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This is the author's manuscript

Original Citation:

Availability:

This version is available <http://hdl.handle.net/2318/1654220> since 2018-10-18T15:04:09Z

Published version:

DOI:10.1016/j.jvolgeores.2017.09.003

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(Article begins on next page)

1 **Measuring effusion rates of obsidian lava flows by means of satellite thermal**
2 **data**

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26 **Abstract**

27

28 Space-based thermal data are increasingly used for monitoring effusive eruptions, especially for
29 calculating lava discharge rates and forecasting hazards related to basaltic lava flows. The
30 application of this methodology to silicic, more viscous lava bodies (such as obsidian lava flows) is
31 much less frequent, with only few examples documented in the last decades. The 2011-2012
32 eruption of Cordón Caulle volcano (Chile) produced a voluminous obsidian lava flow ($\sim 0.6 \text{ km}^3$)
33 and offers an exceptional opportunity to analyze the relationship between heat and volumetric flux
34 for such type of viscous lava bodies. Based on a retrospective analysis of MODIS infrared data
35 (MIROVA system), we found that the energy radiated by the active lava flow is robustly correlated
36 with the erupted lava volume, measured independently. We found that after a transient time of
37 about 15 days, the coefficient of proportionality between radiant and volumetric flux becomes
38 almost steady, and stabilizes around a value of $\sim 5 \times 10^6 \text{ J m}^{-3}$. This coefficient (i.e. radiant density)
39 is much lower than those found for basalts ($\sim 1 \times 10^8 \text{ J m}^{-3}$) and likely reflects the appropriate
40 spreading and cooling properties of the highly-insulated, viscous flows. The effusion rates trend
41 inferred from MODIS data correlates well with the tremor amplitude and with the plume elevation
42 recorded throughout the eruption, thus suggesting a link between the effusive and the coeval
43 explosive activity. Modelling of the eruptive trend indicates that the Cordón Caulle eruption
44 occurred in two stages, either incompletely draining a single magma reservoir or more probably
45 tapping multiple interconnected magmatic compartments.

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47

48 **Keywords:** effusion rates, obsidian lava flow, radiant power, Puyehue-Cordón Caulle, satellite
49 thermal remote sensing

50

51

1. Introduction

The rate at which magma is erupted is a key parameter for understanding and modelling volcanic eruptions. When the magma is effused or extruded, the discharge rate that characterizes a given eruption reveals the pressure changes inside the magma chamber, and its modelling may constrain the location and capacity of magma storage zones (Wadge, 1981; Stasiuk et al., 1993; Melnik and Sparks, 1999). Lava discharge rates are essential for evaluating eruption dynamics (e.g. Harris et al., 2000), and represent one of the key parameter necessary to forecast lava flow paths and evaluate the associated hazards (e.g. Ganci et al., 2012; Harris et al., 2016).

During the past thirty years, several works focused on estimating lava discharge rates by using satellite thermal data (Harris, 2013 and reference therein). This approach, hereby called “thermal proxy”, is essentially based on the relationships between heat and volumetric fluxes of active lava bodies (e.g. Pieri and Baloga, 1986; Harris et al., 1998; Wright et al., 2001; Harris et al., 2007; Dragoni and Tallarico, 2009; Harris and Baloga, 2009; Garel et al., 2012; Coppola et al., 2013; Harris et al., 2016). Notably, most of the literature has been focused on estimation of the effusion rates at basaltic volcanoes, such as Bardarbunga-Holuhraun (Coppola et al., 2017), Etna (Harris et al., 1998, 2011; Harris and Neri, 2002; Gouhier et al., 2012; Ganci et al., 2012), Kilauea (Koeppen et al., 2013), Hekla (Harris et al., 2000), Stromboli (Calvari et al., 2005, 2010; Valade et al., 2016; Zakšek et al., 2015), Piton de la Fournaise (Coppola et al., 2009, 2017), Nyamulagira (Coppola and Cigolini, 2013; Coppola et al., 2016), Ambrym (Coppola et al., 2016), and Okmok (Patrick et al., 2003). In contrast, the number of studies drastically drops when considering viscous lavas bodies such as silicic flows (Harris et al., 2002, 2004) and domes (Harris and Ripepe, 2007; van Manen et al., 2010; Coppola et al., 2016). Studies are limited by a smaller number of eruptions characterized by felsic domes-flows emplacement, with respect to basaltic lava flows (Wright, 2016), but also by the complex relationships between eruption rate, heat balance, morphology and rheology that characterize the emplacement of viscous lava (e.g. Fink et al., 1998; Griffith, 2000; Harris and

77 [Baloga, 2009](#)). The reliability of the thermal approach as a universal method to estimate effusion
78 rates over a broad spectrum of lava bodies, is still matter of debate (i.e. [Dragoni and Tallarico,](#)
79 [2009](#); [Garel et al., 2012](#)). For example, [Harris and Baloga \(2009\)](#) stressed that the relationships
80 between effusion rates, flow planar areas and radiant flux will vary between thermal, rheological,
81 compositional and ambient (e.g. slope and flow bed roughness) conditions, so that a relationship
82 developed for basaltic lavas cannot be directly applied to andesitic lavas or other higher in silica
83 content. Moreover, recent laboratory and analytic models suggest that the relationship between
84 radiated power and effusion rate becomes valid (i.e. stationary) only after a transient time, in which
85 the lava flow reaches a thermal equilibrium ([Garel et al., 2012, 2014](#)). While for basaltic lava flow a
86 transient time of hours to days is now in now well constrained from theory and observations, ([Garel](#)
87 [et al., 2012, 2014](#), [Coppola et al., 2013](#)), for silicic lava domes there is still a lack of measurements,
88 with thermal modelling suggesting transient time of several years ([Garel et al., 2012](#)). The 2011-
89 2012 rhyodacitic eruption of Cordón Caulle (CC) provides an exceptional training opportunity to
90 test the thermal proxy over a voluminous ($\sim 0,6 \text{ km}^3$), long-lasting (~ 1 year) obsidian lava flow
91 ([Bertin et al., 2015](#)).

92 In this paper, we used MODIS (Moderate Resolution Imaging Spectroradiometer) infrared data,
93 automatically processed by the MIROVA (Middle Infrared Observation of Volcanic Activity)
94 system ([Coppola et al., 2016](#)), to analyze and quantify the thermal output related to the Cordón
95 Caulle eruption. Hence, we assess the reliability of the thermal proxy over silicic flows, by
96 comparing the radiant energy emitted by the obsidian CC lava flow with independent and
97 systematic measurements of lava flow volumes, derived from satellite-based topographic mapping
98 ([Bertin et al., 2015](#)). The comparison of satellite-based effusive trend with other geophysical
99 parameters is finally used to interpret the effusive process of CC eruption, in terms of magma
100 discharge models.

101

102 **2. Geological setting and chronology of the 2011-2012 Cordón Caulle eruption**

103 *2.1 Cordón Caulle Volcanic Complex*

104 The Cordón Caulle Volcanic Complex (CCVC) is a 15 km NW-SE elongated corridor of eruptive
105 centres located in the Southern Volcanic Zone (SVZ) of the Andes. This complex (Fig. 1a) is
106 formed by the Cordón Caulle fissure system (CC), which connects the Pleistocene Cordillera
107 Nevada caldera, at the NW tip, with the Puyehue stratovolcano, on SE (Lara et al., 2006a; Singer et
108 al., 2008; Lara and Moreno, 2006). Tectonic setting of CCVC is characterised by the
109 superimposition of the Quaternary tectonic regime (see Cembrano and Lara, 2009 for a review)
110 over a pre-Andean NW striking structure (Lara et al., 2004). This results in complex interactions
111 between the pathways of magmatic ascent system and the structural setting, especially along the
112 Cordón Caulle fissure (Lara et al., 2004, 2006a, 2006b). Holocene eruptions evacuated rhyodacitic
113 to rhyolitic magmas mostly from Cordón Caulle, whereas basaltic to andesitic lavas were erupted
114 exclusively from Puyehue stratovolcano (Lara et al., 2004; Singer et al., 2008). In the latter century,
115 Cordón Caulle showed a remarkable explosive and effusive activity, with the 1921-1922, the 1960
116 and the 2011-2012 eruptions characterised by the emission of large volumes ($> 0.5 \text{ km}^3$ of tephra,
117 comparable to lava volume) of silicic materials (up to 71 wt% in SiO_2 ; Castro et al., 2013) (cf.
118 Singer et al., 2008; Jay et al., 2014). Earthquake-volcano mechanisms may be responsible for the
119 triggering of the latter eruptions due to (i) the occurrence of high-magnitude, subduction-related
120 seismic events prior the eruptive phases (Lara et al., 2004), or (ii) intra-arc tectonics (Wendt et al.,
121 2017).

122 **Figure 1**

123

124 *2.2 Chronology of the 2011 – 2012 eruption*

125 The eruption of the Cordón Caulle began on 4 June 2011, following two months of increasing
126 seismic activity below the CCVC (Silva-Parejas et al., 2012; Bertin et al., 2015; Elissondo et al.,
127 2016). The first explosive stage was characterised by vigorous pyroclastic and gas-vent activity,
128 with eruptive column reaching ~14 km during the first hours of the eruption (Castro et al., 2013).
129 The ash plume rapidly reached the Atlantic coast affecting Buenos Aires and several Argentinean
130 provinces (Collini et al., 2013, Pistolesi et al., 2015). During the 27 h climax phase, ~ 0.25 km³ of
131 rhyodacitic tephra was ejected, releasing about 0.2 Mt of sulphur dioxide (Silva-Parejas et al., 2012;
132 Theys et al., 2013; Farquharson et al., 2015; Jay et al., 2014). Pulses of major explosive activity
133 continued until the 15 June, with a mass flow rate constantly above 10⁶ kg s⁻¹. However, a general
134 decrease in the height of pyroclastic columns from 12 km to 8 km was observed in the following
135 days (Bonadonna et al., 2015). On 15 June 2011, the extrusion of a viscous lava body began from
136 the same vent. The effusive phase was characterised by an initial (~10 days) discharge rate up to 70
137 m³ s⁻¹ (Castro et al., 2013; Bertin et al., 2015), generating an extensive compound flow with
138 structural and textural features typical of obsidian flows (Tuffen et al., 2013). The emplacement was
139 characterised by significant flow inflation and the formation of several breakout lobes that gradually
140 enlarged the flow field. Lava flow facies comprised rubbly, 'A'ā-like lava surface (30–45 m thick)
141 as well as large coherent slabs, spines and tongues of lava (20 m thick), that were essentially
142 localised along *en-échelon* tensional fractures and breakouts (Tuffen et al., 2013; Farquharson et al.,
143 2015). Apparent viscosities of ~10¹⁰ Pa s were estimated on the basis of the eruption temperature
144 (900°C; Castro et al., 2013) and the average advance rate of breakout lobes (1-3 m day⁻¹;
145 Farquharson et al., 2015). Fourteen months after the onset of eruption (August 2012), the lava field
146 covered an area of ~ 7 km² reaching a total volume of ~ 0.6 km³ (Jay et al., 2014; Bertin et al.,
147 2015). Mild explosive activity, characterized by mixed gas and ash jetting punctuated by Vulcanian
148 blasts, accompanied the entire effusive phase (Shipper et al., 2013). Lava transport and lateral
149 spreading was observed for several months after the termination of the eruption due to an efficient
150 thermal insulation of the inner core (Tuffen et al., 2013; Farquharson et al., 2015). A rapid post-

151 eruptive re-inflation was reported by [Delgado et al. \(2016\)](#) for the period 2012-2015 based on
152 ground deformation inferred from InSAR.

153 3. Methods

154 3.1 MIROVA system

155 MIROVA (Middle Infrared Observation of Volcanic Activity; www.mirovaweb.it) is an automated
156 global hot spot detection system ([Coppola et al. 2016](#)) based on near-real time ingestion of
157 Moderate Resolution Imaging Spectroradiometer (MODIS) infrared data. Two MODIS instruments
158 carried on NASA's Terra (EOS-AM) and Aqua (EOS-PM) satellites provide daily global
159 observation at 16 infrared spectral channels. These instruments deliver approximately 4 images per
160 day, with a nominal spatial resolution of ~ 1 km at nadir. The MIROVA system completes
161 automatic detection, location and quantification of high-temperature ($> 200^{\circ}\text{C}$) thermal anomalies
162 related to volcanic activity, within 1 to 4 hours of each satellite overpass (Fig. 2a). MIROVA
163 implements a hybrid algorithm using spectral and spatial principles to identify pixels contaminated
164 by thermally anomalous features (see [Coppola et al., 2016](#) for details). When pixel(s) are detected,
165 MIROVA automatically calculates the total "above background" Middle Infrared Radiance
166 (ΔL_{MIR}):

$$167 \quad \Delta L_{\text{MIR}} = \sum_{i=1}^{i=n} L_{\text{MIR},i} - L_{\text{bk},i} \quad [\text{equation 1}]$$

168 where L_{MIR} is the MIR radiance of the active pixel i (MODIS bands 21 or 22 centered at $3.959 \mu\text{m}$),
169 L_{bk} is the background MIR radiance (obtained from not contaminated, adjacent pixels) and n is the
170 number of detected pixels. The "above background" MIR radiance thus provides a bulk
171 measurement of the spectral radiance ($\text{W m}^{-2} \text{sr}^{-1} \mu\text{m}^{-1}$) emitted by the hot surface(s) in the $4 \mu\text{m}$
172 region of the electromagnetic spectrum (Fig. 2b). This last is then converted into Volcanic Radiative
173 Power (VRP, in Watt) by using the "MIR method" ([Wooster et al., 2003](#)):

$$174 \quad \text{VRP} = 18.9 \times A_{\text{PIX}} \times \Delta L_{\text{MIR}} \quad [\text{equation 2}]$$

175 where A_{PIX} is the pixel's area (1 km² for MODIS), and 18.9 representing the constant
176 proportionality (m⁻² sr μm), derived from the linear relationship (\pm 30%) between spectral radiance
177 and radiant power, for hot target having integrated temperature between 600 and 1500 K (Wooster
178 et al., 2003).

179

180 During the Cordón Caulle eruption (1 June 2011 - 31 August 2012), the MIROVA system detected
181 619 alerts over a total of 2983 MODIS overpasses (\sim 21%), with VRP spanning from less than 1
182 MW to more than 5000 MW (Fig. 2c). However, the visual inspection of all the alerted images
183 facilitated identification of a large number of data acquired in cloudy conditions (143 images),
184 and/or under poor geometrical conditions (i.e. high satellite zenith; 208 images) that strongly
185 deformed and affected the thermal anomaly at ground level. We thus selected 268 high quality
186 images (\sim 43% of alerted images; \sim 9% of the total MODIS overpasses; 1 alert every 2 days, on
187 average) that were used to calculate the total volcanic radiant energy (VRE) produced by the
188 eruption (Fig. 3). As demonstrated by Massimetti et al. (2017), the trapezoidal integration of VRP
189 signal related exclusively to supervised images (rather than to the whole unsupervised dataset),
190 provides a robust and more accurate quantification of the total energy radiated by the eruptive
191 events.

192

193 **Figure 2**

194

195 *3.2 Landsat 7 ETM+ Thermal Analysis*

196 A sequence of thermal images, acquired by the Enhanced Thematic Mapper plus (Landsat 7 ETM+)
197 is also presented in Fig. 3b-i. This sensor provides multispectral images in 7 spectral channels (from
198 visible to infrared) with a spatial resolution of 30 m (plus a pancromatic channel with 10 m
199 resolution). Here we used the channel 6, centered at 11.5 μm, that offered a synoptic view of the

200 thermal state of the CC lava flow during the eruption. By tracing the distinct boundary between the
201 background and the saturation level (about 78 °C; Donegan and Flynn, 2004) the thermally active
202 zones of the flow field were visualized (Fig. 3b-i). Despite a systematic flaw in all the ETM+
203 images acquired after May 2003 (diagonal data gaps prevent a complete view of the flow field), the
204 images shown in Fig. 3b-i facilitates improved tracking of the sources of the VRP identified by
205 MODIS.

206

207 **Figure 3**

208

209 **4. Results**

210 *4.1 Thermal output of the 2011-2012 Cordón Caulle eruption*

211

212 During the initial intense explosive phase (4-15 June 2011; gray field in Fig. 2c), MODIS-
213 MIROVA discontinuous alerts (27 alerts in total), with the radiant power reaching very high values
214 (VRP > 1000 MW) several time. In this first week of strong explosive activity, the detections were
215 mostly plume-contaminated, thus providing only a minimum estimate of the radiant power
216 generated by the active vent(s), where incandescent material was continuously ejected during
217 explosive activity (Silva-Parejas et al., 2012)..

218 In contrast, since 15 June 2011, the thermal output associated with the emplacement of the obsidian
219 lava flow was systematically tracked by MIROVA alerts (Fig. 3a). The radiant power reached the
220 maximum of 915 MW on 21 June, 6 days after the beginning of the lava effusion. Hence, the
221 thermal radiance started to decline gradually until September 2011 when the VRP was reduced to
222 ~100 MW. The slow declining trend was interrupted in early October 2011 when the VRP
223 gradually increase for about 4 weeks reaching again 200 MW in mid November. On 26 November
224 2011, an isolated peak of 634 MW marked a new breakout (Fig. 3a). This peak was followed by a

225 declining pattern similar to the one observed during the beginning of the eruption phase. Small
226 fluctuations were still recorded by February 2012 (between 100-150 MW) but between April and
227 July 2012 the VRP declined definitively to less than 5 MW, likely in correspondence of the end of
228 the effusion. This reduction in output corresponds to the lowering of the alert code from orange to
229 yellow (Fig. 2c).

230 By integrating (trapezoidal integration) the dataset of suitable images, collected during the effusive
231 phase between 15 June 2011 and 31 August 2012 (408 days, excluding the explosive and plume
232 injection phase), we calculated that the Cordón Caulle lava flow radiated approximately 2.75×10^{15}
233 J into the atmosphere, with a mean radiant flux of about 73 MW (Fig. 3a).

234

235 **Figure 3**

236

237 *4.2 Radiant density of CC obsidian lava flow*

238 Satellite-based thermal data have been widely used to estimate heat flux and lava discharge rates
239 during effusive eruptions (see [Harris, 2013](#) and references therein). The basic principle of this
240 method relies on a mutual relationship between effusion rates, active flow area and the thermal flux
241 that has been well documented ([Pieri and Baloga, 1986](#); [Wright et al., 2001](#); [Harris and Baloga](#)
242 [2009](#), for details). In particular, after a certain transient time required to reach thermal equilibrium
243 inside an active lava flow ([Garel et al., 2012, 2016](#)), the thermal energy radiated (VRE) can be
244 related to the erupted lava volume (Vol) throughout a best-fit parameter ([Coppola et al., 2013](#)).

$$245 \quad c_{rad} = \frac{VRE}{Vol} \quad [\text{equation 3}]$$

246 where c_{rad} (hereby called radiant density) is an empirical that embeds the appropriate rheological,
247 insulation and topographic conditions for the studied lava flow. Hence, by knowing the energy

248 radiated during an eruption and the volume of the erupted lava, it is possible to infer the appropriate
249 radiant density that characterizes a lava flow.

250 Systematic measurements of CC lava flow growth (Bertin et al., 2015) throughout the effusive
251 phase, were combined with MIROVA radiance data. These datasets facilitate the calculation of
252 radiant energy (VRE) for each independent volume measurement (Fig. 4a) and assess the evolution
253 of the radiant density (equation 3) during the eruption (Fig. 4b).

254 As a whole, we observe a very high correlation ($R^2 = 0.98$) between VRE and volume, with a best-
255 fit (linear interpolation) c_{rad} equal to $\sim 4.9 \times 10^6 \text{ J m}^{-3}$ (Fig. 4c). Notably, the radiant density
256 increased steadily during the first two weeks of extrusion (inset of Fig. 4b), and then stabilized to
257 the best-fit value reported above ($\pm 18\%$) until the end of the eruption (Fig. 4b).

258 The ratio between radiant and volumetric flux of active lava bodies is mainly controlled by their
259 bulk rheological properties (Coppola et al., 2013), with low-viscosity basaltic lava flows exhibiting
260 the highest value of c_{rad} ($1-4 \times 10^8 \text{ J m}^{-3}$) and viscous silicic flows represented by the lowest end-
261 member ($< 1 \times 10^7 \text{ J m}^{-3}$). Coppola et al. (2013) provided an empirical method to calculate the
262 radiant density of a lava body ($\pm 50\%$), on the basis of the silica content of erupted lavas, the latter
263 being considered a first-order proxy of its bulk rheological properties:

$$264 \quad c_{rad} = 6.45 \times 10^{25} \times (X_{SiO_2})^{-10.4} \quad \text{[equation 4]}$$

265 where X_{SiO_2} is the silica content of the erupted lavas (wt %). For Cordon Caulle lavas ($X_{SiO_2} = 70$
266 wt%; Castro et al., 2013), we calculated a radiant density of $4.2 (\pm 2.1) \times 10^6 \text{ J m}^{-3}$, in strong
267 agreement with those obtained from eq. 3 (Fig. 4d).

268

269 **Figure 4**

270

271 *4.3 Effusive trend of CC 2011-2012 eruption*

272 By using a radiant density equal to $4.9 \times 10^6 \text{ J m}^{-3}$, we calculated the effusion rates trends for the
273 CC 2011-2012 eruption, according to:

$$274 \quad TADR = \frac{VRP}{C_{rad}} \quad \text{[equation 5]}$$

275 where TADR is the time averaged lava discharge rate (Coppola et al., 2013).

276 The evolution of TADR is compared in Fig. 5 with the topographic estimates provided by Bertin et
277 al. (2015) and indicates the robustness of the thermal approach in tracking the variations of the
278 extrusive process, during an ongoing silicic eruption.

279 The MODIS-derived effusion rates show a general waxing-waning trend, typically observed during
280 pressurized eruptions (Wadge, 1981). The exponential decay of effusion rates is particularly evident
281 during the first 4 months of activity when the TADRs declined from $\sim 100 \text{ m}^3 \text{ s}^{-1}$ (on late June
282 2011) to $\sim 15 \text{ m}^3 \text{ s}^{-1}$ (on late September 2011). However, between October and November 2011 our
283 data suggest the occurrence of a slight increase in the effusive activity, that reached TADR values
284 of $40 - 50 \text{ m}^3 \text{ s}^{-1}$ (Fig. 5a). The inspection of ETM+ images suggests that this phase was
285 accompanied by the recrudescence of the surface activity in proximity of the vent, and emplacement
286 of new lava lobe(s) at the northwestern flow margins (Fig. 3d-e). In subsequent months, the
287 extrusion of lava slowed and was characterized by the emplacement at the northeastern lava flow
288 unit and temporary breakouts at the margins of the compound flow field (Fig. 3g-i).

289 Similarities are noted between MODIS data and the amplitude of the quasi-harmonic tremor (Bertin
290 et al., 2015) associated to the effusion of lava (initial waning stage followed by November 2011
291 pulse; Fig. 5a). The altitude of the erupted plumes, tracked by the Buenos Aires Volcanic Ash
292 Advisory Center (VAAC; GVP, 2013), also correlate to MODIS output. The remarkable
293 correspondence between lava discharge rate, seismicity and plume height (Fig. 5), point towards a
294 common origin of the observed trends until November 2011. However, this relationship requires

295 exogenous growth of the flow field (i.e. lateral expansion of the flow margin rather than inflation).
296 Accordingly, we suggest the MODIS-derived lava discharge rate pulse recorded on late November
297 2011, represents the buffered response of the compound flow field to a real increase in magma flux
298 at the vent, occurring 2-3 weeks earlier (25 October – 1 November). This would correspond to
299 increases in the harmonic tremor and plume injection altitude indicative of enhanced activity (Fig.
300 5b-c).

301

302 **Figure 5**

303

304 **5. Discussion**

305 The results presented here outline that during the CC eruption, the MIR-derived radiant power
306 (VRP) was strongly correlated to the effusive process (Fig. 5). In particular, we found that the
307 timescale over which the thermal proxy becomes stable (i.e. the c_{rad} reaches a quasi-steady value)
308 was approximately two weeks (Fig 4c). This transient time likely reflects a buffered thermal
309 response of the bulk, viscous lava flow to shorter variations of lava emission at the vent, as inferred
310 for example by variations of tremor amplitude (Fig. 5b) or plume elevation (Fig. 5c).

311 The timescale calculated here is much shorter than the one modelled by [Garel et al. \(2012\)](#),
312 according which thermal equilibrium during emplacement of lava domes or silicic domes is reached
313 only after several years. [Garel et al. \(2016\)](#) effectively suggest that the heat radiated by the whole
314 lava surface of viscous lava bodies may not reach a steady state emission (very long transient time),
315 and should not be used to calculate TADR. However, the same author also suggest that if only the
316 hottest and younger portions of the flow field is considered, the ratio between radiant flux and
317 discharge rate (i.e. c_{rad}) reaches a steady-state more rapidly, and the transient time becomes shorter.
318 This important conclusion implies that the wavelength (and method) used to calculate the radiant

319 flux is fundamental in determining what portion of the lava field is considered active, and what is
320 the appropriate transient time.

321 As argued by [Coppola et al. \(2016\)](#), the use of the MIR method ([Woster et al., 2003](#)) to calculate the
322 radiant power of active lava flows (characterized by a continuum of surface temperatures, from the
323 effusion temperature to the background), relies on the notion that the flow surfaces at temperatures
324 below 500-600 K do not contribute substantially to the pixel-integrated MIR radiance. Accordingly,
325 the VRP estimated using equation 2 does not provide the radiant power of the whole lava surface,
326 but more likely of smaller, younger and hotter portions of the flow field (cf. [Coppola et al., 2010](#)).
327 As shown by the Landsat images (Fig. 3c), in the case of long-lived viscous lava bodies, these areas
328 are restricted to the vent area, where the magma is extruded, as well as to active lobes at the flow
329 margins, shear fracture zones and hot cracks within the upper surface ([Bernstein et al., 2013](#)).

330 The timescale of ~ 15 days represents the temporal window over which the discharge rate should be
331 averaged (TADR), when applying the MIR-derived thermal proxy to viscous lava bodies similar to
332 the CC lava flow. Accordingly, a single TADR measurement does not necessarily indicates sharp
333 variations in the effusion rate at the vent, but may only reflect local surge(s) of lava, such as from
334 temporary breakouts at the margins of the compound flow field. Such a mechanism of emplacement
335 was recurrently observed during the CC eruption, with a significant evolution of breakout lava
336 observed over a period of only 6 days ([Tuffen et al., 2013](#)). This is possibly a common feature of
337 the obsidian flows that may affect the short-term thermal emission of the flow surface, thus causing
338 a higher variability of a single measurement (gray circles in Fig. 5a). On the other hand, variations
339 in viewing geometry of sun-synchronous satellites, such as MODIS, have been recognized by
340 [Flower et al. \(2016\)](#) as a possible source of high frequency (< 8 days) non-geophysical radiance
341 cycles. Despite selected only images with good viewing geometry, this non-volcanic cyclic
342 behavior may still partially affect our analysis. Therefore, when the TADR values are averaged over
343 an appropriate timescale, the local and temporary flow dynamics, as well as the noise created by
344 sensor viewing geometry, become smoothed, and the MODIS-derived trend better reflects the

345 effusive process at broad-scale (blue line in Fig. 5a). This suggests that a radiant density of $\sim 5 \times$
346 10^6 J m^{-3} , probably characterizes also the spreading and cooling processes of other obsidian lava
347 flows, whose bulk viscosities ($\sim 10^{10} \text{ Pa s}$) and emplacement styles are similar to the ones observed
348 at CC (Farquharson et al., 2015).

349

350 One of the most interesting feature of the CC eruption was the hybrid and coeval explosive-effusive
351 activity that led to emplacement of the large obsidian lava flow during sustained vulcanian ash
352 emissions (Castro et al., 2013; Shipper et al., 2013). As stressed by Shipper et al., (2013), this
353 coupled activity contrasts with most of the models that assume the explosive-effusive transition, as
354 composed by two end-member styles, separated in time and resulting from contrasting mechanisms.
355 The correlation observed between the tremor amplitude, plume height and the buffered MODIS-
356 derived effusion rates (Fig. 5), provides new elements to interpret this hybrid activity and suggests
357 that the explosive and effusive processes were fed by a common pressure source.

358 The overall effusive trend of the CC eruption (exponential decay) is indicative of a pressurized
359 eruption (Wadge, 1981) and is consistent with the sin-eruptive deformation data suggesting magma
360 withdrawal from one or multiple sources located between $\sim 4 \text{ km}$ and $\sim 6 \text{ km}$ depth (Jay et al.,
361 2014).

362 Here we modeled the MODIS-derived cumulative volume according to an asymptotic exponential
363 growth curve that is typical for such type of eruptions (Wadge, 1981):

364
$$V(t) = V_e \left(1 - \exp\left(-\frac{t}{\tau}\right) \right) \quad \text{[equation 6]}$$

365 where V_e is the excess magma volume inside the pressurized chamber, and τ is the decay time
366 constant. Our best-fit regression ($R^2 = 0.984$) indicates an initial excess volume V_e , equal to ~ 0.75
367 km^3 , and a time constant of ~ 205 days (Fig. 6a). Since the final volume of erupted lava was only \sim
368 0.6 km^3 , the above model suggests that $\sim 0.15 \text{ km}^3$ of magma remained unerupted after the end of
369 the eruption. This value is very similar to the volume of magma ($\sim 0.125 \text{ km}^3$) that was intruded

370 inside the edifice after the eruption (between March 2012 and April 2015), as documented by the
371 post-eruptive deformation data (Delgado et al., 2016). According to this model, a blockage in the
372 conduit, or in the magma path, would provide a possible link between the termination of the
373 eruptive phase and the subsequent rapid reinflation of the volcano. Analysis of residuals (Fig. 6a)
374 suggests that two monthly-long cycles were overprinted on the whole exponential trend. Although
375 we may not exclude that these cycles could be linked to exogenous growing process of the lava
376 field, cyclic behaviors are frequently observed at silicic volcanoes and could be related to non-linear
377 processes (degassing, crystallization, rheological stiffening) acting within the plumbing system
378 (Denlinger and Hoblitt, 1999).

379 On the other hand, and on the basis of the correlation with tremor amplitude and plume height (Fig.
380 5), the effusive trend of CC eruption could be considered as composed by two distinct stages (*Stage*
381 *1*: 15 June- 31 October 2014; *Stage 2*: 1 November 2014 –31 August 2012, respectively; Fig. 6). In
382 this case, our best-fit regression analysis indicates two very similar stages, characterized by V_e of
383 ~ 0.360 and ~ 0.265 km³ and time constants τ equal to ~ 62 and ~ 73 days, respectively (Fig. 6b).
384 This two-stages model would be consistent with tapping of distinct but interconnected melt bodies,
385 which is in agreement with a large and complex plumbing system, eventually composed by distinct
386 but hydraulically connected compartments (e.g., Gudmundsson, 2012), as suggested by deformation
387 pattern (Jay et al., 2014; Delgado et al., 2016) and geochemical analysis of erupted magmas
388 (Alloway et al., 2015).

389

390 **6. Conclusions**

391 The excellent correlation ($R^2 = 0.98$) between cumulative volumes and radiant energy, provides a
392 significant validation of the thermal proxy for the Cordón Caulle obsidian flow. We found that the
393 equilibrium between volumetric and radiant fluxes was reached after the first 15-20 days of activity,
394 and remained relatively stable ($c_{\text{rad}} = 4.9 \pm 0.9 \times 10^6$ J m⁻³) throughout the rest of the eruption. We

395 suggest that during long-lived silicic eruptions, the satellite-derived effusion rates must be time-
396 averaged over periods of at least 15 days. This correction allows filtering the cyclic noise related to
397 satellite viewing conditions (Flower et al., 2016), and permits to smooth local surge(s) of lava from
398 temporary breakouts unrelated to the magma flux at the vent. The absolute value of the coefficient
399 of proportionality (i.e. c_{rad} , radiant density) is significantly lower than those found for basaltic lava
400 flows ($1-4 \times 10^8 \text{ J m}^{-3}$) and falls well in the range of typical c_{rad} of viscous silicic flows ($2-10 \times 10^6$
401 J m^{-3} ; Coppola et al., 2013). This suggests that the spreading and cooling processes of lava bodies
402 can be essentially expressed through a viscosity-dependent relationship, linking the radiant energy
403 and the erupted lava volume. For any type of lava flow, if the timescale and the appropriate range of
404 c_{rad} values can be reasonably inferred (from direct measurements (eq. 3) or empirically (eq.4)), the
405 effusion rates and the erupted lava volumes can be calculated in near real time from satellite thermal
406 data. This methodology provides a new tool for tracking effusion rate trends not only for basaltic
407 eruptions (e.g. Coppola et al., 2017) but also for silicic, flow-forming eruptions whose effusive
408 dynamic is much less studied. In the case of CC eruption, the MODIS-derived effusion rates
409 correlate with the seismic tremor and the plume height, thus suggesting a common pressure source
410 for the coeval effusive and explosive activities. Modelling of the effusion rate trend allow us to
411 suggest alternatively an exponential discharge dynamic, typical of pressurized eruptions (Wadge,
412 1981), characterized by incomplete magma withdrawal and linked to the post-eruptive inflation
413 phase; or a more complex magma discharge, probably tapping different interconnected magma
414 bodies.

415

416 **Acknowledgments**

417

418 We acknowledge the Editor and the anonymous reviewer for the constructive comments and
419 suggestions on the original manuscript. LANCE-MODIS system ([http://lance-
420 modis.eosdis.nasa.gov/](http://lance-modis.eosdis.nasa.gov/)) provided Level 1B MODIS data. Tremor signal was processed at the

421 Observatorio Volcanológico de los Andes del Sur in SERNAGEOMIN and lava effusion rate from
422 computed volumes were obtained as part of response during the Cordón Caulle 2011 eruption.
423 MIROVA is a collaborative project between the Universities of Turin and Florence (Italy), and is
424 supported by the Italian Civil Protection Department

425

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626

627 **Figure Captions**

628

629 **Figure 1** – (a) Location of Cordón Caulle Volcanic Complex (CCVC) on the Southern Andes
630 Volcanic Zone, with the Cordón Caulle (CC) fissures system lying between the Cordillera Nevada,
631 on the NW, and the Puyehue stratovolcano on the SE. The lava flow related to the 2011-2012
632 eruption is shown in red (shaded relief map from Google Maps). (b) A detailed view of the 7.2 km²
633 obsidian lava flow emplaced throughout the eruption (image from Google Maps).

634

635 **Figure 2** – (a) Example of MODIS images (in transparency on Google Earth) acquired over CCVC
636 on 30 December 2011 (05:55 UTC) and elaborated by the MIROVA system (Coppola et al., 2016).
637 The image displays the MIR radiance recorded at 3.959 μm (MODIS band 22), over an area of
638 50x50 km centered on Puyehue volcano (pixel resolution is 1 km). Thermal anomalies appear as
639 bright pixels at the eruptive site, while the ash volcanic plume, directed toward north, is represented
640 by black (cold) pixels. (b) Zoomed view over the CC lava flow (dashed black line) showing the
641 distributions of alerted pixels detected by the MIROVA algorithm (black contour line). The total
642 “above background” MIR radiance, ΔL_{MIR} , is equal to 7.78 W m⁻² sr⁻¹ μm⁻¹, corresponding to a
643 VRP of 147 MW (eq. 2). (c) Volcanic Radiative Power (VRP) recorded during the 2011-2012 CC
644 eruption (on log scale). The horizontal color bar refers to the alert codes provided during the

645 eruptive crisis by the *Servicio Nacional de Geología y Minería*. The gray field outlines the initial
646 phase (4 – 15 June 2011) characterized by intense explosive activity and ash emission.

647

648 **Figure 3** – (a) Radiant power (left axis) and energy (right axis) recorded by MODIS during the
649 effusive phase of the 2011-2012 Cordón Caulle eruption (15 June 2011 – 31 August 2012). Gray
650 circles represent VRP of single selected images. The blue line (7 points moving average) describes
651 the long-term pattern of the radiant flux between June 2011 and August 2012. (b) Sequence of
652 thermal images derived from Enhanced Thematic Mapper Plus (Landsat 7 ETM+) acquired on 2011-
653 2012 over CC lava flow. Colors represent the “above background” radiance, recorded at 11.5 μm
654 (ETM+ Channel 6). White pixels represent flow surfaces at temperature higher than the saturation
655 level (pixel integrated temperature of $\sim 78^\circ\text{C}$). Black contours outline the final extension of the lava
656 flow. Diagonal data gaps are due to failure of Landsat 7 Scan Line Corrector (SLC) occurred in
657 May 2003.

658

659 **Figure 4** – (a) Temporal evolution of CC lava flow volume (black squares) and radiant energy
660 measured by MODIS (red circles). Volume data from [Bertin et al. 2015](#); (b) Temporal evolution of
661 the radiant density (eq. 3) throughout the eruption. After the first 15 days of activity the radiant
662 density reaches an almost steady value ($\sim 4.9 \times 10^6 \text{ J m}^{-3}$), representing by best-fit relationship (c)
663 between Volcanic Radiant Energy (VRE) and erupted volume (Vol); (d) Relationship between
664 radiant density and silica content of 28 lava bodies (gray squares) analyzed by [Coppola et al. 2013](#).
665 The high silica Cordón Caulle obsidian flow (yellow star) is characterized by a radiant density in
666 excellent agreement with those predicted by the eq. 4 (black dashed line).

667

668 **Figure 5** - (a) Time-Averaged Discharge Rate (TADR) estimated for the effusive stage of 2011-
669 2012 Cordón Caulle eruption by using a radiant density $c_{\text{rad}} = 4.5 \times 10^6 \text{ J m}^{-3}$ (eq. 3). Single MODIS-
670 derived measurements (gray circles) likely tracked the continuous occurrence of short-lived, local

671 breakouts that produced the sharp fluctuations of the TADR. The weekly average (blue line) better
672 reproduces the extrusive process as measured by the volumes estimates of Bertin et al., 2015 (black
673 squares). (b) Quasi-harmonic tremor (maximum reduced displacement; RD) recorded between June
674 and December 2011 (modified from Bertin et al., 2015); (c) Plume height recorded during the 2011-
675 2012 CC eruption (data from the Buenos Aires Volcanic Ash Advisory Center VAAC and
676 SERNAGEOMIN; source Global Volcanism Program, 2013).

677

678 **Figure 6** – MODIS-derived cumulative volume (blue circles) measured during the CC eruption. (a)
679 A single stage magma discharge model (red line, equation 6) indicates an excess magma volume V_e
680 $= 0.75 \text{ km}^3$ and decay constant $\tau = 205$ days. At the end of the eruption (31 august 2012) only 0.6
681 km^3 were erupted, thus implying a volume of unerupted magma equal to 0.15 km^3 . The occurrence
682 of two sub-cycles (as evidenced by residuals) may result from non-linear processes inside the
683 plumbing system. (b) Two-stages magma discharge model (red lines). Stage 1 (15 June- 31 October
684 2014) is characterized by $V_{e1} = 0,360 \text{ km}^3$ and $\tau_1=62$ days; Stage 2 (1 November 2014 –31 August
685 2012), is characterized by $V_{e2}= 0,266 \text{ km}^3$ and $\tau_2=73$ days. The consecutive eruptions of two
686 interconnected magmatic sources (compartments), could be at the origin of this trend.