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Extended Hamiltonians and shift, ladder functions and operators

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Abstract

In recent years, many natural Hamiltonian systems, classical and quantum, with constants of motion of high degree, or symmetry operators of high order, have been found and studied. Most of these Hamiltonians, in the classical case, can be included in the family of extended Hamiltonians, geometrically characterized by the structure of warped manifold of their configuration manifold. For the extended manifolds, the characteristic constants of motion of high degree are polynomial in the momenta of determined form. We consider here a different form of the constants of motion, based on the factorization procedure developed by S. Kuru, J. Negro and others. We show that an important subclass of the extended Hamiltonians admits factorized constants of motion and we determine their expression. The classical constants may be non-polynomial in the momenta, but the factorization procedure allows, in a type of extended Hamiltonians, their quantization via shift and ladder operators, for systems of any finite dimension.

1 Introduction

In the past few years, the systematic study of classical Hamiltonian systems admitting polynomial constants of motion of high degree, and quantum Hamiltonian operators with symmetry operators of high order, produced several articles [5, 4, 15, 16, 18, 19, 20, 21, 22, 24], many of them focused on the superintegrability of the systems considered. In several articles [6, 7, 8, 9, 10, 11, 12, 13], we, together with Luca Degiovanni, were able to unify most of those Hamiltonian systems under the common structure of *extended Hamiltonians*. For the case of natural Hamiltonians, characterized geometrically in [9], to be an extended Hamiltonian implies that the configuration manifold is a particular *warped Riemannian manifold* [23], equipped with a conformal Killing vector field of special type. Extended Hamiltonians allow the existence of privileged coordinate systems adapted to the warped geometry, with the property of splitting a given $(N+1)$ -dimensional system into one-dimensional and N -dimensional parts. The

most important consequence of the extended Hamiltonian structure, in the classical case, is the existence of constants of motion, that we call *characteristic*, determined by the power of an operator applied to a well determined function, in correspondence of a rational parameter employed in the construction of the extended Hamiltonian itself. The classical constants of motion of high degree in the Tremblay-Turbiner-Winternitz (TTW) system (and in its particular cases: three-particles Jacobi-Calogero and Wolfes systems), in the Post-Winternitz system, and in their multi-dimensional and non-zero-curvature manifold generalizations, as shown in [12], are exactly of that type. Until now, we did not have a similar unified procedure for determining the symmetry operators of the quantum Hamiltonian corresponding to an extended Hamiltonian. In the cases considered in literature, essentially the TTW system and its generalizations, the symmetry operators are determined case by case, making use of the complete separation of variables of the Hamiltonians. The symmetry operators are therefore always obtained by some factorization procedure. One of these factorization techniques is that developed by S. Kuru, J. Negro and co-authors [17, 1, 2, 5]. The procedure is based on the well known factorization in creation-annihilation, or shift and ladder, operators, developed firstly for the harmonic oscillator. This technique is applied to the classical case also, obtaining the factorization of classical constants of motion in shift and ladder functions. As S. Kuru and J. Negro pointed out in a private communication, the basic ansatz of their factorization technique (the existence of particular ladder functions) implies the structure of an extended manifold for the configuration manifold of the system. By assuming this ansatz, we are able to show here how the classical characteristic constants of motion of extended Hamiltonian systems factorize in shift and ladder functions. Moreover, we show that, in some important cases, the extended Hamiltonian structure implies the validity of the ansatz, that it is then no longer a restriction. Most importantly, in analogy with these results we are able to build a general quantization procedure of the characteristic constant of motion, for one of the three possible forms of extensions, in any finite dimension, the model being the shift-ladder factorization of above. The content of the article, Section by Section, is the following. In Section 2 we recall the basics of the theory of extended Hamiltonians; in Section 3 we show that the existence of certain ladder functions implies the factorization of the characteristic first integrals of the extended Hamiltonians (Theorem 6); in Section 4, on the contrary, we show that a large class of extended Hamiltonians admit a factorized first integral, according to Theorem 6. In Section 5, we make a comparison between the Kuru-Negro factorization and the one proposed here. In Section 6, we consider the quantum version of the procedure and we propose a quantized extension procedure, holding for a subclass of the extended Hamiltonians, based on the Kuru-Negro technique. Given a type of extended Hamiltonian, we obtain an extended quantum Hamiltonian and its factorized characteristic symmetry operator (Theorem 10), the construction of which is described in subsections 6.1-6.3.

2 Extensions: definitions and relevant properties

Let $L(q^i, p_i)$ be a Hamiltonian with N degrees of freedom, that is defined on the cotangent bundle of an N -dimensional manifold. We say that L *admits extensions*, if there exists $(c, c_0) \in \mathbb{R}^2 - \{(0, 0)\}$ such that there exists a non null solution $G(q^i, p_i)$ of

$$X_L^2(G) = -2(cL + c_0)G, \quad (1)$$

where X_L is the Hamiltonian vector field of L . If L admits extensions, then, for any $\gamma(u)$ solution of the ODE

$$\gamma' + c\gamma^2 + C = 0, \quad (2)$$

depending on the arbitrary constant parameter C , we say that any Hamiltonian $H(u, q^i, p_u, p_i)$ with $N + 1$ degree of freedom of the form

$$H = \frac{1}{2}p_u^2 - k^2\gamma'L + k^2c_0\gamma^2 + \frac{\Omega}{\gamma^2}, \quad k \in \mathbb{Q} - \{0\}, \quad \Omega \in \mathbb{R} \quad (3)$$

is an *extension of L* .

Extensions of Hamiltonians where introduced in [8] and studied because they admit polynomial in the momenta first integrals generated via a recursive algorithm. Moreover, the degree of the first integrals is related with the (choice of) numerator and denominator of the rational parameter $k = m/n$, $m, n \in \mathbb{N}^*$. Indeed, for any $m, n \in \mathbb{N} - \{0\}$, let us consider the operator

$$U_{m,n} = p_u + \frac{m}{n^2}\gamma X_L \quad (4)$$

Proposition 1. [11] For $\Omega = 0$, the Hamiltonian (3) is in involution with the function

$$K_{m,n} = U_{m,n}^m(G_n) = \left(p_u + \frac{m}{n^2}\gamma(u)X_L\right)^m(G_n) \quad (5)$$

where $m/n = k$ and G_n is the n -th term of the recursion

$$G_1 = G, \quad G_{n+1} = X_L(G)G_n + \frac{1}{n}G X_L(G_n), \quad (6)$$

starting from any solution G of (1).

Remark 1. The functions G_n satisfy

$$X_L^2 G_n = -2n^2(cL + c_0)G_n.$$

For $\Omega \neq 0$, the recursive construction of a first integral is more complicated: we consider the following function, depending on two non-zero integers $2s, r$ (the first one has to be even)

$$\bar{K}_{2s,r} = (U_{2s,r}^2 + 2\Omega\gamma^{-2})^s(G_r). \quad (7)$$

where the operator $U_{2s,r}^2$ is defined according to (4) as

$$U_{2s,r}^2 = \left(p_u + \frac{2s}{r^2}\gamma(u)X_L\right)^2,$$

and G_r is, as in (5), the r -th term of the recursion (6). with $G_1 = G$ solution of (1). For $\Omega = 0$ the functions (7) reduces to (5) and thus can be computed also when the first of the indices is odd.

We have

Theorem 2. [12] For any $\Omega \in \mathbb{R}$, the Hamiltonian (3) with $k = m/n$ satisfies for $m = 2s$,

$$\{H, \bar{K}_{m,n}\} = 0, \quad (8)$$

for $m = 2s + 1$,

$$\{H, \bar{K}_{2m,2n}\} = 0. \quad (9)$$

We call K and \bar{K} , of the form (5) and (7) respectively, *characteristic first integrals* of the corresponding extensions

Remark 2. In [9] it is proven that the ODE (2) defining γ is a necessary condition in order to get a characteristic first integral of the form (5) or (7). According to the value of c , the explicit form of $\gamma(u)$ is given (up to constant translations of u) by

$$\gamma = \begin{cases} -Cu & c = 0 \\ \frac{1}{T_\kappa(cu)} = \frac{C_\kappa(cu)}{S_\kappa(cu)} & c \neq 0 \end{cases} \quad (10)$$

where $\kappa = C/c$ is the ratio of the constant parameters appearing in (2) and T_κ , S_κ and C_κ are the trigonometric tagged functions [21]

$$S_\kappa(x) = \begin{cases} \frac{\sin \sqrt{\kappa}x}{\sqrt{\kappa}} & \kappa > 0 \\ x & \kappa = 0 \\ \frac{\sinh \sqrt{|\kappa|}x}{\sqrt{|\kappa|}} & \kappa < 0 \end{cases} \quad C_\kappa(x) = \begin{cases} \cos \sqrt{\kappa}x & \kappa > 0 \\ 1 & \kappa = 0 \\ \cosh \sqrt{|\kappa|}x & \kappa < 0 \end{cases},$$

$$T_\kappa(x) = \frac{S_\kappa(x)}{C_\kappa(x)},$$

(see also [11] for a summary of their properties).

It is proved in [8, 12] that the characteristic first integrals K or \bar{K} are functionally independent from H , L , and from any first integral $I(p_i, q^i)$ of L . This means that the extensions of (maximally) superintegrable Hamiltonians are (maximally) superintegrable Hamiltonians with one additional degree of freedom.

Remark 3. The equation (1) is the fundamental condition for the extension procedure. When L is a N -dimensional natural Hamiltonian with metric tensor $\mathbf{g} = (g_{ij})$ and G does not depend on the momenta, (1) splits into several equations, the coefficients of monomials in the momenta, whose leading term is the Hessian equation

$$\nabla_i \nabla_j G + cG g_{ij} = 0, \quad i, j = 1, \dots, N. \quad (11)$$

This equation is of paramount importance in theory of warped manifolds, see for example [23], where it is proved, for manifolds of positive definite signature, the strict connection between the constant c and the curvature of the metric g

(Theorem 2, where it is also proved that, if the metric is geodesically complete and $c \neq 0$, then g is of constant curvature c) and the warped structure of g whenever the equation (11) admits solutions $G(q^k)$ (Lemma 1.2 of [23]).

However, by relaxing the hypothesis of the these statements (for example allowing the non geodesic completeness of the manifold in the case of $c \neq 0$), solutions of (11) can be found also in non-constant curvature manifolds, [14]. Moreover, the structure of the extensions H of natural Hamiltonians shows that the extended metric is again a warped metric.

We want to study other properties of extensible and extended natural Hamiltonians, related with the existence of ladder and shift functions as well as ladder and shift operators generalizing what done by Kuru Negro as instance in [5]. These properties seem to be strictly connected with separability of the dynamics of the additional degree of freedom, parametrized by the variable u , but they appear to be independent of the separability within L .

3 Ladder functions and factorized first integrals of extended Hamiltonians

Let H be a Hamiltonian of the form

$$H = \frac{1}{2}p_u^2 + \alpha(u)L(p_i, q^i) + \beta(u), \quad (12)$$

and let X_H , X_L be the Hamiltonian vector fields of H and L , respectively.

Definition 1. *The function $F(p_i, q^j)$ is a ladder function of L if there exists a function $f(L)$ such that*

$$X_L F = f(L)F.$$

Proposition 3. *A ladder function F of the Hamiltonian L verifies (1), i.e.,*

$$X_L^2 F = -2(cL + c_0)F, \quad (13)$$

if and only if

$$f(L) = \pm\sqrt{-2(cL + c_0)}. \quad (14)$$

Hence,

Theorem 4. *If L admits a ladder function G with factor $f(L)$ given by (14), then L admits extensions of the form (3). Moreover, for*

$$\alpha = -k^2\gamma', \quad \beta = k^2c_0\gamma^2 + \frac{\Omega}{\gamma^2},$$

the Hamiltonian (12) is an extension of L , and the ladder function G generates the first integrals (5), (7).

The existence of a ladder function of L allows the factorization of the characteristic first integrals K or \bar{K} of any extension of L .

Lemma 5. *If $X_L G = f(L)G$, then, for any $n \in \mathbb{N} - \{0\}$,*

$$X_L G_n = n f G_n,$$

$$G_n = (2f)^{n-1} G^n,$$

hence, both G^n and G_n are additional ladder functions of L .

Proof. Recall that

$$G_1 = G, \quad G_{n+1} = X_L GG_n + \frac{1}{n} GX_L G_n,$$

then, $G_2 = 2GX_LG = 2f(L)G^2$, and

$$X_L G_2 = 2X_L G X_L G + 2GX_L^2 G = 4G^2 f^2 = 2f(2G^2 f) = 2fG_2.$$

By induction on n , we assume that $X_L G_n = nfG_n$ and $G_n = (2f)^{n-1}G^n$, then

$$G_{n+1} = (fG)G_n + \frac{1}{n}(nfG_n)G = (2fG)G_n = (2f)^n G^{n+1},$$

$$X_L G_{n+1} = X_L(2fGG_n) = 2(n+1)f^2GG_n = (n+1)fG_{n+1}.$$

□

Theorem 6. Let G be a ladder function of L with factor $f(L)$ given by (14), and H be an extension (3) of L admitting first integrals (5) (for $\Omega = 0$) or (7) (for $\Omega \neq 0$). Then,

i) for any (m, n) , the functions (5) factorize as

$$K_{m,n} = (2f)^{n-1}G^n F,$$

where

$$F = \left(p_u + \frac{m}{n}\gamma f \right)^m.$$

ii) for any (s, r) , the functions (7) factorize as

$$\bar{K}_{2s,r} = (2f)^{r-1}G^r F,$$

where

$$\begin{aligned} F &= \left[\left(p_u + \frac{2s}{r}\gamma f \right)^2 + 2\Omega\gamma^{-2} \right]^s \\ &= \left(p_u + \frac{2s}{r}\gamma f + i\frac{\sqrt{2\Omega}}{\gamma} \right)^s \left(p_u + \frac{2s}{r}\gamma f - i\frac{\sqrt{2\Omega}}{\gamma} \right)^s. \end{aligned} \quad (15)$$

In both cases, F satisfies the relation $X_H F = nk^2\gamma' f F$.

Proof. By (5) and Lemma 5 we get

$$K_{m,n} = \left(p_u + \frac{m}{n^2}\gamma nf \right)^m G_n,$$

that, after simplification, gives the statement. Since $K_{m,n}$ is a first integral of H , we have

$$0 = X_H K_{m,n} = X_H(FG_n) = X_H FG_n + FX_H G_n.$$

Then, because for any function $\phi(q^i, p_i)$ we have $X_H \phi = -k^2\gamma' X_L \phi$, we find that F satisfies

$$X_H F = -\frac{F}{G_n} X_H G_n = k^2\gamma' \frac{F}{G_n} X_L G_n = k^2 n \gamma' f F.$$

The same reasoning applies to (7). □

In [5] the functions F in Theorem 6 are called *shift functions*. Clearly, the factor $(2f)^{n-1}$, or $(2f)^{r-1}$, in the factorization of the first integrals (7) and (5) given in Theorem 6 is irrelevant and can be dropped down. Thus, we have that the first integrals of the extended Hamiltonian (3) are product of shift and ladder functions. This fact makes easier the computation of the algebra of symmetries of H [5].

Remark 4. Due to the expression of the function f , the characteristic first integrals of Theorem 6 are not polynomial in the momenta.

4 Extended Hamiltonians and ladder functions

Let us assume that L is a N -dimensional natural Hamiltonian

$$L = \frac{1}{2}g^{ij}p_ip_j + V(q^i). \quad (16)$$

In the previous section we have shown that a function G satisfying condition (1) for the Hamiltonian L which is also a ladder function for L (or equivalently a ladder function of L with special factor (14)) allows the factorization of the characteristic first integrals $K_{m,n}$ of any extension H of L . All the extensions studied until now are based on the existence of solutions G of (1) that are polynomial in the momenta. Consequently, the characteristic first integrals of extensions are also polynomial in the momenta. These G , however, are not ladder functions of L . In this section we show how to obtain ladder functions of L from some of these solutions G and, therefore, an alternative factorization of the first integrals is presented. Inspired by [5], we search for ladder functions of the form

$$G^\pm = A_\pm^i p_i + B_\pm f + \frac{1}{f}C_\pm, \quad (17)$$

with $f = \sqrt{-2(cL + c_0)}$ and $\mathbf{A}_\pm, B_\pm, C_\pm, V$ depending on the (q^i) only. If we are able to determine $\mathbf{A}_\pm, B_\pm, C_\pm, V$ such that

$$X_L G^\pm = \pm f G^\pm, \quad (18)$$

then we obtain ladder functions G^\pm of L verifying the hypothesis of Theorem 6.

Proposition 7. *If we assume $C_\pm = c_1$ constant, then a solution in form (17) of (18) is given by*

$$G^\pm = \pm(\nabla G)^i p_i + Gf + \frac{1}{f}c_1, \quad f = \sqrt{-2(cL + c_0)}, \quad (19)$$

with the function G and the constant c_1 satisfying

$$[\mathbf{g}, \nabla G] + c\mathbf{g}G = 0, \quad (20)$$

$$\nabla V \cdot \nabla G - 2(cV + c_0)G + c_1 = 0, \quad (21)$$

where $[\cdot, \cdot]$ are the Schouten brackets.

Proof. By expanding equation (18) with G^\pm and L given by (17) and (16) respectively, we have, if C_\pm is constant,

$$\begin{aligned} X_L G^\pm \mp f G^\pm &= \\ &= p_i p_j (\nabla^i A_\pm^j \pm c g^{ij} B_\pm) + p_i f (\nabla^i B_\pm \mp A_\pm^i) - A_\pm^i \nabla_i V \pm 2(cV - c_0) B_\pm \mp C_\pm. \end{aligned}$$

By equating to zero the coefficients of the monomials in the momenta, we obtain easily

$$\begin{aligned} [\mathbf{g}, \mathbf{A}_\pm] &= \mp c \mathbf{g} B_\pm, \\ \nabla B_\pm &= \pm \mathbf{A}_\pm, \\ \nabla V \cdot \mathbf{A}_\pm \mp 2(cV + c_0) B_\pm \pm C_\pm &= 0, \end{aligned}$$

or, equivalently,

$$[\mathbf{g}, \nabla B_\pm] + c \mathbf{g} B_\pm = 0, \quad (22)$$

$$\nabla V \cdot \nabla B_\pm - 2(cV + c_0) B_\pm + C_\pm = 0, \quad (23)$$

$$\nabla B_\pm = \pm \mathbf{A}_\pm. \quad (24)$$

Equations (22), (23) are the same for (B_+, C_+) and (B_-, C_-) , thus, their solutions coincide with the solutions (G, c_1) of (20) and (21). Therefore, by setting $\mathbf{A}_\pm = \pm \nabla G$, $B_\pm = G$, and $C_\pm = c_1$, we get (19). \square

Remark 5. If for a given V there exists G satisfying (21) for some constants c, c_0, c_1 , then the same G is a solution of

$$\nabla \tilde{V} \cdot \nabla G - 2(\tilde{c} \tilde{V} + \tilde{c}_0) G + \tilde{c}_1 = 0,$$

for the potential $\tilde{V} = a_1 V + a_2$ ($a_1, a_2 \in \mathbb{R}$) by choosing $\tilde{c} = c$, $\tilde{c}_1 = a_1 c_1$, $\tilde{c}_0 = a_1 c_0 + c a_2$. This means that if a solution G is known for $c_0 = 0$, $c \neq 0$, then, by adding the inessential constant \tilde{c}_0/c to the potential V we get a solution G of (21) for any value of \tilde{c}_0 .

Equations (20) and (21) must be compared with equation (1)

$$X_L^2(G) = -2(cL + c_0)G,$$

which determines the function G necessary to build the extensions of L . By assuming $G(q^i)$, equation (1) splits into [8]

$$\nabla_i \nabla_j G + c g_{ij} G = 0, \quad (25)$$

$$\nabla V \cdot \nabla G - 2(cV + c_0) G = 0. \quad (26)$$

It is clear that (20) and (21) are equivalent to

$$X_L^2(G) = -2(cL + c_0)G + c_1, \quad (27)$$

since $[g, \nabla G] = [g, [g, G]]$. Thus, G must verify the hessian equation (25). Hence, for $c_1 = 0$, we get (25) and (26). This means that, given any function G solution of (25), (26), we can immediately obtain solutions G^\pm (19) just by putting $c_1 = 0$:

Proposition 8. Let $G(q^i)$ be any solution of (25), (26) (i.e. of $X_L^2 G = f^2 G$), then, the functions

$$G^\pm = \pm \nabla^i G p_i + G f,$$

are ladder functions of L , i.e. solutions of

$$X_L G^\pm = \{L, G^\pm\} = \pm f G^\pm.$$

Moreover, the extended Hamiltonian H of L admits the factorized first integrals described in Theorem 6.

The solutions G of (25) have been computed for several Riemannian and pseudo-Riemannian manifolds in [7], [8], [9], [10]. Moreover, in [8] it is proved that the constant c is related to the curvature, or the sectional curvatures, of the base manifold of L . Furthermore, equation (26) can be seen as a compatibility condition that the potential V has to fulfil. By admitting $c_1 \neq 0$, (i.e., considering (23) instead of (26) as compatibility condition of the potential) we enlarge the class of compatible potentials V , which would depend on c_1 and which reduces for $c_1 = 0$ to the class of potential verifying (26).

However, this technique does not imply that we are enlarging the set of Hamiltonians that admit extensions. Indeed, more general functions $G(q^i, p_i)$ which are polynomial in the momenta and solutions of $X_L^2 G = f^2 G$ on the same manifold, can be compatible with these potentials V , as is the case of the TTW system (see [12] and [5]), where G is linear and homogeneous in the momenta. In this case, for example, we have $V = (C_1 + C_2 \cos q) \sin^{-2} q$, and we know the solution $G = p \sin q$ for $c = 1$ and $c_0 = 0$. However, for the same values of c and c_0 , the (20), (21) can be integrated, yielding the solution $G = \cos q$ and $c_1 = -C_2$.

Remark 6. Since any ladder function of L is also a solution of (1), the solutions G^\pm of the form (19), for a given L , allow the extension of the Hamiltonian L by putting $G = G^\pm$. We observe that the conditions on c , i.e. constraints on the curvature or sectional curvature of the base manifold of L , remain unchanged.

Example 1. In [8] are computed the extensions for natural Hamiltonian L on \mathbb{S}^3 with Hopf coordinates ($q^1 = \eta, q^2 = \xi_1, q^3 = \xi_2$). In these coordinates the non-zero metric components are $g_{11} = 1$, $g_{22} = \sin^2 \eta$ and $g_{33} = \cos^2 \eta$. The function G , obtained for $c = 1$ equal to the curvature of the sphere and $c_0 = 0$, is

$$G = (a_3 \sin \xi_1 + a_4 \cos \xi_1) \sin \eta + (a_1 \sin \xi_2 + a_2 \cos \xi_2) \cos \eta,$$

where a_i are arbitrary constants. The corresponding ladder functions are

$$\begin{aligned} G^\pm &= \pm \left\{ [(a_3 \sin \xi_1 + a_4 \cos \xi_1) \cos \eta - (a_1 \sin \xi_2 + a_2 \cos \xi_2) \sin \eta] p_\eta + \right. \\ &\quad \left. \left[\frac{a_3 \cos \xi_1 - a_4 \sin \xi_1}{\sin \eta} \right] p_{\xi_1} + \left[\frac{a_1 \cos \xi_2 - a_2 \sin \xi_2}{\cos \eta} \right] p_{\xi_2} \right\} + G f + \frac{c_1}{f}. \end{aligned}$$

In [8], it is provided a solution V of (26) for $a_1 = 1, a_2 = a_3 = a_4 = 0$, as

$$V = \frac{1}{\sin^2 \eta} F \left(\xi_1, \frac{\tan \eta}{\cos \xi_2} \right),$$

where F is any smooth function, potential that can be separable or not in the chosen coordinates. With the same assumptions, the integration of (21) for $c_1 \neq 0$ gives

$$V = \frac{1}{\sin^2 \eta} F \left(\xi_1, \frac{\tan \eta}{\cos \xi_2} \right) + \frac{c_1 \cos \eta \sin \xi_2}{1 - \cos^2 \eta \sin^2 \xi_2},$$

in accordance to the fact that, for $c_0 = 0$, equation (26) is a linear homogeneous differential equation in V and a solution of (21) is the sum of the solution of the homogeneous equation with a particular solution of the non-homogeneous case.

5 Comparison with Kuru-Negro factorization

In [5] the authors determine for the classical Hamiltonian of the TTW system in polar coordinates (r, θ) shift and ladder functions, from which constants of the motion X^\pm are obtained depending on a rational parameter $k = m/n$:

$$X^\pm = (\Xi^\pm)^n (A^\mp)^{2m} (D^\pm)^m.$$

Their method is: (i) to split the Hamiltonian into a radial part combined together with a one-dimensional Hamiltonian H_θ (the Posch-Teller Hamiltonian); (ii) to determine ladder functions Ξ^\pm of H_θ with factor $\mp\sqrt{-8H_\theta}$; (iii) to construct the functions A and D interconnecting the radial part with the constant of the motion $k\sqrt{H_\theta}$.

We show that this construction of a factorized first integral of H can be straightforwardly generalized to any H which is an extension (3) with $c \neq 0$ of a Hamiltonian L admitting ladder functions G^\pm such that

$$X_L G^\pm = \pm \sqrt{-2(cL + c_0)} G^\pm. \quad (28)$$

However, while Theorem 6 holds for any extension (also for $c = 0$), in the following construction the hypothesis that $c \neq 0$ seems to be crucial.

By using the property (2) of the γ , we can write any extension (3) of L with $c \neq 0$ as

$$H = \frac{1}{2} p_u^2 + \frac{1}{2} \left(\gamma^2 + \frac{C}{c} \right) M^2 + \frac{\omega^2}{2\gamma^2} - \frac{C}{c} k^2 c_0, \quad (29)$$

where, in order to underline the similarity with the formulas in [5], we set $M = k\sqrt{2(cL + c_0)}$ and $\omega^2 = 2\Omega$.

Therefore, we have

Proposition 9. *Given a Hamiltonian L such that equation (28) admits a solution G^\pm for some constants c and c_0 not both zero, then, (i) L admits extensions, (ii) for any extension H (29) of L , the functions*

$$X^\pm = (G^\pm)^{2n} (A^\mp)^{2m} (D^\pm)^m,$$

with

$$\begin{aligned} A^\pm &= \mp \frac{i}{\sqrt{2}} p_u + \frac{\omega}{\sqrt{2}\gamma} - \frac{M}{\sqrt{2}} \gamma, \\ D^\pm &= \mp \frac{i}{\gamma} p_u + \frac{\omega}{\gamma^2} - \frac{H}{\omega} - \frac{C}{c\omega} \left(c_0 k^2 - \frac{M^2}{2} \right), \end{aligned}$$

are factorized first integrals of H .

Proof. The Hamiltonian L admits an extension H because of Proposition 3. Moreover, we have, for any solution G^\pm of (28)

$$\begin{aligned}\{H, G^\pm\} &= k^2(c\gamma^2 + C)\{L, G^\pm\} \\ &= \pm k^2(c\gamma^2 + C)\sqrt{-2(cL + c_0)}G^\pm \\ &= \mp ik\gamma' MG^\pm,\end{aligned}$$

and, since $k = m/n$,

$$\{H, (G^\pm)^{2n}\} = \mp i2m\gamma' M(G^\pm)^{2n}. \quad (30)$$

The functions A^\pm and D^\pm satisfy the following factorization and commutation formulas

$$\begin{aligned}H &= A^+A^- + \omega M + \frac{C}{c}\left(\frac{M^2}{2} - k^2d\right), \\ D^+D^- - \left(\frac{H}{\omega} + \frac{C}{c\omega}\left(c_0k^2 - \frac{M^2}{2}\right)\right)^2 &= -M^2, \\ \{H, A^\pm\} &= \mp i\gamma'\left(M + \frac{\omega}{\gamma^2}\right)A^\pm, \quad (31) \\ \{H, D^\pm\} &= \mp 2i\gamma'\frac{\omega}{\gamma^2}D^\pm. \quad (32)\end{aligned}$$

The statement follows by computing directly $\{H, X^\pm\}$ taking into account (31), (32) and (30). \square

We recall that, from Theorem 6 and Proposition 8, it follows that the factorization of the characteristic first integral is, for $k = \frac{m}{n}$, of the form $(G^\pm)^n F^m$ (up to the factor f^{n-1} , inessential because itself first integral of H). Therefore, the factorization determined by Theorem 6 appears to be simpler than the one given in Proposition 9.

Example 2. The ladder functions Ξ^\pm of [5] for the Tremblay-Turbiner-Winternitz Hamiltonian satisfy (13) and guarantee the extensibility of the Hamiltonian H_θ . Indeed, given (26) of [5]

$$\{H_\theta, \Xi^\pm\} = \pm 4i\sqrt{H_\theta}\Xi^\pm,$$

with

$$\begin{aligned}\Xi^\pm &= \pm ip_\theta \sin 2\theta + \sqrt{H_\theta} \cos 2\theta + \frac{\beta^2 - \alpha^2}{\sqrt{H_\theta}}, \\ H_\theta &= p_\theta^2 + \frac{\alpha^2}{\cos^2 \theta} + \frac{\beta^2}{\sin^2 \theta},\end{aligned}$$

and

$$\{\theta, p_\theta\} = -1,$$

as we usually assume, we have that the $G^\pm = \Xi^\pm$ are solutions of

$$\{L, G\} = \pm \sqrt{-2(cL + c_0)}G,$$

and therefore of

$$\{L, \{L, G\}\} = -2(cL + c_0)G,$$

with $L = H_\theta$, $c = 8$ and $c_0 = 0$ (sec 3.2 of [5]). Up to a factor 1/2, the radial term of H in [5] coincides with the (u, p_u) -dependent term in the extensions (29), in the case $\gamma = 1/r$ ($C = 0, c = 1$) and $\omega^2 = 2\Omega$.

6 Shift and Ladder Operators for Quantum Extended Hamiltonians

Our aim is to give a general method for constructing a $(N+1)$ quantum Hamiltonian operator admitting a factorized symmetry, starting from a natural Hamiltonian L with N degrees of freedom which admits a ladder function of the form (17), i.e., such that there exists a non trivial solution of (20) and (21). As a first step, we look for shift and ladder operators similar to those given in [5]. Finally, as in [5], we use them in order to construct a symmetry operator for the extended quantum Hamiltonian. The structure of this Hamiltonian operator is strictly related with the classical extension of L , but does not exactly coincide with the Laplace-Beltrami quantization of an extension of L [14].

Definition 2. Let \hat{H}_M be an operator depending on a parameter M . We say that the M dependent operator $\hat{\mathcal{S}}_M$ is a shift operator for \hat{H}_M if there exists a function $\bar{M} = \bar{M}(M)$ such that

$$\hat{\mathcal{S}}_M \hat{H}_M = \hat{H}_{\bar{M}} \hat{\mathcal{S}}_M.$$

Let \hat{H} be an operator and the function ψ_λ be an eigenfunction of \hat{H} with eigenvalue λ . We say that $\hat{\mathcal{L}}_\lambda$ is a ladder operator for the eigenfunctions of \hat{H} if there exists a function $\bar{\lambda} = \bar{\lambda}(\lambda)$ such that

$$\hat{H}\psi_\lambda = \lambda\psi_\lambda \Rightarrow \hat{H}(\hat{\mathcal{L}}_\lambda\psi_\lambda) = \bar{\lambda}\hat{\mathcal{L}}_\lambda\psi_\lambda.$$

If $\bar{\lambda}(\lambda) = \lambda$ for all eigenvalues λ we say that $\hat{\mathcal{L}}_\lambda$ is a symmetry for \hat{H} .

We consider the Schrödinger operator associated with the natural Hamiltonian L (16),

$$\hat{L} = -\frac{\hbar^2}{2}\Delta + V. \quad (33)$$

where Δ is the Laplace-Beltrami operator associated to the metric \mathbf{g} . Let $\bar{\mathbf{g}}$ be the metric tensor of the extension of (16) which is of the form

$$\bar{g}^{ij} = \begin{pmatrix} 1 & 0 \\ 0 & \alpha(u)g^{ij} \end{pmatrix}, \quad \alpha(u) = \frac{m^2}{n^2}(c\gamma^2 + C) = -k^2\gamma', \quad (34)$$

since γ satisfies (2). We choose $\bar{\mathbf{g}}$ as metric tensor for the quantum extended Hamiltonian of (33). Therefore, the Laplace-Beltrami operator $\bar{\Delta}$ associated with the metric $\bar{\mathbf{g}}$ is

$$\bar{\Delta}\psi = \partial_u^2\psi + Nc\gamma\partial_u\psi - k^2\gamma'\Delta\psi, \quad (35)$$

where N is the dimension of the base manifold of L . We consider an extended Schrödinger operator of (33) of the form

$$\hat{H} = -\frac{\hbar^2}{2}(\partial_u^2\psi + Nc\gamma\partial_u\psi - k^2\gamma'\Delta\psi) - k^2\gamma'V + W(u). \quad (36)$$

where k is a constant and the function W will be suitably determined later. If we consider eigenfunctions ψ_λ of the operator \hat{L}

$$\hat{L}\psi_\lambda = \lambda\psi_\lambda,$$

then, for any $\phi = \phi(u)$,

$$\hat{H}(\phi\psi_\lambda) = -\frac{\hbar^2}{2} (\partial_u^2 + Nc\gamma\partial_u) (\phi)\psi_\lambda + (W - k^2\gamma'\lambda)\phi\psi_\lambda.$$

Since by (2) $\gamma' = -c\gamma^2 - C$, we may write

$$\hat{H}(\phi\psi_\lambda) = -\frac{\hbar^2}{2} (\partial_u^2 + Nc\gamma\partial_u) (\phi)\psi_\lambda + (W + ck^2\lambda\gamma^2 + Ck^2\lambda)\phi\psi_\lambda. \quad (37)$$

Hence, we have that $\hat{H}(\phi\psi_\lambda) = E\phi\psi_\lambda$ (that is $\phi\psi_\lambda$ is an eigenfunction of (36) associated with the eigenvalue E) if and only if $\phi(u)$ is an eigenfunction of the M -dependent operator

$$\hat{H}^M = -\frac{\hbar^2}{2} (\partial_u^2 + Nc\gamma(u)\partial_u) + M(c\gamma^2(u) + C) + W(u), \quad (38)$$

with $M = k^2\lambda$. We denote such a ϕ by $\phi_E^{k^2\lambda}$ in order to underline its dependence on the spectral parameters E and λ .

Definition 3. We say that \hat{X} is a *warped symmetry operator* of (36) if, for all partially separated eigenfunctions $f_E = \phi_E^{k^2\lambda}\psi_\lambda$, $\hat{H}\hat{X}(f_E) = E\hat{X}(f_E)$, or equivalently, that $[\hat{H}, \hat{X}](f_E) = 0$. We allow \hat{X} to depend on both λ and E .

In the following, we construct operators \hat{G} , \hat{A} , \hat{D} which are the factors of a warped symmetry \hat{X} of \hat{H} for $k \in \mathbb{Q}$. In particular, with Theorem 10 below, we produce a quantum version of the factorized characteristic first integrals of extended Hamiltonians of Proposition 9, valid only for the case $c \neq 0$, $C = 0$. The proof of the Theorem is given in Sec. 6.4, while in Sec. 6.1-6.3 are studied the relevant properties of the operators \hat{G} , \hat{A} and \hat{D} , respectively.

Given any function $G(q^1, \dots, q^N)$, we define the operators

$$\hat{G}_\epsilon^\pm = a_0 \nabla^i G \nabla_i + a_1^\pm G + a_2^\pm \quad (39)$$

where the constants a_i verify the relations

$$a_1^\pm = a_0 \left(\frac{c(1-N)}{2} \pm \epsilon \frac{\sqrt{2|c|}}{\hbar} \right), \quad a_2^\pm = \frac{-2c_1 a_0}{-c\hbar^2 \pm 2\hbar\epsilon\sqrt{2|c|}}, \quad (40)$$

with $\epsilon, c_1 \in \mathbb{R}$, $c \neq 0$. Moreover, for any integer $n > 0$, we construct the operators $(\hat{G}_\epsilon^\pm)^n$ as follows

$$(\hat{G}_\epsilon^\pm)^n = \hat{G}_{\epsilon\pm(n-1)s}^\pm \circ \hat{G}_{\epsilon\pm(n-2)s}^\pm \circ \cdots \circ \hat{G}_{\epsilon\pm s}^\pm \circ \hat{G}_\epsilon^\pm \quad (41)$$

where $s = \sqrt{\frac{|c|}{2}}\hbar$, and \hat{G}_ϵ^\pm is defined in (39).

We define the operators

$$\hat{A}_\mu^{\sigma, \tau} = \partial_u + \frac{1}{\hbar} \left[\sigma\omega cu + \frac{\hbar c(N-1) + \tau\sqrt{2|c|}\mu}{2cu} \right], \quad (\sigma, \tau = \pm 1), \quad (42)$$

$$\hat{D}_E^\pm = \frac{\hbar}{\sqrt{2}}u\partial_u + \frac{\hbar(N+1)}{2\sqrt{2}} \pm \left(\frac{E}{2c\omega} - \frac{\omega}{\sqrt{2}}cu^2 \right), \quad (43)$$

where $E, \mu, \omega \in \mathbb{R}$, $c \neq 0$. Moreover, for any integer $m > 0$, we define the operators

$$(\hat{A}_\mu^{\sigma,\tau})^m = \hat{A}_{\mu+\tau(m-1)s}^{\sigma,\tau} \circ \hat{A}_{\mu+\tau(m-2)s}^{\sigma,\tau} \circ \cdots \circ \hat{A}_{\mu+\tau s}^{\sigma,\tau} \circ \hat{A}_\mu^{\sigma,\tau}, \quad (44)$$

$$(\hat{D}_E^\pm)^m = \hat{D}_{E\pm 2(m-1)\delta}^\pm \circ \hat{D}_{E\pm 2(m-2)\delta}^\pm \circ \cdots \circ \hat{D}_{E\pm 2\delta}^\pm \circ \hat{D}_E^\pm, \quad (45)$$

with $s = \sqrt{\frac{|c|}{2}}\hbar$, $\delta = \hbar\omega$, $\sigma, \tau = \pm 1$.

In (41) and (45), the signs in upper and lower indices must be the same (all + or all -).

We are able to state the theorem as follows

Theorem 10. Let $L_0 = \frac{1}{2}g^{ij}p_ip_j + V_0(q^i)$ be a classical natural Hamiltonian with N degrees of freedom and G a function of (q^i) satisfying (20), (21) for $c \neq 0$ and $c_0 = 0$ i.e., generating a ladder function of the form (19) for L_0 . For any $k \in \mathbb{Q}^+$, let us construct the operator

$$\hat{X}_{\epsilon,E}^k = (\hat{G}_\epsilon^+)^{2n} \circ (\hat{A}_{k\epsilon}^{1,1})^{2m} \circ (\hat{D}_E^+)^m \quad (46)$$

where m, n are positive integers such that $k = m/n$, the operators $(\hat{G}_\epsilon^+)^{2n}$, $(\hat{A}_\mu^{\sigma,\tau})^{2m}$ (with $\mu = k\epsilon$, $\sigma = \tau = 1$) and $(\hat{D}_E^+)^m$ are defined in (41) (44) (45) and where \hat{G}_ϵ^\pm are constructed from the function G . Then, the quantum Hamiltonian

$$\hat{L}_0 = -\frac{\hbar^2}{2}\Delta + V_0, \quad (47)$$

associated with L_0 , can be extended to the quantum Hamiltonian

$$\hat{H} = -\frac{\hbar^2}{2} \left(\partial_{uu}^2 + \frac{N}{u} \partial_u \right) + \frac{k^2}{cu^2} \hat{L}_0 + \frac{\omega^2}{2} c^2 u^2 + \frac{c_N(k^2 + 1)}{cu^2}, \quad (48)$$

with

$$c_N = -\frac{\hbar^2 c(N-1)^2}{8}, \quad (49)$$

which admits the operator $\hat{X}_{\epsilon,E}^k$ as warped symmetry: i.e.,

$$\hat{H}(\hat{X}_{\epsilon,E}^k f_E) = E \hat{X}_{\epsilon,E}^k(f_E),$$

for all (partially) multiplicatively separated eigenfunctions $f_E = \phi_E^M(u)\psi_\lambda(q^i)$ of \hat{H} such that ψ_λ is an eigenfunction of \hat{L}_0 with eigenvalue λ , and

$$M = k^2(\lambda + c_N), \quad \epsilon = \sqrt{|\lambda + c_N|}, \quad c(\lambda + c_N) \geq 0.$$

Remark 7. For $N = 1$, the quantum extended Hamiltonian (48) is the quantum operator associated with the extension of L_0 with $\gamma = 1/(cu)$ (that is for $c \neq 0$, $C = 0$). However, for $N > 1$ (that is when L_0 has more than one degree of freedom) this is no more true, due to the presence of the last addendum in (48). The addendum could be absorbed in L_0 , by replacing L_0 with $L_1 = L_0 + c_N \frac{k^2+1}{k^2}$. Nevertheless, in the proof of Theorem 10 it is essential to consider the eigenvalues of \hat{L}_0 , not those of the shifted operator \hat{L}_1 . From a different view point, the additional term proportional to c_N added to the quantization of the classical extension of L_0 suggests a link with the introduction of a kind of quantum correction, used as instance in [14, 3] to allow quantization for classical quadratic in the momenta first integrals.

Remark 8. Other choices of the signs in \hat{G}_ϵ^\pm , $\hat{A}_\mu^{\sigma,\tau}$, and \hat{D}_E^\pm (as well as the commutation between $(\hat{G}_\epsilon^\pm)^{2n}$ and $(\hat{A}_{k\epsilon}^{\sigma,\tau})^{2m} \circ (\hat{D}_E^+)^m$) allow us to construct other warped symmetry operators for (48): by applying the same kind of shown given in Section 6.4, it is easy to check that the operators

$$(\hat{G}_\epsilon^-)^{2n}(\hat{A}_{k\epsilon}^{-1,-1})^{2m}(\hat{D}_E^-)^m, \quad (\hat{G}_\epsilon^-)^{2n}(\hat{A}_{k\epsilon}^{1,-1})^{2m}(\hat{D}_E^+)^m, \quad (\hat{G}_\epsilon^+)^{2n}(\hat{A}_{k\epsilon}^{-1,1})^{2m}(\hat{D}_E^-)^m.$$

commute with (75), on separated eigenfunctions.

Remark 9. Shift and ladder operators arise usually from the factorization, complete or partial, of the Hamiltonian operator. While this property is useful for computations, we did not make use of it in the construction of the symmetry operators or of the first integrals.

6.1 Ladder operators for the eigenfunctions of \hat{L}

Proposition 11. Let G be solution of (20) and (21), with $c \neq 0$ and

$$c_0 = -\frac{c^2 \hbar^2}{8}(N-1)^2. \quad (50)$$

By setting, for all λ such that $c\lambda \geq 0$,

$$\epsilon = \sqrt{|\lambda|} \quad (\text{i.e., } \lambda = \check{c}\epsilon^2 \text{ with } \check{c} = \text{sgn}(c)),$$

then, for any $a_0 \neq 0$, the operators \hat{G}_ϵ^\pm (39) are ladder operators for the eigenfunctions of the Hamiltonian operator (33): i.e., for any ψ_λ such that $\hat{L}\psi_\lambda = \lambda\psi_\lambda$

$$\hat{L}\hat{G}_\epsilon^-\psi_\lambda = \check{c}(\epsilon + s)^2 \hat{G}_\epsilon^-\psi_\lambda, \quad \hat{L}\hat{G}_\epsilon^+\psi_\lambda = \check{c}(\epsilon - s)^2 \hat{G}_\epsilon^+\psi_\lambda, \quad (51)$$

with $s = \hbar\sqrt{\frac{|c|}{2}}$.

The proof of the above proposition is based on several lemmas.

Lemma 12. Let $\hat{G} = a_0\nabla^i G \nabla_i + a_1 G + a_2$, where G is solution of (20), and a_0, a_1, a_2 are constants. We have for any function ψ

$$\begin{aligned} \Delta\hat{G}(\psi) &= a_0\nabla^i G g^{lm} \nabla_l \nabla_m \nabla_i \psi + (2a_1 - a_0c)\nabla^i G \nabla_i \psi \\ &\quad + G(a_1 - 2a_0c)\Delta\psi - a_1 c N G \psi + a_2 \Delta\psi. \end{aligned} \quad (52)$$

Proof. Direct computation, taking into account that solutions G of (20) verify (25) and its contraction with the contravariant metric tensor: $\Delta G = -cNG$. \square

Lemma 13. We have for any function ψ

$$(g^{jk} \nabla_k \nabla_j \nabla_i - \nabla_i \Delta)\psi = -R^l_i \nabla_l \psi, \quad (53)$$

where R^l_i are the components of the Ricci tensor.

Proof. It follows from the direct computation, using the definition of R^l_{ijk} as

$$R^l_{ijk} = \partial_j \Gamma^l_{ik} - \partial_k \Gamma^l_{ij} + \Gamma^m_{ik} \Gamma^l_{mj} - \Gamma^m_{ij} \Gamma^l_{mk},$$

the relation

$$\partial_i g^{jk} = -g^{jl} \Gamma^k_{il} - g^{lk} \Gamma^j_{il},$$

and the equations $R_{ri} = R^j_{rij}$, $R^l_{ijk} = -R^l_{ijk}$. \square

Lemma 14. *If G is a solution of (20) then*

$$R_{ij}\nabla^j G = c(1 - N)\nabla_i G.$$

Proof. By differentiating $\nabla_i \nabla_j G = -cg_{ij}G$ we have (see also [8])

$$c(g_{jl}\nabla_i G - g_{il}\nabla_j G) = \nabla_k GR_{lij}^k = \nabla^k GR_{klkj} = -\nabla^k GR_{lkij},$$

by contracting the first and the last terms with g^{lj} we have the statement. \square

Then, it follows

Lemma 15. *Let us consider $\hat{G} = a_0\nabla^i G \nabla_i + a_1 G + a_2$, where G is solution of (20) and a_0, a_1, a_2 are constants. For any function ψ we have*

$$\begin{aligned} \Delta \hat{G}(\psi) &= a_0 \nabla^i G \nabla_i \Delta \psi + (2a_1 - a_0 c(2 - N)) \nabla^i G \nabla_i \psi - a_1 c N G \psi \\ &\quad + G(a_1 - 2a_0 c) \Delta \psi + a_2 \Delta \psi. \end{aligned} \quad (54)$$

Under the assumptions of Lemma 15 we can easily evaluate the operator $\Delta \hat{G}$ on an eigenfunction ψ_λ of $-\frac{\hbar^2}{2}\Delta + V$ with eigenvalue λ i.e., such that

$$\Delta \psi_\lambda = \frac{2}{\hbar^2}(V - \lambda)\psi_\lambda.$$

Lemma 16. *Let be $\hat{G} = a_0\nabla^i G \nabla_i + a_1 G + a_2$, where G is solution of (20)-(21), with a_0, a_1, a_2 constants. We have*

$$\begin{aligned} \left(-\frac{\hbar^2}{2}\Delta + V\right) \hat{G}(\psi_\lambda) &= \left[\frac{\hbar^2}{2}(a_0 c(2 - N) - 2a_1) + a_0 \lambda\right] \nabla^i G \nabla_i \psi_\lambda + \\ &\quad + \left[\lambda(a_1 - 2a_0 c) + \frac{\hbar^2}{2}ca_2 N - 2a_0 c_0\right] G \psi_\lambda + (a_0 c_1 + a_2 \lambda) \psi_\lambda. \end{aligned} \quad (55)$$

We now require that $\hat{G}\psi_\lambda$ is an eigenfunction of $\hat{L} = -\frac{\hbar^2}{2}\Delta + V$. By (55), the condition is clearly equivalent to the equations

$$\begin{aligned} \left[\frac{\hbar^2}{2}(a_0 c(2 - N) - 2a_1) + a_0 \lambda\right] a_1 &= \left[\lambda(a_1 - 2a_0 c) + \frac{\hbar^2}{2}ca_1 N - 2a_0 c_0\right] a_0, \\ \left[\frac{\hbar^2}{2}(a_0 c(2 - N) - 2a_1) + a_0 \lambda\right] a_2 &= a_0^2 c_1 + a_0 a_2 \lambda. \end{aligned}$$

Therefore,

Lemma 17. *The function $\hat{G}\psi_\lambda = (a_0\nabla^i G \nabla_i + a_1 G + a_2)\psi_\lambda$ is an eigenfunction of $\hat{L} = -\frac{\hbar^2}{2}\Delta + V$ if and only if the constants a_0, a_1, a_2 satisfy*

$$2(c\lambda + c_0)a_0^2 + c\hbar^2(1 - N)a_0 a_1 - \hbar^2 a_1^2 = 0, \quad (56)$$

$$\frac{\hbar^2}{2}c(N - 2)a_0 a_2 + \hbar^2 a_1 a_2 + a_0^2 c_1 = 0. \quad (57)$$

If $a_0 = 0$, then (56) implies $a_1 = 0$, while equation (57) is satisfied for all a_2 . However in this case the operator \hat{G} reduces to multiplication by a constant (a trivial symmetry operator). Hence, we consider $a_0 \neq 0$.

For $c \neq 0$, from (56), we deduce

$$a_1 = a_0 \left[\frac{c}{2}(1-N) \pm \frac{1}{2\hbar} \sqrt{c^2\hbar^2(N-1)^2 + 8(c\lambda + c_0)} \right].$$

which is real if $c\lambda + c_0 \geq -c^2\hbar^2(N-1)^2/8$. If in (57) the coefficient of a_2 vanishes, then the solution of (56–57) is $a_0 = a_1 = 0$ already considered. Otherwise, the value for the parameter a_2 satisfying (57) is

$$a_2 = \frac{2c_1 a_0}{c\hbar^2 \mp \hbar \sqrt{c^2\hbar^2(N-1)^2 + 8(c\lambda + c_0)}}$$

and the corresponding eigenvalue for the eigenfunction $\hat{G}\psi_\lambda$ is

$$\bar{\lambda} = \lambda + \frac{\hbar^2 c}{2} \mp \frac{\hbar}{2} \sqrt{c^2\hbar^2(N-1)^2 + 8(c\lambda + c_0)}.$$

Proof of Proposition 11. In [5] the ladder operators depend on a parameter ϵ and transform the eigenvalue $\lambda = \epsilon^2$ into $\bar{\lambda} = (\epsilon \pm 2)^2$. If we want that the eigenvalue associated with $\hat{G}\psi_\lambda$ is of the form

$$\bar{\lambda} = \pm(\sqrt{|\lambda|} + s)^2,$$

for a constant s , then (after Lemmas 12–17) the value of c_0 has to be (50), the condition $c\lambda \geq 0$ has to be satisfied and the values of a_1 , a_2 and $\bar{\lambda}$ are those appearing in (40) and (51), respectively. \square

Remark 10. It is easy to check what happens to \hat{G}_ϵ^\pm in the particular case $N = 1$, $V = c_0 = c_1 = 0$: in accordance with the first equation (9) and relations (10) in [5], for $c = 4$ and $G = \cos 2\theta$, after setting $\hbar^2 = 2$, we obtain the solution $a_1 = \pm 2\epsilon a_0$, $a_2 = 0$ with the eigenvalue $\bar{\lambda} = (\epsilon \pm 2)^2$.

Remark 11. By choosing

$$a_0 = \hbar \sqrt{\frac{|c|}{2}} \left(\hbar \sqrt{\frac{|c|}{2}} \mp 2\epsilon \right),$$

the operators \hat{G}_ϵ^\pm become polynomials in ϵ and the relations (40) reduce to

$$a_1 = \frac{c}{2} \left(\hbar \sqrt{\frac{|c|}{2}} \mp 2\epsilon \right) \left(\hbar(1-N) \sqrt{\frac{|c|}{2}} \pm 2\epsilon \right), \quad a_2 = c_1.$$

Remark 12. For $c = 0$ equations (56) and (57) reduce to

$$a_1^2 = \frac{2}{\hbar^2} c_0 a_0^2, \quad a_0^2 c_1 = -\hbar^2 a_1 a_2.$$

If $c_0 > 0$ we have non trivial solutions for $a_0 \neq 0$: $a_1 = \pm \frac{\sqrt{2c_0}}{\hbar} a_0$, $a_2 = \mp \frac{c_1}{\hbar \sqrt{2c_0}} a_0$. Moreover, the eigenvalue of $\hat{G}\psi_\lambda$ is $\bar{\lambda} = \lambda - \hbar^2 a_1/a_0 = \lambda \mp \hbar \sqrt{2c_0}$. This ladder operator is then independent of λ and does not seem useful to generalize the approach used in [5].

Since the eigenvalue of $\hat{G}_\epsilon\psi_\lambda$ is of the form $(\epsilon + s)^2$ we have that

Corollary 18. *The operators (41) map the eigenfunctions of \hat{L} associated with the eigenvalue $\lambda = \check{c}\epsilon^2$ to eigenfunctions of \hat{L} with eigenvalue $\bar{\lambda} = \check{c}(\epsilon + ns)^2$ for any integer $n > 0$. Moreover, for the choice of a_0 proposed in Remark 11 these operators are polynomials in ϵ whose coefficients can be easily determined explicitly.*

Proposition 11 provides a method for constructing ladder operators \hat{G}_ϵ starting from the spectrum of \hat{L} by setting $\epsilon = \sqrt{|\lambda|}$ for all eigenvalues λ of \hat{L} with suitable sign (recall that we need that $c\lambda \geq 0$). In the following proposition, we show how to avoid the condition (50). By Remark 5, we know that if a solution G of (20)-(21) is known for $c \neq 0$ and arbitrary c_0 , then the same G satisfies the equations with the same c and $c_0 = c_N c$ with c_N given as in (49), if we add the constant term $c_N - c_0/c$ to V . Indeed, we have

Corollary 19. *Let \hat{L}_0 defined in (47) be a Schrödinger operator such that equations (20)-(21) admit a solution G for $c \neq 0$ and $c_0 = 0$. Then, the operators \hat{G}_ϵ^\pm defined in (39) are ladder operators for the eigenfunctions of*

$$\hat{L}_N = -\frac{\hbar^2}{2}\Delta + V_0 - \frac{\hbar^2 c(N-1)^2}{8}. \quad (58)$$

Proof. The potential V of (58) is of the form $V_0 + c_N$. Hence, \hat{L}_N verifies the hypotheses of Proposition 11. \square

Since eigenfunctions of (47) and (58) are the same and

$$\hat{L}_0\psi_\lambda = \lambda\psi_\lambda \Rightarrow \hat{L}_N\psi_\lambda = (\lambda + c_N)\psi_\lambda$$

that is the eigenvalues are shifted by the constant c_N , we try to construct ladder operators for \hat{L}_0 :

Proposition 20. *Let G be solution of (20) and (21) for $c \neq 0$ and for $c_0 = 0$. By setting, for all λ such that $c(\lambda + c_N) > 0$,*

$$\epsilon = \sqrt{|\lambda + c_N|} = \sqrt{\left|\lambda - \frac{\hbar^2 c(N-1)^2}{8}\right|}, \quad (\text{i.e., } \lambda = \check{c}\epsilon^2 - c_N), \quad (59)$$

then, the operators \hat{G}_ϵ^\pm (39) are ladder operators for the eigenfunctions of the Hamiltonian operator (47) i.e., for any ψ_λ such that $\hat{L}_0\psi_\lambda = \lambda\psi_\lambda$, we have

$$\hat{L}_0\hat{G}_\epsilon^- \psi_\lambda = (\check{c}(\epsilon + s)^2 - c_N) \hat{G}_\epsilon^- \psi_\lambda, \quad (60)$$

$$\hat{L}_0\hat{G}_\epsilon^+ \psi_\lambda = (\check{c}(\epsilon - s)^2 - c_N) \hat{G}_\epsilon^+ \psi_\lambda, \quad (61)$$

with $s = \sqrt{\frac{|c|}{2}}\hbar$.

Proof. For any ψ_λ such that $\hat{L}_0\psi_\lambda = \lambda\psi_\lambda$ we have $\hat{L}_N\psi_\lambda = \lambda_N\psi_\lambda$ with $\lambda_N = \lambda + c_N$ and \hat{L}_N as in (58). Hence, since $c(\lambda + c_N) > 0$ means $c\lambda_N > 0$, \hat{L}_N satisfies the hypotheses of Proposition 11, then the operators \hat{G}_ϵ^\pm defined by (40) for $\epsilon = \sqrt{|\lambda_N|}$ verify (51), where \hat{L} is given by (58). Thus, operator (47) verifies (60) and the relation between ϵ and λ is given by (59). \square

Corollary 21. *The operators $(\hat{G}_\epsilon^\pm)^n$ defined in (41) are ladder operators for \hat{L}_0 : that is, for any ψ_λ such that $\hat{L}_0\psi_\lambda = \lambda\psi_\lambda$, with $c(\lambda + c_N) > 0$, and for ϵ as in (59), we have*

$$\hat{L}_0(\hat{G}_\epsilon^-)^n\psi_\lambda = (\check{c}(\epsilon + ns)^2 - c_N)(\hat{G}_\epsilon^-)^n\psi_\lambda, \quad (62)$$

$$\hat{L}_0(\hat{G}_\epsilon^+)^n\psi_\lambda = (\check{c}(\epsilon - ns)^2 - c_N)(\hat{G}_\epsilon^+)^n\psi_\lambda. \quad (63)$$

6.2 Quantum extended Hamiltonians \hat{H} with Shift-ladder operators

As in [17] we look for Shift-Ladder operators of the M -dependent radial operator (38), that is we search for operators that modify the value of M as well as the energy (they add a constant multiple of the identity to \hat{H}^M). We use an ansatz on the form of $W(u)$ and of $\gamma(u)$, which allows us to exactly generalize the shift-ladder operators of the radial operator analysed in [17] for $N = 1$ to an arbitrary N . A study of a more general W and γ is not yet complete.

Proposition 22. *If the operators \hat{H}^M (38) verify $c \neq 0$, $C = 0$ (that is $\gamma(u) = \frac{1}{cu}$) and*

$$W(u) = \frac{\omega^2}{2}c^2u^2 - \frac{\hbar^2(N-1)^2}{8u^2} = \frac{\omega^2}{2}c^2u^2 - \frac{c_N}{cu^2}, \quad \omega \in \mathbb{R}, \quad (64)$$

then, for all M such that $cM > 0$, by setting

$$\mu = \sqrt{|M|},$$

the operators (42) satisfy

$$\hat{H}^{\bar{M}}\hat{A}_\mu^{\sigma,\tau} = \hat{A}_\mu^{\sigma,\tau}(\hat{H}^M - \sigma\hbar c\omega),$$

with

$$\bar{M} = \check{c}(\mu - \tau s)^2 = \check{c}\left(\mu - \tau\hbar\sqrt{\frac{|c|}{2}}\right)^2. \quad (65)$$

Proof. We assume that the operators \hat{A} are of the form

$$\hat{A} = \partial_u + \left(a\gamma(u) + b\frac{1}{\gamma(u)}\right), \quad a, b \in \mathbb{R},$$

then we find that a necessary condition for having

$$\hat{H}^{\bar{M}}\hat{A} = \hat{A}(\hat{H}^M - \delta), \quad \delta \in \mathbb{R}, \quad (66)$$

is $A = 0$, that is $\gamma(u) = cu$, and

$$\hat{A} = \partial_u + \frac{1}{\hbar^2} \left[\delta u + \frac{1}{2cu}(\hbar^2 cN - 2(\bar{M} - M)) \right].$$

Moreover, for $W(u)$ of the form

$$W = \frac{\omega^2}{2}\gamma^{-2} + \beta\gamma^2, \quad \omega, \beta \in \mathbb{R}$$

the following conditions are sufficient in order to get the identity (66):

$$\delta^2 = \hbar^2 c^2 \omega^2,$$

(i.e., $\delta = \pm \hbar c \omega$) and

$$\bar{M}^2 - \bar{M}(2M + \hbar^2 c^2) + \frac{(2M - \hbar^2 c^2)^2 - \hbar^4 c^4 (N-1)^2 - 8\hbar^2 c^2 \beta}{4} = 0$$

that is

$$\bar{M} = M + \frac{\hbar^2 c}{2} \pm \frac{\hbar}{2} \sqrt{\hbar^2 c^2 (N-1)^2 + 8\beta + 8cM}.$$

For $cM > 0$, by setting $\mu = \sqrt{|M|}$, if

$$\beta = -\frac{\hbar^2 c^2 (N-1)^2}{8}. \quad (67)$$

then we have (65). Furthermore, W becomes (64) and we get the operators $\hat{A}_\mu^{\sigma,\tau}$ of the statement. \square

Remark 13. Operators $\hat{A}_\mu^{\sigma,\tau}$ are a mixture of shift and ladder operators, as in [5], since both the parameter M and the energy of the Hamiltonian \hat{H}^M are changed by them. Since we are not requiring that \hat{H}^M can be written (up to a multiple of the identity) as a composition of $\hat{A}_\mu^{\sigma,\tau}$, we have more freedom in their form than in [5] (e.g. the $\hat{A}_\mu^{\sigma,\tau}$ are defined up to a constant factor).

Remark 14. The value of β in (67) coincides with the value of c_0 given by (50). Moreover, the shift s of M^2 coincides with the shift of ϵ^2 produced by \hat{G}_ϵ .

Corollary 23. If the radial operator \hat{H}^M defined in (38) is

$$\hat{H}^M = -\frac{\hbar^2}{2} \left(\partial_u^2 + \frac{N}{u} \partial_u \right) + \frac{8M - \hbar^2 c(N-1)^2}{8cu^2} + \frac{\omega^2}{2} c^2 u^2, \quad M = k^2 \lambda. \quad (68)$$

then, for all M such that $cM > 0$, by setting $\mu = \sqrt{|M|}$, $M = \check{c}\mu^2$, for any positive integer m , the operators $(\hat{A}_\mu^{\sigma,\tau})^m$ (44) verify

$$(\hat{A}_\mu^{\sigma,\tau})^m (\hat{H}^M - \sigma m \delta) = \hat{H}^{\check{c}(\mu - \tau ms)^2} (\hat{A}_\mu^{\sigma,\tau})^m$$

with $s = \sqrt{\frac{|c|}{2}} \hbar$, $\delta = \hbar c \omega$, $\sigma, \tau = \pm 1$.

Remark 15. A direct attempt to reabsorb the constant c_N in M , as we did for the ladder operators \hat{G} , does not seem to work. Indeed, we define the operator

$$\hat{H}^0 = -\frac{\hbar^2}{2} \left(\partial_{uu}^2 + \frac{N}{u} \partial_u \right) + \frac{\omega^2}{2} c^2 u^2. \quad (69)$$

Let \hat{H}_0^ν be

$$\hat{H}_0^\nu = \hat{H}_0 + \frac{\nu}{cu^2}, \quad (70)$$

then, we have the relations

$$\hat{H}^M = \hat{H}_0^{M+c_N}, \quad \hat{H}_0^\nu = \hat{H}^{\nu-c_N},$$

with c_N as in (49). We look for shift-ladder operator of \hat{H}_0^ν . By setting

$$\mu = \sqrt{|\nu - c_N|},$$

and considering the operators defined in Proposition 22, we get

$$\hat{A}_\mu^{\sigma,\tau} \hat{H}_0^\nu = (\hat{H}_0^{\bar{\nu}} + \sigma\delta) \hat{A}_\mu^{\sigma,\tau},$$

where $\bar{\nu} = \check{c}(\sqrt{|\nu - c_N|} - \tau s)^2 + c_N$, $s = \sqrt{\frac{|c|}{2}}\hbar$, $\delta = \hbar c\omega$, and $\sigma, \tau = \pm 1$. The form of $\bar{\nu}$ allows to apply also the operator $(\hat{A}_\mu^{\sigma,\tau})^m$, which for μ as above gives

$$(\hat{A}_\mu^{\sigma,\tau})^m \hat{H}_0^\nu = (\hat{H}_0^{\bar{\nu}} + m\sigma\delta)(\hat{A}_\mu^{\sigma,\tau})^m, \quad (71)$$

with $\bar{\nu} = \check{c}(\mu - m\tau s)^2 + c_N$. For any function $\phi(u)$ and for any eigenfunction ψ_λ of \hat{L}_0 (47) the operator

$$\hat{H}_1 = \hat{H}_0 - \frac{k^2}{cu^2} \hat{L}_0$$

acts on $(\hat{G}_\epsilon^-)^{2n}(\hat{A}_\mu^{\sigma,\tau})^{2m}(\phi\psi_\lambda)$ (where ϵ is given by (59), and μ has to be determined) as

$$\hat{H}_1((\hat{G}_\epsilon^-)^{2n} \circ (\hat{A}_\mu^{\sigma,\tau})^{2m}(\phi\psi_\lambda)) = \hat{H}_1[(\hat{G}_\epsilon^-)^{2n}(\psi_\lambda)(\hat{A}_\mu^{\sigma,\tau})^{2m}(\phi)]. \quad (72)$$

Since $(\hat{G}_\epsilon^-)^{2n}(\psi_\lambda)$ is an eigenfunction of \hat{L}_0 associated with the eigenvalue $\bar{\lambda} = \check{c}(\epsilon + 2ns)^2 - c_N$, we get that (72) equals

$$\hat{H}_0^{\bar{\nu}}[(\hat{A}_\mu^{\sigma,\tau})^{2m}(\phi)](\hat{G}_\epsilon^-)^{2n}(\psi_\lambda) = \hat{H}^{\bar{\nu}-c_N}[(\hat{A}_\mu^{\sigma,\tau})^{2m}(\phi)](\hat{G}_\epsilon^-)^{2n}(\psi_\lambda)$$

with $\bar{\nu} = \check{c}(k\epsilon + 2ms)^2 - k^2c_N$. But it does not seem possible to go further and apply (71) for any μ if $c_N \neq 0$.

6.3 Shift operators for the eigenfunctions of \hat{H}

Let $\phi_E^M(u)$ be an eigenfunction of (68) with eigenvalue E : $\hat{H}^M(\phi_E^M) = E\phi_E^M$. We look for shift operators for these eigenfunctions.

Proposition 24. *The operators (43) satisfy*

$$\hat{H}^M \hat{D}_E^\pm \phi_E^M = (E \pm 2\hbar c\omega) \hat{D}_E^\pm \phi_E^M, \quad (73)$$

i.e., they map eigenfunctions of \hat{H}^M to eigenfunction of the same operator with a different eigenvalue:

$$\hat{D}_E^+ : \phi_E^M \mapsto \phi_{E+2\hbar c\omega}^M, \quad \hat{D}_E^- : \phi_E^M \mapsto \phi_{E-2\hbar c\omega}^M,$$

Proof. By using the ansatz

$$\hat{D}_E \phi = a_0 cu \partial_u \phi + (a_1 + a_2 c^2 u^2) \phi$$

($a_i \in \mathbb{R}$, $i = 0, 1, 2$) and imposing the condition that ϕ has to be an eigenfunction of \hat{H}^M , we compute $\hat{H}^M(\hat{D}_E(\phi_E^M))$ and we determine the values of $a_1 = \mp \frac{c}{\hbar} \omega a_0$, $a_2 = (c \frac{N+1}{2} \pm \frac{E}{\hbar\omega}) a_0$ and the shift of the eigenvalue (73). For $a_0 = \hbar/(\sqrt{2}c)$ we obtain (43). \square

Up to a constant factor, the relations (73) coincide with the analogous ones in [5] for $\hbar^2 = 2$.

The expressions of the operators \hat{D}_E^\pm and the shift of the eigenvalue do not depend on M . Hence, the same operators are ladder operators for the eigenfunctions of the operator (70).

The operators (45) verify the shift property: if $\hat{H}^M \phi_E^M = E \phi_E^M$, then

$$\hat{H}^M (\hat{D}_E^+)^m \phi_E^M = (E + 2m\delta) (\hat{D}_E^+)^m \phi_E^M, \quad \hat{H}^M (\hat{D}_E^-)^m \phi_E^M = (E - 2m\delta) (\hat{D}_E^-)^m \phi_E^M, \quad (74)$$

with $\delta = \hbar c\omega$.

6.4 Proof of Theorem 10

Proof of Theorem 10. Let \hat{L}_0 be the Schrödinger operator defined in (47) such that (20) and (21) admit a non trivial solution G with $c \neq 0$ and $c_0 = 0$. Let \hat{H}^0 be the operator defined in (69), then we can write the operator (48) as

$$\hat{H} = \hat{H}^0 + \frac{k^2}{cu^2} \hat{L}_0 + (k^2 + 1) \frac{c_N}{cu^2}, \quad (75)$$

with $k \in \mathbb{Q}^+$ and c_N given by (49). Moreover, we can write the M -dependent operator (68) as

$$\hat{H}^M = \hat{H}^0 + \frac{M + c_N}{cu^2}.$$

Let ψ_λ be an eigenfunction of \hat{L}_0 associated with an eigenvalue λ such that $c(\lambda + c_N) > 0$. Then, the function $\phi(u)\psi_\lambda$ is an eigenfunction of 75 with eigenvalue E , if and only if $\phi_E(u) = \phi_E^M$ with $\hat{H}^M(\phi_E^M) = E\phi_E^M$ for

$$M = k^2(\lambda + c_N).$$

Obviously, $cM > 0$ if and only if $c(\lambda + c_N) > 0$. We set ϵ as in (59), so that

$$\sqrt{|M|} = k\epsilon,$$

or, equivalently,

$$M = \check{c}k^2\epsilon^2.$$

According to Definition 3, we prove that

$$\hat{H}(\hat{X}_{\epsilon,E}^k(\phi_E^M \psi_\lambda)) = E \hat{X}_{\epsilon,E}^k(\phi_E^M \psi_\lambda),$$

for all multiplicatively separated eigenfunctions $\phi_E^M \psi_\lambda$ of \hat{H} such that $c(\lambda + c_N) > 0$. We have the splitting

$$\hat{X}_{\epsilon,E}^k(\phi_E^M \psi_\lambda) = (\hat{A}_{k\epsilon}^{1,1})^{2m} \circ (\hat{D}_E^+)^m(\phi_E^M) \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda).$$

Hence, by applying Propositions 11, 22 and 24, we have

$$\begin{aligned}
& \hat{H}(\hat{X}_{\epsilon,E}^k(\phi_E^M \psi_\lambda)) = \\
&= \hat{H}^0(\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^M) \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) + \\
&\quad + (\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^M) \cdot \frac{k^2}{cu^2} \hat{L}_0(\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) + \\
&\quad + (\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^M) \cdot \frac{(k^2+1)c_N}{cu^2} (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= \hat{H}^0(\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^{k\epsilon}) \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) + \\
&\quad + (\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^M) \cdot \frac{\check{c}k^2(\epsilon+2ns)^2 + c_N}{cu^2} (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= [\hat{H}^{\check{c}k^2(\epsilon-2ns)^2}(\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^M)] \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= [\hat{H}^{\check{c}(k\epsilon-2ms)^2}(\hat{A}_{k\epsilon}^{1,1})^{2m}(\hat{D}_E^+)^m(\phi_E^M)] \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= [(\hat{A}_{k\epsilon}^{1,1})^{2m} \circ (\hat{H}^M - 2m\delta) \circ (\hat{D}_E^+)^m(\phi_E^M)] \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= (\hat{A}_{k\epsilon}^{1,1})^{2m}[(E + 2m\delta)(\hat{D}_E^+)^m(\phi_E^M) - 2m\delta(\hat{D}_E^+)^m(\phi_E^M)] \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= E(\hat{A}_{k\epsilon}^{1,1})^{2m} \circ (\hat{D}_E^+)^m(\phi_E^M) \cdot (\hat{G}_\epsilon^+)^{2n}(\psi_\lambda) = \\
&= E\hat{X}(\phi_E^M \psi_\lambda).
\end{aligned}$$

This proves the statement. \square

7 Conclusions

We analysed the Kuru-Negro ansatz about the existence of ladder functions in the context of extended Hamiltonian systems. We showed that the ansatz fits naturally in this theory, leading always to a factorized form of the characteristic first integrals of extended systems. Moreover, we proved that for a large subclass of extended Hamiltonians the ansatz is not restrictive.

In the quantum context, instead of a deduction of the classical ladder functions and factorized first integrals from the quantum factorized symmetry operators constructed through shift and ladder operators, by a sort of semiclassical limit procedure, we worked in the opposite direction: from the properties of the classical Hamiltonian L we constructed a quantum operator \hat{H} admitting a factorized symmetry operator. Indeed, we have considered the quantum Hamiltonian associated with a classical Hamiltonian L admitting a ladder function. We have explicitly constructed an extended quantum Hamiltonian and four factorized symmetry operators, by using shift and ladder operators, as in the Kuru-Negro theory. The quantum extension is a Schrödinger operator, whose classical counterpart is similar to a particular case (case $c \neq 0, C = 0$) of an extension of L . However, the two Hamiltonians coincide only if L has one degree of freedom.

These results show the intimate connection between the theory of extended Hamiltonians and the procedure of factorization in shift-ladder operators. The connection reveals itself in essentially two properties: the possibility of a partial separation of variables (u, q^i) and the existence of solutions of the equation $X_L^2 G = -2(cL + c_0)G$. Both the properties are consequences of the structure of warped manifold: the first one, of the configuration manifold of H , the second

one, of the configuration manifold of L [23]. The relations between factorization and warped manifolds owe to be studied in greater detail. The opportunity of several further studies emerges clearly:

- (i) the construction of factorized symmetry operators should be carried on also for the general extended metric tensor (case $C \neq 0, c \neq 0$);
- (ii) the shift and ladder operators are here built on the model of [5], but different, more general, forms of them should be investigated: in particular, a quantization based on the factorization results of Theorem 6 should be studied;
- (iii) it is an open problem if the ansatz is actually a restrictive condition for the more general $G(q^i, p_j)$.

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