



UNIVERSITÀ DEGLI STUDI DI TORINO

AperTO - Archivio Istituzionale Open Access dell'Università di Torino

Segregated configurations involving the square root of the laplacian and their free boundaries

This is the author's manuscript		
Original Citation:		
Availability:		
This version is available http://hdl.handle.net/2318/1705346 since 2019-07-05T22:41:13Z		
Published version:		
DOI:10.1007/s00526-019-1529-9		
Terms of use:		
Open Access		
Anyone can freely access the full text of works made available as "Open Access". Works made available under a Creative Commons license can be used according to the terms and conditions of said license. Use of all other works requires consent of the right holder (author or publisher) if not exempted from copyright protection by the applicable law.		

(Article begins on next page)

SEGREGATED CONFIGURATIONS INVOLVING THE SQUARE ROOT OF THE LAPLACIAN AND THEIR FREE BOUNDARIES

DANIELA DE SILVA AND SUSANNA TERRACINI

ABSTRACT. We study the local structure and the regularity of free boundaries of segregated minimal configurations involving the square root of the laplacian. We develop an improvement of flatness theory and, as a consequence of this and Almgren's monotonicity formula, we obtain partial regularity (up to a small dimensional set) of the nodal set, thus extending the known results in [4, 21] for the standard diffusion to some anomalous case.

1. INTRODUCTION

The analysis of the nodal sets of segregated stationary configurations for systems of elliptic equations has been the subject of an intense study in the last decade, starting from the works [9, 10, 11, 6, 4, 3, 17, 12, 13, 14]. The present paper is concerned with the geometric structure of the nodal set, when the creation of a free boundary is triggered by the interplay between *fractional diffusion* and *competitive interaction*. A prototypical example comes the following system of fractional Gross-Pitaevskii equations

(1.1)
$$\begin{cases} (-\Delta + m_i^2)^{1/2} u_i + V(x) u_i = \omega_i u_i^3 - \beta u_i \sum_{j \neq i} a_{ij} u_j^2 \\ u_i \in H^{1/2}(\mathbb{R}^N), \end{cases}$$

with $a_{ij} = a_{ji} > 0$, which is the relativistic version of the Hartree-Fock approximation theory for mixtures of Bose-Einstein condensates in different hyperfine states which overlap in space. The sign of ω_i reflects the type of interaction of the particles within each single state. If ω_i is positive, the self interaction is attractive (focusing problems defocusing otherwise). *V* represents an external potential. The sign of β , on the other hand, accounts for the interaction of particles in different states. This interaction is attractive when negative and repulsive otherwise. If the condensates repel, and the competition rate tends to infinity, the densities eventually separate spatially, giving rise to a free boundary: the common nodal set of the components u_i 's. This phenomenon is called phase separation and has been described in the recent literature, both physical and mathematical, in the case of standard diffusion. It is by now a well-established fact that in the case of elliptic systems with standard diffusion this nodal set is comparable, as regards to the qualitative properties, to that of the scalar solutions. The main reason can be attributed to the validity of a weak reflection law (see [21]) which constitutes the condition of extremality at the

Date: November 28, 2018.

²⁰¹⁰ Mathematics Subject Classification. 35J70, 35J75, 35R11, 35B40, 35B44, 35B53,

Key words and phrases. Free boundary problems, nonlocal diffusion, improvement of flatness, monotonicity formulas, blow-up classification.

Work partially supported by the ERC Advanced Grant 2013 n. 339958 Complex Patterns for Strongly Interacting Dynamical Systems - COMPAT .

common interface. Relevant connections have been established with optimal partition problems involving spectral functionals (cfr [21, 18]).

We consider the following model with fractional diffusion: according to [5], the *n*-dimensional half laplacian can be interpreted as a (nonlinear) Dirichlet-to-Neumann operator. For this reason we shall state all our results for harmonic functions with non-linear Neumann boundary conditions involving strong competition terms. Precisely, the following uniform-in- β estimates have been derived in [23]. Here we denote by $B_r^+ := B_r \cap \{z > 0\} \subset \mathbb{R}^{n+1}$ and by $\mathcal{B}_r := B_r \cap \{z = 0\}$ (where *z* is the (n+1)-th coordinate).

Theorem 1.1 (Local uniform Hölder bounds, [23]). Let the functions $f_{i,\beta}$ be continuous and uniformly bounded (w.r.t. β) on bounded sets, and let $\{\mathbf{u}_{\beta} = (u_{i,\beta})_{1 \le i \le k}\}_{\beta}$ be a family of $H^1(B_1^+)$ solutions to the problems

$$(P_{\beta}) \qquad \begin{cases} -\Delta u_{i} = 0 & \text{in } B_{1}^{+} \\ \partial_{\nu} u_{i} = f_{i,\beta}(u_{i}) - \beta u_{i} \sum_{j \neq i} u_{j}^{2} & \text{on } \mathcal{B}_{1}. \end{cases}$$

Let us assume that

$$\|\mathbf{u}_{\beta}\|_{L^{\infty}(B_{1}^{+})} \leq M,$$

for a constant M independent of β . Then for every $\alpha \in (0, 1/2)$ there exists a constant $C = C(M, \alpha)$, not depending on β , such that

$$\|\mathbf{u}_{\beta}\|_{\mathcal{C}^{0,\alpha}\left(\overline{B_{1/2}^{+}}\right)} \leq C(M,\alpha).$$

Furthermore, $\{\mathbf{u}_{\beta}\}_{\beta}$ is relatively compact in $H^{1}(B^{+}_{1/2}) \cap \mathcal{C}^{0,\alpha}\left(\overline{B^{+}_{1/2}}\right)$ for every $\alpha < 1/2$.

As a byproduct, up to subsequences, we have convergence of the solutions to (P_{β}) to some limiting profile, whose components are segregated on the boundary \mathcal{B}_1 . If furthermore $f_{i,\beta} \to f_i$, uniformly on compact sets, we can prove that this limiting profile satisfies

(1.2)
$$\begin{cases} -\Delta u_i = 0 & \text{in } B_1^+ \\ u_i \partial_\nu u_i = f_i(u_i) u_i & \text{on } \mathcal{B}_1. \end{cases}$$

One can see that, for solutions of this type of equation, the highest possible regularity corresponds indeed to the Hölder exponent $\alpha = 1/2$. As a matter of fact, it has been proved that the limiting profiles do enjoy such optimal regularity.

Theorem 1.2 (Optimal regularity of limiting profiles, [23]). Under the assumptions above, assume moreover that the locally Lipschitz continuous functions f_i satisfy $f_i(s) = f'_i(0)s + O(|s|^{1+\varepsilon})$ as $s \to 0$, for some $\varepsilon > 0$. Then $\mathbf{u} \in C^{0,1/2}\left(\overline{B^+_{1/2}}\right)$.

It is worthwhile noticing that these result apply to the (local) minimizers of the functionals

(1.3)
$$J_{\beta}(U) = \sum_{i=1}^{k} \int_{B_{1}^{+}} \frac{1}{2} \nabla u_{i}(x,z) |^{2} dx \, dz + \beta \sum_{1 \le i < j \le k} \int_{\mathcal{B}_{1}} u_{i}^{2}(x,0) u_{j}^{2}(x,0) \, dx$$

in the set of all configurations with fixed boundary data. Taking the singular limit as $\beta \rightarrow +\infty$ we are naturally lead to consider the energy minimizing profiles which segregate only at the characteristic hyperplane $\{z = 0\}$.

Our main goal is to describe, from differential and geometric measure theoretical points of view, the structure of the trace on the characteristic hyperplane $\{z = 0\}$ of the common nodal set of these limiting profiles. From now on, for the sake of simplicity we shall

assume that the reactions f_i 's are identically zero and we reflect the components u_i 's through the hyperplane $\{z = 0\}$. It is worthwhile noticing that we cannot deduce from the system (1.2) alone any regularity property of the common nodal set $\mathcal{N}(\mathbf{u}) = \{x \in \Omega \cap \mathbb{R}^n \times \{0\} : \mathbf{u}(x, 0) = 0\}$, as the equations can be independently solved for arbitrary, though mutually disjoint, nodal sets $\mathcal{N}(u_i) = \{(x, 0) \in \Omega \cap \mathbb{R}^n \times \{0\} : u_i(x, 0) = 0\}$ on the characteristic hyperplane $\{z = 0\}$.

Definition 1.3 (Segregated minimal configurations, class $\mathfrak{M}(\Omega)$). For an open, z-symmetric $\Omega \subset \mathbb{R}^{n+1}$, we define the class $\mathfrak{M}(\Omega)$ of the segregated minimal configurations as the set of all the even-in-z vector valued functions $\mathbf{u} = (u_1, \ldots, u_k) \in (H^1(\Omega))^k \cap \mathcal{C}^{0,1/2}(\Omega)$, whose components are all nonnegative and achieving the minimal the energy among configurations segregating only at the characteristic hyperplane $\{z = 0\}$, that is solutions to

$$\min\left\{\sum_{i=1}^{k} \int_{\Omega} |\nabla u_i|^2 : \begin{array}{c} u_i(x,0) \cdot u_j(x,0) \equiv 0 \ \mathbb{R}^n \text{-a.e. for } i \neq j, \\ u_i = \varphi_i, \text{ on } \partial\Omega \text{ for } i = 1, \dots, k, \end{array}\right\}$$

where the φ_i 's are nonnegative $H^{1/2}$ -boundary data which are even-in-z and segregated on the hyperplane $\{z = 0\}$.

For such class of solutions, we are going to prove a theorem on the structure of the nodal set $\mathcal{N}(\mathbf{u})$, which is the perfect counterpart of the results in [4, 21].

Theorem 1.4 (Structure of the nodal set of segregated minimal configurations). Let $\Omega \subset \mathbb{R}^{n+1}$, with $n \geq 2$, **u** be a segregated minimal configuration and let $\mathcal{N}(\mathbf{u}) = \{x \in \tilde{\Omega} : \mathbf{u}(x,0) = 0\}$. Then, $\mathcal{N}(\mathbf{u})$ is the union of a relatively open regular part $\Sigma_{\mathbf{u}}$ and a relatively closed singular part $\mathcal{N}(\mathbf{u}) \setminus \Sigma_{\mathbf{u}}$ with the following properties:

- (i) $\Sigma_{\mathbf{u}}$ is a locally finite collection of hyper-surfaces of class $C^{1,\alpha}$ (for some $0 < \alpha < 1$).
- (ii) $\mathcal{H}_{dim}(\mathcal{N}(\mathbf{u}) \setminus \Sigma_{\mathbf{u}}) \leq n-2$ for any $n \geq 2$. Moreover, for n = 2, $\mathcal{N}(\mathbf{u}) \setminus \Sigma_{\mathbf{u}}$ is a locally finite set.

Remark 1.5. (a) In the light of the extension facts related to the half-laplacian, our theory applies, among others, to segregated minimizing configurations involving non local energies, like, for instance the solutions to the following problem (when s = 1/2):

$$\min\left\{\sum_{i=1}^{k} \int_{\mathbb{R}^{2n}} \frac{|u_i(x) - u_i(y)|^2}{|x - y|^{n+2s}} : u_i(x, 0) \cdot u_j(x, 0) \equiv 0 \text{ a.e. in } \mathbb{R}^n \text{ for } i \neq j, \\ u_i \equiv \varphi_i, \text{ on } \mathbb{R}^n \setminus \tilde{\Omega} \text{ for } i = 1, \dots, k, \right\}$$

where φ_i are nonnegative $H^{1/2}(\mathbb{R}^n)$ data which are segregated themselves. In this regard, our results extend those of [4] to the fractional case, or, equivalently, to the case when the phase segregation takes place only on the characteristic hyperplane.

(b) It is worthwhile noticing that, in case of the standard diffusion, the nodal set of the segregated minimal configurations shares the same measure theoretical features with the nodal set of harmonic functions; this is not the case of the fractional diffusion; indeed, as shown in [20], the stratified structure of the nodal set of *s*-harmonic functions is far more complex than that of the segregated minimal configurations. The asymptotics and properties of limiting profiles of competition diffusion systems with quadratic (Lotka-Volterra) mutual interactions have been investigated in [25]; as discussed there, the free boundary, in the Lotka-Volterra case, resembles the nodal set of *s*-harmonic functions with some important differences however, enlightened in that paper.

(c) The theory developed in this paper is suitable to extend in order to cover the limiting cases, as $\beta \to \infty$, of the variational problems associated with (1.1).

In order to prove our main result, we will actually consider two other different notions of solution. The first is the class of solutions in the variational sense - concept that we have expressed through the validity of the domain variation formula in Definition 6.1 - and a second notion of solution, expressed in terms of viscosity solutions in Definitions 2.3 (for two components) and 7.1 (for k components). As we will see in Theorem 7.3, the latter notion is weaker (although they are probably equivalent) but carries precious information on the regularity of the nodal set. In particular, both notions encode a reflection rule about the free boundary (see e.g. Proposition 6.3), which will be the ultimate reason for the regularity of the nodal set.

A major achievement toward the proof of our structure theorem will be an improvement of flatness argument for the case of two components, which was inspired by the work in [15]. As a byproduct, it yields the following local regularity result: in two dimensions, let U(t, z) be defined in polar coordinates as

(1.4)
$$U(t,z) = r^{1/2} \cos \frac{\theta}{2},$$
$$t = r \cos \theta, \quad z = r \sin \theta, \quad r \ge 0, \quad -\pi \le \theta \le \pi,$$

and let

$$\bar{U}(t,z) := U(-t,z).$$

By abuse of notation we denote,

$$U(x,z) := U(x_n,z), \quad \overline{U}(x,z) := \overline{U}(x_n,z).$$

Theorem 1.6 (Local regularity of the free boundary). There exists $\bar{\varepsilon} > 0$ small depending only on n, such that if $\mathbf{u} = (u_1, u_2)$ is a viscosity solution in B_1 in the sense of Definition 2.3 satisfying

(1.5)
$$\|u_1 - U\|_{\infty} \le \bar{\varepsilon}, \quad \|u_2 - \bar{U}\|_{\infty} \le \bar{\varepsilon}$$

then $\mathcal{N}(\mathbf{u})$ is a $C^{1,\alpha}$ graph in $B_{\frac{1}{2}} \cap \mathbb{R}^n \times \{0\}$ for every $\alpha \in (0,1)$ with $C^{1,\alpha}$ norm bounded by a constant depending on α and n.

The paper is organized as follows. Section 2 contains the definition of viscosity solution for 2-components systems, and the basic related facts. Section 3 is devoted to the study of the linearized problem associated with ε -domain variations around the fundamental solution of the system defined in (1.4). Next, Section 4 contains Harnack estimates for viscosity solutions to the free boundary problem for two components. Such estimates will be crucial tools in proving the improvement of flatness result, in Section 5, which concludes the analysis of the regular part of the free boundary of 2-systems and proves Theorem 1.6. The rest of the paper concerns k-vector systems: Section 6 is devoted to the study of k-vector solutions in the variational sense (class \mathcal{G}) and the consequences of the associated Almgren's monotonicity formula, with a focus on the existence and classification of blow-ups and conic entire solutions. In Section 7 we introduce the notion of k-vector solutions in the viscosity sense and we connect it to the variational one, in order to prove Theorem 1.4. The two appendices contain some ancillary known results.

CONTENTS

1.	Introduction	1
2.	Viscosity solutions for two component systems	5
2.1.	Notations and definitions.	5
2.2.	Comparison principle.	6

SEGREGATED CONFIGURATIONS INVOLVING THE SQUARE ROOT OF THE LAPLACIAN	5
2.3. Renormalization.	8
2.4. Domain Variations	8
3. The linearized problem.	11
4. Harnack Inequality	19
5. Improvement of Flatness	25
6. Segregated critical profiles	29
6.1. Compactness and convergence of blowup sequences	32
6.2. Segregated entire homogenous profiles	33
7. Viscosity solutions to <i>k</i> -component systems	35
7.1. Structure of the nodal set	35
Appendix A.	37
Appendix B.	
References	

2. VISCOSITY SOLUTIONS FOR TWO COMPONENT SYSTEMS

2.1. Notations and definitions. First, we introduce some notations.

A point $X \in \mathbb{R}^{n+1}$ will be denoted by $X = (x, z) \in \mathbb{R}^n \times \mathbb{R}$, and sometimes $x = (x', x_n)$ with $x' = (x_1, \ldots, x_{n-1})$.

A ball in \mathbb{R}^{n+1} with radius r and center X is denoted by $B_r(X)$ and for simplicity $B_r = B_r(0)$. Also, for brevity, \mathcal{B}_r denotes the n-dimensional ball $B_r \cap \{z = 0\}$ (previously denoted with $\partial^0 B_r^+$).

Let v(X) be a continuous non-negative function in B_1 . We associate with v the following sets:

$$\begin{split} B_{1}^{+}(v) &:= B_{1} \setminus \mathcal{N}(v), \quad \mathcal{N}(v) := \{(x,0) : v(x,0) = 0\}; \\ \mathcal{B}_{1}^{+}(v) &:= B_{1}^{+}(v) \cap \mathcal{B}_{1}; \\ F(v) &:= \partial_{\mathbb{R}^{n}} \mathcal{N}^{o}(v) \cap \mathcal{B}_{1}, \quad \mathcal{N}^{o}(v) := Int_{\mathbb{R}^{n}}(\mathcal{N}(v)). \end{split}$$

We now introduce the definition of viscosity solutions for a problem with two components. Let $u_1(x, z), u_2(x, z)$ be non-negative continuous functions in the ball $B_1 \subset \mathbb{R}^{n+1} = \mathbb{R}^n \times \mathbb{R}$, which vanish on complementary subsets of $\mathbb{R}^n \times \{0\}$ and are even in the *z* variable. We consider the following free boundary problem

(2.1)
$$\begin{cases} \Delta u_i = 0, & \text{in } B_1^+(u_i), i = 1, 2, \\ \mathcal{B}_1 = \mathcal{N}(u_1) \cup \mathcal{N}(u_2), & \mathcal{N}^o(u_1) \cap \mathcal{N}^o(u_2) = \emptyset, \\ \frac{\partial u_1}{\partial \sqrt{t}} = \frac{\partial u_2}{\partial \sqrt{t}}, & \text{on } F(u_1, u_2) := \partial_{\mathbb{R}^n} \mathcal{N}^o(u_1) \cap \mathcal{B}_1 = \partial_{\mathbb{R}^n} \mathcal{N}^o(u_2) \cap \mathcal{B}_1, \end{cases}$$

where

(2.2)
$$\frac{\partial u_i}{\partial \sqrt{t}}(x_0) := \lim_{t \to 0^+} \frac{u_i(x_0 + t\nu_i(x_0), 0)}{\sqrt{t}}, \quad x_0 \in F(u_1, u_2)$$

with $\nu_i(x_0)$ the normal to $F(u_1, u_2)$ at x_0 pointing toward $\{u_i(x, 0) > 0\}$.

Definition 2.1. Given g, v continuous in B_1 , we say that v touches g by below (resp. above) at $X_0 \in B_1$ if $g(X_0) = v(X_0)$, and

$$g(X) \ge v(X)$$
 (resp. $g(X) \le v(X)$) in a neighborhood O of X_0 .

If this inequality is strict in $O \setminus \{X_0\}$, we say that v touches g strictly by below (resp. above).

Definition 2.2. We say that the ordered pair (v_1, v_2) is a (strict) comparison subsolution to (2.1) if the $v_i \in C(B_1)$ are non-negative functions even-in-*z* that satisfy

- (i) v_1 is C^2 and $\Delta v_1 \ge 0$, in $B_1^+(v_1)$,
- (ii) v_2 is C^2 and $\Delta v_2 \le 0$, in $B_1^+(v_2)$,

(iii)
$$\mathcal{B}_1 = \mathcal{N}(v_1) \cup \mathcal{N}(v_2), \quad \mathcal{N}^o(v_1) \cap \mathcal{N}^o(v_2) = \emptyset, \quad \mathcal{N}^o(v_i) \neq \emptyset;$$

(iv) $F(v_1, v_2) := F(v_1) = F(v_2)$ is C^2 and if $x_0 \in F(v_1, v_2)$ we have

$$v_i(x_0 + t\nu_i(x_0), 0) = \alpha_i(x_0)\sqrt{t} + o(\sqrt{t}), \text{ as } t \to 0^+,$$

with

$$\alpha_1(x_0) \ge \alpha_2(x_0),$$

where $\nu_i(x_0)$ denotes the unit normal at x_0 to $F(v_1, v_2)$ pointing toward $\{v_i(x, 0) > 0\}$;

(v) Either the v_i are both not harmonic in $B_1^+(v_i)$ or $\alpha_1(x_0) > \alpha_2(x_0)$ at all $x_0 \in F(v_1, v_2)$.

Similarly the ordered pair (w_1, w_2) is a (strict) comparison supersolution if (w_2, w_1) is a (strict) comparison subsolution.

Definition 2.3. We say that (u_1, u_2) is a viscosity supersolution to (2.1) if $u_i \ge 0$ is a continuous function in B_1 which is even-in-z and it satisfies

- (i) $\Delta u_i = 0$ in $B_1^+(u_i)$;
- (ii) $\mathcal{B}_1 = \mathcal{N}(u_1) \cup \mathcal{N}(u_2), \quad \mathcal{N}^o(u_1) \cap \mathcal{N}^o(u_2) = \emptyset;$
- (iii) If (v_1, v_2) is a (strict) comparison subsolution then v_1 and v_2 cannot touch u_1 and u_2 respectively by below and above at a point $X_0 = (x_0, 0) \in F(u_1, u_2) := \partial \mathcal{N}^o(u_1) \cap \mathcal{B}_1$.

Respectively, we say that (u_1, u_2) is a viscosity subsolution to (2.1) if the conditions above hold with (iii) replaced by

(iv) If (w_1, w_2) is a (strict) comparison supersolution then w_1 and w_2 cannot touch u_1 and u_2 respectively by above and below at a point $X_0 = (x_0, 0) \in F(u_1, u_2)$.

We say that (u_1, u_2) is a viscosity solution if it is both a super and a subsolution.

2.2. Comparison principle. We now derive a basic comparison principle.

Definition 2.4. Let (v_1, v_2) be a comparison subsolution to (2.1). We say that (v_1, v_2) is monotone in the e_n direction whenever v_1 is monotone increasing and v_2 is monotone decreasing in the e_n direction.

Lemma 2.5 (Comparison principle). Let $(u_1, u_2), (v_1^t, v_2^t) \in C(\overline{B}_1)$ be respectively a solution and a family of comparison subsolutions to (2.1), $t \in [0, 1]$. Assume that

 $\begin{array}{ll} \text{(i)} & v_1^0 \leq u_1, & v_2^0 \geq u_2 \text{ in } \overline{B}_1; \\ \text{(ii)} & v_1^t \leq u_1, & v_2^t \geq u_2 \text{ on } \partial B_1 \text{ for all } t \in [0,1]; \end{array}$

- (iii) $v_1^t < u_1 \text{ on } \mathcal{F}(v_1^t)$ which is the boundary in ∂B_1 of the set $\partial \mathcal{B}_1^+(v_1^t) \cap \partial \mathcal{B}_1$, for all $t \in [0, 1]$;
- (iv) $v_2^t > u_2$ on $\mathcal{F}(u_2)$ which is the boundary in ∂B_1 of the set $\partial \mathcal{B}_1^+(u_2) \cap \partial \mathcal{B}_1$, for all $t \in [0, 1]$;
- (v) v_i^t is continuous in $(x, t) \in \overline{B}_1 \times [0, 1]$ and $\overline{\mathcal{B}_1^+(v_i^t)}$ is continuous in the Hausdorff metric.

Then

$$v_1^t \le u_1, \quad v_2^t \ge u_2, \quad \text{in } B_1, \text{ for all } t \in [0, 1].$$

Proof. Let

$$A := \{ t \in [0,1] : v_1^t \le u_1, v_2^t \ge u_2 \text{ on } \overline{B}_1 \}.$$

In view of (i) and (v) A is closed and non-empty. Our claim will follow if we show that A is open. Let $t_0 \in A$, then $v_1^{t_0} \leq u_1, v_2^{t_0} \geq u_2$ on \overline{B}_1 and by the definition of viscosity solution

$$F(v_i^{t_0}) \cap F(u_i) = \emptyset, \quad i = 1, 2$$

For i = 1 we argue as follows. Together with (iii) the identity above implies that

$$\mathcal{B}_{1}^{+}(v_{1}^{t_{0}}) \subset \mathcal{B}_{1}^{+}(u_{1}), \quad F(v_{1}^{t_{0}}) \cup \mathcal{F}(v_{1}^{t_{0}}) \subset \{x \in \overline{\mathcal{B}}_{1} : u_{1}(x,0) > 0\}.$$

By (v) this gives that for t close to t_0

(2.3)
$$\mathcal{B}_1^+(v_1^t) \subset \mathcal{B}_1^+(u_1), \quad F(v_1^t) \cup \mathcal{F}(v_1^t) \subset \{x \in \overline{\mathcal{B}}_1 : u_1(x,0) > 0\}.$$

Call $D := B_1 \setminus (\mathcal{B}_1^0(v_1^t) \cup F(v_1^t))$. Combining (2.3) with assumption (ii) we obtain

$$v_1^t \leq u_1 \quad \text{on } \partial D_1$$

and by the maximum principle the inequality holds also in D. Hence

$$v_1^t \le u_1 \quad \text{in } \overline{B}_1.$$

Similarly for i = 2, combining (iv) and (v) we get that for t close to t_0 ,

$$v_2^t \ge u_2 \quad \text{in } \overline{B}_1$$

Hence $t \in A$ which shows that A is open.

The corollary below is now a straightforward consequence of Lemma 2.5.

Corollary 2.6. Let (u_1, u_2) be a solution to (2.1) and let (v_1, v_2) be a comparison subsolution to (2.1) in B_2 which is strictly monotone in the e_n direction in $B_2^+(v_i)$. Call

$$v_i^t(X) := v_i(X + te_n), \quad X \in B_1.$$

Assume that for $-1 \le t_0 < t_1 \le 1$

$$v_1^{t_0} \le u_1, \quad v_2^{t_0} \ge u_2 \quad \text{in } B_1,$$

$$v_1^{\iota_1} \le u_1 \quad \text{on } \partial B_1, \quad v_1^{\iota_1} < u_1 \quad \text{on } \mathcal{F}(v_1^{\iota_1}),$$

and

$$v_2^{t_1} \geq u_2 \quad \textit{on } \partial B_1, \quad v_2^{t_1} > u_2 \quad \textit{on } \mathcal{F}(u_2).$$

Then

$$v_1^{t_1} \leq u_1, \quad v_2^{t_1} \geq u_2 \quad \text{ in } \overline{B}_1.$$

2.3. **Renormalization**. The following result allows us to replace the flatness assumption (1.5) with the property that u_1 and u_2 are trapped between two nearby translates of U and \overline{U} respectively. Precisely, we have the following lemma.

Lemma 2.7. Let (u_1, u_2) be non negative continuous functions in \overline{B}_2 satisfying

 $\Delta u_i = 0, \quad in \, B_2^+(u_i),$

(2.4)
$$\mathcal{B}_1 = \mathcal{N}(u_1) \cup \mathcal{N}(u_2), \quad \mathcal{N}^o(u_1) \cap \mathcal{N}^o(u_2) = \emptyset$$

and the flatness assumption

$$(2.5) ||u_1 - U||_{\infty} \le \delta, ||u_2 - \overline{U}||_{\infty} \le \delta$$

with $\delta > 0$ small universal. Then in B_1 ,

(2.6)
$$U(X - \varepsilon e_n) \le u_1(X) \le U(X + \varepsilon e_n), \quad U(X + \varepsilon e_n) \le u_2(X) \le U(X - \varepsilon e_n),$$

for some $\varepsilon = K\delta$, K universal.

Lemma 2.7 follows immediately from Lemma A.3 in the Appendix. Indeed, (2.4)-(2.5) guarantee that u_i satisfies the assumption (A.5).

2.4. **Domain Variations.** We recall the definition of ε -domain variation corresponding to U and some basic lemmas from [15]. We also introduce in a similar fashion the ε -domain variation corresponding to \overline{U} and deduce its properties.

Let $\varepsilon > 0$ and let g be a continuous non-negative function in \overline{B}_{ρ} . Here and henceforth we denote by P^{\pm} the half-hyperplanes

$$P^{+} := \{ X \in \mathbb{R}^{n+1} : x_n \ge 0, z = 0 \},\$$
$$P^{-} := \{ X \in \mathbb{R}^{n+1} : x_n \le 0, z = 0 \},\$$

and by

$$L := \{ X \in \mathbb{R}^{n+1} : x_n = 0, z = 0 \}.$$

To each $X\in \mathbb{R}^{n+1}\setminus P^-$ we associate $\tilde{g}_\varepsilon(X)\subset \mathbb{R}$ via the formula

(2.7)
$$U(X) = g(X - \varepsilon w e_n), \quad \forall w \in \tilde{g}_{\varepsilon}(X).$$

Similarly, to each $X \in \mathbb{R}^{n+1} \setminus P^+$ we associate $\tilde{\bar{g}}_{\varepsilon}(X) \subset \mathbb{R}$ via the formula

(2.8)
$$\overline{U}(X) = g(X - \varepsilon w e_n), \quad \forall w \in \tilde{\overline{g}}_{\varepsilon}(X).$$

By abuse of notation, we write $\tilde{g}_{\varepsilon}(X)$, $\tilde{\tilde{g}}_{\varepsilon}(X)$ to denote any of the values in this set. As observed in [15], if g satisfies

(2.9)
$$U(X - \varepsilon e_n) \le g(X) \le U(X + \varepsilon e_n) \quad \text{in } B_{\rho},$$

then for all $\varepsilon > 0$ we can associate with g a possibly multi-valued function \tilde{g}_{ε} defined at least on $B_{\rho-\varepsilon} \setminus P^-$ and taking values in [-1,1] which satisfies

(2.10)
$$U(X) = g(X - \varepsilon \tilde{g}_{\varepsilon}(X)e_n).$$

Moreover if g is strictly monotone in the e_n direction in $B^+_{\rho}(g)$, then \tilde{g}_{ε} is single-valued. A similar statement holds for \tilde{g}_{ε} , when g satisfies the flatness assumption

(2.11)
$$U(X + \varepsilon e_n) \le g(X) \le U(X - \varepsilon e_n) \quad \text{in } B_{\rho}.$$

The following elementary lemmas hold. The proof of the first one can be found in [15]. The second one can be obtained similarly.

Lemma 2.8. Let g, v be non-negative continuous functions in B_{ρ} . Assume that g satisfies the flatness condition (2.9) in B_{ρ} and that v is strictly increasing in the e_n direction in $B_{\rho}^+(v)$. Then if

$$v \leq g$$
 in B_{ρ} ,

and \tilde{v}_{ε} is defined on $B_{\rho-\varepsilon}\setminus P^-$ we have that

$$\tilde{v}_{\varepsilon} \leq \tilde{g}_{\varepsilon} \quad \text{on } B_{\rho-\varepsilon} \setminus P^-.$$

Viceversa, if \tilde{v}_{ε} is defined on $B_s \setminus P^-$ and

$$\tilde{v}_{\varepsilon} \leq \tilde{g}_{\varepsilon}$$
 on $B_s \setminus P^-$,

then

$$v \leq g$$
 on $B_{s-\varepsilon}$.

Lemma 2.9. Let g, v be non-negative continuous functions in B_{ρ} . Assume that g satisfies the flatness condition (2.11) in B_{ρ} and that v is strictly decreasing in the e_n direction in $B_{\rho}^+(v)$. Then if

$$v \geq g$$
 in B_{ρ} ,

and $\tilde{\bar{v}}_{\varepsilon}$ is defined on $B_{\rho-\varepsilon} \setminus P^+$ we have that

$$\tilde{\bar{v}}_{\varepsilon} \leq \tilde{\bar{g}}_{\varepsilon}$$
 on $B_{\rho-\varepsilon} \setminus P^+$.

Viceversa, if \tilde{v}_{ε} is defined on $B_s \setminus P^+$ and

$$\tilde{\bar{v}}_{\varepsilon} \leq \tilde{\bar{g}}_{\varepsilon} \quad \text{on } B_s \setminus P^+,$$

then

$$v \ge g$$
 on $B_{s-\varepsilon}$.

We now state and prove a key comparison principle, which will follow immediately from the lemmas above and Corollary 2.6.

Lemma 2.10. Let (u_1, u_2) , (v_1, v_2) be respectively a solution and a comparison subsolution to (2.1) in B_2 , with (v_1, v_2) strictly monotone in the e_n direction in $B_2^+(v_i)$. Assume that u_1 and u_2 satisfy respectively the flatness assumptions (2.9)-(2.11) in B_2 for $\varepsilon > 0$ small and that $(v_1)_{\varepsilon}, (v_2)_{\varepsilon}$ are defined in $B_{2-\varepsilon} \setminus P^{\mp}$ and satisfy

$$|(\widetilde{v_1})_{\varepsilon}| \le C, \quad |(\widetilde{\overline{v_2}})_{\varepsilon}| \le C.$$

If

(2.12)
$$\widetilde{(v_1)}_{\varepsilon} + c \leq \widetilde{(u_1)}_{\varepsilon} \quad in \left(B_{3/2} \setminus \overline{B}_{1/2}\right) \setminus P^-,$$

(2.13)
$$(\widetilde{v_2})_{\varepsilon} + c \leq (\widetilde{\overline{u_2}})_{\varepsilon} \quad in \left(B_{3/2} \setminus \overline{B}_{1/2}\right) \setminus P^+$$

then

(2.14)
$$\widetilde{(v_1)}_{\varepsilon} + c \leq \widetilde{(u_1)}_{\varepsilon} \quad \text{in } B_{3/2} \setminus P^-,$$

and

(2.15) $\widetilde{(\overline{v_2})}_{\varepsilon} + c \leq \widetilde{(\overline{u_2})}_{\varepsilon} \quad in \ B_{3/2} \setminus P^+.$

Proof. We wish to apply Corollary 2.6 to the functions (u_1, u_2) and

 $v_{i,t} = v_i(X + \varepsilon t e_n).$

We need to verify that for some $t_0 < t_1 = c$

 $(2.16) v_{1,t_0} \le u_1 \quad v_{2,t_0} \ge u_2 \quad \text{in } \overline{B}_1,$

and for all $\delta > 0$ and small

 $(2.17) v_{1,t_1-\delta} \le u_1 \quad \text{on } \partial B_1, \quad v_{1,t_1-\delta} < u_1 \quad \text{on } \mathcal{F}(v_{1,t_1-\delta}),$

$$(2.18) v_{2,t_1-\delta} \ge u_2 \quad \text{on } \partial B_1, \quad v_{2,t_1-\delta} > u_2 \quad \text{on } \mathcal{F}(u_2).$$

Then our Corollary implies

$$\begin{aligned} v_{1,t_1-\delta} &\leq u_1 \quad \text{in } \overline{B}_1, \\ v_{2,t_1-\delta} &\geq u_2 \quad \text{in } \overline{B}_1. \end{aligned}$$

By letting δ go to o, we obtain that

$$v_{1,t_1} \leq u_1, \quad v_{2,t_1} \geq u_2 \quad \text{in } \overline{B}_1,$$

which in view of Lemma 2.8 gives

$$\begin{split} \widetilde{(v_{1,t_{1}})_{\varepsilon}} &\leq \widetilde{(u_{1})}_{\varepsilon} \quad \text{in } B_{1-\varepsilon} \setminus P^{-}, \\ \widetilde{(v_{2,t_{1}})}_{\varepsilon} &\leq \widetilde{(u_{2})}_{\varepsilon} \quad \text{in } B_{1-\varepsilon} \setminus P^{+}, \end{split}$$

assuming that the ε -domain variations on the left hand side exist on $B_{1-\varepsilon} \setminus P^{\mp}$. On the other hand, it is easy to verify that on such set

(2.19)
$$\widetilde{(v_{1,t})}_{\varepsilon}(X) = (\widetilde{v_1})_{\varepsilon}(X) + t, \quad \widetilde{(\overline{v_{2,t}})}_{\varepsilon}(X) = \widetilde{(v_2)}_{\varepsilon}(X) + t,$$

and hence we have for $t_1 = c/2$,

$$\widetilde{(v_1)}_{\varepsilon} + t_1 \leq \widetilde{(v_1)}_{\varepsilon} + c \leq \widetilde{(u_1)}_{\varepsilon} \quad \text{in } B_{1-\varepsilon} \setminus P^-,$$

$$\widetilde{(v_2)}_{\varepsilon} + t_1 \leq \widetilde{(v_2)}_{\varepsilon} + c \leq \widetilde{(u_2)}_{\varepsilon} \quad \text{in } B_{1-\varepsilon} \setminus P^+,$$

which gives the desired conclusion.

We are left with the proof of (2.16)-(2.17)-(2.18).

In view of Lemmas 2.8-2.9, in order to obtain (2.16) it suffices to show that

$$\widetilde{(v_{1,t_0})}_{\varepsilon} \leq \widetilde{(u_1)}_{\varepsilon}, \quad \text{in } B_{1+\varepsilon} \setminus P^-,$$
$$\widetilde{(v_{2,t_0})}_{\varepsilon} \leq \widetilde{(u_2)}_{\varepsilon}, \quad \text{in } B_{1+\varepsilon} \setminus P^+,$$

which by (2.19) becomes

$$\widetilde{(v_1)}_{\varepsilon} + t_0 \le \widetilde{(u_1)}_{\varepsilon}, \quad \text{in } B_{1+\varepsilon} \setminus P^-,$$

$$\widetilde{(v_2)}_{\varepsilon} + t_0 \le \widetilde{(u_2)}_{\varepsilon}, \quad \text{in } B_{1+\varepsilon} \setminus P^+.$$

These last inequalities hold trivially for an appropriate choice of t_0 since the functions involved are bounded.

For (2.17)-(2.18), notice that the first inequality follows easily from our assumption (2.12)-(2.13) together with (2.19) and Lemmas 2.8-2.9. More precisely we have that for all $\delta \leq c/2$, the stronger statement

$$(2.20) v_{1,t_1} \le u_1 \quad v_{2,t_1+\delta} \ge u_2 \quad \text{in } B_{\frac{3}{2}-\varepsilon} \setminus B_{\frac{1}{2}+\varepsilon},$$

holds.

In particular, from the strict monotonicity of v_1 in the e_n -direction in $B_2^+(v_1)$ we have that

$$v_{1,t_1} > 0 \quad \text{on } \mathcal{F}(v_{1,t_1-\delta}),$$

which combined with the previous inequality gives that

$$\iota_1 > 0 \quad \text{on } \mathcal{F}(v_{1,t_1-\delta}),$$

that is the second condition in (2.17). Similarly the second inequality in (2.17) follows from the second inequality in (2.20) and the strict monotonicity of v_2 in $B_2^+(v_2)$.

Finally, given $\varepsilon > 0$ small and Lipschitz functions $\tilde{\varphi}, \tilde{\tilde{\varphi}}$ defined on $B_{\rho}(\bar{X})$, with values in [-1, 1], then there exists a unique pair $(\varphi_{\varepsilon}, \bar{\varphi}_{\varepsilon})$ defined at least on $B_{\rho-\varepsilon}(\bar{X})$ such that

- (2.21) $U(X) = \varphi_{\varepsilon}(X \varepsilon \tilde{\varphi}(X)e_n), \quad X \in B_{\rho}(\bar{X}),$
- (2.22) $\overline{U}(X) = \overline{\varphi}_{\varepsilon}(X \varepsilon \widetilde{\overline{\varphi}}(X)e_n), \quad X \in B_{\rho}(\overline{X}).$

It is readily seen that if (u_1, u_2) satisfies the flatness assumption (2.6) in B_1 then (say $\rho, \varepsilon < 1/4, \bar{X} \in B_{1/2}$,)

(2.24)
$$\tilde{\varphi} \leq (\tilde{u}_2)_{\varepsilon} \quad \text{in } B_{\rho}(\bar{X}) \setminus P^+ \Rightarrow \bar{\varphi}_{\varepsilon} \geq u_2 \quad \text{in } B_{\rho-\varepsilon}(\bar{X})_{\varepsilon}$$

The following proposition is contained in [16]. An analogous statement can be clearly obtained for $\bar{\varphi}_{\varepsilon}$.

Proposition 2.11. Let φ be a smooth function in $B_{\lambda}(\bar{X}) \subset \mathbb{R}^{n+1} \setminus P^-$. Define (for $\varepsilon > 0$ small) the function φ_{ε} as above by

(2.25)
$$U(X) = \varphi_{\varepsilon}(X - \varepsilon \varphi(X)e_n).$$

Then,

(2.26)
$$\Delta \varphi_{\varepsilon} = \varepsilon \Delta (U_n \varphi) + O(\varepsilon^2), \quad \text{in } B_{\lambda/2}(\bar{X})$$

with the function in $O(\varepsilon^2)$ depending on $\|\varphi\|_{C^5}$ and λ .

3. The linearized problem.

We introduce here the linearized problem associated with (2.1). Given $g \in C(B_1)$ and $X_0 = (x'_0, 0, 0) \in B_1 \cap L$, we call

$$g_r(X_0) := \lim_{(x_n, z) \to (0, 0)} \frac{g(x'_0, x_n, z) - g(x'_0, 0, 0)}{r}, \quad r^2 = x_n^2 + z^2.$$

Once the change of unknowns (2.7)-(2.8) has been done, the linearized problem associated with (2.1) is

(3.1)
$$\begin{cases} \Delta(U_n g_1) = 0, & \text{in } B_1 \setminus P^-, \\ \Delta(-\bar{U}_n g_2) = 0, & \text{in } B_1 \setminus P^+, \\ g_1 = g_2, & (g_1)_r + (g_2)_r = 0, & \text{on } B_1 \cap L \end{cases}$$

Definition 3.1. We say that (g_1, g_2) is a viscosity solution to (3.1) if $g_i \in C(B_1)$, g_i is even-in-z and it satisfies

(i) $\Delta(U_n g_1) = 0$ in $B_1 \setminus P^-$; (ii) $\Delta(-\bar{U}_n g_2) = 0$ in $B_1 \setminus P^+$; (iii) $g_1 = g_2$ on $B_1 \cap L$;

(iv) Let ϕ_i be continuous around $X_0 = (x'_0, 0, 0) \in B_1 \cap L$ and satisfy

$$\phi_1(X) = l(x') + b_1(X_0)r + O(|x' - x'_0|^2 + r^{3/2}),$$

$$\phi_2(X) = l(x') + b_2(X_0)r + O(|x' - x'_0|^2 + r^{3/2}),$$

with $l(x') = a_0 + a_1 \cdot (x' - x'_0)$. If $b_1(X_0) + b_2(X_0) > 0$ then (ϕ_1, ϕ_2) cannot touch (g_1, g_2) by below at X_0 , and if $b_1(X_0) + b_2(X_0) < 0$ then (ϕ_1, ϕ_2) cannot touch (g_1, g_2) by above at X_0 .

We wish to prove the following existence and regularity result.

Theorem 3.2. Given a boundary data (h_1, h_2) with $h_i \in C(\overline{B}_1), |h_i| \leq 1, h_i$ even-in-z and $h_1 = h_2$ on $B_1 \cap L$, there exists a unique classical solution (g_1, g_2) to (3.1) such that $g_i \in C(\overline{B}_1), g_i = h_i$ on $\partial B_1, g_i$ is even-in-z and it satisfies

(3.2)
$$|g_i(X) - a_0 - a' \cdot (x' - x'_0) - b_i r| \le C(|x' - x'_0|^2 + r^{3/2}), \quad X_0 \in B_{1/2} \cap L,$$

for a universal constants C , and $a' \in \mathbb{R}^{n-1}, a_0, b_i \in \mathbb{R}$ depending on X_0 and satisfying

$$b_1 + b_2 = 0.$$

As a corollary of the theorem above we obtain the following regularity result.

Theorem 3.3 (Improvement of flatness). There exists a universal constant C such that if (g_1, g_2) is a viscosity solution to (3.1) in B_1 with

$$-1 \le g_i(X) \le 1 \quad \text{in } B_1,$$

then

$$(3.3) a_0 + a' \cdot x' + b_i r - C|X|^{3/2} \le g_i(X) \le a_0 + a' \cdot x' + b_i r + C|X|^{3/2},$$
 for some $a' \in \mathbb{R}^{n-1}$, $a_0, b_i \in \mathbb{R}$ such that

 $b_1 + b_2 = 0.$

Proof. Let (w_1, w_2) be the unique classical solution to (3.1) in $B_{1/2}$ with boundary data (g_1, g_2) . We will prove that $w_i = g_i$ in $B_{1/2}$ and hence it satisfies the desired estimate in view of (3.2). Denote by

$$\bar{w}^i_{\varepsilon} := w_i - \varepsilon + \varepsilon^2 r.$$

Then, for ε small

$$\bar{w}_{\varepsilon}^i < g_i \quad \text{on } \partial B_{1/2}.$$

We wish to prove that

Now, notice that $(\bar{w}^1_{\varepsilon}, \bar{w}^2_{\varepsilon})$ (and all its translations) is a classical strict subsolution to (3.1) that is

(3.5)
$$\begin{cases} \Delta(U_n \bar{w}_{\varepsilon}^1) = 0, & \text{in } B_{1/2} \setminus P^-, \\ \Delta(-\bar{U}_n \bar{w}_{\varepsilon}^2) = 0, & \text{in } B_{1/2} \setminus P^+, \\ \bar{w}_{\varepsilon}^1 = \bar{w}_{\varepsilon}^2, & (\bar{w}_{\varepsilon}^1)_r + (\bar{w}_{\varepsilon}^2)_r > 0, & \text{on } B_{1/2} \cap L. \end{cases}$$

Since g_i is bounded, for t large enough $\bar{w}_{\varepsilon}^i - t$ lies strictly below g_i . We let $t \to 0$ and show that the first contact point cannot occur for $t \ge 0$. Indeed since $\bar{w}_{\varepsilon}^i - t$ is a strict subsolution which is strictly below g_i on $\partial B_{1/2}$ then no touching can occur either in

 $B_{1/2} \setminus P^{\mp}$ or on $B_{1/2} \cap L$. We only need to check that no touching occurs on $P^{\mp} \setminus L$. This follows from Lemma A.4 in the Appendix.

Thus (3.4) holds. Passing to the limit as $\varepsilon \to 0$ we deduce that

$$w_i \leq g_i \quad \text{in } B_{1/2}.$$

Similarly we also infer that

$$w_i \ge g_i \quad \text{in } B_{1/2},$$

and the desired equality holds.

We remark that there is no theory readily available for this class of degenerate problems with a Neumann type boundary condition on the thin boundary L, which in this setting has positive capacity (see the classical work on degenerate elliptic problems [?]). Thus Theorem 3.3 calls for a direct proof.

The existence of the classical solution of Theorem 3.2 will be achieved via a variational approach in the appropriate functional space. Precisely, define the weighted Sobolev space

$$H^{1}_{\omega}(B_{1}) := H^{1}(B_{1}, U^{2}_{n}dX) \times H^{1}(B_{1}, \bar{U}^{2}_{n}dX)$$

endowed with the norm

$$\|(h_1,h_2)\| := \int_{B_1} (h_1^2 + |\nabla h_1|^2) U_n^2 dX + \int_{B_1} (h_2^2 + |\nabla h_2|^2) \bar{U}_n^2 dX$$

and its subspace

$$V_0(B_1) := \{(h_1, h_2) : h_i \in C_0^{\infty}(B_1), h_1 = h_2 \text{ on } B_1 \cap L\}.$$

Let $\mathcal{H}_{0,\omega}^1(B_1) := \overline{V_0(B_1)}$ be the completion of $V_0(B_1)$ in $H_{\omega}^1(B_1)$. Consider the energy functional

$$J(h_1, h_2) := \int_{B_1} (U_n^2 |\nabla h_1|^2 + \bar{U}_n^2 |\nabla h_2|^2) dX, \quad (h_1, h_2) \in H^1_{\omega}(B_1).$$

Given a boundary data (h_1, h_2) with $h_i \in C^{\infty}(\bar{B}_1)$, and $h_1 = h_2$ on $\bar{B}_1 \cap L$, we minimize $J(h_1 + \varphi_1, h_2 + \varphi_2)$

among all $(\varphi_1, \varphi_2) \in \mathcal{H}^1_{0,\omega}(B_1)$. Equivalently, a minimizing pair (g_1, g_2) must satisfy

$$J(g_1, g_2) \le J(g_1 + \varphi_1, g_2 + \varphi_2), \quad \forall \varphi_i \in C_0^{\infty}(B_1), i = 1, 2,$$

with

(3.6)
$$\varphi_1 = \varphi_2 \quad \text{on } B_1 \cap L.$$

Since J is strictly convex, we conclude that (g_1, g_2) is a minimizer if and only if

$$\lim_{\varepsilon \to 0} \frac{J(g_1, g_2) - J(g_1 + \varepsilon \varphi_1, g_2 + \varepsilon \varphi_2)}{\varepsilon} = 0, \quad \forall \varphi_i \in C_0^\infty(B_1), i = 1, 2, \ldots$$

satisfying (3.6), or equivalently,

(3.7)
$$\int_{B_1} (U_n^2 \nabla g_1 \cdot \nabla \varphi_1 + \bar{U}_n^2 \nabla g_2 \cdot \nabla \varphi_2) dx = 0, \quad \forall \varphi_i \in C_0^\infty(B_1), i = 1, 2,$$

satisfying (3.6).

Lemma 3.4. Let (g_1, g_2) be a minimizer to J in B_1 , then

$$\Delta(U_n g_1) = 0 \quad in B_1 \setminus P^-,$$

$$\Delta(-\bar{U}_n g_2) = 0 \quad in B_1 \setminus P^+.$$

Proof. By (3.7)

$$\int_{B_1} (U_n^2 \nabla g_1 \cdot \nabla \varphi^- + \bar{U}_n^2 \nabla g_2 \cdot \nabla \varphi^+) dx = 0, \quad \forall \varphi^{\mp} \in C_0^{\infty}(B_1 \setminus P^{\mp}).$$

After integration by parts we obtain,

$$\begin{split} &\operatorname{div}(U_n^2 \nabla g_1) = 0 \quad \text{in } B_1 \setminus P^-, \\ &\operatorname{div}(\bar{U}_n^2 \nabla g_2) = 0 \quad \text{in } B_1 \setminus P^+. \end{split}$$

Since
$$U_n, \overline{U}_n$$
 are bounded in $B_1 \setminus P^-, B_1 \setminus P^+$ respectively, we conclude by elliptic regularity that

$$g_1 \in C^{\infty}(B_1 \setminus P^-), g_2 \in C^{\infty}(B_1 \setminus P^+).$$

Now a simple computation concludes the proof. Precisely,

$$\operatorname{div}(U_n^2 \nabla g_1) = U_n^2 \Delta g_1 + 2 \sum_{i=1}^{n+1} U_n U_{ni}(g_1)_i = 0 \quad \text{in } B_1 \setminus P^-.$$

Since $U_n>0$ and $\Delta U=0$ in $B_1\setminus P^-$ the identity above is equivalent to

$$\Delta(U_n g_1) = U_n \Delta g_1 + 2\nabla U_n \cdot \nabla g_1 = 0 \quad \text{in } B_1 \setminus P^-,$$

as desired. The computation for g_2 is obtained analogously.

Lemma 3.5. Let $(g_1, g_2) \in C(B_1)$ be a solution to

$$\Delta(U_n g_1) = 0 \quad in \ B_1 \setminus P^-,$$

$$\Delta(-\bar{U}_n g_2) = 0 \quad in \ B_1 \setminus P^+,$$

with $g_1 = g_2$ on $B_1 \cap L$ and assume that

$$\lim_{r \to 0} (g_i)_r(x', x_n, z) = b_i(x'), \quad (g_i)_r(X) = \frac{x_n}{r} g_{x_n}(X) + \frac{z}{r} g_z(X);$$

with $b_i(x')$ a continuous function. Then (g_1, g_2) is a minimizer to J in B_1 if and only if $b_1 + b_2 \equiv 0$.

Proof. The pair (g_1, g_2) minimizes J if and only if

$$\int_{B_1} (U_n^2 \nabla g_1 \cdot \nabla \varphi_1 + \bar{U}_n^2 \nabla g_2 \cdot \nabla \varphi_2) dx = 0,$$

for all $\varphi_i \in C_0^\infty(B_1)$ which coincide on L. We compute that

$$(3.8) \quad \int_{B_1} U_n^2 \nabla g_1 \cdot \nabla \varphi_1 dx = \int_{B_1 \setminus P^-} \operatorname{div}(U_n^2 g_1) \varphi_1 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} U_n^2 \varphi_1 \nabla g_1 \cdot \nu d\sigma$$

where $C_{\delta} = \{r \leq \delta\}$ and ν is the inward normal to C_{δ} , and similarly

$$(3.9) \quad \int_{B_1} \bar{U}_n^2 \nabla g_2 \cdot \nabla \varphi_2 dx = \int_{B_1 \setminus P^+} \operatorname{div}(\bar{U}_n^2 g_2) \varphi_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma_2 dx + \lim_{\delta \to 0} \int_{\partial C_\delta \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nabla g_2 \nabla$$

Indeed, say for g_1 , given $\delta > 0$, for P_{ε} a strip of width ε around P^-

$$\int_{B_1 \setminus C_{\delta}} U_n^2 \nabla g_1 \cdot \nabla \varphi_1 dx = \lim_{\varepsilon \to 0} \int_{(B_1 \setminus C_{\delta}) \cap (B_1 \setminus P_{\varepsilon})} U_n^2 \nabla g_1 \cdot \nabla \varphi_1 dx$$

and integrating by parts

$$\int_{B_1 \setminus C_{\delta}} U_n^2 \nabla g_1 \cdot \nabla \varphi_1 dx = \int_{(B_1 \setminus C_{\delta}) \setminus P^-} \operatorname{div}(U_n^2 g_1) \varphi_1 dx + \int_{\partial C_{\delta} \cap B_1} U_n^2 \varphi_1 \nabla g_1 \cdot \nu d\sigma$$

14

$$+\lim_{\varepsilon\to 0}\int_{\partial P_{\varepsilon}\cap (B_1\setminus C_{\delta})}U_n^2\varphi_1\nabla g_1\cdot\nu d\sigma.$$

It is easily seen that the last limit goes to zero, hence our claim is proved. In fact, in the region of integration we have

$$U_n \leq C\varepsilon,$$

and

$$|\nabla(U_n g_1)|, |\nabla U_n| \le C,$$

from which it follows that

$$|U_n|\nabla g_1| \le C.$$

Finally using the formula for U we can compute that

(3.10)
$$\lim_{\delta \to 0} \int_{\partial C_{\delta} \cap B_1} U_n^2 \varphi_1 \nabla g_1 \cdot \nu d\sigma = \pi \int_L b_1(x') \varphi_1(x', 0, 0) dx',$$

(3.11)
$$\lim_{\delta \to 0} \int_{\partial C_{\delta} \cap B_1} \bar{U}_n^2 \varphi_2 \nabla g_2 \cdot \nu d\sigma = \pi \int_L b_2(x') \varphi_1(x',0,0) dx'.$$

Precisely, say for g_1 ,

$$\int_{\partial C_{\delta} \cap B_{1}} U_{n}^{2} \varphi_{1} \nabla g_{1} \cdot \nu d\sigma = \frac{1}{\delta} \int_{\partial C_{\delta} \cap B_{1}} \cos^{2}(\frac{\theta}{2})(g_{1})_{r} \varphi_{1} d\sigma$$
$$= \int_{\partial C_{1} \cap B_{1}} \cos^{2}(\frac{\theta}{2})((g_{1})_{r} \varphi_{1})(X', \delta \cos \theta, \delta \sin \theta) dx' d\theta,$$

and the desired equality follows.

Combining (3.8)-(3.9)-(3.10)-(3.11) with the fact that $\varphi_2 = \varphi_2$ on L, and the computation at the end of Lemma 3.4, we conclude the proof.

The key ingredient in showing the regularity of the minimizing pair is the following Harnack inequality.

Lemma 3.6 (Harnack inequality). Let (g_1, g_2) be a minimizer to J in B_1 among competitors which are even-in-z, and assume $|g_i| \leq 1$. Then $g_i \in C^{\alpha}(B_{1/2})$ and

$$[g_i]_{C^{0,\alpha}(B_{1/2})} \le C,$$

with C universal.

Indeed, since our linear problem is invariant under translations in the x'-direction, we see that discrete differences of the form

$$g_i(X+\tau) - g_i(X),$$

with τ in the x'-direction are also minimizers. Thus by standard arguments (see [2]) we obtain the following corollary.

Corollary 3.7. Let (g_1, g_2) be as in Lemma 3.6. Then $D_{x'}^{\beta}g_i \in C^{0,\alpha}(B_{1/2})$ and

$$[D_{x'}^{\beta}g_i]_{C^{\alpha}(B_{1/2})} \le C,$$

with C depending on β .

The proof of Harnack inequality relies on the following comparison principle for minimizers.

Lemma 3.8 (Comparison Principle). Let (g_1, g_2) and (h_1, h_2) be minimizers to J in B_1 . If

$$g_i \ge h_i$$
 a.e in $B_1 \setminus B_\rho$, $i = 1, 2$,

then

$$g_i \ge h_i$$
 a.e. in B_1 , $i = 1, 2$.

Proof. Call

$$M_i := \max\{g_i, h_i\}, \quad m_i := \min\{g_i, h_i\}, \quad i = 1, 2.$$

We have that

$$(3.12) J(M_1, M_2) + J(m_1, m_2) = J(g_1, g_2) + J(h_1, h_2).$$

On the other hand, since (g_1, g_2) and (h_1, h_2) are minimizing pairs, and under our assumptions $(M_1, M_2), (m_1, m_2)$ are admissible competitors,

$$J(M_1, M_2) \ge J(g_1, g_2), \quad J(m_1, m_2) \ge J(h_1, h_2).$$

We conclude that

$$J(M_1, M_2) = J(g_1, g_2), \quad J(m_1, m_2) = J(h_1, h_2),$$

and by uniqueness,

$$M_i = g_i, \quad m_i = h_i, \quad i = 1, 2.$$

Our desired result follows.

Proof of Lemma 3.6. The key step consists in proving the following claim. The remaining ingredients are the standard Harnack inequality and Boundary Harnack inequality for harmonic functions.

Claim: There exist universal constants η, c such that if $g_i \ge 0$ a.e. in B_1 and

$$g_1(\frac{1}{4}e_n) \ge 1,$$

then

$$a_i \ge c$$
 a.e. in B_η .

By boundary Harnack inequality, since U_ng_1 and U_n are non-negative and harmonic in $B_1 \setminus P^-$ and vanish on P^- we conclude that

$$g_1 \ge c$$
 a.e. on $C := \{ |x'| \le 3/4, 1/8 \le r \le 1/2 \}$

. Let

$$v_1(X) := -\frac{|x'|^2}{n-1} + 2(x_n+1)r, \quad v_2(X) := -\frac{|x'|^2}{n-1} - 2(x_n+1)r$$

We show that (v_1, v_2) is a minimizer to J in B_1 . To do so we prove that (v_1, v_2) satisfies Lemma 3.5.

To prove that

$$\Delta(U_n v_1) = 0 \quad \text{in } B_1 \setminus P^-,$$

we use that $2rU_n=U$ and that U,U_n are harmonic outside of P^- and do not depend on $x^\prime.$ Thus

$$\Delta(U_n v_1) = -\Delta(\frac{|x'|^2}{n-1}U_n) + \Delta((x_n+1)U) = -2U_n + 2U_n = 0.$$

We argue similarly for v_2 .

Finally the fact that

$$\lim_{r \to 0} (v_1)_r + (v_2)_r = 0,$$

follows immediately from the definition of v_i , i = 1, 2.

Now by choosing δ small depending on c we can guarantee that

$$\begin{split} v_1 + \delta &\leq c \leq g_1, \quad v_2 + \delta \leq 0 \leq g_2 \quad \text{a.e. on } \{ |x'| \leq 3/4, \delta/2 \leq r \leq \delta \} \\ v_i + \delta &\leq 0 \leq g_i, \quad i = 1, 2 \quad \text{on } \{ |x'| = 3/4, r \leq \delta \}. \end{split}$$

By our comparison principle, Lemma 3.8 we conclude that

$$v_i + \delta \leq g_i$$
 a.e in $\{|x'| \leq 3/4, r \leq \delta\}$

Since $v_i + \delta$ is a continuous function positive at zero, our claim follows immediately. \Box

Proof of Theorem 3.2. Let (g_1, g_2) be a minimizer to J in which achieves the boundary data (h_1, h_2) . By approximation, we can assume that h_i is smooth. Then by Lemma 3.4, (g_1, g_2) satisfy

$$\Delta(U_n g_1) = 0 \quad \text{in } B_1 \setminus P^-,$$

$$\Delta(-\bar{U}_n g_2) = 0 \quad \text{in } B_1 \setminus P^+,$$

and in view of Lemma 3.6 and Corollary 3.7, $\Delta_{x'}g_i \in C^{0,\alpha}(B_1)$ with universally bounded norm in say $B_{3/4}$. Thus, it is shown in Theorem 8.1 in [15] that (after reflecting g_2)

(3.13)
$$|g_i(x', x_n, z) - g_i(x', 0, 0) - b_i(x')r| = O(r^{3/2}), \quad (x', 0, 0) \in B_{1/2} \cap L,$$

with b_i Lipschitz and moreover,

$$(3.14) |(g_i)_r(x',x_n,z) - b_i(x')| \le Cr^{1/2}, \quad (x',0,0) \in B_{1/2} \cap L,$$

with C universal. Hence by Lemma 3.5,

$$(g_1)_r(x') + (g_2)_r(x') = b_1(x') + b_2(x') = 0$$
 on $B_{1/2} \cap L$.

For completeness, we sketch the argument to obtain (3.13)-(3.14), say for g_1 . Since g_1 solves

$$\Delta(U_n g_1) = 0 \quad \text{in } B_1 \setminus P^-$$

and U_n is independent on x^\prime we can rewrite this equation as

$$\Delta_{x_n,z}(U_ng_1) = -U_n\Delta_{x'}g_1$$

and according to Corollary 3.7 we have that

$$\Delta_{x'}g_1 \in C^{\alpha}(B_{1/2}),$$

with universal bound. Thus, for each fixed x', we need to investigate the 2-dimensional problem

$$\Delta(U_t g_1) = U_t f, \quad \text{in } B_{1/2} \setminus \{t \le 0, z = 0\}$$

with

$$f \in C^{\alpha}(B_{1/2}),$$

and g_1 bounded. Without loss of generality, for a fixed x' we may assume $g_1(x', 0, 0) = 0$. Let H(t, z) be the solution to the problem

$$\Delta H = U_t f, \quad \text{in } B_{1/2} \setminus \{t \le 0, z = 0\},$$

such that

$$H = U_t h$$
 on $\partial B_{1/2}$, $H = 0$ on $B_{1/2} \cap \{t \le 0, z = 0\}$.

We wish to prove that

$$U_t h = H.$$

First notice that

$$\Delta(H - U_t h) = 0 \quad \text{in } B_{1/2} \setminus \{ t \le 0, z = 0 \},\$$

and

(3.16)

$$H = 0 \quad \text{on } \partial B_{1/2} \cup (B_{1/2} \cap \{t < 0, z = 0\}).$$

We claim that

(3.17)
$$\lim_{(t,z)\to(0,0)}\frac{H-U_th}{U_t} = \lim_{(t,z)\to(0,0)}\frac{H}{U_t} = 0.$$

If the claims holds, then given any $\varepsilon > 0$

$$-\varepsilon U_t \le H - U_t h \le \varepsilon U_t, \quad \text{in } B_\delta$$

with $\delta = \delta(\varepsilon)$. Then by the maximum principle the inequality above holds in the whole $B_{1/2}$ and by letting $\varepsilon \to 0$ we obtain (3.16).

To prove the claim (3.17) we show that H satisfies the following

(3.18)
$$|H(t,z) - aU(t,z)| \le C_0 r^{1/2} U(t,z), \quad r^2 = t^2 + z^2, a \in \mathbb{R},$$

with C_0 universal.

To do so, we consider the holomorphic transformation

$$\Phi: (s, y) \to (t, z) = (\frac{1}{2}(s^2 - y^2), sy)$$

which maps $B_1 \cap \{s > 0\}$ into $B_{1/2} \setminus \{t \le 0, z = 0\}$ and call

$$\tilde{H}(s,y)=H(t,z),\quad \tilde{f}(s,y)=f(t,z).$$

Then, easy computations show that

$$\Delta \tilde{H} = s \tilde{f} \quad \text{in } B_1 \cap \{s > 0\}, \quad \tilde{H} = 0 \quad \text{on } B_1 \cap \{s = 0\}.$$

Since the right-hand side is C^{α} we conclude that $\tilde{H} \in C^{2,\alpha}$. In particular \tilde{H}_s satisfies

$$|\tilde{H}_s(s,y) - a| \le C_0 |(s,y)|, \quad a = \tilde{H}_s(0,0)$$

with C_0 universal. Integrating this inequality between 0 and s and using that $\tilde{H} = 0$ on $B_1 \cap \{s = 0\}$ we obtain that

$$|\tilde{H}(s,y) - as| \le C_0 s |(s,y)|.$$

In terms of H, this equation gives us

$$(3.19) |H - aU| \le C_0 r^{1/2} U$$

as desired.

Thus (3.16) and (3.18) hold and by combining them and using that $U/U_t=2r$ we deduce

$$|h - 2ar| \le 2C_0 r^{3/2}$$

which is the desired estimate i.e. (recall that above we assumed $g_1(x', 0, 0) = 0$)

$$|g_1(x', x_n, x) - g_1(x', 0, 0) - b_1(x')r| \le 2C_0 r^{3/2},$$

with

$$b_1(x') = 2\tilde{H}_s(x', 0, 0).$$

We remark that b(x') is Lipschitz. Indeed, notice that the derivatives $(g_1)_i$, i = 1, ..., n-1 still satisfy the same equation (3.15) as g_1 , where the C^{α} norm of the right-hand side has a universal bound. Thus, we can argue as above to conclude that

$$\left|\partial_i \tilde{H}_s(x', 0, 0)\right| \le C,$$

which together with the formula for $b_1(x')$ shows that $b_1(x')$ is a Lipschitz function.

To obtain the second of our estimates (3.14) we proceed similarly as above. Since $U_t h = H$ one can compute easily that

$$U_t h_r = H_r + \frac{1}{2} \frac{H}{r}$$

Moreover, after our holomorphic transformation

(3.21)
$$2rH_r(t,z) = sH_s(s,y) + yH_y(s,y)$$

As observed above,

$$|\tilde{H}_s - a(x')| \le C|(s, y)|,$$

and similarly since $\tilde{H} = 0$ on $B_1 \cap \{s = 0\}$

$$|\tilde{H}_y| \le Cs$$

These two inequalities combined with (3.21) give us

(3.22)
$$|2rH_r - a(x')U| \le Cr^{1/2}U.$$

Combining (3.20) with (3.19)-(3.22) we obtain (3.14) as desired.

Finally, since g_i is C^{∞} in the x'-direction and $g_1 = g_2$ on L, we conclude that for a fixed $X_0 \in B_{1/4} \cap L$

$$|g_i(x',0,0) - g_i(x'_0,0,0) - a' \cdot (x' - x'_0)| \le C|x' - x'_0|^2$$

from which (3.2) follows, using that b_i is Lipschitz continuous.

We are left with the proof that the boundary data is achieved continuously. Indeed this follows by classical elliptic theory if we restrict to $\partial B_1 \setminus P^{\mp}$ (for g_1 and g_2 respectively.)

If $X_0 \in \partial B_1 \cap (P^- \setminus L)$ then in a small neighborhood of X_0 intersected with $B_1 \cap \{z > 0\}$ the function $U_n g_1$ is harmonic continuous up to the boundary and vanishes continuously on $\{z = 0\}$ (since h is bounded). The continuity of g_1 at X_0 then follows from standard boundary regularity for the harmonic function $U_n g_1$ after reflecting oddly across z = 0 and using that the boundary data is smooth. We argue similarly if $X_0 \in \partial B_1 \cap (P^+ \setminus L)$ and we must show the continuity of g_2 at X_0 .

Finally, on the set $\partial B_1 \cap L$ as in the case of Laplace equation, it suffices to construct at each point X_0 a local barrier pair (minimizing pair) for (g_1, g_2) which is zero at X_0 and strictly negative in a neighborhood of X_0 . Such barrier is given by a multiple of (see Lemma 3.5)

$$((x' - x'_0) \cdot x'_0, (x' - x'_0) \cdot x'_0).$$

4. HARNACK INEQUALITY

This section is devoted to determine a Harnack type inequality for solutions to our free boundary problem (2.1).

Theorem 4.1 (Harnack inequality). There exists $\bar{\varepsilon} > 0$ such that if (u_1, u_2) solves (2.1) and it satisfies in $B_{\rho}(X^*)$,

$$\begin{split} U(X+\varepsilon a_0e_n) \leq u_1(X) \leq U(X+\varepsilon b_0e_n), \quad \bar{U}(X+\varepsilon \bar{a}_0e_n) \leq u_2(X) \leq \bar{U}(X+\varepsilon \bar{b}_0e_n), \\ \text{with} \end{split}$$

$$\varepsilon(b_0 - a_0) \le \overline{\varepsilon}\rho, \quad \varepsilon(\overline{a}_0 - \overline{b}_0) \le \overline{\varepsilon}\rho,$$

then

$$U(X + \varepsilon a_1 e_n) \le u_1(X) \le U(X + \varepsilon b_1 e_n) \quad \text{in } B_{\eta\rho}(X^*),$$

$$\bar{U}(X + \varepsilon \bar{a}_1 e_n) \le u_2(X) \le \bar{U}(X + \varepsilon \bar{b}_1 e_n) \quad \text{in } B_{\eta\rho}(X^*),$$

with

$$a_0 \le a_1 \le b_1 \le b_0, \quad (b_1 - a_1) \le (1 - \eta)(b_0 - a_0), \\ \bar{b}_0 \le \bar{b}_1 \le \bar{a}_1 \le \bar{a}_0, \quad (\bar{a}_1 - \bar{b}_1) \le (1 - \eta)(\bar{a}_0 - \bar{b}_0),$$

for a small universal constant η . Moreover, if $B_{\rho}(X^*) \cap F(u_1, u_2) \neq \emptyset$, and $a_0 = \bar{b}_0, b_0 = \bar{a}_0$ then the conclusion holds with $a_1 = \bar{b}_1, b_1 = \bar{a}_1$.

Let (u_1, u_2) be as in the Theorem above, and denote A_{ε}^{\mp} be the following sets

(4.1)
$$A_{\varepsilon}^{-} := \{ (X, (\tilde{u}_{1})_{\varepsilon}(X)) : X \in B_{1-\varepsilon} \setminus P^{-} \} \subset \mathbb{R}^{n+1} \times [a_{0}, b_{0}].$$

(4.2)
$$A_{\varepsilon}^{+} := \{ (X, (\tilde{\tilde{u}}_{2})_{\varepsilon}(X)) : X \in B_{1-\varepsilon} \setminus P^{+} \} \subset \mathbb{R}^{n+1} \times [\bar{b}_{0}, \bar{a}_{0}] \}$$

Since ε -domain variations may be multivalued, we mean that given X all pairs (X, Z) with Z an element of the ε -variation belong to A_{ε}^{\mp} . Applying Theorem 4.1 iteratively we obtain

(4.3)
$$A_{\varepsilon}^{-} \cap (B_{\frac{1}{2}\eta^{m}-\varepsilon}(X^{*}) \times [a_{0},b_{0}]) \subset B_{\frac{1}{2}\eta^{m}-\varepsilon}(X^{*}) \times [a_{m},b_{m}],$$

(4.4)
$$A_{\varepsilon}^{+} \cap (B_{\frac{1}{2}\eta^{m}-\varepsilon}(X^{*}) \times [\bar{b}_{0},\bar{a}_{0}]) \subset B_{\frac{1}{2}\eta^{m}-\varepsilon}(X^{*}) \times [\bar{b}_{m},\bar{a}_{m}],$$

with

(4.5)
$$b_m - a_m \le (b_0 - a_0)(1 - \eta)^m, \quad \bar{a}_m - \bar{b}_m \le (\bar{a}_0 - \bar{b}_0)(1 - \eta)^m$$

for all m's such that both inequalities are satisfied

(4.6)
$$2\varepsilon(1-\eta)^m\eta^{-m}(b_0-a_0) \le \bar{\varepsilon}, \quad 2\varepsilon(1-\eta)^m\eta^{-m}(\bar{a}_0-\bar{b}_0) \le \bar{\varepsilon}.$$

Thus we obtain the following corollary, which will be the key of the improvement of flatness arguments of Section 5.

Corollary 4.2. If

$$\begin{split} U(X-\varepsilon e_n) &\leq u_1(X) \leq U(X+\varepsilon e_n) \quad \text{in } B_1, \\ \bar{U}(X+\varepsilon e_n) &\leq u_2(X) \leq \bar{U}(X-\varepsilon e_n) \quad \text{in } B_1, \end{split}$$

with $\varepsilon \leq \overline{\varepsilon}/2$, given $m_0 > 0$ such that

$$2\varepsilon(1-\eta)^{m_0}\eta^{-m_0} \le \bar{\varepsilon},$$

then the set $A_{\varepsilon}^{\mp} \cap (B_{1/2} \times [-1, 1])$ is above the graph of a function $y = a_{\varepsilon}^{\mp}(X)$ and it is below the graph of a function $y = b_{\varepsilon}^{\mp}(X)$ with

$$b_{\varepsilon}^{\mp} - a_{\varepsilon}^{\mp} \le 2(1-\eta)^{m_0-1},$$

and $a_{\varepsilon}^{\mp}, b_{\varepsilon}^{\mp}$ having a modulus of continuity bounded by the Hölder function αt^{β} for α, β depending only on η . Moreover, $a_{\varepsilon}^{-} = a_{\varepsilon}^{+}, b_{\varepsilon}^{-} = b_{\varepsilon}^{-}$ on $B_{1/2} \cap F(u_1, u_2)$.

The following lemma is the key ingredient in the proof of Theorem 4.1.

Lemma 4.3. There exists $\bar{\varepsilon} > 0$ such that for all $0 < \varepsilon \leq \bar{\varepsilon}$ if (u_1, u_2) is a solution to (2.1) in B_1 such that

(4.7)
$$u_1(X) \ge U(X), \quad u_2(X) \le \overline{U}(X) \quad \text{in } B_{1/2},$$

and at $\bar{X} \in B_{1/8}(\frac{1}{4}e_n)$

(4.8)
$$u_1(\bar{X}) \ge U(\bar{X} + \varepsilon e_n),$$

then

(4.9)
$$u_1(X) \ge U(X + \tau \varepsilon e_n), \quad u_2 \le \overline{U}(X + \tau \varepsilon e_n) \quad \text{in } B_{\delta_2}$$

for universal constants τ , δ .

Lemma 4.3 relies on building an appropriate family of radial subsolutions. Precisely, let R > 0 and denote by

$$V_R(t,z) = U(t,z)((n-1)\frac{t}{R}+1), \quad \bar{V}_R(t,z) = \bar{U}(t,z)((n-1)\frac{t}{R}+1).$$

Then set,

$$v_R(X) = V_R(R - \sqrt{|x'|^2 + (x_n - R)^2}, z)$$

$$\bar{v}_R(X) = \bar{V}_R(R - \sqrt{|x'|^2 + (x_n - R)^2}, z).$$

Finally, for $\beta > 0$, define (for notational simplicity we drop the dependence on β ,)

(4.10)
$$v_1^R = (1 + \frac{\beta}{R})v_R, \quad v_2^R = (1 + \frac{\beta}{R})\bar{v}_R$$

Proposition 4.4. If R is large enough, the ordered pair (v_1^R, v_2^R) is a comparison subsolution to (2.1) in B_2 which is strictly monotone in the e_n -direction outside $\{v_i^R(x, 0) = 0\}$. Moreover, there exist functions $\tilde{v}_1^R, \tilde{v}_2^R$ such that (4.11)

$$U(X) = v_1^R(X - \tilde{v}_1^R(X)e_n), \quad \text{in } B_1 \setminus P^-, \quad \bar{U}(X) = v_2^R(X - \tilde{\tilde{v}}_2^R(X)e_n), \quad \text{in } B_1 \setminus P^+$$

and

(4.12)
$$|\tilde{v}_1^R(X) - \gamma_R(X)| \le \frac{C}{R^2}, \quad |\tilde{\bar{v}}_2^R(X) - \bar{\gamma}_R(X)| \le \frac{C}{R^2},$$

with (4.13)

where

$$\begin{split} &\gamma_R(X) = -\frac{|x'|^2}{2R} + 2(n-1)\frac{x_n r}{R} + 2\beta \frac{r}{R}, \quad \bar{\gamma}_R(X) = -\frac{|x'|^2}{2R} - 2(n-1)\frac{x_n r}{R} - 2\beta \frac{r}{R} \\ &\text{for } r = \sqrt{x_n^2 + z^2} \text{ and } C \text{ depending on } \beta. \end{split}$$

Proof. **Step 1.** We need to show that v_1^R is strictly subharmonic in $B_2^+(v_1^R)$ while v_2^R is strictly superharmonic in $B_2^+(v_2^R)$. We sketch the proof of the claim about v_2^R . The claim about v_1^R can be obtained by similar computations which are already contained in Proposition 6.4 in [15] to which we refer for further details.

On the set $B_2^+(v_2^R)$,

$$\Delta \bar{v}_R(X) = ((\partial_{tt} + \partial_{zz}))\bar{V}_R(R - \rho, z) - \frac{n-1}{\rho}\partial_t \bar{V}_R(R - \rho, z),$$

 $\rho := \sqrt{|x'|^2 + (x_n - R)^2}.$

Also for (t, z) outside the set $\{(t, 0) : t \ge 0\}$

$$\begin{aligned} \Delta_{t,z}\bar{V}_R(t,z) &= (\partial_{tt} + \partial_{zz})\bar{V}_R(t,z) = \frac{2(n-1)}{R}\partial_t\bar{U}(t,z) + (1+(n-1)\frac{t}{R})\Delta_{t,z}\bar{U}(t,z) \\ &= \frac{2(n-1)}{R}\partial_t\bar{U}(t,z), \end{aligned}$$

and

(4.14)
$$\partial_t \bar{V}_R(t,z) = (1 + (n-1)\frac{t}{R})\partial_t \bar{U}(t,z) + \frac{n-1}{R}\bar{U}(t,z).$$

Thus we need to prove that in $B_2^+(v_2^R)$

$$\frac{2(n-1)}{R}\partial_t \bar{U} - \frac{n-1}{\rho} [(1+(n-1)\frac{R-\rho}{R})\partial_t \bar{U} + \frac{n-1}{R}\bar{U}] < 0,$$

where \overline{U} and $\partial_t \overline{U}$ are evaluated at $(R - \rho, z)$.

Set $t = R - \rho$, then straightforward computations reduce the inequality above to

$$[2(R-t) - R - (n-1)t]\partial_t \bar{U}(t,z) - (n-1)\bar{U}(t,z) < 0$$

Using that $\partial_t \bar{U}(t,z) = -\bar{U}(t,z)/(2r)$ with $r^2 = t^2 + z^2$, this inequality is immediately satisfied for R large.

Now we prove that (v_1^R, v_2^R) satisfies the free boundary condition. Observe that

$$F(v_1^R, v_2^R) = \partial \mathcal{B}_R(Re_n, 0) \cap \mathcal{B}_2,$$

and hence it is smooth. By the radial symmetry it is enough to show that the free boundary condition is satisfied at 0. Precisely we have,

(4.15)
$$v_R(x,z) = U(x_n,z) + o(|(x,z)|^{1/2}), \text{ as } (x,z) \to (0,0),$$

and

(4.16)
$$\bar{v}_R(x,z) = \bar{U}(x_n,z) + o(|(x,z)|^{1/2}), \text{ as } (x,z) \to (0,0).$$

Indeed, let us prove the expansion for \bar{v}_R . Since \bar{U} is Hölder continuous with exponent 1/2, it follows that

$$|\bar{V}_R(t,z) - \bar{V}_R(t_0,z)| \le C|t - t_0|^{1/2}$$
 for $|t - t_0| \le 1$.

Thus for $(x, z) \in B_s$, s small

 $|\bar{v}_R(x,z) - \bar{V}_R(x_n,z)| = |\bar{V}_R(R-\rho,z) - \bar{V}_R(x_n,z)| \le C|R-\rho-x_n|^{1/2} \le Cs.$ Hence for $(x,z) \in B_s$

$$|\bar{v}_R(x,z) - \bar{U}(x_n,z)| \le Cs + |\bar{V}_R(x_n,z) - \bar{U}(x_n,z)|.$$

Thus from the formula for \bar{V}_R

$$|\bar{v}_R(x,z) - \bar{U}(x_n,z)| \le Cs + (n-1)\frac{|x_n|}{R}\bar{U}(x_n,z) \le C's, \quad (x,z) \in B_s$$

which gives the desired expansion (4.16).

Now, we show that v_2^R is strictly monotone decreasing in the e_n -direction in $B_2^+(v_2^R)$. Outside of its zero plate,

$$\partial_{x_n} \bar{v}_R(x) = -\frac{x_n - R}{\rho} \partial_t \bar{V}_R(R - \rho, z)$$

Thus we only need to show that $\bar{V}_R(t,z)$ is strictly monotone decreasing in t outside $\{(t,0): t \ge 0\}$. This follows immediately from (4.14) and the formula for \bar{U} .

Step 2. We show the existence of \tilde{v}_1^R satisfying (4.11) and (4.12). The claim for v_2^R can be proved with similar arguments.

Following Proposition 3.5 in [16] it suffices to show the 2-dimensional claim:

$$U(t + \gamma_R(t, z) - \frac{C}{R^2}, z) \le v_{\beta, R}(t, z) \le U(t + \gamma_R(t, z) + \frac{C}{R^2}, z)$$

with

$$v_{\beta,R}(t,z) := (1+\frac{\beta}{R})V_R(t,z),$$

and

$$\gamma_R(t,z) = 2\frac{(n-1)}{R}rt + 2\beta\frac{r}{R}, \quad r = t^2 + z^2.$$

Indeed since (see the properties listed at the beginning of the Appendix,)

$$|U_{tt}| \le Cr^{-1}U_t$$

we have that if $|\mu| \leq r/2$ then

$$|U(t+\mu,z) - (U(t,z) + \mu U_t(t,z))| \le \mu^2 |U_{tt}(t',z)| \le C\mu^2 r^{-1} U_t(t,z),$$

where in the last inequality we used (A.1). Thus, since $U_t = U/(2r)$,

$$(1 + \frac{\mu}{2r} + C\frac{\mu^2}{r^2})U(t,s) \ge U(t + \mu, z) \ge (1 + \frac{\mu}{2r} - C\frac{\mu^2}{r^2})U(t,z).$$

Choosing

$$\mu = \tilde{\mu} \pm 4C \frac{\tilde{\mu}^2}{r}$$

we obtain that

(4.17)
$$U(t + \tilde{\mu} + 4C\frac{\tilde{\mu}^2}{r}, z) \ge (1 + \frac{\tilde{\mu}}{2r})U(t, z) \ge U(t + \tilde{\mu} - 4C\frac{\tilde{\mu}^2}{r}, z),$$

provided that $|\tilde{\mu}|/r \leq c$, with c sufficiently small. Since

$$v_{\beta,R} = (1 + \frac{n-1}{R}t + \frac{\beta}{R} + \frac{(n-1)\beta}{R^2}t)U$$

we can apply the inequality above with

$$\tilde{\mu} = 2r(\frac{n-1}{R}t + \frac{\beta}{R} + \frac{(n-1)\beta}{R^2}t)$$

hence $|\tilde{\mu}|/r \leq \frac{C}{R}$ and obtain the claim.

Furthermore, these radial subsolutions satisfy the following estimates.

Corollary 4.5. There exist δ , c_0 , C_0 , C_1 depending on β , such that (4.18)

$$v_{1}^{R}(X + \frac{c_{0}}{R}e_{n}) \geq U(X + \frac{c_{0}}{2R}e_{n}), \quad v_{2}^{R}(X + \frac{c_{0}}{R}e_{n}) \leq \bar{U}(X + \frac{c_{0}}{2R}e_{n}), \quad \text{in } B_{\delta},$$
(4.19)

$$v_{1}^{R}(X - \frac{C_{1}}{R}e_{n}) \leq U(X), \quad v_{2}^{R}(X - \frac{C_{1}}{R}e_{n}) \geq \bar{U}(X), \quad \text{in } \overline{B}_{1},$$
(4.20)

$$v_{1}^{R}(X + \frac{c_{0}}{R}e_{n}) \leq (1 + \frac{C_{0}}{R})U(X), \quad \text{in } \overline{B}_{1} \setminus B_{1/4},$$

with strict inequality on $F(v_1^R(X + \frac{c_0}{R}e_n)) \cap (\overline{B}_1 \setminus B_{1/4})$, and

(4.21)
$$v_2^R(X + \frac{c_0}{R}e_n) \ge (1 - \frac{C_0}{R})\bar{U}(X), \quad \text{in } \overline{B}_1 \setminus B_{1/4},$$

with strict inequality on $L \cap (\overline{B}_1 \setminus B_{1/4})$.

Proof. Estimates (4.18)-(4.19) follow from (4.11) and Lemmas 2.8-2.9. Let us prove (4.21). The proof of (4.20) can be obtained similarly.

Using again (4.11) and Lemma 2.9 we get that

$$v_2^R(X + \frac{c_0}{R}e_n) \ge \bar{U}(X) \quad \text{in } B_1 \cap \{|x'| \ge \frac{1}{8}, |x_n| \le \bar{\delta}\},$$

with $c_0, \bar{\delta}$ depending on β and R large (with strict inequality on L.) Thus we need to show that (4.21) holds on the complement of $(\bar{B}_1 \setminus B_{1/4}) \cap \{r \leq \bar{\delta}\}$. On the other hand in B_1 ,

$$v_2^R(X + \frac{c_0}{R}e_n) \ge \bar{U}(X + \frac{\bar{C}}{R}e_n)$$

for some \overline{C} depending on β (combining again (4.11) and Lemma 2.9.) From the version of Lemma A.2 in the Appendix for the function \overline{U} , we conclude that

$$\bar{U}(X + \frac{\bar{C}}{R}e_n) \ge (1 - C\frac{\bar{C}}{R})\bar{U}(X),$$

as long as $r > \overline{\delta}$ and with $C = C(\overline{\delta})$. The desired conclusion immediately follows. \Box

Having established the Proposition and the Corollary above, we are ready to present the proof of Lemma 4.3.

Proof of Lemma 4.3. In view of (4.8), we have

$$u_1(\bar{X}) - U(\bar{X}) \ge U(\bar{X} + \varepsilon e_n) - U(\bar{X}) = \partial_t U(\bar{X} + \lambda e_n) \varepsilon \ge c\varepsilon, \quad \lambda \in (0, \varepsilon).$$

Hence by Lemma A.1 in the Appendix it follows that

(4.22)
$$u_1(X) \ge (1 + c'\varepsilon)U(X) \quad \text{in } B_{1/4}$$

Now set,

$$w_i^{R,t}(X) := (1 + c''\varepsilon)v_i^R(X + te_n), \quad X \in B_1$$

with (v_1^R, v_2^R) the subsolution pair defined in (4.10) for $\beta = 0, c'' \le c'$ and R large to be chosen later.

If $t_0 = -C_1/R$, then it follows from (4.19) and the inequality above that

$$w_1^{R,t_0} \le (1 + c'\varepsilon)U(X) \le u_1, \quad w_2^{R,t_0} \ge \bar{U} \ge u_2 \quad \text{in } \bar{B}_{1/4}.$$

If $t_1 = c_0/R$, then choosing $R = 2C_0/c''\varepsilon$, and ε small (depending on c''), (4.20) gives

$$w_1^{R,t_1} \leq (1+c'\varepsilon)U \leq u_1 \quad \text{on } \partial B_{1/4}$$

while (4.21) gives

$$w_2^{R,t_1} \ge (1 + c''\varepsilon)(1 - \frac{C_0}{R})\bar{U} \ge \bar{U} \ge u_2 \quad \text{on } \partial B_{1/4},$$

with the inequality being strict on the necessary sets to apply Corollary 2.6. Thus, we conclude that

$$w_1^{R,t_1} \le u_1, \quad w_{R,t_2}^2 \ge u_2 \quad \text{in } B_{1/4}.$$

From (4.18) it follows that in B_{ρ} , for $\rho \leq \delta$,

$$u_1 \ge U(X + \frac{c_0}{2R}e_n)$$

while

$$u_2 \le (1 + c''\varepsilon)\overline{U}(X + \frac{c_0}{2R}e_n) \le \overline{U}(X + \frac{c_0}{2R}e_n - Cc''\rho\varepsilon e_n).$$

The last inequality follows from a variant of Lemma A.1 for \overline{U} , rescaled in the ball B_{ρ} . Therefore, for ρ small (depending on c'', C.) we get that

$$u_2 \le \bar{U}(X + \frac{c_0}{4R}) \quad \text{in } B_\rho$$

and the desired gain is achieved.

5. Improvement of Flatness

In this section we show our main improvement of flatness theorem, from which the desired regularity Theorem 1.6 follows with standard arguments (see for example [2].)

Theorem 5.1 (Improvement of flatness). There exist $\bar{\varepsilon} > 0$ and $\rho > 0$ universal constants such that for all $0 < \varepsilon \leq \bar{\varepsilon}$ if (u_1, u_2) solves (2.1) with $0 \in F(u_1, u_2)$ and it satisfies

(5.1)
$$U(X - \varepsilon e_n) \le u_1(X) \le U(X + \varepsilon e_n) \quad \text{in } B_1,$$

(5.2)
$$U(X + \varepsilon e_n) \le u_2(X) \le U(X - \varepsilon e_n) \quad \text{in } B_1,$$

then

(5.3)
$$\alpha U(x \cdot \nu - \frac{\varepsilon}{2}\rho, z) \le u_1(X) \le \alpha U(x \cdot \nu + \frac{\varepsilon}{2}\rho, z) \quad \text{in } B_\rho,$$

(5.4)
$$\alpha \bar{U}(x \cdot \nu + \frac{\varepsilon}{2}\rho, z) \le u_2(X) \le \alpha \bar{U}(x \cdot \nu - \frac{\varepsilon}{2}\rho, z) \quad \text{in } B_\rho,$$

for some direction $\nu \in \mathbb{R}^n$, $|\nu| = 1$, and $|\alpha - 1| \leq \varepsilon$.

The proof of Theorem 5.1 relies on the next two lemmas. The first Lemma is an immediate consequence of Corollary 4.2.

Lemma 5.2. Let $\varepsilon_k \to 0$ and let (u_1^k, u_2^k) be a sequence of solutions to (2.1) with $0 \in F(u_1^k, u_2^k)$ satisfying

(5.5)
$$U(X - \varepsilon_k e_n) \le u_1^k(X) \le U(X + \varepsilon_k e_n) \quad \text{in } B_1,$$

(5.6)
$$\overline{U}(X - \varepsilon_k e_n) \le u_2^k(X) \le \overline{U}(X + \varepsilon_k e_n)$$
 in B_1 .

Denote by $\tilde{u}_1^k, \tilde{\tilde{u}}_2^k$ the ε_k -domain variation of (u_1^k, u_2^k) with respect to U, \bar{U} respectively. Then the sequences of sets

$$A_1^k := \{ (X, \tilde{u}_1^k(X)) : X \in B_{1-\varepsilon_k} \setminus P^- \},$$

$$A_2^k := \{ (X, \tilde{u}_2^k(X)) : X \in B_{1-\varepsilon_k} \setminus P^+ \},$$

have a subsequence that converge uniformly (in Hausdorff distance) in $B_{1/2} \setminus P^{\mp}$ to the graph

$$A_{\infty}^{+} := \{ (X, \tilde{u}_{\infty}^{+}(X)) : X \in B_{1/2} \setminus P^{+} \},\$$

where \tilde{u}_{∞}^{\mp} is a Hölder continuous function in $B_{1/2}$ and

$$\tilde{u}_{\infty}^{-} = \tilde{u}_{\infty}^{+}$$
 on $B_{1/2} \cap L$

25

Lemma 5.3. The pair $(\tilde{u}_{\infty}^{-}, \tilde{u}_{\infty}^{+})$ in Lemma 5.2 solves the linearized problem (3.1) in $B_{1/2}$. In particular, in $B_{2\rho} \setminus P^{\mp}$ (ρ small universal)

(5.7)
$$a' \cdot x' + b_{\mp}r - \frac{1}{8}\rho \le \tilde{u}_{\infty}^{\mp}(X) \le a' \cdot x' + b_{\mp}r + \frac{1}{8}\rho,$$

for some $a' \in \mathbb{R}^{n-1}, b_{\mp} \in \mathbb{R}$ (universally bounded) such that

$$b_{-} + b_{+} = 0.$$

Proof. We start by showing that $U_n \tilde{u}_{\infty}^-$ is harmonic in $B_{1/2} \setminus P^-$. Similarly, we can show that $-\bar{U}_n \tilde{u}_{\infty}^+$ is harmonic in $B_{1/2} \setminus P^+$.

Let $\tilde{\varphi}$ be a smooth function which touches \tilde{u}_{∞}^- strictly by below at $X_0 \in B_{1/2} \setminus P^-$. We need to show that

(5.8)
$$\Delta(U_n\tilde{\varphi})(X_0) \le 0.$$

Since by the previous lemma, the sequence A_1^k converges uniformly to A_{∞}^- in $B_{1/2} \setminus P^$ we conclude that there exist a sequence of constants $c_k \to 0$ and a sequence of points $X_k \in B_{1/2} \setminus P^-$, $X_k \to X_0$ such that $\tilde{\varphi}_k := \tilde{\varphi} + c_k$ touches \tilde{u}_1^k by below at X_k for all k large enough.

Define the function φ_k by the following identity

(5.9)
$$\varphi_k(X - \varepsilon_k \tilde{\varphi_k}(X) e_n) = U(X).$$

Then according to (2.23) φ_k touches u_1^k by below at $Y_k = X_k - \varepsilon_k \tilde{\varphi}_k(X_k) e_n \in B_1^+(u_1^k)$, for k large enough. Thus, since u_1^k satisfies (2.1) in B_1 it follows that

$$(5.10) \qquad \qquad \Delta \varphi_k(Y_k) \le 0.$$

By Proposition 2.11,

$$\Delta \varphi_k(Y_k) = \varepsilon_k \Delta(U_n \tilde{\varphi})(Y_k) + O(\varepsilon_k^2).$$

Thus, as $k \to \infty$

$$\Delta(U_n\tilde{\varphi})(X_0) \le 0,$$

as desired.

Next we need to show that

$$(\tilde{u}_{\infty}^{-})_r(X_0) + (\tilde{u}_{\infty}^{+})_r(X_0) = 0, \quad X_0 = (x'_0, 0, 0) \in B_{1/2} \cap L,$$

in the viscosity sense of Definition 3.1.

Assume by contradiction that there exist ϕ_i continuous around X_0 and satisfying

$$\phi_1(X) = l(x') + b_1(X_0)r + O(|x' - x'_0|^2 + r^{3/2}),$$

$$\phi_2(X) = l(x') + b_2(X_0)r + O(|x' - x'_0|^2 + r^{3/2}),$$

with $l(x') = a_0 + a_1 \cdot (x' - x'_0)$, and

$$b_1(X_0) + b_2(X_0) > 0$$

with (ϕ_1, ϕ_2) which touch $(\tilde{u}_{\infty}^-, \tilde{u}_{\infty}^+)$ by below at X_0 .

Then we can find constants $\alpha, \beta, \delta, \bar{r}$ and a point $Y' = (y'_0, 0, 0) \in B_2$ depending on ϕ_1, ϕ_2 such that the polynomials

$$q_1(X) = a_0 - \frac{\alpha}{2} |x' - y'_0|^2 + 2\alpha(n-1)x_n r + 2\alpha\beta r,$$

$$q_2(X) = a_0 - \frac{\alpha}{2} |x' - y'_0|^2 - 2\alpha(n-1)x_n r - 2\alpha\beta r,$$

touch ϕ_1, ϕ_2 by below at X_0 in a tubular neighborhood $N_{\bar{r}} = \{|x' - x'_0| \le \bar{r}, r \le \bar{r}\}$ of X_0 , with

$$\phi_i - q_i \ge \delta > 0, \quad \text{on } N_{\bar{r}} \setminus N_{\bar{r}/2}.$$

This implies that

(5.11)
$$\tilde{u}_{\infty}^{-} - q_1 \ge \delta > 0, \quad \tilde{u}_{\infty}^{+} - q_2 \ge \delta > 0 \quad \text{on } N_{\bar{r}} \setminus N_{\bar{r}/2},$$

and

(5.12)
$$\tilde{u}_{\infty}^{-}(X_0) - q_1(X_0) = 0, \quad \tilde{u}_{\infty}^{+}(X_0) - q_2(X_0) = 0.$$

In particular,

(5.13)

$$\begin{split} &|\tilde{u}_{\infty}^{-}(X_{k}^{-}) - q_{1}(X_{k}^{-})| \to 0, \quad |\tilde{u}_{\infty}^{+}(X_{k}^{+}) - q_{2}(X_{k}^{+})| \to 0, \quad X_{k}^{\mp} \in N_{\bar{r}} \setminus P^{\mp}, X_{k}^{\mp} \to X_{0}. \end{split}$$
 Now, let us choose $R_{k} = 1/(\alpha \varepsilon_{k})$ and let us define

$$\begin{split} & w_1^k(X) = v_1^{R_k}(X - Y' + \varepsilon_k a_0 e_n), \quad Y' = (y'_0, 0, 0), \\ & w_2^k(X) = v_2^{R_k}(X - Y' + \varepsilon_k a_0 e_n), \quad Y' = (y'_0, 0, 0), \end{split}$$

with v_i^R the functions defined in Proposition 4.4 for the choice of β above. Then the ε_k -domain variation of w_1^k, w_2^k , which we call $\tilde{w}_1^k, \tilde{w}_2^k$, can be easily computed from the definition

$$w_1^k(X - \varepsilon_k \tilde{w}_1^k(X)e_n) = U(X), \quad w_2^k(X - \varepsilon_k \tilde{w}_2^k(X)e_n) = \bar{U}(X)$$

Indeed, since U and \bar{U} are constant in the x' -direction, this identity is equivalent to

$$v_1^{R_k}(X - Y' + \varepsilon_k a_0 e_n - \varepsilon_k \tilde{w}_1^k(X) e_n) = U(X - Y'),$$

$$v_2^{R_k}(X - Y' + \varepsilon_k a_0 e_n - \varepsilon_k \tilde{w}_2^k(X) e_n) = \bar{U}(X - Y'),$$

which in view of Proposition 4.4 gives us

$$\tilde{v}_{1}^{R_{k}}(X - Y') = \varepsilon_{k}(\tilde{w}_{1}^{k}(X) - a_{0}),$$

$$\tilde{v}_{2}^{R_{k}}(X - Y') = \varepsilon_{k}(\tilde{w}_{2}^{k}(X) - a_{0}).$$

From the choice of R_k , the formula for q_i and (4.12), we then conclude that

$$\begin{split} \tilde{w}_1^k(X) &= q_1(X) + \alpha^2 O(\varepsilon_k), \\ \tilde{w}_2^k(X) &= q_2(X) + \alpha^2 O(\varepsilon_k), \end{split}$$

and hence

$$(5.14) \qquad |\tilde{w}_1^k - q_1| \le C\varepsilon_k, \quad |\tilde{\tilde{w}}_2^k - q_2| \le C\varepsilon_k, \quad \text{in } N_{\bar{r}} \setminus P^{\mp} \text{ respectively.}$$

Thus, from the uniform convergence of A_k^\mp to A_∞^\mp and (5.11)-(5.14) we get that for all k large enough

(5.15)
$$\tilde{u}_k^1 - \tilde{w}_1^k \ge \frac{\delta}{2} \quad \text{in} \left(N_{\bar{r}} \setminus N_{\bar{r}/2} \right) \setminus P^-,$$

(5.16)
$$\tilde{u}_k^2 - \tilde{\bar{w}}_2^k \ge \frac{\delta}{2} \quad \text{in} \left(N_{\bar{r}} \setminus N_{\bar{r}/2}\right) \setminus P^+$$

Similarly, from the uniform convergence of A_k^{\mp} to A_{∞}^{\mp} and (5.14)-(5.13) we get that for k large

 $(5.17) \quad \tilde{u}_k^1(X_k^-) - \tilde{w}_1^k(X_k^-) \le \frac{\delta}{4}, \quad \text{ for some sequence } X_k^- \in N_{\bar{r}} \setminus P^-, X_k^- \to X_0.$

(5.18)
$$\tilde{u}_k^2(X_k^+) - \tilde{w}_2^k(X_k^+) \le \frac{\delta}{4}$$
, for some sequence $X_k^+ \in N_{\bar{r}} \setminus P^+, X_k^+ \to X_0$.

On the other hand, it follows from Lemma 2.10 and (5.15)-(5.16) that

$$\begin{split} \tilde{u}_k^1 - \tilde{w}_1^k &\geq \frac{\delta}{2}, \quad \text{in } N_{\bar{r}} \setminus P^-, \\ \tilde{u}_k^2 - \tilde{w}_2^k &\geq \frac{\delta}{2}, \text{in } N_{\bar{r}} \setminus P^+. \end{split}$$

We have reached a contradiction.

We are now ready to present the proof of Theorem 5.1.

Proof of Theorem 5.1. Let ρ be the universal constant in Lemma 5.3 and assume by contradiction that there exists a sequence $\varepsilon_k \to 0$ and a sequence of solutions (u_1^k, u_2^k) to (2.1) in B_1 which satisfies the flatness assumption but not the conclusion of the theorem. Then by Lemmas 5.2-5.3, for k large, we have in $B_{2\rho} \setminus P^{\mp}$,

(5.19)
$$a' \cdot x' + b_{-}r - \frac{1}{4}\rho \le \tilde{u}_{1}^{k}(X) \le a' \cdot x' + b_{-}r + \frac{1}{4}\rho.$$

(5.20)
$$a' \cdot x' + b_{+}r - \frac{1}{4}\rho \le \tilde{\bar{u}}_{2}^{k}(X) \le a' \cdot x' + b_{+}r + \frac{1}{4}\rho.$$

Let

$$\alpha = 1 + \frac{1}{2}b_{-}\varepsilon = 1 - \frac{1}{2}b_{+}\varepsilon.$$

Set

$$v(X) := \alpha U(x \cdot \nu - \frac{\varepsilon_k}{2}\rho, z), \quad \nu = (\nu', \nu_n) := \frac{(0, 1) + \varepsilon_k(a', 0)}{\sqrt{1 + \varepsilon_k^2 |a'|^2}}.$$

Since $\nu_n > 0$, v is strictly monotone increasing in the e_n -direction say in $B_{2\rho}^+(v)$. We wish to show that

(5.21)
$$\tilde{v}_{\varepsilon_k} \leq \tilde{u}_1^k \quad \text{in } B_{\rho+\varepsilon_k} \setminus P^-.$$

Then by Lemma 2.8 we can conclude that

$$v \leq u_1^k \quad \text{in } B_{\rho}.$$

With similar arguments we obtain that u_1^k, u_2^k satisfy all the bounds in the conclusion of the theorem and reach a contradiction.

Notice that in view of Lemma A.1 in the Appendix,

 $|\tilde{v}_{\varepsilon_k}| \leq C$, *C* universal.

According to (4.17), we have

(5.22)
$$\alpha U(t,z) = U(t + \varepsilon_k b_- r + O(\varepsilon_k^2), z), \quad r^2 = t^2 + z^2.$$

Thus we can estimate $\tilde{v}_{\varepsilon_k}$ using the definition:

$$v(X - \varepsilon_k \tilde{v}_{\varepsilon_k}(X)e_n) = U(X), \quad X \notin P^-.$$

Indeed,

$$v(X - \varepsilon \tilde{v}_{\varepsilon_k}(X)e_n) = \alpha U(t(X), z), \quad t(X) := x \cdot \nu - \varepsilon_k \tilde{v}_{\varepsilon}(X)\nu_n - \frac{\varepsilon_k}{2}\rho,$$

hence by the claim,

$$U(X) = U(t(X) + \varepsilon_k b_- \bar{r} + O(\varepsilon_k^2)\bar{r}^2, z), \quad \bar{r}^2 = t(X)^2 + z^2.$$

We conclude that, say in $B_1 \setminus P^-$

$$x_n = x' \cdot \nu' + x_n \nu_n - \varepsilon \tilde{v}_{\varepsilon_k}(X) \nu_n - \frac{\varepsilon_k}{2} \rho + \varepsilon_k b_- \bar{r} + O(\varepsilon_k^2).$$

28

Moreover,

$$\bar{r} = r\nu_n + O(\varepsilon_k).$$

Indeed, since $\nu_n \ge 1 - C \varepsilon_k^2$

$$|\bar{r} - r\nu_n| \le |t(X) - x_n| + C\varepsilon_k^2$$

and also $\nu' = O(\varepsilon_k)$ and $\tilde{v}_{\varepsilon_k}$ is bounded, thus

$$t(X) - x_n = x' \cdot \nu' + x_n(\nu_n - 1) - \varepsilon \tilde{v}_{\varepsilon_k} \nu_n - \frac{\varepsilon_k}{2} \rho = O(\varepsilon_k).$$

Hence,

$$\tilde{v}_{\varepsilon_k}(X) = \frac{1}{\varepsilon_k \nu_n} [x' \cdot \nu' + (\nu_n - 1)x_n - \frac{\varepsilon_k}{2}\rho + \varepsilon_k b_- r\nu_n + O(\varepsilon_k^2)]$$

and by the definition of $\nu,$

$$\tilde{v}_{\varepsilon_k}(x) \le x' \cdot a' + b_- r - \frac{\rho}{2} + O(\varepsilon_k).$$

The inequality in (5.21) now follows immediately for k large using (5.19).

6. Segregated critical profiles

In order to effectively describe the nodal set of our segregated minimal profiles, we need to take properly into account the extremality features of such configurations. To this aim, we consider the following definition.

Definition 6.1 (Class \mathcal{G}). For an open, z-symmetric $\Omega \subset \mathbb{R}^{n+1}$, we define the class $\mathcal{G}(\Omega)$ of the segregated critical profiles as the set of even-in-z vector valued functions $\mathbf{u} = (u_1, \ldots, u_k) \in (H^1(\Omega))^k \cap C^{0,1/2}(\Omega)$, whose components are all nonnegative and such that

(G1)
$$-\Delta u_i = 0$$
 in $\Omega \setminus \mathcal{N}(u_i), i = 1, \dots, k,$

(G2)
$$\tilde{\Omega} =: \Omega \cap \mathbb{R}^n \times \{0\} = \bigcup_{i=1}^k \mathcal{N}(u_i), \quad \mathcal{N}^o(u_i) \cap \mathcal{N}^o(\sum_{j \neq i} u_j) = \emptyset \; \forall i = 1, \dots, k \;,$$

(G₃)
$$\int_{\Omega} \left\{ dY(X) \nabla \mathbf{u}(X) \cdot \nabla \mathbf{u}(X) - \frac{1}{2} \mathrm{div} Y(X) |\nabla \mathbf{u}(X)|^2 \right\} dX = 0 ,$$

for every vector field $Y \in C_0^{\infty}(\Omega; \mathbb{R}^{n+1})$ whose z-component is odd-in-z (the odd-in-z domain variation formula).

As customary, we say that $\mathbf{u} \in \mathcal{G}_{loc}(\mathbb{R}^{n+1})$ if $\mathbf{u} \in \mathcal{G}(B_R)$ for every R > 0.

Remark 6.2. It is immediate to check, using the penalized functionals (1.3), the a priori estimates of Theorem 1.1 and the maximal regularity result in Theorem 1.2, that the minimizing critical profiles given by Definition 1.3 are indeed segregated critical profiles in the sense above. A word of caution must be entered concerning the second condition in (G2). Let indeed $\mathbf{u} \neq 0$ be a minimal energy configuration in a ball with all component identically vanishing at the characteristic hyperplane. Assume that $u_i \neq 0$ then, as all the other components vanish on \Re^n , it is an *unconstrained even-in-z energy minimizer* and, as such, it can not identically vanish on $\{z = 0\}$.

Now we proceed to the study of the segregated critical profiles in the sense of Definition 6.1. Once more, we shall take advantage of domain variations, though in a more general framework. In what follows the domain variations formula (G₃) will play a key role in two complementary contests. On one hand, we will exploit it in order to prove Almgren's monotonicity formula; furthermore we will show that it implies a fundamental weak reflection law which will allow us to prove that every element of the class is in fact a viscosity solution in the appropriate sense. Such a weak reflection property is expressed by the following proposition.

Proposition 6.3. Let $\mathbf{u} = (u_1, u_2) = (a_1U, a_2\overline{U})$ for some constants, U being defined in (1.4). Then, the validity of the domain variation formula (G3) is equivalent to $a_1^2 = a_2^2$.

Proof. To see this we use the identity, which is equivalent to (G3) in the specific case,

(6.1)
$$\int_{\omega} \left\{ dY(x) \nabla \mathbf{u}(x) \cdot \nabla \mathbf{u}(x) - \frac{1}{2} \operatorname{div} Y(x) |\nabla \mathbf{u}(x)|^2 \right\} dx \\ = \int_{\partial \omega} \left\{ Y(x) \cdot \nabla \mathbf{u}(x) \ \nu(x) \cdot \nabla \mathbf{u}(x) - \frac{1}{2} \nu(x) \cdot Y(x) |\nabla \mathbf{u}(x)|^2 \right\} d\sigma ,$$

where ω is any piecewise smooth domain in \mathbb{R}^{n+1} , z-symmetric. Now we take $Y \equiv e_n = (0, \ldots, 0, 1, 0)$ in \mathbb{R}^{n+1} and the cylinder

$$\omega = C_{l,r} = Q_l \times B_r,$$

where, $B_r \subset \mathbb{R}^2$ denotes a ball of radius r in the plane (x_n, z) , and $Q_l \subset \mathbb{R}^{n-1}$ a cube of edge length equal to 2l, both centered at zero.

We obtain, since ν is orthogonal to e_n and ∇u_i on $\partial Q_l \times B_r$,

$$0 = \int_{\partial(Q_l \times B_r)} \left\{ e_n \cdot \nabla \mathbf{u}(x) \,\nu(x) \cdot \nabla \mathbf{u}(x) - \frac{1}{2}\nu(x) \cdot e_n |\nabla \mathbf{u}(x)|^2 \right\} d\sigma$$

$$= \int_{Q_l \times \partial B_r} \left\{ e_n \cdot \nabla \mathbf{u}(x) \,\nu(x) \cdot \nabla \mathbf{u}(x) - \frac{1}{2}\nu(x) \cdot e_n |\nabla \mathbf{u}(x)|^2 \right\} d\sigma$$

$$= (2l)^{n-1} \int_0^{2\pi} \left\{ \frac{a_1^2}{4} \cos^2(\theta/2) - \frac{a_1^2}{8} \cos\theta - \frac{a_2^2}{4} \sin^2(\theta/2) - \frac{a_2^2}{8} \cos\theta \right\} d\theta .$$

In order to prove the Almgren monotonicity formula, arguing as in [23, Theorem 3.3], we can take advantage of the local Pohožaev identity, a direct consequence of the domain variation formula (6.1) for spherical domains and the particular choice of the vector field $Y = X - X_0$. Indeed, for $\mathbf{u} \in \mathcal{G}(B_R((x_0, 0)))$ we have

(6.2)
$$(1-n) \int_{B_r} \sum_i |\nabla u_i|^2 \, \mathrm{d}x \mathrm{d}z + r \int_{\partial B_r} \sum_i |\nabla u_i|^2 \, \mathrm{d}\sigma = 2r \int_{\partial B_r} \sum_i |\partial_\nu u_i|^2 \, \mathrm{d}\sigma$$

for a.e. $r \in (0, R)$.

Let us define, for every $x_0 \in \mathbb{R}^n$, $X_0 = (x_0, 0) \in \mathbb{R}^{n+1}$, and r > 0,

$$\begin{split} E(r) &= E(x_0, r) = E(x_0, \mathbf{u}, r) := \frac{1}{r^{n-1}} \int_{B_r(x_0, 0)} \sum_i |\nabla u_i|^2 \, \mathrm{d}x \mathrm{d}z \\ H(r) &= H(x_0, r) = H(x_0, \mathbf{u}, r) := \frac{1}{r^n} \int_{\partial B_r(x_0, 0)} \sum_i u_i^2 \, \mathrm{d}\sigma. \end{split}$$

Since $\mathbf{u} \in H^1_{\text{loc}}(\mathbb{R}^{n+1}, \mathbb{R}^k)$, both E and H are locally absolutely continuous functions on $(0, +\infty)$, that is, both E' and H' are $L^1_{\text{loc}}(0, \infty)$ (here, ' = d/dr).

Theorem 6.4. Let $\mathbf{u} \in \mathcal{G}(B_{r_0}((x_0, 0)), \mathbf{u} \neq 0$. For every $x_0 \in \mathbb{R}^n$ the function (Almgren frequency function)

$$N(x_0, r) := \frac{E(x_0, r)}{H(x_0, r)}$$

is well defined on $(0, r_0)$, absolutely continuous, non decreasing, and it satisfies the identity

(6.3)
$$\frac{\mathrm{d}}{\mathrm{d}r}\log H(r) = \frac{2N(r)}{r}$$

Moreover, if $N(r) \equiv \gamma$ on an open interval, then $N \equiv \gamma$ for every r, and \mathbf{u} is a homogeneous function of degree γ .

Proof. There holds

$$H^{2}(r)N'(r) = E'(r)H(r) - E(r)H'(r)$$
 for $r \in (r_{1}, r_{2})$.

The Pohožaev identity (6.2) can be used to compute the derivative of *E*:

$$E'(r) = \frac{1-n}{r^n} \int_{B_r} \sum_{i} |\nabla u_i|^2 \, \mathrm{d}x \mathrm{d}y + \frac{1}{r^{n-1}} \int_{\partial B_r} \sum_{i} |\nabla u_i|^2 \, \mathrm{d}\sigma$$
$$= \frac{2}{r^{n-1}} \int_{\partial B_r} \sum_{i} |\partial_{\nu} u_i|^2 \, \mathrm{d}\sigma,$$

While for H we find

(6.4)

$$H'(r) = \frac{2}{r^n} \int\limits_{\partial B_r} \sum_i u_i \partial_\nu u_i \,\mathrm{d}\sigma = \frac{2}{r} E(r).$$

Thus the Cauchy-Schwarz inequality yields, for $r \in (r_1, r_2)$,

$$\frac{r^{2n-1}}{2} H^2(r) N'(r) =$$

$$= \int_{\partial B_r} \sum_i |\partial_\nu u_i|^2 \, \mathrm{d}\sigma \int_{\partial B_r} \sum_i u_i^2 \, \mathrm{d}\sigma - \left(\int_{\partial B_r} \sum_i u_i \partial_\nu u_i \, \mathrm{d}\sigma \right)^2 \ge 0.$$

To conclude this section, we observe that the monotonicity of N(x, r) implies that both for r small and for r large the corresponding limits are well defined. Consequently we have (see, e.g. [21, Corollary 2.6]):

Proposition 6.5. Let $\mathbf{u} \in \mathcal{G}(\Omega)$. Then

(i) N(x,0⁺) is a non negative upper semicontinuous function on Ω;
(ii) If Ω = ℝⁿ⁺¹, then N(x,∞) is constant on ℝⁿ (possibly ∞).

Another relevant consequence of the monotonicity formula is the following comparison property which, with $r_2 = 2r_1$, entails the doubling property and gives a unique continuation principle. **Proposition 6.6.** Given $\mathbf{u} \in \mathcal{G}(\Omega)$ and $\tilde{\omega} \in \tilde{\Omega}$, there exist $\tilde{C} > 0$ and $\tilde{r} > 0$ such that

$$H(x_0, \mathbf{u}, r_2) \le H(x_0, \mathbf{u}, r_1) \left(\frac{r_2}{r_1}\right)^{2\tilde{C}}$$

for every $x_0 \in \tilde{\omega}$, $0 < r_1 < r_2 \leq \tilde{r}$.

6.1. Compactness and convergence of blowup sequences. In what follows we will frequently happen to study the local behavior of solutions, which will be analyzed via a blowup procedure. We will therefore need a compactness theorem, in order to obtain convergence of sequences. As a main results of [23, Proposition 6.15], we have uniform bounds in spaces of Hölder functions

Theorem 6.7 (A priori bounds in Hölder spaces). Let $\{\mathbf{u}_j\}_j$ be a sequence of elements of $\mathcal{G}(B_R)$ such that

$$\|\mathbf{u}_j\|_{L^\infty(B_R)} \le \bar{m}_j$$

with \bar{m} independent of j. Then for every $\alpha \in (0, 1/2)$ there exists a constant $C = C(\bar{m}, \alpha, R)$, not depending on j, such that

$$\|\mathbf{u}_j\|_{\mathcal{C}^{0,\alpha}(\overline{B_R/2})} \le C.$$

Furthermore, $\{\mathbf{u}_j\}_j$ is relatively compact in $H^1(B_{R/2}) \cap \mathcal{C}^{0,\alpha}(\overline{B_{R/2}})$ for every $\alpha < 1/2$.

This result is the key to the study of convergence of blow-up sequences in the class $\mathcal{G}(\Omega)$. Take $\mathbf{u} \in \mathcal{G}(\Omega)$ and define, for every fixed $\rho, t > 0$ and $x_0 \in \Omega$, the rescaled function

$$\mathbf{v}(x) = \mathbf{v}_{x_0,t}(x) = \frac{1}{\sqrt{H(x_0,t)}} \mathbf{u}_{x_0,t}(x) = \frac{\mathbf{u}(x_0 + tx)}{\sqrt{H(x_0,t)}}, \quad \text{for } x \in \tilde{\Omega} := \frac{\Omega - x_0}{t}.$$

It is immediate to check that $\mathbf{v} \in \mathcal{G}(\tilde{\Omega})$. Now we consider the convergence of such blowup sequences with fixed or variable centers. Let $\omega \in \Omega$ and take some sequences $x_m \in \omega, t_m \downarrow 0$. We define a normalized blowup sequence as

$$\mathbf{u}_m(x) = \mathbf{u}_{x_m, t_m} = \frac{\mathbf{u}(x_m + t_m x)}{\sqrt{H(x_m, t_m)}}, \qquad \text{for } x \in \frac{\Omega - x_m}{t_m},$$

Thanks to the a priori bounds in Hölder spaces, the blow-up sequences do possess limits that, in two different situations, have the notable property of being homogeneous functions. This property will give us the idea for the classification of nodal points.

Theorem 6.8 (Compactness of blowup sequences). Under the previous notations there exists a function $\mathbf{\bar{u}} \in \mathcal{G}_{loc}(\mathbb{R}^{n+1})$ such that, up to a subsequence, $\mathbf{u}_m \to \mathbf{\bar{u}}$ in $\mathcal{C}_{loc}^{0,\alpha}(\mathbb{R}^{n+1})$ for every $0 < \alpha < 1/2$ and strongly in $H^1_{loc}(\mathbb{R}^n)$. Furthermore we have $\mathbf{\bar{u}} \in \mathcal{C}_{loc}^{0,1/2}(\mathbb{R}^{n+1})$.

Corollary 6.9. Under the previous notations, suppose that one of these situations occurs:

(i) $x_m = x_0 \in \mathcal{N}(\mathbf{u})$ for every m, (ii) $x_m \in \mathcal{N}(\mathbf{u})$ and $x_m \to x_0 \in \mathcal{N}(\mathbf{u})$ with $N(x_0, \mathbf{u}, 0^+) = 1/2$.

Then there are h > 2 and $\nu > 1/2$ such that $\bar{\mathbf{u}}$ is homogeneous of degree ν with h nontrivial components.

In other words, $\mathbf{u} \in \mathcal{H}(\nu, n, h)$, the class of functions introduced in the following definition.

Definition 6.10 (Segregated entire homogenous profiles). For $\nu > 0$ and for every dimension *n*, we define the class

$$\mathcal{H}(\nu, n, h) := \left\{ \mathbf{u} \in \mathcal{G}_{\mathrm{loc}}(\mathbb{R}^{n+1}) : \begin{array}{ll} \mathbf{u} \text{ is } \nu \text{-homogeneous,} \\ \text{ with } h \text{ nontrivial components} \end{array} \right\}$$

6.2. **Segregated entire homogenous profiles.** Now we turn to the study of the entire segregated homogeneous profiles, as in Definition 6.10. For such profiles, the Almgren's quotient centered at zero does not depend on r. The following result has been proved in [23, Lemmata 7.6 and 7.7].

Lemma 6.11. If $h \ge 2$, we have

$$u(n,h) = \inf\{
u > 0: \mathcal{H}(
u,n,h) ext{ is non empty}\} = rac{1}{2}$$

We will need to refine the characterization of this extremal frequency. Toward this aim, we will apply the improvement of flatness Theorem 1.6 and thus at this point we need to connect the two definitions of solution, namely segregated critical profiles, that is elements of the class \mathcal{G} when k = 2, , and viscosity solutions in the sense of Definition 2.3.

Theorem 6.12. Let $\mathbf{u} \in \mathcal{G}(B_1)$ with k = 2; then \mathbf{u} is a viscosity solution in B_1 in the sense of Definition 2.3.

Proof. First, comparing with Definition 6.1, (i) is trivially satisfied. As to (ii), the two conditions follow from the segregation property (G2) and the doubling Proposition 6.6. Next we prove (iii). Indeed, assume not; then, there is a strict comparison subsolution $(v_1, v_2) =: \mathbf{v}$ touching u_1 and u_2 by below and above respectively at $X_0 = (x_0, 0) \in F(v_1, v_2)$. With no loss of generality, we can assume that v_1 and v_2 are harmonic where not vanishing, and thus $\alpha_1(x_0) > \alpha_2(x_0)$ in Definition 2.2.

As $u_1 \ge v_1$ and $F(v_1, v_2) \in C^2$, we have $H(x_0, \mathbf{u}, r) \ge H(x_0, \mathbf{v}, r) \ge Cr$ as $r \to 0$, thus $N(x_0, 0^+) \le 1/2$. Since by Lemma 6.11, 1/2 is the smallest possible critical frequency, we actually infer $N(x_0, 0^+) = 1/2$. Now, according to Theorem 6.8, let us take a converging blowup sequence for \mathbf{u} with fixed centre x_0 , such that $H(x_0, \mathbf{u}, r) \simeq r$ and call its limit $\bar{\mathbf{u}}$. We know that the two components, say \bar{u}_1 and \bar{u}_2 are not vanishing and, up to a rotation, they are positive multiples of (U, \bar{U}) , so $(\bar{u}_1, \bar{u}_2) = (aU, a\bar{U})$ for a positive constant a. On the other hand, let us consider the similarly rescaled sequence

$$\tilde{\mathbf{v}}_m(x) = \tilde{\mathbf{v}}_{x_0, t_m}(x) = \frac{\mathbf{v}(x_0 + t_m x)}{\sqrt{H(x_0, t_m)}} ,$$

and its uniform limit $\tilde{\mathbf{v}} = (\tilde{\alpha}_1(x_0)U, \tilde{\alpha}_2(x_0)\overline{U})$. As it still touches \overline{u}_1 and \overline{u}_2 by below, and above respectively, we deduce

$$a \ge \tilde{\alpha}_1 > \tilde{\alpha}_2 \ge a$$

a contradiction. The proof of (iv) is symmetrical.

In particular, we can apply Theorem 1.6 to segregated critical profiles with two components. This will allow us to prove the following refinement of the Lemma above.

Theorem 6.13. Let $\mathbf{u} \in \mathcal{H}(\nu, n, h)$; then we have

- (i) (uniqueness) $\nu = 1/2$ if and only if h = 2 and, up to orthogonal transformations, $\mathbf{u} = (aU, a\overline{U}, 0, \dots, 0)$, for some positive constant a;
- (ii) (gap) there is $\delta_n > 0$ such that, if $h \ge 2$ and $\nu > 1/2$, then $\nu > \delta_n + 1/2$.

Proof. The characterization of the homogeneous critical profiles depends on some facts related to some spectral properties on the sphere (see also [24, 22, 1]). In our situation, the spectral problem we need to face takes the following form.

Definition 6.14 (The spectral problem). For each open subset ω of $\mathbb{S}^{n-1} := \mathbb{S}^n \cap \mathbb{R}^n \times \{0\}$ we define the first eigenvalue associated with ω as

$$\lambda_1(\omega) := \inf \left\{ \frac{\int_{\mathbb{S}^n} |\nabla_T u|^2 \,\mathrm{d}\sigma}{\int_{\mathbb{S}^n} u^2 \,\mathrm{d}\sigma} : u \in H^1(\mathbb{S}^n), \text{ even and } u \equiv 0 \text{ on } \mathbb{S}^{n-1} \setminus \omega \right\}.$$

Here $\nabla_T u$ stands for the (tangential) gradient of u on \mathbb{S}^n . Furthermore, we define its characteristic exponent as follows

$$\gamma(\omega) = \sqrt{\left(rac{n-1}{2}
ight)^2 + \lambda_1(\omega) - rac{n-1}{2}}$$

As it is well known, u achieves $\lambda_1(\omega)$ if and only if it is one signed, and its $\gamma(\lambda_1(\omega))$ homogeneous extension to \mathbb{R}^{n+1} is harmonic (see [22, 1]). Furthermore, there is a unique minimizer, up to multiplicative constants, and the eigenvalue is strictly monotone with respect to inclusion. Going back to the proof of Theorem 6.13 (i), as in [24], by a rearrangement argument (the foliated Schwarz symmetrization), one may prove that, among sets having the same measure, the minimal eigenvalue, and therefore characteristic exponent, is achieved if and only if ω is a spherical cap. In such a situation, writing $\Gamma(\vartheta) :=$ $\gamma(\lambda_1(\omega_\vartheta))$ for the spherical cap ω_ϑ with opening ϑ , we have

$$\Gamma\left(\frac{\pi}{2}\right) = \frac{1}{2} \; .$$

If $\hat{\omega}$ is the spherical cup of measure one third of that of the sphere, we have

$$\Gamma(\hat{\omega}) = \frac{1}{2} + \hat{\delta}_n$$

for some positive $\hat{\delta}_n$. If $h \ge 3$, at least one of the components supports will intersect the sphere \mathbb{S}^n in a set measuring less than one third of the measure of the whole sphere. Consequently, there holds

$$\nu(n,h) = \inf\{\nu > 0: \mathcal{H}(\nu,n,h) \text{ is non empty}\} \ge \frac{1}{2} + \hat{\delta}_n \ , \qquad \forall h \ge 3 \ ,$$

In other words,

$$\mathbf{u} \in \mathcal{H}(1/2, n, h) \Longrightarrow h = 2$$
.

Now, let $\mathbf{u} = (u_1, u_2) \in \mathcal{H}(1/2, n, 2)$ and $\omega_i = \{x \in \mathbb{S}^{n-1} : u_i(x, 0) \neq 0\}$; then necessarily the ω_i 's have the same measure and are two complementary half spheres. Thus $\mathbf{u} = (u_1, u_2) = (a_1U, a_2\overline{U})$ for some positive constants. On the other hand, by Proposition 6.3 the validity of the domain variation formula (G3) yields $a_1 = a_2$. This proves (i).

In order to prove (ii), in light of the previous discussion we only need to examine the case of two components. Assume by contradiction that there exist sequences $\nu_m \downarrow 1/2$ and $\mathbf{u}_m \in \mathcal{H}(\nu_m, n, 2)$. Thanks to the global compactness Theorem 6.7, we can normalize in such a way that, up to a subsequence, $\mathbf{u}_m = (u_{1,m}, u_{2,m}) \rightarrow \bar{\mathbf{u}} = (U, \bar{U})$ in $C_{\text{loc}}^{0,\alpha}(\mathbb{R}^{n+1})$ for every $0 < \alpha < 1/2$ and strongly in $H_{\text{loc}}^1(\mathbb{R}^n)$. Thus, in the unit ball we get that for m large and $\bar{\varepsilon}$ as in Theorem 1.6,

$$||u_{1,m} - U||_{\infty} \le \bar{\varepsilon}, \quad ||u_{2,m} - U||_{\infty} \le \bar{\varepsilon}.$$

As already remarked, $(u_{1,m}, u_{2,m})$ is a viscosity solution in the sense of Definition 2.3. In view of Theorem 1.6 we infer that $F(u_{1,m}, u_{2,m})$ is a $C^{1,\alpha}$ graph in $\mathcal{B}_{\frac{1}{2}}$ for every $\alpha \in (0,1)$ with $C^{1,\alpha}$ norm bounded by a constant depending on α and n. On the other hand, the nodal set of \mathbf{u}_m is a cone for every m. We then deduce that $F(u_{1,m}, u_{2,m})$ is an hyperplane and $\nu_m = 1/2$ for a sufficiently large m.

7. Viscosity solutions to k-component systems

This section is devoted to the discussion of a suitable notion of viscosity solution to the segregated system, in the spirit of the class S defined in [11] and used in [3].

Definition 7.1. We say that $\mathbf{u} = (u_1, u_2, \dots, u_k)$ is a k-vector viscosity solution in Ω if each $u_i \ge 0$ is a continuous function in Ω which is even with respect to z = 0 and it satisfies

(V1)
$$\Delta u_i = 0 \quad \text{in } \Omega \setminus \mathcal{N}(u_i) ,$$

(V2)
$$\tilde{\Omega} = \bigcup_{i=1}^{k} \mathcal{N}(u_i), \quad \mathcal{N}^o(u_i) \cap \mathcal{N}^o(\sum_{j \neq i} u_j) = \emptyset \; \forall i \; ,$$

(V3) For every $i \in \{1, ..., k\}$, if (v_1, v_2) is a (strict) comparison subsolution then v_1 and v_2 cannot touch u_i and $\sum_{j \neq i} u_j$ respectively by below and above at a point $X_0 = (x_0, 0) \in F(u_i) := \partial \mathcal{N}^o(u_i) \cap \mathcal{B}_1$.

Remark 7.2. In the case k = 2 the above definition coincides with Definition 2.3.

Theorem 7.3. Let $\mathbf{u} \in \mathcal{G}(\Omega)$; then \mathbf{u} is a k-vector viscosity solution in Ω .

Proof. The proof proceeds exactly with the same arguments used in Theorem 6.12, applied repeatedly to the pairs $(u_i, \sum_{j \neq i} u_j)$. Just notice that, in the end of the blowup process we find only two nontrivial components.

7.1. Structure of the nodal set. Our main interest is the study of the free boundary $\mathcal{N}(\mathbf{u}) = \{x \in \tilde{\Omega} : \mathbf{u}(x, 0) = 0\}$ for every element of $\mathfrak{M}(\Omega)$.

Definition 7.4. Given $\mathbf{u} \in \mathfrak{M}(\Omega)$ we define its regular and singular nodal sets respectively by

$$\begin{split} \Sigma_{\mathbf{u}} &= \{ x \in \mathcal{N}(\mathbf{u}) : \ N(x, U, 0^+) = 1/2 \}, \quad \text{and} \\ S_{\mathbf{u}} &= \mathcal{N}(\mathbf{u}) \setminus \Sigma_{\mathbf{u}} = \{ x \in \mathcal{N}(\mathbf{u}) : \ N(x, U, 0^+) > 1/2 \}. \end{split}$$

As a consequence of Theorem 6.13 and Proposition 6.5, the first set is relatively open while the second is relatively closed in $\mathcal{N}(\mathbf{u})$.

Theorem 7.5. Let $\mathbf{u} \in \mathfrak{M}(\Omega)$. Then

- (i) $\Sigma_{\mathbf{u}}$ is a locally finite collection of hyper-surfaces of class $C^{1,\alpha}$ (for some $0 < \alpha < 1$).
- (ii) $\mathcal{H}_{dim}(S_{\mathbf{u}}) \leq n-2$ for any $n \geq 2$. Moreover, for n = 2, $S_{\mathbf{u}}$ is a locally finite set.

Proof of Theorem 7.5 (ii). We apply the Federer's Reduction Principle, Theorem B.1, to the following class of functions

$$\mathcal{F} = \left\{ \mathbf{u}|_{\mathbb{R}^n} \in \left(L^{\infty}_{\text{loc}}(\mathbb{R}^n) \right)^k : \text{there exists some domain } \Omega \subset \mathbb{R}^{n+1} \\ \text{such that } B_2(0) \Subset \Omega \text{ and } \mathbf{u}|_{\Omega} \in \mathfrak{M}(\Omega) \right\},$$

where \mathfrak{M} is the class of minimizing segregated configuration of Definition 1.3. Indeed (H1) is obvious and (H2) holds by Theorem 6.8. In order to check (H3), let us define $\mathcal{S}: \mathcal{F} \to \mathcal{C}$ by $\mathcal{S}(\mathbf{u}) = S_{\mathbf{u}}$ (which is closed by Theorem 6.13 (ii). It is immediate to check (H3)-(i). As for (H3)-(ii), take $\mathbf{u}_m, \mathbf{u} \in \mathcal{F}$ as stated. Then in particular $\mathbf{u}_m \to \mathbf{u}$ uniformly in $B_2(0)$ and by using Theorem 6.7 it is easy to obtain strong convergence in $H^1(B_{3/2}(0))$. Suppose now that (H3)-(ii) does not hold; then there exists a sequence $x_m \in B_1(0)$ ($x_m \to x$, up to a subsequence, for some x) and $\bar{\varepsilon} > 0$ such that $N(x_m, \mathbf{u}_m, 0^+) \ge 1/2 + \delta$ and dist $(x_m, \mathcal{S}(\mathbf{u})) \ge \bar{\varepsilon}$. But then for any small r we obtain

$$N(x_m, \mathbf{u}_m, r) \ge 1/2 + \delta$$

and hence (since $N(x_m, \mathbf{u}_m, r) \to N(x, \mathbf{u}, r)$ in *m* for any *r*) $N(x, \mathbf{u}, 0^+) \ge 1/2 + \delta$, a contradiction. It is a standard fact that **u** inherits the minimizing property from the approximating sequence.

Finally let us prove that $d := \mathcal{H}_{\text{dim}}(S_{\mathbf{u}}) \leq n-2$. At first, we notice that the full d = n dimension is not allowed for energy minimizing configurations, as already explained in Remark 6.2. Next, let us assume that d = n-1: in such a case, we would find a function \mathbf{v} , ν -homogeneous with respect to every point $y \in \mathbb{R}^{n-1} \times \{0\}$, such that $S_{\mathbf{v}} \neq \emptyset$. In such a case, by performing a sequence of blow-ups centered at the elements of the canonical basis of $\mathbb{R}^{n-1} \times \{0\}$, we obtain the existence of a ν -homogeneous element $\mathbf{u} \in \mathcal{G}$ whose full nodal set $\mathcal{N}(\mathbf{u})$ is invariant under translations in the same linear space. But this yields $\nu = 1/2$, a contradiction.

Proof of Theorem 7.5 (i). To prove point (i) we need to establish some preliminary issues relative to flatness and separation properties of the regular part of the free boundary. Our next Lemma 7.6 gives that $\Sigma_{\mathbf{u}}$ verifies the (n-1)-dimensional (ε, R) -Reifenberg flatness condition for every $0 < \varepsilon < 1$ and some $R = R(\varepsilon) > 0$. It is a straightforward consequence of the compactness of the blowup sequences with variable centers, Corollary 6.9, and the characterization of the entire profiles at the lowest frequency $\nu = 1/2$, Theorem 6.13 (see the very similar statement and proof in [21, Lemma 5.3]). Subsequently, in Lemma 7.7, we will show that, in a neighborhood of a regular point, only two components are not identically vanishing, and they are a viscosity solution to the two-component system in the sense of Definition 2.3. Moreover, as argued before, we can arrange for their normalized blow-up to converge locally uniformly to (U, \overline{U}) . This concludes the proof in light of Theorem 1.6.

Lemma 7.6. For $\tilde{\omega} \Subset \Omega$ and any given $0 < \varepsilon < 1$, there exists R > 0 such that for every $x \in \Sigma_{\mathbf{u}} \cap \tilde{\omega}$ and 0 < r < R there exists a hyper-plane $H = H_{x,r}$ containing x such that ¹

(7.1)
$$d_{\mathcal{H}}(\mathcal{N}(\mathbf{u}) \cap B_r(x), H \cap B_r(x)) \le \varepsilon r.$$

Lemma 7.7 (Local Separation Property). Given $x_0 \in \Sigma_{\mathbf{u}}$ there exists a radius $R_0 > 0$ such that $\mathcal{B}_{R_0}(x_0) \subset \tilde{\Omega}$ and in \mathcal{B}_{R_0} there are only two nontrivial components, say u_i and

¹Here $d_{\mathcal{H}}(A, B) := \max\{\sup_{a \in A} \operatorname{dist}(a, B), \sup_{b \in B} \operatorname{dist}(A, b)\}$ denotes the Hausdorff distance. Notice that $d_{\mathcal{H}}(A, B) \leq \varepsilon$ if and only if $A \subseteq N_{\varepsilon}(B)$ and $B \subseteq N_{\varepsilon}(A)$, where $N_{\varepsilon}(\cdot)$ is the closed ε -neighborhood of a set: for $A \subseteq \mathbb{R}^n$ and $\varepsilon > 0$, $N_{\varepsilon}(A) = \{x \in \mathbb{R}^n : \operatorname{dist}(x, A) \leq \varepsilon\}$.

 u_j , so that $\mathcal{B}_{R_0}(x_0) \setminus \mathcal{N}(\mathbf{u}) = \mathcal{B}_{R_0}(x_0) \cap \{|\mathbf{u}| > 0\}$ has exactly two (possibly disconnected) components $\tilde{D}_i = \{u_i > 0\}$, $\tilde{D}_j = \{u_j > 0\}$. Consequently, the pair (u_i, u_j) is a viscosity solution in B_{R_0} in the sense of Definition 2.3.

Proof. Indeed, in view of the previous lemma, for sufficiently small $\varepsilon > 0$, we can uniquely associate with any $y \in \mathcal{N}(\mathbf{u}) \cap \mathcal{B}_{R_0}(x_0)$ and $0 < r < R_0 - |y - x_0|$, and ordered pair of indices (i, j) so that there exist a hyper-plane $H_{y,r}$ (passing through y) and a unitary vector $\nu_{y,r}$ (orthogonal to $H_{y,r}$) such that

$$\{ x + t\nu_{y,r} \in B_r(y) : x \in H_{y,r}, t \ge \varepsilon r \} \subset D_i,$$

$$\{ x - t\nu_{y,r} \in B_r(y) : x \in H_{y,r}, t \ge \varepsilon r \} \subset \tilde{D}_j.$$

Now it is easy to check that the pair (i, j) depends continuously on y and r, thus it is locally constant. Finally, thanks to Theorem 7.3 it is immediate to check that the pair (u_i, u_j) is a viscosity solution, whenever so is **u** and all the other components vanish. \Box

Appendix A

We collect here some know results from [15, 16], which we used throughout the paper. The analogue claims for \bar{U} can be easily obtained and the details are left to the reader.

- First of all, we use the following straightforward properties of the function U:
 - (i) $\Delta U = 0$, U > 0 in $\mathbb{R}^{n+1} \setminus P$.

(ii)
$$U_t = \frac{1}{2}r^{-1/2}\cos\frac{\theta}{2} = \frac{1}{2r}U$$
 and $U_t > 0$ in $\mathbb{R}^{n+1} \setminus P$.

Since U_t is positive harmonic in $\mathbb{R}^2 \setminus \{(t,0), t \leq 0\}$, homogenous of degree -1/2and vanishes continuously on $\{(t,0), t < 0\}$ one can see from boundary Harnack inequality (or by direct computation) that values of U_t at nearby points with the same second coordinate are comparable in dyadic rings. Precisely we have

(A.1)
$$\frac{U_t(t_1,s)}{U_t(t_2,s)} \le C \quad \text{if} \quad |t_1 - t_2| \le \frac{1}{2} |(t_2,s)|.$$

Lemma A.1. Let $g \in C(B_2)$, $g \ge 0$ be a harmonic function in $B_2^+(g)$ and let $\overline{X} = \frac{3}{2}e_n$. Assume that

$$g \ge U$$
 in B_2 , $g(\bar{X}) - U(\bar{X}) \ge \delta_0$

for some $\delta_0 > 0$, then

(A.2)
$$g \ge (1 + c\delta_0)U \quad in B_1$$

for a small universal constant c. In particular, for any $0 < \varepsilon < 2$

(A.3)
$$U(X + \varepsilon e_n) \ge (1 + c\varepsilon)U(X) \quad in B_1,$$

with \boldsymbol{c} small universal.

Lemma A.2. For any $\varepsilon > 0$ small, given $2\varepsilon < \overline{\delta} < 1$, there exists a constant C > 0 depending on $\overline{\delta}$ such that

$$U(t+\varepsilon,z) \leq (1+C\varepsilon)U(t,z)$$
 in $\overline{B}_1 \setminus B_{\overline{\delta}} \subset \mathbb{R}^2$.

Lemma A.3. Let $g \in C(\overline{B}_2)$, $g \ge 0$ be a harmonic function in $B_2^+(g)$ satisfying

$$\|g - U\|_{L^{\infty}(\overline{B}_2)} \le \delta_2$$

and

(A.5)
$$\{x \in \mathcal{B}_2 : x_n \le -\delta\} \subset \{x \in \mathcal{B}_2 : g(x,0) = 0\} \subset \{x \in \mathcal{B}_2 : x_n \le \delta\},\$$

with $\delta > 0$ small universal. Then

(A.6)
$$U(X - \varepsilon e_n) \le g(X) \le U(X + \varepsilon e_n)$$
 in B_1 ,

for some $\varepsilon = K\delta$, K universal.

Lemma A.4. Let $w_1, w_2 \in C(B_1)$ satisfy

$$\Delta(U_n w_i) = 0, \quad in B_1 \setminus P^-, i=1,2.$$

Then w_1 and w_2 cannot touch (either by above or below) on $P^- \setminus L$, unless they coincide.

Appendix B

Here we state Federer's Reduction Principle in a form suited to our needs. Take a class of functions \mathcal{F} invariant under rescaling and translation, and consider a map \mathcal{S} which associates to each function $\Phi \in \mathcal{F}$ a subset of \mathbb{R}^n . For us $\mathcal{S}(\Phi)$ will be the singular set associated with the segregated configuration Φ . This principle establishes rather elementary conditions on \mathcal{F} and \mathcal{S} which imply that, the Hausdorff dimension of $\mathcal{S}(\Phi)$ for every $\Phi \in \mathcal{F}$ can be controlled by the Hausdorff dimension of $\mathcal{S}(\Phi)$ for elements which are homogeneous of some degree.

Theorem B.1 (Federer's Reduction Principle). Let $\mathcal{F} \subset L^{\infty}_{loc}(\mathbb{R}^n)$, and define, for any given $\mathbf{u} \in \mathcal{F}, x_0 \in \mathbb{R}^n$ and t > 0, the rescaled and translated function

$$\mathbf{u}_{x_0,t} := \mathbf{u}(x_0 + t \cdot).$$

We say that $\mathbf{u}_m \to \mathbf{u}$ in \mathcal{F} iff $\mathbf{u}_m \to \mathbf{u}$ in $L^{\infty}_{loc}(\mathbb{R}^n)$. Assume that \mathcal{F} satisfies the following conditions:

- (H1) (Closure under appropriate scalings and translations) Given any $|x_0| \le 1 t, 0 < t < 1, \rho > 0$ and $\mathbf{u} \in \mathcal{F}$, we have that also $\rho \cdot \mathbf{u}_{x_0,t} \in \mathcal{F}$.
- (H2) (Existence of a homogeneous "blow-up") Given $|x_0| < 1, t_m \downarrow 0$ and $\mathbf{u} \in \mathcal{F}$, there exists a sequence $\rho_m \in (0, +\infty)$, a real number $\nu \ge 0$ and a function $\bar{\mathbf{u}} \in \mathcal{F}$ homogeneous of degree ν such that, if we define $\mathbf{u}_m(x) = \mathbf{u}(x_0 + t_m x)/\rho_m$, then

$$\mathbf{u}_m \to \bar{\mathbf{u}}$$
 in \mathcal{F} , up to a subsequence.

- (H3) (Singular Set hypotheses) There exists a map $S : \mathcal{F} \to \mathcal{C}$ (where $\mathcal{C} := \{A \subset \mathbb{R}^n : A \cap B_1(0) \text{ is relatively closed in } B_1(0)\}$) such that
 - (i) Given $|x_0| \le 1 t$, 0 < t < 1 and $\rho > 0$, it holds

$$\mathcal{S}(\rho \cdot \mathbf{u}_{x_0,t}) = (\mathcal{S}(\mathbf{u}))_{x_0,t} := \frac{\mathcal{S}(\mathbf{u}) - x_0}{t}.$$

(ii) Given $|x_0| < 1$, $t_m \downarrow 0$ and $\mathbf{u}, \bar{\mathbf{u}} \in \mathcal{F}$ such that there exists $\rho_m > 0$ satisfying $\mathbf{u}_m := \rho_m \mathbf{u}_{x_0, t_m} \to \bar{\mathbf{u}}$ in \mathcal{F} , the following "continuity" property holds:

$$\forall \varepsilon > 0 \ \exists k(\epsilon) > 0 : \ k \ge k(\varepsilon) \Rightarrow \mathcal{S}(\mathbf{u}_m) \cap B_1(0) \subseteq \{ x \in \mathbb{R}^n : \ dist(x, \mathcal{S}(\bar{\mathbf{u}})) < \varepsilon \}.$$

Then, if we define

(B.1)
$$d = \max \{ \dim L : L \text{ is a vector subspace of } \mathbb{R}^n \text{ and there exist } \mathbf{u} \in \mathcal{F}$$

and $\nu \ge 0$ such that $\mathcal{S}(\mathbf{u}) \neq \emptyset$ and $\mathbf{u}_{y,t} = t^{\nu} \mathbf{u} \; \forall y \in L, \; t > 0 \}$

either $S(\mathbf{u}) \cap B_1(0) = \emptyset$ for every $\mathbf{u} \in \mathcal{F}$, or else $\dim_{\mathcal{H}}(S(\mathbf{u}) \cap B_1(0)) \leq d$ for every $\mathbf{u} \in \mathcal{F}$. Moreover in the latter case there exist a function $\mathbf{v} \in \mathcal{F}$, a d-dimensional subspace $L \leq \mathbb{R}^n$ and a real number $\nu \geq 0$ such that

 $\mathbf{v}_{y,t} = t^{\nu} \mathbf{v} \qquad \forall y \in L, \ t > 0, \qquad and \qquad \mathcal{S}(\mathbf{v}) \cap B_1(0) = L \cap B_1(0).$

If d = 0 then $S(\mathbf{u}) \cap B_{\rho}(0)$ is a finite set for each $\mathbf{u} \in \mathcal{F}$ and $0 < \rho < 1$.

This is the readjusted version of the Federer principle as it appears in Simon's book [19, Appendix A]. The version we present here can be seen as a particular case of a generalization made by Chen (see [8, Theorem 8.5] and [7, Proposition 4.5]).

References

- [1] Krzysztof Bogdan and Tomasz Byczkowski. Potential theory for the α -stable Schrödinger operator on bounded Lipschitz domains. *Studia Math.*, 133(1):53–92, 1999.
- [2] Luis A. Caffarelli and Xavier Cabré. Fully nonlinear elliptic equations, volume 43 of American Mathematical Society Colloquium Publications. American Mathematical Society, Providence, RI, 1995.
- [3] Luis A. Caffarelli, Aram L. Karakhanyan, and Fang-Hua Lin. The geometry of solutions to a segregation problem for nondivergence systems. J. Fixed Point Theory Appl., 5(2):319–351, 2009.
- [4] Luis A. Caffarelli and Fang-Hua Lin. Singularly perturbed elliptic systems and multi-valued harmonic functions with free boundaries. J. Amer. Math. Soc., 21(3):847–862, 2008.
- [5] Luis A. Caffarelli and Luis Silvestre. An extension problem related to the fractional Laplacian. Comm. Partial Differential Equations, 32(7-9):1245-1260, 2007.
- [6] Shu-Ming Chang, Chang-Shou Lin, Tai-Chia Lin, and Wen-Wei Lin. Segregated nodal domains of twodimensional multispecies Bose-Einstein condensates. *Phys. D*, 196(3-4):341–361, 2004.
- [7] Xu-Yan Chen. On the scaling limits at zeros of solutions of parabolic equations. J. Differential Equations, 147(2):355–382, 1998.
- [8] Xu-Yan Chen. A strong unique continuation theorem for parabolic equations. *Math. Ann.*, 311(4):603–630, 1998.
- [9] Monica Conti, Susanna Terracini, and Gianmaria Verzini. Nehari's problem and competing species systems. Ann. Inst. H. Poincaré Anal. Non Linéaire, 19(6):871–888, 2002.
- [10] Monica Conti, Susanna Terracini, and Gianmaria Verzini. An optimal partition problem related to nonlinear eigenvalues. J. Funct. Anal., 198(1):160–196, 2003.
- [11] Monica Conti, Susanna Terracini, and Gianmaria Verzini. A variational problem for the spatial segregation of reaction-diffusion systems. *Indiana Univ. Math. J.*, 54(3):779–815, 2005.
- [12] E. Norman Dancer, Kelei Wang, and Zhitao Zhang. Uniform Hölder estimate for singularly perturbed parabolic systems of Bose-Einstein condensates and competing species. J. Differential Equations, 251(10):2737– 2769, 2011.
- [13] E. Norman Dancer, Kelei Wang, and Zhitao Zhang. Dynamics of strongly competing systems with many species. Trans. Amer. Math. Soc., 364(2):961-1005, 2012.
- [14] E. Norman Dancer, Kelei Wang, and Zhitao Zhang. The limit equation for the Gross-Pitaevskii equations and S. Terracini's conjecture. J. Funct. Anal., 262(3):1087–1131, 2012.
- [15] Daniela De Silva and Jean Michel Roquejoffre. Regularity in a one-phase free boundary problem for the fractional Laplacian. Ann. Inst. H. Poincaré Anal. Non Linéaire, 29(3):335–367, 2012.
- [16] Daniela De Silva and Ovidiu Savin. C^{2,α} regularity of flat free boundaries for the thin one-phase problem.
 J. Differential Equations, 253(8):2420–2459, 2012.
- [17] Benedetta Noris, Hugo Tavares, Susanna Terracini, and Gianmaria Verzini. Uniform Hölder bounds for nonlinear Schrödinger systems with strong competition. *Comm. Pure Appl. Math.*, 63(3):267–302, 2010.
- [18] Miguel Ramos, Hugo Tavares, and Susanna Terracini. Extremality conditions and regularity of solutions to optimal partition problems involving Laplacian eigenvalues. Arch. Ration. Mech. Anal., 220(1):363–443, 2016.
- [19] Leon Simon. Lectures on geometric measure theory, volume 3 of Proceedings of the Centre for Mathematical Analysis, Australian National University. Australian National University, Centre for Mathematical Analysis, Canberra, 1983.
- [20] Yannick Sire, Susanna Terracini, and Giorgio Tortone. On the nodal set of solutions to degenerate or singular elliptic equations with an application to *s*-harmonic functions. *preprint*, 2018, https://arxiv.org/abs/1808.01851.

D. DE SILVA AND S. TERRACINI

- [21] Hugo Tavares and Susanna Terracini. Regularity of the nodal set of segregated critical configurations under a weak reflection law. *Calc. Var. Partial Differential Equations*, 45(3-4):273–317, 2012.
- [22] Susanna Terracini, Giorgio Tortone, and Stefano Vita. On s-harmonic functions on cones. Anal. PDE, 11(7):1653-1691, 2018.
- [23] Susanna Terracini, Gianmaria Verzini, and Alessandro Zilio. Uniform Hölder bounds for strongly competing systems involving the square root of the laplacian. J. Eur. Math. Soc. (JEMS), 18(12):2865–2924, 2016.
- [24] Susanna Terracini and Stefano Vita. On the asymptotic growth of positive solutions to a nonlocal elliptic blow-up system involving strong competition. Ann. Inst. H. Poincaré Anal. Non Linéaire, 35(3):831–858, 2018.
- [25] Gianmaria Verzini and Alessandro Zilio. Strong competition versus fractional diffusion: the case of Lotka-Volterra interaction. Comm. Partial Differential Equations, 39(12):2284–2313, 2014.

DANIELA DE SILVA BARNARD COLLEGE, COLUMBIA UNIVERSITY DEPARTMENT OF MATHEMATICS, NEW YORK, NY 10027, USA *E-mail address*: desilva@math.columbia.edu

Susanna Terracini

DIPARTIMENTO DI MATEMATICA "GIUSEPPE PEANO", UNIVERSITÀ DI TORINO, VIA CARLO ALBERTO, 10, 10123 TORINO, ITALY *E-mail address*: susanna.terracini@unito.it