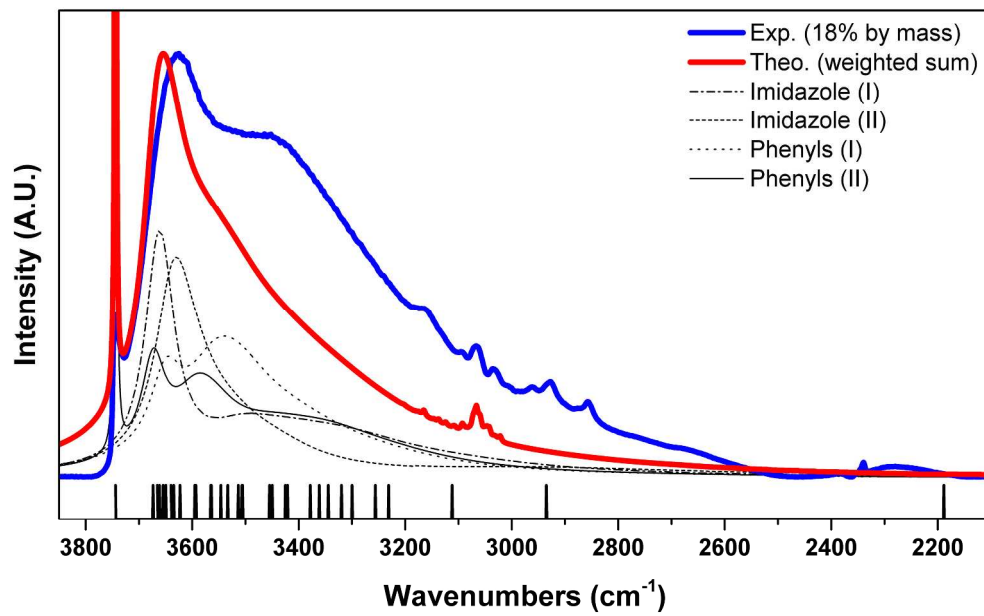
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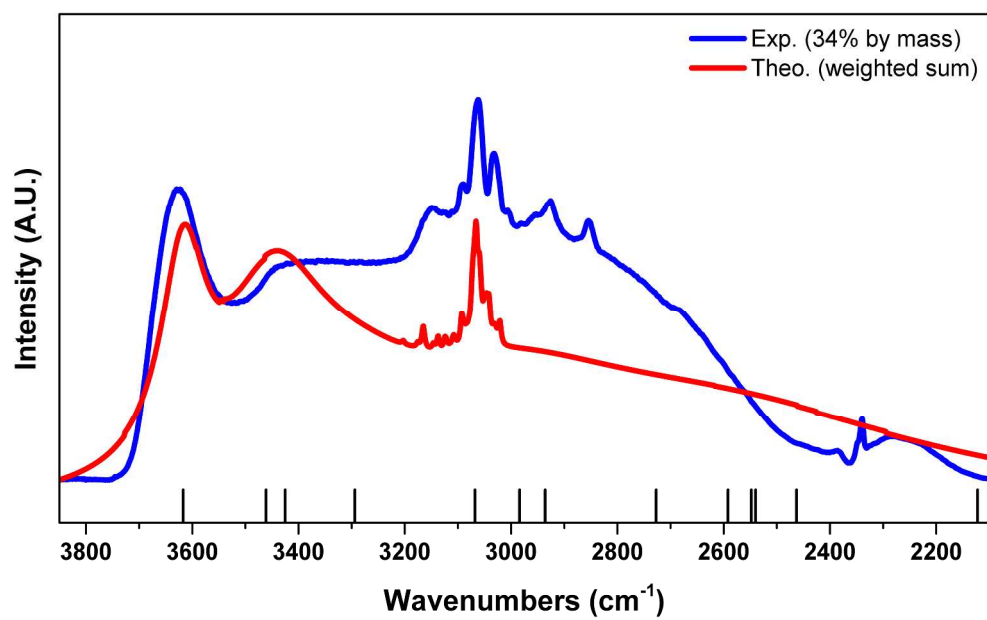
305x149mm (300 x 300 DPI)



225x183mm (300 x 300 DPI)



268x168mm (300 x 300 DPI)



270x169mm (300 x 300 DPI)