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Evolutionary Algorithm-based Crystal Structure Prediction for Copper(I) Fluoride

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Abstract: Despite numerous experimental studies since 1824, the binary copper(I) fluoride still remains unknown. We have carried out a crystal structure prediction for CuF using the USPEX evolutionary algorithm and a dispersion-corrected hybrid density functional method. In total about 5000 hypothetical structures were investigated. The energetics of the predicted structures were also counter-checked with local second-order Møller-Plesset perturbation theory. We report 39 new hypothetical copper(I) fluoride structures that are lower in energy compared to the previously predicted cinnabar-type structure. Cuprophilic Cu–Cu interactions are present in all the low-energy structures, leading to ordered Cu substructures such as helical or zig-zag-type Cu–Cu motifs. The lowest-energy structure adopts a trigonal crystal structure with space group $P3_121$. From the electronic point of view, the predicted CuF modification is a semiconductor with an indirect band gap of 2.3 eV.

Introduction

Binary copper(I) halides CuCl, CuBr, and CuI are known to be semiconductors possessing the cubic zincblende structure (space group *F*-43*m*).¹ In contrast, copper(I) fluoride (CuF) has not been synthesized unambiguously and is not structurally characterized to this date. In the case of CuF, a major challenge is the disproportionation of Cu(I) to Cu(II) and Cu metal in aqueous solution: CuF \rightarrow CuF₂ + Cu. The first report on the synthesis of CuF dates back to Berzelius in 1824, and since then a series of attempts have been made to obtain CuF.^{2–7} However, to date there has not been a reproducible successful synthesis of CuF and its structure remains a mystery.^{5,8–11} An early experimental attempt to describe the crystal structure of CuF was made in 1933,³ suggesting that CuF possesses the

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same zincblende structure as the other Cu(I) halides. This observation is still being reported in some textbooks and crystal structure databases even though Haendler *et al.* later showed that the interplanar spacings of the CuF crystal structure reported in 1933 are practically identical to those in Cu₂O.¹⁰ In several other studies, the characterization of the product assumed to be CuF was not possible due to the disproportionation to CuF₂ and Cu.^{4,5,12}

Due to the experimental challenges in the synthesis of CuF, several theoretical and computational studies towards improved understanding of CuF and its structural chemistry have been carried out. Barber et al. calculated lattice energies for monofluorides of the 3d metals using simple empirical formula and reported all of them to be unstable towards disproportionation.¹³ Schwerdtfeger et al. carried out pioneering computational studies on the possible structure of CuF by comparing rocksalt, zincblende, and cluster models using density functional theory (DFT).14,15 Walsh et al. used DFT methods to investigate six hypothetical CuF structures based on known binary structure types: cinnabar, graphite, NiAs, sphalerite, rocksalt, and wurtzite.¹⁶ They identified cinnabar as the lowest-energy structure for CuF by using a hybrid DFT method (Figure 1a). The structures of the binary compounds investigated by Walsh et al. were also included in a later study where 22 hypothetical CuF structures were investigated using hybrid DFT and Local-MP2 (LMP2) methods.¹¹ Also in this study, the cinnabar structure was found to be the lowest-energy structure for CuF. Besides the cinnabar structure, a low-energy structure derived by distorting a wurtzite-type structure was found (Figure 1b).



Figure 1. The lowest-energy crystal structures of CuF identified in previous computational studies: (a) cinnabar-type structure $(P3_121)^{16}$ and (b) distorted wurtzite-type structure (*Cmcm*).¹¹ Reddish brown: Cu, light blue: F.

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So far, all hypothetical CuF crystal structures considered in computational studies have been derived from known binary structure types, which is a very strong limitation on the search space. To account for the possibility that CuF adopts its own new structure type, one may utilize crystal structure prediction algorithms that enable screening of a vast number of hypothetical structures.^{17–21} Evolutionary algorithm-based methods (EA) have been proved to be one of the most successful techniques to search for global minima and to predict unknown crystal structures.^{22–27} An example of a robust EA-based method is the USPEX code (Universal Structure Predictor: Evolutionary Xtallography), which has been successfully used for numerous different materials.^{24–33}

Here, we describe the first evolutionary algorithm-based crystal structure prediction study for copper(I) fluoride. By combining the USPEX method with dispersion-corrected hybrid DFT method and *ab initio* LMP2 calculations, we are able to report 39 hypothetical CuF structures that have lower energy than the previously reported cinnabar-type structure.

Results and Discussion

Overview of the CuF structures predicted by USPEX.

To span the configuration space of an unknown crystal structure such as CuF, different numbers of formula units (*Z*) have to be considered. For CuF, we investigated compositions from Z = 2 to Z = 8. For each case, except for Z = 7 and Z = 8, we ran several distinct USPEX searches to improve the sampling of the configuration space (see below for details). In some cases, a symmetry analysis of the structural candidate obtained from USPEX resulted in a final structure with smaller *Z* than in the original USPEX search. For example, an USPEX search with Z = 2, Z = 4, Z = 6, Z = 8 could yield a final structure with Z = 1, *Z*

= 2, Z = 3, or Z = 4 respectively. We did not run USPEX for Z = 1 as this composition proved to be too small in our tests. Furthermore, the simulations with Z = 2–8 yielded many final structures with Z = 1, which were found to be high-energy structures not discussed here. In general, the number of formula units used in the USPEX simulations is limited by the available computational resources and searches with Z = 7 and Z = 8 proved already to be rather demanding with the hybrid DFT-PBE0 method.

As the cinnabar-type structure was previously obtained as the lowest-energy crystal structure for CuF, we used it as a reference for our energy comparisons. Thus, all relative energies (ΔE) are reported with respect to the cinnabar-type (Z = 3): $\Delta E = E(CuF \text{ structure})/Z(CuF \text{ structure}) - E(cinnabar-type CuF)/Z(cinnabar-type CuF).$

Figure 2 illustrates the relative energies of the predicted lowest-energy CuF structures and Table 1 lists detailed information for them. The structures are labeled using a scheme that is explained in the caption of Figure 2. The found structures as well as the differences between the DFT-PBE0-D3/TZVP and LMP2/TZVPP levels of theory are discussed in more detail in the following sections. The unit cell parameters and atom coordinates of the predicted CuF structures are given as SI in CIF format.



Figure 2. Relative DFT-PBE0-D3/TZVP energies of the lowest-energy CuF structures predicted by USPEX. The energies are given with respect to the cinnabar-type CuF ($\Delta E = 0$). The structures are identified by a label like **Z2a**. The label tells the number of formula units used in the USPEX search that produced the structure (Z = 2 for **Z2a**). The letter (a, b, c, ...) tells the energy ranking of the structure within this class of structures, a being the lowest-energy structure. Note that the final Z may be different from the Z used in the original USPEX simulation (see Table 1).

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Structure ^a	Space group	Ζ	ΔE (kJ mol ⁻¹ per	Band gap (eV)	Density (g cm ⁻³)	Coordination number of	Cu–F distances (Å) ^c	Cu–Cu distances
			CuF unit) ^b			the Cu ^c		(Å) ^e
72a	$P_{-1}(2)$	2	_0 1	19	5 70	4	1 97/2 44	2 66 2 70
72h	$P_{-1}(2)$	2	_0.1	23	5.43	2	1.9/ 1.96	3.0
Z3a	P3121	3	-7.9	2.3	5.71	4	2.00/2.22	2.64, 2.96
73h	(152) C2 (5)	3	-0.6	19	5.81	3.5	Cu-3 ^d · 2 00/2 13	2 65-2 95
200	02 (0)	0	0.0	1.0	0.01	3, 3	Cu-5: 2.09, 2.10/2.21–2.36	2.00 2.00
Z3c	<i>P</i> 1 (1)	3	-0.4	1.9	5.80	3, 4, 5	Cu-3: 2.00/2.13 Cu-4: 2.03, 2.07/2.19, 2.34 Cu-5: 2.11, 2.16/2.20–2.29	2.65–2.92
Z4a	Pnma (62)	4	-4.9	2.1	5.62	4	1.92, 1.95/2.37	2.70
Z4b	P2 ₁ (4)	4	-3.9	2.2	5.66	4	1.98-2.04/2.20-2.36	2.58-2.91
Z4c	P-1 (2)	4	-3.3	2.1	5.65	3	1.91-1.98/2.31, 2.43	2.70-2.91
Z4d	P21(4)	4	-3.1	2.0	5.55	3, 4	Cu-3: 1.90, 1.92/2.35	2.62, 2.89
Z4e	$P2_{1}(4)$	4	-1.7	2.2	5.58	4	1.98-2.07/2.21-2.40	2.61.2.68
Z4f	P-1 (2)	4	-1.6	2.2	5.58	3	1.95-1.99/2.33 2.37	2.72-2.91
Z4a	$P2_{1}/c(14)$	4	-1.4	2.0	5.56	4	1.97. 1.99/2.29. 2.30	2.63, 2.88
Z4h	$P_{2_12_12_1}$ (19)	4	-1.3	2.1	5.72	4	1.96, 1.98/2.37, 2.41	2.82, 2.87
Z4i	Pnma (62)	4	-1.3	2.3	5.56	2	1.96, 1.99	2.83
Z4j	P1 (1)	4	-0.9	2.1	5.71	3, 4, 5	Cu-3: 1.98, 2.00/2.16 Cu-4: 1.97–2.07/2.18–2.38 Cu-5: 2.11, 2.12/2.18–2.36	2.62–2.77
Z4k	P-1 (2)	4	-0.8	2.0	5.67	4	1.95–1.98/2.38, 2.44	2.65, 2.68
Z4I	Cc (9)	4	-0.4	1.9	5.83	4	1.99-2.05/2.25-2.46	2.73-2.86
Z4m	C2/c (15)	2	-0.3	2.0	5.26	4	1.95/2.46	2.80
Z4n	<i>P</i> 1 (1)	4	-0.2	2.0	5.72	3, 4	Cu-2: 1.97, 1.99 Cu-3: 1.96–2.03/2.36 Cu-4: 1.98, 2.00/2.37	2.72–2.93
Z4o	C2/c (15)	2	-0.2	2.3	5.47	2	1.95	3.00
Z4p	<i>P</i> 1 (1)	2	-0.1	1.9	5.53	3, 5	Cu-3: 1.93, 1.96/2.39 Cu-5: 2.08, 2.14/2.20–2.27	2.79, 2.97
Z5a	<i>P</i> 1 (1)	5	-4.3	2.1	5.66	3, 4	Cu-3: 1.96, 2.00/2.20 Cu-4: 1.98–2.08/2.17–2.39	2.64–2.94
Z5b	P1 (1)	5	-3.4	2.0	5.73	4	1.99-2.11/2.17-2.31	2.59-2.93
Z5c	C2 (5)	5	-3.3	2.2	5.63	4	1.97-2.07/2.18-2.41	2.63-2.77
Z6a	C2/c (15)	6	-3.9	2.2	5.51	3, 4	Cu-3: 1.92/2.37 Cu-4: 1.96, 1.99/2.31, 2.44	2.63–2.99
Z6b	<i>P</i> 1 (1)	6	-3.3	2.1	5.76	3, 4, 5	Cu-3: 2.05, 2.07/2.10 Cu-4: 1.97–2.02/2.21–2.31 Cu-5: 2.09, 2.14/2.20–2.28	2.59–2.90
Z6c	Cc (9)	6	-3.1	2.4	5.25	3	1.92-1.94/2.26-2.29	2.66, 2.67
Z6d	C2/c (15)	6	-2.3	2.1	5.46	3, 4	Cu-3: 1.91, 1.94/2.42 Cu-4: 1.95/2.33	2.71–2.92
Z6e	<i>P</i> -1 (2)	6	-2.1	2.1	5.63	4	1.97–2.10/2.20–2.49	2.70–2.80
Z6f	Cmc2 ₁ (36)	6	-2.0	2.0	5.83	3, 4	Cu-3: 2.03/2.11 Cu-4: 2.03/2.19, 2.25	2.69, 2.74
Z6g	<i>P</i> 1 (1)	6	-2.0	2.1	5.58	2, 4	Cu-2: 1.97 Cu-4: 1.99–2.13/2.16–2.34	2.80, 2.87
Z6h	Cc (9)	6	-1.7	2.1	5.78	4	1.97-2.16/2.18-2.44	2.58-2.95
Z6i	P1 (1)	6	-1.7	2.3	5.54	2, 3, 4	Cu-2: 1.93, 1.96 Cu-3: 1.92–1.99/2.44	2.75–2.89
							Cu-4: 1.93–2.00/2.33–2.42	
Z6j	Cmc2 ₁ (36)	6	-1.4	2.1	5.50	3, 4	Cu-3: 1.98/2.19 Cu-4: 2.05/2.11, 2.35	2.65, 2.69
Z6k	P-1 (2)	6	-0.6	2.2	5.58	2, 3, 5	Cu-2: 1.94 Cu-3: 1.92, 1.93/2.46 Cu-5: 2.13, 2.14/2.30–2.41	2.70–2.89
Z8a	Pbca (61)	8	-2.5	2.0	5.79	4	2.00, 2.10/2.17, 2.30	2.67
Z8b	P2 ₁ /c (14)	8	-1.8	2.3	5.64	2	1.95, 1.96	2.67-2.97
Z8c	P1 (1)	8	-1.0	2.2	5.59	3, 4	Cu-3: 1.91–1.99/2.33, 2.38 Cu-4: 1.97–2.09/2.22–2.45	2.66–2.94
Z8d	$P2_{1}/c(14)$	4	-0.3	2.2	5.75	3	1.97, 2.00/2.40	2.76, 2.92

 $^{\rm a}$ Label of the structure. See caption of Figure 2 for an explanation of the labeling.

^b Relative energy compared to the previously reported cinnabar-type CuF.

 $^{\circ}$ Cu–F distances smaller than 2.49 Å were used to determine the coordination number of Cu.

^d When there are Cu atoms with different coordination numbers, these are denoted as Cu-3, Cu-4, *etc*.

^e Cu-Cu distances longer than 3 Å are not listed.

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CuF structures from USPEX simulations with composition $\ensuremath{\text{Cu}_2\text{F}_2}$

We started the structure prediction with a composition possessing only two copper and two fluorine atoms in the unit cell (Cu2F2). Overall, five USPEX simulations were carried out, screening 1090 hypothetical Cu2F2 structures in total. However, only two structures were found to possess relative energy that is similar to the cinnabar-type structure: **Z2a** (space group *P*-1) and **Z2b** (space group *P*-1). Both structures are triclinic, whereas the cinnabar-type CuF should adopt the trigonal space group *P*3121. Band gaps of **Z2a** and **Z2b** are rather similar at about 2 eV, while the density of the **Z2a** structure (5.70 g cm–3) is larger than that of **Z2b** (5.43 g cm–3).

The USPEX simulation for Z = 2 also produced the distorted wurtzite structure discovered in a previous study to be a lowenergy structure (Figure 1b).¹¹ Compared to the previous PBE0/TZVP results, we found that taking dispersion interactions into account clearly changes the energetics of the distorted wurtzite structure. Previously the structure was estimated to be 1.2 kJ mol⁻¹ per CuF unit higher in energy than the cinnabartype structure, but inclusion of the dispersion corrections increases the relative energy to 2.8 kJ mol⁻¹. Structural parameters of the distorted wurtzite also change clearly in case of DFT-PBE0-D3/TZVP, resulting in a = 2.90 Å, b = 5.38 Å, and c = 6.25 Å, while the corresponding DFT-PBE0/TZVP values were 3.09 Å, 5.50 Å, and 6.38 Å. Cu–F distances are almost identical, 1.95 Å and 1.96 Å, for DFT-PBE0/TZVP and DFT-PBE0-D3/TZVP, respectively.

CuF structures from USPEX simulations with composition Cu_3F_3

We ran three USPEX simulations with composition Cu_3F_3 , resulting in total in 591 hypothetical CuF structures. Three structures were found to possess lower energy than the cinnabar-type structure (Table 3). All *Z* = 3 structures produced by USPEX are lower in energy than the structures obtained with the starting composition Cu_2F_2 . Among all predicted CuF structures, structure **Z3a** was found to be the lowest-energy structure with $\Delta E = -7.9$ kJ mol⁻¹ per CuF.

Z3a adopts a trigonal crystal structure with space group P3₁21. This is in fact the same space group as for the cinnabar-type CuF and the structure is closely related (Figure 3) as also the Pearson code and Wyckoff sequence are similar. If only the shortest Cu-F distances are considered (2.0 Å), the coordination number of both the Cu and F atoms is two. While the Cu atom is almost linearly coordinated by two F atoms (162.5°), the F atom interconnects two Cu atoms in an angle of 82.3°. The 3₁ screw axis of the space group leads to onedimensional infinite helices described with the Niggli formula 1_{∞} [Cu_{2/2}F_{2/2}] parallel to the c-axis as is the case for α -HgS. Due to the 31 screw axis, one turn of the helix is finished after the length of the c-axis. One turn of the Cu-F-helix is finished after the c-axis length of 2.96 Å, whereas in α-HgS one turn is finished after circa 9.44 Å. While the coordination numbers of the Hg and S atoms in α -HgS can be described as Hg^[2+2+2]S^[2+2+2], both reminiscent of octahedra as in galena (PbS), the coordination numbers of the respective atoms in

CuF are much better described as Cu^[2+2]F^[2+2]. In α -HgS, the Hg–S distances within the octahedron-like coordination polyhedron are 2.2, 3.1, and 3.5 Å. The respective Cu–F distances are 2.0 and 2.2 Å, while the next-nearest Cu–F distances are 3.1 Å. So, a coordination number of 4 in CuF is much more pronounced in comparison to α -HgS.

Based on single-bond covalent radii from the literature, Cu–F single bonds would be expected to be about 1.76 Å.⁶³ However, in CuF₂ with a Jahn-Teller distorted octahedral coordination, there are four Cu–F distances of 1.91 Å and two distances of 2.31 Å. The CuF structures predicted here have Cu–F distances that are rather similar to CuF₂, a typical pattern being that there are two distances slightly below 2.0 Å and two distances clearly longer than this (about 2.2–2.4 Å). The coordination numbers in Table 1 have been reported in such way that both the "short" and "long" Cu–F distances are counted in the coordination number.

Table 1 shows that the predicted CuF structures have Cu–Cu distances shorter than 3.0 Å, suggesting weak cuprophilic Cu–Cu interactions. The shortest Cu–Cu distance in **Z3a** is 2.64 Å. Each Cu atom has two Cu–Cu distances of 2.64 Å and two longer Cu–Cu distances of 2.96 Å that are equal to the lattice parameter *c*. Due to the 3₁ screw axis, also the closer Cu–Cu contacts form helical chains along the *c*-axis (Figure 3b). The density and band gap of **Z3a** are 5.71 g cm⁻³ and 2.3 eV, respectively. The band gap is indirect and the valence bands are dominated by contributions from Cu atoms. The band structure and the density of states of **Z3a** are illustrated in the SI (Figure S1).

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Figure 3. Predicted low-energy CuF structure **Z3a** ($P3_121$, Z = 3). Reddish brown atoms are Cu, light blue atoms F. a) Primitive unit cell, b) Illustration of the Cu atom substructure and the short Cu–Cu distances of 2.64 Å, c) View that highlights the Cu–F helices. Cu–F distances of 2.0 Å shown in red and the interhelical Cu–F distances of 2.2 Å shown as dotted lines.

CuF structures from USPEX simulations with composition Cu_4F_4

We carried out five USPEX simulations with composition Cu_4F_4 , resulting in total in 1441 hypothetical CuF structures. 16 structures were found to possess lower energy than the cinnabar-type structure (Table 1). Three of the low-energy structures obtained from a starting composition Cu_4F_4 actually have Z = 2 in the final crystal structure (**Z4m**, **Z4o**, and **Z4p**).

As in case of Z = 3, the lowest-energy Z = 4 structures produced by the USPEX simulation with the starting composition Cu₄F₄ are lower in energy than the structures obtained with starting composition Cu₂F₂, indicating that Z = 2is not large enough to cover the CuF configuration space. Among all predicted CuF structures with Z = 4, structure **Z4a** was found to be the lowest-energy structure with $\Delta E = -4.9$ kJ mol⁻¹ per CuF (Figure 4).



Figure 4. Predicted low-energy CuF structure Z4a (*Pnma*, Z = 4). Reddish brown atoms are Cu, light blue atoms F. a) Primitive unit cell, b) Illustration of the Cu atom substructure and the short Cu–Cu distances of 2.70 Å, c) $2\times 2\times 2$ supercell of the structure.

The structure **Z4a** adopts an orthorhombic crystal structure with space group *Pnma*. The coordination number of both Cu and F atoms in **Z4a** is 4, with two shorter Cu–F distances of 1.92 and 1.95 Å and two longer distances of 2.36 and 2.37 Å (Figure 4a). The F–Cu–F angle is 168.9° for the shorter and 76.0° for the longer Cu–F distances, respectively. Each Cu atom has two short Cu–Cu distances of 2.70 Å, forming a zig-

zag-type chain of such interactions along the *b*-axis of the structure (Figure 4b). The Cu–Cu distance in **Z3a** is slightly shorter (2.64 Å). Each Cu atom also possesses two next-nearest neighbors at a Cu–Cu distance of 2.91 Å along the *b*-axis. The band gap of about 2 eV and the density of about 5.6 g cm⁻³ for the **Z4a** structure are rather typical among the **Z4** structures.

Overall, as in the case of the Z2 and Z3 structures, the Cu– F distances in Z4 structures are about 1.9–2.0 Å for the shorter and about 2.2–2.4 Å for the longer ones. The Cu–Cu distances vary from 2.6 to 2.9 Å, the only exception being Z4o with 3.0 Å. The coordination numbers of Cu in the Z4 structures vary from 2 to 5 and some structures possess Cu atoms with different coordination numbers (Z4d, Z4j, Z4n, and Z4p).

Besides the structure **Z4a**, the structures **Z4b** (*P*2₁), **Z4c** (*P*-1), and **Z4d** (*P*2₁) were found to be low-energy structures with ΔE between -3 and -4 kJ mol⁻¹ (Figure 5). As the structures have rather low symmetry, they possess a range of different Cu–F distances, which are overall rather similar: 1.9–2.0 Å for the short distances and 2.2–2.4 Å for the long distances. Cu–Cu distances vary between 2.6 and 2.9 Å.

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Figure 5. Low-energy CuF structures Z4b, Z4c, and Z4d, obtained from USPEX simulations with composition Cu₄F₄. Reddish brown atoms are Cu, light blue atoms F.

Other **Z4** structures that are not illustrated here are available as SI in CIF format. One of the predicted **Z4** structures, **Z4g** ($P2_1/c$), looks similar to the distorted MnP structure investigated also in a previous study on hypothetical CuF structures.¹¹ Also, structures **Z4a** (Pnma), **Z4d** ($P2_1$), **Z4g** ($P2_1/c$), and **Z4k** (P-1) seem to have a very similar framework; however, the energy difference between **Z4a** and **Z4k** is 4.1 kJ

mol⁻¹ per CuF unit (Table 1). This illustrates that even small structural modifications lead to significant changes in the energetics of the CuF structures studied here.

CuF structures from USPEX simulations with composition $\text{Cu}_{5}\text{F}_{5}$

We carried out two USPEX simulations with composition Cu_5F_5 , producing 440 hypothetical structures. Only three predicted Z = 5 structures were found to have lower energy than the cinnabar-type structure (Table 1). However, all of them possess rather low relative energy (Figure 6). All Z = 5 candidates are low-symmetry structures, **Z5a** and **Z5b** being triclinic and **Z5c** monoclinic (*C*2). **Z5a** includes two types of Cu atoms with different coordination numbers (three and four), whereas all Cu atoms in **Z5b** and **Z5c** are surrounded by four fluorine atoms. The Cu substructure in **Z5a** combines helical and zig-zag motifs seen in **Z3** and **Z2** structures.

CuF structures from USPEX simulations with composition Cu_6F_6

We ran three USPEX simulations with composition Cu_6F_6 , resulting in total 1068 structures and 11 hypothetical structures that have lower energy than cinnabar-type CuF (Table 1). The USPEX simulation with composition Cu_6F_6 also reproduced the lowest-energy structure **Z3a**, but here we focus on structures with *Z* = 6. The structures **Z6a** (*C2/c*), **Z6b** (*P*1), and **Z6c** (*Cc*) were found to be among the lowest-energy hypothetical CuF structures (Figure 7), their relative energies being rather similar to structures **Z4b**, **Z4c**, and **Z4d** (**Z6a**: -3.9 kJ mol⁻¹, **Z6b**: -3.3kJ mol⁻¹, **Z6c**: -3.1 kJ mol⁻¹). The structure **Z6c** has the lowest density among all studied CuF structures (5.25 g cm⁻³). Cu–F and Cu–Cu distances of all **Z6** structures are rather similar: short Cu–F distances are about 2.0 Å, and longer distances are about 2.2–2.4 Å. Cu–Cu distances are about 2.7–3.0 Å. These ranges are also similar to **Z2** and **Z4** structures.



Figure 6. Low-energy CuF structures Z5a, Z5b, and Z5c, obtained from USPEX simulations with composition $Cu_5F_5.$ Reddish brown atoms are Cu, light blue atoms F.



Figure 7. Low-energy CuF structures Z6a, Z6b, and Z6c, obtained from USPEX simulations with composition Cu_6F_6 . Reddish brown atoms are Cu, light blue atoms F.

CuF structures from USPEX simulations with composition $\ensuremath{\text{Cu}_7\text{F}_7}$

We ran one USPEX simulation with the composition Cu_7F_7 , producing 231 structures in total. No structures were found to possess lower energy in comparison to cinnabar-type CuF. The lowest energy Cu_7F_7 structure is 0.2 kJ mol⁻¹ per CuF higher in energy in comparison to cinnabar-type CuF.

CuF structures from USPEX simulation with composition Cu_8F_8

Finally, we ran one USPEX simulation with the composition Cu₈F₈, producing 313 structures in total. Only four structures were found to possess lower energy than cinnabar-type CuF (Table 1). The USPEX simulations with Z = 8 are already computationally rather demanding with the DFT-PBE0-D3 method and the small number of low-energy structures suggests that the configurational space has been rather wellspanned already by the USPEX simulations with smaller Z. The structures Z8a-Z8d have lower relative energy than the lowest-energy Z2 structures, but overall the structure Z8a is not among the 10 lowest energy structures ($\Delta E = -2.5 \text{ kJ mol}^{-1}$ per CuF). To summarize, the composition Cu₂F₂ appears to be too small to span essential parts of the configurational space, while Cu₈F₈ is already computationally rather demanding and does not produce low-energy structures with as good hit rate as the compositions $Cu_3F_3-Cu_6F_6$.

Relative energies of the predicted CuF structures studied with the *ab initio* LMP2 method

The DFT-PBE0-D3/TZVP energy comparisons show that there are three hypothetical CuF structures for which the relative energy is lower than -4.0 kJ mol⁻¹ per CuF (**Z3a**, **Z4a**, and **Z5a** with $\Delta E = -7.9$, -4.9, and -4.3 kJ mol⁻¹ per CuF unit, respectively). In addition, there are many structural candidates within a rather narrow energy range from -3 to -4 kJ mol⁻¹ per CuF unit. Even the lowest-energy structures found here do not appear to be stable with respect to disproportionation to CuF₂ and Cu metal at ambient pressure. The disproportionation reaction CuF(Z3a) \rightarrow CuF₂ + Cu is excenergetic by -37 kJ mol⁻¹, whereas previous studies reported -34 kJ mol⁻¹ and -49 kJ mol⁻¹.^{11,16} We also investigated the disproportionation reaction at 10 GPa pressure, where the reaction is still excenergetic by -37 kJ mol⁻¹.

To further investigate the energetics of the predicted CuF structures, in particular from the point of view of the weak Cu–Cu interactions, we also carried out single-point energy calculations at the LMP2/TZVPP level of theory (Table 2).⁶⁴ While the D3 dispersion correction is empirical, the LMP2 method offers an *ab initio* approach for evaluating the energetics of the CuF structures with weak Cu–Cu interactions.

Table 2. Relative energies of the predicted CuF at the DFT-PBE0-D3/TZV	Ρ
and LMP2/TZVPP levels of theory.	

Structure	ΔE (kJ mol ⁻¹ p	er CuF unit)
	DFT-PBE0-	LMP2/TZVPP
	D3/TZVP	
Z3a	-7.9	-4.8 (-5.4) ^a
Z4a	-4.8	-7.0 (-4.6)
Z5a	-4.3	-1.9 (-0.4)
Z6a	-3.9	-3.8 (-3.4)
Z4b	-3.9	-0.8 (-0.3)
Z5b	-3.4	0.2
Z6b	-3.3	0.5
Z4c	-3.3	-3.2
Z5c	-3.3	1.4 (1.9)
Z6c	-3.1	-5.6 (-3.9)
Z4d	-3.1	-3.0
Z8a	-2.5	1.7
Z6d	-2.3	-4.4 (-2.1)
Z6e	-2.1	-0.4
Z6f	-2.0	3.2
Z6g	-2.0	-1.0
Z8b	-1.8	-0.5
Z6h	-1.7	2.8
Z6i	-1.7	-1.3
Z4e	-1.7	1.2
Z4f	-1.6	-0.5
Z6j	-1.4	1.9
Z4g	-1.4	0.4
Z4h	-1.3	1.5
Z4i	-1.3	-1.5
Z8c	-1.0	0.7
Z4j	-0.9	2.2
Z4k	-0.8	2.1
Z6k	-0.6	-0.4
Z3b	-0.6	3.5
Z3c	-0.4	3.6
Z4I	-0.4	3.3
Z8d	-0.3	2.9
Z4m	-0.3	-2.8
Z4n	-0.2	2.7
Z4o	-0.2	-0.6
Z4p	-0.1	1.9
Z2a	-0.1	2.5
Z2b	-0.1	-0.7

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^a Relative energies in parentheses have been obtained with the extended TZVPP+f basis set (see text).

The DFT-PBE0-D3/TZVP and LMP2/TZVPP relative energies of the hypothetical CuF structures are illustrated in Figure 8. Altogether 39 crystal structures were predicted by DFT to have lower energy than the cinnabar-type CuF. Single-point energy calculations of these structures with LMP2 suggest that 19 of the structures have lower energy than cinnabar-type CuF.

The two lowest-energy structures predicted by DFT retain a low relative energy also with LMP2, but their order is reversed. The relative energy of the structure **Z3a** deteriorates from -7.9 to -4.8 kJ mol⁻¹ per CuF unit, whereas the relative energy of the structure **Z4a** improves from -4.9 to -7.0 kJ mol⁻¹ per CuF unit. The relative energy of **Z5a** deteriorates from -4.3 to -1.9 kJ mol⁻¹ per CuF unit.

Concerning the DFT and LMP2 relative energies of other low-energy structures, the two methods are in relatively good agreement for the following cases (energies in kJ mol⁻¹ per CuF unit): -3.3 vs. -3.2 for **Z4c**, -3.1 vs. -3.0 for **Z4d**, -1.3 vs. -1.5 for **Z4i**, -0.2 vs. -0.6 for **Z4o**, -3.9 vs. -3.8 for **Z6a**, -1.7 vs. -1.7 vs. -1.5

1.3 for **Z6i**, -0.6 vs. -0.4 for **Z6k**. In some cases, though, the DFT and LMP2 energies differed significantly. The largest deviations between the DFT and LMP2 energies were found to be about 5 kJ mol⁻¹ (-3.3 kJ vs. 1.4 kJ mol⁻¹ per CuF unit for **Z5c**, -2.0 kJ vs. 3.2 kJ mol⁻¹ per CuF unit for **Z6f**, -1.7 kJ vs. 2.8 kJ mol⁻¹ per CuF unit for **Z6h**).

It appears that the CuF structures where the shorter Cu-F distances are larger than 1.96 Å become energetically less favorable with LMP2/TZVPP (structures Z2a, Z3a, Z3b, Z3c, Z4b, Z4e-Z4h, Z4j-Z4l, Z4n, Z5a, Z4p, Z5b, Z5c, Z6b, Z6e-Z6h, Z6j, Z8a, Z8c, Z8d). For comparison, CuF structures where the shorter Cu-F distances are less than 1.96 Å appear to have similar relative energies with DFT-PBE0-D3/TZVP and LMP2/TZVPP (Z4o, Z8b) or possess even lower relative energy with LMP2 (Z2b, Z4a, Z4m, Z6c, Z6d). Since all the structures have been obtained as a result of DFT structural optimizations, it can be that not all of the structures are energetically that favorable with LMP2. USPEX simulations with LMP2 are not yet feasible since analytical LMP2 gradients are not available and the HF+LMP2 calculations are significantly more demanding computationally compared to hybrid DFT.

The LMP2/TZVPP single-point energies also highlighted two structures, where the relative energy improved so much that they became low-energy structures (**Z6c** with improvement from -3.1 to -5.6 kJ mol⁻¹ per CuF unit and **Z6d** with improvement from -2.3 to -4.4 kJ mol⁻¹ per CuF unit). Thus, **Z6c** is the structure with the second-lowest relative energy at the LMP2/TZVPP level of theory.

For the five lowest-energy structures found by DFT-D3 (**Z3a**, **Z4a**, **Z5a**, **Z6a**, **Z4b**), as well as low-energy LMP2/TZVPP structures **Z6c**, and **Z6d**, we carried out LMP2 calculations with extended TZVPP+f basis set to get a better understanding of the basis set effects (see Computational details). In these calculations, the relative LMP2 energy of **Z3a** changed from -4.8 to -5.4 kJ mol⁻¹ per CuF unit and the relative energy of **Z4a** changed from -7.0 to -4.6 kJ mol⁻¹ per CuF unit. This energy ordering of **Z3a** and **Z4a** is in agreement with the DFT-PBE0-D3/TZVP results, but one should keep in mind that the geometries were not optimized with the LMP2 method and this

could further change the relative ordering. The main finding is that the structures **Z3a** and **Z4a** are the lowest-energy structures at both DFT-PBE0-D3/TZVP and LMP2/TZVPP+f levels of theory. For the three other structures **Z5a**, **Z6a**, and **Z4b**, the LMP2 relative energies change as follows: from -3.8 to -3.4 kJ mol⁻¹ per CuF unit (**Z6a**), from -1.9 to -0.4 kJ mol⁻¹ per CuF unit (**Z5a**), and from -0.8 to 0.3 kJ mol⁻¹ per CuF unit (**Z4b**). In the case of **Z6c** and **Z6d**, the relative energies change from - 5.6 to -3.9 kJ mol⁻¹ per CuF unit (**Z6c**) and from -4.4 to -2.1 kJ mol⁻¹ per CuF unit, bringing them close to the DFT-PBE0/D3 relative energies. We also carried out extended basis set calculation for **Z5c**, which has the largest deviation between the DFT-D3 and the LMP2 results. The relative LMP2 energy of the **Z5c** structure changed from 1.4 to 1.9 kJ mol⁻¹ per CuF unit (DFT-D3: -3.3 kJ mol⁻¹ per CuF unit).

Conclusions

We have carried out crystal structure predictions for copper(I) fluoride by using the evolutionary algorithm-based USPEX code and a recently developed CRYSTAL interface. We screened about 5000 hypothetical crystal structures by using a dispersioncorrected hybrid DFT method and also carried out single-point energy calculations with ab initio LMP2 method. We identified 39 hypothetical CuF structures that are lower in energy in comparison to the previously reported cinnabar-type CuF structure (19 structures at the LMP/TZVPP level of theory). The predicted low-energy structures are true local minima that could exist as metastable species, but this remains to be proven by definitive experimental validation. Relatively short Cu-Cu distances suggesting cuprophilic Cu–Cu interactions are present in all the low-energy structures, leading in ordered Cu substructures. Minor structural modifications can lead to significant changes in the relative energies of the CuF structures. The lowest-energy CuF structures predicted here can be used to identify unknown phases in the experimental investigations towards copper(I) fluoride.



Figure 8. Relative DFT-PBE0-D3/TZVP and LMP2/TZVPP energies of the lowest-energy CuF structures predicted by USPEX. The energies are given with respect to the cinnabar-type CuF (ΔE = 0). For the labeling of the structures, please see the caption of Figure 2.

Computational details

Crystal structure predictions were carried out by using USPEX 9.4.4 code (USPEX input file is included as Supporting Information, SI).25,26,28 All quantum chemical calculations within the USPEX simulations were performed using the CRYSTAL17 code.34 We used a recently developed CRYSTAL interface for USPEX.27 Hybrid DFT-PBE0 functional with 25% Hartree-Fock exchange was utilized.35,36 For a 3d metal such as Cu, the use of hybrid DFT over GGA or GGA + U is expected to increase the accuracy of the predictions.³⁷⁻⁴³ All-electron, Gaussian-type triple- ζ valence + polarization (TZVP) and split-valence + polarization (SVP) basis sets based on Karlsruhe def2 basis sets were used within the crystal structure predictions (a list with of all used basis sets is provided as SI).⁴⁴ Cu(I) ions are expected to show weak "cuprophilic" $d^{10}-d^{10}$ interactions,45-48 which were taken into account in the USPEX search using Grimme's D3 dispersion correction with zero-damping (ZD).^{49,50} To accelerate the evolutionary searches, the local geometry optimizations within USPEX were carried out by using relatively weak convergence criteria and, in some cases, the smaller SVP basis set. Input files for CRYSTAL geometry optimizations within the USPEX search are provided as SI.

Low-energy structures produced by USPEX were re-optimized at the DFT-PBE0-D3(ZD)/TZVP level of theory using tighter convergence criteria. For the re-optimization, the reciprocal space *k*-point meshes were chosen depending on the magnitude of the corresponding direct space lattice parameter *d*: d < 4 Å $\rightarrow 12$ *k*-points along *d*; 4 Å < d < 6 Å $\rightarrow 8$, 6 Å < d < 8 Å $\rightarrow 6$; 8 Å < d < 12 Å $\rightarrow 4$, d > 12 Å $\rightarrow 2$. Tightened tolerance factors (TOLINTEG) of 8, 8, 8, 8, and 16 were used for the evaluation of the Coulomb and exchange integrals. All re-optimized structures that were found to be at least 3 kJ mol⁻¹ per CuF lower in energy compared to the cinnabar-type structure were confirmed to be true local minima by means of a harmonic frequency calculation.^{51,52}

To evaluate the performance of the applied USPEX/DFT-PBE0-D3(ZD) crystal structure prediction methodology for d^{10} -metal monofluorides, we confirmed that rocksalt structure (*Fm*-3*m*) was correctly found as the lowest-energy structure for the known AgF.⁵³

For investigating the energetics of the disproportionation reaction $CuF \rightarrow CuF_2 + Cu$, we also optimized the structures of CuF_2 and Cu at the DFT-PBE0-D3/TZVP level of theory. CuF_2 was studied using an antiferromagnetic ground state.⁵⁴ 12×8×8 and 16×16×16 *k*-point meshes were used for CuF_2 and Cu metal, respectively. For metallic Cu, Fermi surface smearing with T = 316 K (0.001 a.u.) was applied and a double-density *k*-mesh was used for the determination of the Fermi energy. The optimized structures of CuF_2 and Cu metal are included as SI. In the USPEX simulations on CuF, we do not expect to see any disproportionation to CuF_2 and Cu as the used unit cells are still relatively small (up to Z = 8) and we are running closed-shell calculations with no unpaired electrons.

To confirm the presence of weak Cu–Cu interactions, we also carried out periodic local second-order Møller-Plesset perturbation theory (LMP2) single-point energy calculations for the CuF structures which were found to possess lower energy than the previously reported cinnabar-type structure at the DFT-PBE0-D3(ZD)/TZVP level of theory. The LMP2 calculations were carried out with a development version of the CRYSCOR software,⁵⁵ which implements orbital specific virtuals (OSVs) to represent the truncated pair-specific virtual space.⁵⁶ In the OSV-LMP2 formalism, it is not necessary to manually define excitation domains for the virtual space as in the previous implementation based on projected atomic orbitals (PAO-LMP2). The OSV-LMP2 straightforwardly enables

the calculation of smooth potential energy surfaces and relative energies of structural frameworks with different topologies. 57

The HF reference wavefunction and the localized valence-space Wannier functions (WFs) necessary for the LMP2 procedure were obtained with CRYSTAL17. Very tight TOLINTEG tolerance factors of 10, 10, 10, 20 and 50 were used in the Hartree-Fock (HF) part. All-electron, and triple- ζ valence + double polarization (TZVPP) basis set based on Karlsruhe def2-TZVPP basis set was used for Cu, while TZVP basis set was used for F. For the five CuF structures with the lowest energy according to DFT-D3, we also calculated LMP2 single-point energies using even larger basis set that added an f-function for F and a third f-function for Cu. These calculations were carried our using a dual basis set scheme.58,59 In the LMP2 calculations, we utilized the direct-space density-fitting technique for computing the two-electron four-index integrals. A Poisson/Gaussian-type auxiliary basis set of triple-zeta-valence quality was employed for the density-fitting.60-62 From a practical point of view, the calculation of the reference wavefunction with HF can be computationally even more expensive than the actual LMP2 calculation. An input file example for the CRYSCOR calculations is available as SI.

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Evolutionary crystal structure prediction algorithms combined with dispersion-corrected hybrid density functional calculations yield a large number of new lowenergy structural candidates for the unknown copper(I) fluoride.



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Evolutionary Algorithm-based Crystal Structure Prediction for Copper(I) Fluoride