Precision measurement of the integrated luminosity of the data taken by BESIII at center-of-mass energies between 3.810 GeV and 4.600 GeV

This content has been downloaded from IOPscience. Please scroll down to see the full text.
2015 Chinese Phys. C 39 093001
(http://iopscience.iop.org/1674-1137/39/9/093001)

View the table of contents for this issue, or go to the journal homepage for more

Download details:
IP Address: 192.135.19.160
This content was downloaded on 07/06/2016 at 12:06

Please note that terms and conditions apply.
Precision measurement of the integrated luminosity of the data taken by BESIII at center-of-mass energies between 3.810 GeV and 4.600 GeV*

M. Ablikim(麦迪娜)1, M. N. Achasov9,9,a X. C. Ai(艾小聪)1, O. Albayrak5, M. Albrecht4, D. J. Ambroise44, A. Amoroso48.4,48c F. F. An(安芬芬)1, Q. An(安琪)45, J. Z. Bao(白晶)1, R. Baldini Ferroli20,4,31, Y. Ban(班勇)31, D. W. Bennett19, J. V. Bennett5, M. Bertani20,4,34, D. Bettoni21,4, J. M. Bian(边瀚鸣)43, F. Bianchi48.4,48c E. Boger23,4,9, O. Bondarenko25, I. Boyko23, R. A. Briere6, H. Cai(蔡浩)50, X. Cai(蔡啸)1, O. Bakr30,4,4, A. Calcaterra20,4, G. F. Cao(曹国富)1, S. A. Cetin40,4, J. F. Chang(常劲帆)1, G. Chelkov23,4, G. Chen(陈刚)1, H. S. Chen(陈和生)1, H. Y. Chen(陈海云)2, J. C. Chen(陈江川)1, M. L. Chen(陈玛丽)1, S. J. Chen(陈奕宁)29, X. Chen(陈晓荣)26, Y. Chen(陈学鹏)26, H. P. Cheng(程和平)17, X. K. Chu(朱新坤)31, G. Cibinetto21,4, D. Cronin-Hennessy43, H. L. Dai(代洪亮)1, J. P. Dai(代建平)34, A. Dbeysi14, D. Dedovich23, Z. Y. Deng(邓子艳)1, A. Denig22, I. Denysenko23, M. Destefanis48.4,48c, F. De Mori48.4,48c, Y. Ding(丁勇)27, C. Dong(董超)30, J. Dong(董静)1, L. Y. Dong(董耀原)1, M. Y. Dong(董明义)1, S. X. Du(杜书先)52, P. F. Duan(段鹏飞)1, J. Z. Fan(范荆州)38, F. Fang(方建)1, S. S. Fang(房双第)1, X. Fang(方啸)45, Y. Fang(方易)1, L. Fava48.4,48c, F. Feldbauer22, G. Felici40,4, C. Q. Feng(封常青)45, E. Fioravanti21,4, M. Fritsch14,22, C. D. Fu(傅成栋)1, Q. Gao(高倩)41, Y. Gao(高原)39, Z. Gao(高棪)45, I. Garzia21,4, C. Geng(耿聪)45, K. Goetzen10, W. X. Gong(龚文斌)1, W. Gradl22, M. Greco48.4,48c, M. H. Gu(顾晓东)21, Y. T. Gu(顾运宇)12, Y. H. Guan(管慧)1, A. Q. Guo(郭爱强)1, L. B. Guo(郭立波)28, Y. Guo(郭珂)1, Y. P. Guo22, Z. Haddad25, A. Hafner22, S. Hau(韩爽)30, Y. L. Han(韩艳龙)4, X. Q. Hao(郝喜庆)15, F. A. Harris42, K. L. He(何康林)1, Z. Y. He(何振亚)30, T. Held4, Y. K. Heng(衡昆)1, Z. L. Hou(侯治龙)1, C. Hu(胡琛)28, H. M. Hu(胡海明)1, J. F. Hu(胡继斐)48.4,48c, T. Hu(胡海)1, Y. Hu(胡莉)1, G. M. Huang(黄光明)6, G. S. Huang(黄光顺)45, H. P. Huang(黄海鹏)50, J. S. Huang(黄子川)15, X. T. Huang(黄思瑜)33, Y. Huang(黄勇)29, T. Hussain47, Q. Ji(纪全)41, P. Q. Ji(姬清平)40, X. B. Ji(李晓斌)1, X. L. Ji(李静宜)1, L. L. Jiang(姜丽丽)1, L. W. Jiang(姜武军)26, X. S. Jiang(江晓山)1, J. B. Jiao(焦健斌)33, Z. Jiao(焦信)17, D. P. Jin(金大鹏)1, S. Jin(金山)1, T. Johansson49, A. Julin41, N. Kalantar-Nayestanaki25, X. L. Kang(康晓思)1, X. S. Kang(康晓思)40, M. Kavatsyuk25, B. C. Ke5, R. Kliemt14, B. Klose22, O. B. Koley48.4,4,6, B. Kopf4, M. Kornicer22, W. Kuehn24, A. Kupsch9, W. Lai(赖南)1, J. S. Lange24, M. Lara19, P. Larin14, C. Leng48c, C. H. Li(李春花)26, Cheng Li(李澄)45, D. M. Li(李德民)22, F. Li(李飞)1, G. Li(李刚)1, H. B. Li(李海波)1, J. C. Li(李家才)1, Jin Li(李瑞)32, K. Li(李康)13, K. Li(李科)33, Lei Li(李蕾)33, P. R. Li(李培盛)41, T. Li(李腾飞)33, W. D. Li(李卫东)1, W. G. Li(李卫国)1, X. L. Li(李晓玲)33, X. M. Li(李小梅)12, X. N. Li(李小勇)1, X. Q. Li(李学刚)30, Z. B. Li(李志兵)38, H. Liang(梁晃)15, Y. F. Liang(梁勇飞)36, Y. T. Liang(梁羽飞)24, G. R. Liao(廖敬春)1, D. X. Liu(刘春)14, B. J. Liu(刘北江)1, C. X. Liu(刘春秀)4, F. H. Liu(刘福辉)35, Fang Liu(刘方)1, Feng Liu(刘峰)6, H. B. Liu(刘宏博)12, H. H. Liu(刘慧慧)16, H. H. Liu(刘欢欢)1, H. M. Liu(刘怀民)1, J. Liu(刘健)1, J. P. Liu(刘岩)50, J. Y. Liu(刘晶晶)41, K. Liu(刘凯)39, K. Y. Liu(刘有才)27, L. D. Liu(刘兰雕)31, P. L. Liu(刘佩莲)1, Q. Liu(刘倩)41, S. B. Liu(刘树彬)45

Received 13 March 2015

*Supported by National Key Basic Research Program of China (2015CB856700), National Natural Science Foundation of China (NSFC) (11125525, 11232544, 11335008, 11425524), Chinese Academy of Sciences (CAS) Large-Scale Scientific Facility Program, Joint Large-Scale Scientific Facility Funds of the NSFC and CAS (11179007, U1232001, U1332001) CAS (KJCX2-YW-N29, KJCX2-YW-N45), 100 Talents Program of CAS, INPAC and Shanghai Key Laboratory for Particle Physics and Cosmology, German Research Foundation DFG (Collaborative Research Center CRC-1044), Istituto Nazionale di Fisica Nucleare, Italy; Ministry of Development of China (NSFC) (11125525, 11235011, 11322544, 11335008, 11425524), Chinese Academy of Sciences (CAS) Large-Scale Scientific Facility Program, Joint Large-Scale Scientific Facility Funds of the NSFC and CAS (11179007, U1232001, U1332001) CAS (KJCX2-YW-N29, KJCX2-YW-N45), 100 Talents Program of CAS, INPAC and Shanghai Key Laboratory for Particle Physics and Cosmology, German Research Foundation DFG (Collaborative Research Center CRC-1044), Istituto Nazionale di Fisica Nucleare, Italy; Ministry of Development of China.
Abstract: From December 2011 to May 2014, about 5 fb$^{-1}$ of data were taken with the BESIII detector at center-of-mass energies between 3.810 GeV and 4.600 GeV to study the charmonium-like states and higher excited charmonium states. The time-integrated luminosity of the collected data sample is measured to a precision of 1% by analyzing events produced by the large-angle Bhabha scattering process.

Key words: precision measurement, luminosity, Bhabha scattering, charmonium

1 Introduction

As a τ-charm factory, the BESIII experiment has collected the world’s largest sample of e+e− collision data at center-of-mass (CM) energies between 3.810 GeV and 4.600 GeV. In this energy region, the charmonium-like states and higher excited charmonium states are produced copiously, which makes comprehensive studies possible.

The charmonium-like states discovered in recent years have drawn great attention from both theorists and experimentalists for their exotic properties, as reviewed e.g. in Ref. [1]. Being well above the open charm threshold, the strong coupling of these states to hidden charm processes makes their interpretation as conventional charmonium states very difficult. On the other hand, the theory of the strong interaction, Quantum Chromodynamics (QCD), does not prohibit the existence of exotic states beyond the quark model, e.g. molecular states, tetraquark states, hybrid states, etc. Either the verification or the exclusion of the existence of such states will help to evaluate the quark model and better understand QCD. Even though some states have been identified as higher excited charmonium states, such as the ψ(4040), ψ(4160), and ψ(4415), their large widths and interference with each other make their precise study complicated. In addition, the relationship between the charmonium-like states and higher excited charmonium states is still not clear. Precise knowledge of the time-integrated luminosity is essential for quantitative analysis of these states.

In this paper, we present a measurement of the integrated luminosity based on analysis of the Bhabha scattering process e+e− → (γ)e+e−. A similar method has been used in the luminosity measurement of ψ(3770) data at BESIII [2]. The process has a simple and clean signature and a large production cross section, which allows for a small systematic and a negligible statistical uncertainty. A cross check of the result is performed by analyzing the di-gamma process e+e− → γγ.

2 The detector

BESIII is a general purpose detector which covers 93% of the solid angle and operates at the e+e− collider BEPCII. A detailed description of the facilities is given in Ref. [3]. The detector consists of four main components: (a) A small-cell, helium-based main drift chamber (MDC) with 43 layers provides an average single-hit resolution of 135 µm, and a momentum resolution of 0.5% for charged tracks at 1 GeV/c in a 1 T magnetic field; (b) An electro-magnetic calorimeter (EMC), consisting of 6240 CsI(Tl) crystals in a cylindrical structure (barrel and two endcaps). The energy resolution for 1.0 GeV photons is 2.5% (5%) in the barrel (endcaps); (c) A time-of-flight system (TOF), constructed of 5 cm thick plastic scintillators, arranged in 88 detectors of 2.4 m length in two layers in the barrel and 96 fan-shaped detectors in the endcaps. The barrel (endcap) time resolution of 80 ps (110 ps) provides 2σ K/π separation for momenta up to about 1.0 GeV/c; (d) A muon counter (MUC), consisting of nine layers of resistive plate chambers in the barrel and eight layers for each endcap. It is incorporated in the iron return yoke of the superconducting magnet. Its position resolution is about 2 cm. A GEANT4 [4, 5] based detector simulation package has been developed to model the detector response. Due to the crossing angle of the beams at the interaction point, the e+e− CM system is slightly boosted with respect to the laboratory frame.

3 Data sample and Monte Carlo simulation

Twenty-one data samples have been taken at CM energies between 3.810 GeV and 4.600 GeV. Six of the data sets exceed the others in accumulated statistics by an order of magnitude. These samples were taken on the peaks of charmonium-like states, like the Y(4260), Y(4360), and Y(4630), or higher excited charmonium states, like ψ(4040), and ψ(4415), in order to study these resonances and their decays in great detail. The data samples taken at the other CM energies serve as scan points to study the behavior of the cross section around these resonances. All individual data samples are listed in Table 1.

At each energy point, one million Bhabha events were generated using the BABAYAGA3.5 generator with the options presented in Table 2. For the BABAYAGA3.5 generator, the uncertainty in calculating the cross section is 0.5%, which meets the demand of the total uncertainty of luminosity measurement. The kinematic distributions of the final state particles from the BABAYAGA3.5 generator are consistent with those from data. In the simulation, the scattering angles of the final state particles were limited to a range from 20° to 160°, which slightly exceeds the angular acceptance of the detector, in order to save on computing resources. An energy threshold of 0.04 GeV was applied to the final state particles. The acollinearity of the events has not been constrained. Finally, the generation was taking into account the running of the electromagnetic coupling constant and final state radiation (FSR).

To study the background and optimize the event selection criteria, an inclusive Monte Carlo (MC) sample corresponding to a luminosity of 500 pb−1 at CM energy of 4.260 GeV was generated, in which the Quantum Electrodynamics (QED) processes e+e− → e+e−,
\[ e^+e^- \rightarrow \mu^+\mu^- \text{ and } e^+e^- \rightarrow \gamma\gamma, \text{ the continuum production of hadrons, and the initial state radiation (ISR) to } J/\psi \text{ and } \psi(3686) \text{ resonance process were included. The BABAYAGA3.5 generator was used to simulate the relevant QED processes. Other processes, such as the decays of the } J/\psi, \text{ were generated with specialized models that have been packaged and customized for the BESIII Offline Software System (BOSS) (see [7] for an overview).} \]

Table 1. Center-of-mass energy, luminosity obtained from the nominal measurement (\( L \)), cross check results (\( L_{ck} \)), and relative differences between the two results. The uncertainties are statistical only. Superscripts indicate separate samples acquired at the same CM energy.

<table>
<thead>
<tr>
<th>CM energy/GeV</th>
<th>( L/\text{pb}^{-1} )</th>
<th>( L_{ck}/\text{pb}^{-1} )</th>
<th>relative difference (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>3.810</td>
<td>50.54\pm0.03</td>
<td>50.11\pm0.08</td>
<td>-0.85\pm0.17</td>
</tr>
<tr>
<td>3.900</td>
<td>52.61\pm0.03</td>
<td>52.57\pm0.08</td>
<td>-0.08\pm0.17</td>
</tr>
<tr>
<td>4.009</td>
<td>481.96\pm0.01</td>
<td>480.54\pm0.23</td>
<td>-0.30\pm0.05</td>
</tr>
<tr>
<td>4.096</td>
<td>52.63\pm0.03</td>
<td>52.32\pm0.08</td>
<td>-0.40\pm0.17</td>
</tr>
<tr>
<td>4.190</td>
<td>43.09\pm0.03</td>
<td>43.08\pm0.08</td>
<td>-0.03\pm0.20</td>
</tr>
<tr>
<td>4.210</td>
<td>54.55\pm0.03</td>
<td>54.27\pm0.09</td>
<td>-0.62\pm0.18</td>
</tr>
<tr>
<td>4.220</td>
<td>54.13\pm0.03</td>
<td>54.22\pm0.09</td>
<td>+0.17\pm0.18</td>
</tr>
<tr>
<td>4.230</td>
<td>44.40\pm0.03</td>
<td>44.64\pm0.08</td>
<td>+0.54\pm0.20</td>
</tr>
<tr>
<td>4.230\textsuperscript{2}</td>
<td>1047.34\pm0.14</td>
<td>1041.56\pm0.37</td>
<td>-0.56\pm0.04</td>
</tr>
<tr>
<td>4.245</td>
<td>55.59\pm0.04</td>
<td>55.52\pm0.09</td>
<td>-0.13\pm0.18</td>
</tr>
<tr>
<td>4.260\textsuperscript{1}</td>
<td>523.74\pm0.10</td>
<td>524.57\pm0.26</td>
<td>+0.16\pm0.06</td>
</tr>
<tr>
<td>4.260\textsuperscript{2}</td>
<td>301.93\pm0.08</td>
<td>301.11\pm0.20</td>
<td>-0.28\pm0.08</td>
</tr>
<tr>
<td>4.310</td>
<td>44.90\pm0.03</td>
<td>45.29\pm0.08</td>
<td>+0.87\pm0.19</td>
</tr>
<tr>
<td>4.360</td>
<td>539.84\pm0.10</td>
<td>541.38\pm0.28</td>
<td>+0.29\pm0.06</td>
</tr>
<tr>
<td>4.390</td>
<td>55.18\pm0.04</td>
<td>55.27\pm0.09</td>
<td>+0.16\pm0.18</td>
</tr>
<tr>
<td>4.420\textsuperscript{1}</td>
<td>44.67\pm0.03</td>
<td>44.77\pm0.08</td>
<td>+0.22\pm0.20</td>
</tr>
<tr>
<td>4.420\textsuperscript{2}</td>
<td>1028.89\pm0.13</td>
<td>1029.63\pm0.37</td>
<td>+0.07\pm0.04</td>
</tr>
<tr>
<td>4.470</td>
<td>109.94\pm0.04</td>
<td>109.51\pm0.13</td>
<td>-0.39\pm0.13</td>
</tr>
<tr>
<td>4.530</td>
<td>109.98\pm0.04</td>
<td>109.47\pm0.13</td>
<td>-0.46\pm0.13</td>
</tr>
<tr>
<td>4.575</td>
<td>47.67\pm0.03</td>
<td>47.57\pm0.08</td>
<td>-0.21\pm0.18</td>
</tr>
<tr>
<td>4.600</td>
<td>566.93\pm0.11</td>
<td>563.45\pm0.28</td>
<td>-0.62\pm0.06</td>
</tr>
</tbody>
</table>

Table 2. Options for the BABAYAGA3.5 generator used to generate the simulated MC data samples.

| parameters       | value
|------------------|-----------------
| Ebeam            | 2.130 GeV or others
| MinThetaAngle    | 20º
| MaxThetaAngle    | 160º
| MinimumEnergy    | 0.04 GeV
| MaximumAcollinearity | 180º
| RunningAlpha     | 1
| FSR switch       | 1

4 Event selection and results

Signal candidates are required to have exactly two oppositely charged tracks. The tracks must originate from a cylindrical volume, centered around the interaction point, which is defined by a radius of 1 cm perpendicular to the beam axis and a length of ±10 cm along the beam axis. In addition, the charged tracks are required to be within \(|\cos\theta| < 0.8\), where \(\theta\) is the polar angle, measured by the MDC. Without applying further particle identification, the tracks are assigned as electrons and positrons depending on their charge. The deposited energies of electron and positron in EMC must be larger than \(\sqrt{s}/4.26 \times 1.55\) (GeV) to remove the di-muon background, where \(\sqrt{s}\) is the CM energy in GeV; the momenta of electron and positron are required to be larger than \(\sqrt{s}/4.26 \times 2\) (GeV/c), to suppress background events from lighter vector resonances produced in the ISR process, such as \(J/\psi, \psi(3686)\) and other resonances, decaying into \(e^+e^-\) pairs. For the data sample with a CM energy of 3.810 or 3.910 GeV, the effect of the remaining \(\psi(3686)\) events is studied by applying a 20% larger momentum requirement, and is found to be negligible. The requirements on the deposited energies and momenta are not optimized in detail, as the number of the signal events in such an analysis is large enough. All the variables mentioned above are determined in the initial \(e^+e^-\) CM frame. The ratio of the number of remaining background events to the number of signal events, estimated from the inclusive MC sample, is found to be less than \(2 \times 10^{-4}\), which is negligible. Thus all the selected events are taken as Bhabha events.

Figure 1 shows the comparisons between data and MC simulation for the kinematic variables of the leptons by taking data at the CM energy of 4.260 GeV as an example. Reasonable agreement is observed in the angular and momentum distributions. The striking difference between data and simulation found in the distributions of energies deposited by the leptons in the EMC emerges from imperfections in the simulation of the energy response of individual detector channels. At the CM energies analyzed in this work, a single shower in the calorimeter can be so energetic that the deposited energy per crystal exceeds the dynamic range of the analog-to-digital converter (ADC), causing individual ADC channels to saturate. In the analysis presented here, the very loose requirements on the energy deposits will not cause any bias, since they have been applied in regions of reasonable agreement between data and simulation. Relevant deviations between data and MC are considered as contributions to the systematic uncertainties.

The integrated luminosity is calculated with

\[
L = \frac{N_{\text{obs}}}{\sigma_{\text{Bhabha}} \times \epsilon},
\]

where \(N_{\text{obs}}\) is the number of observed Bhabha events, \(\sigma_{\text{Bhabha}}\) is the cross section of the Bhabha process, and \(\epsilon\) is the efficiency determined by analyzing the signal MC sample. The cross sections are calculated with the BABAYAGA3.5 generator using the parameters listed in Table 2 and decrease with increasing energies. The efficiencies are almost independent of the CM energy, as intended by the choice of relative conditions on lepton momenta and deposited energies. The luminosity results
Fig. 1. (color online) Comparison between data and MC simulation at the CM energy of 4.260 GeV. The top row is for positron and the bottom row for electron. From left to right, the plots show the distribution of deposited energy in EMC, the distribution of the cosine of the polar angle measured by the MDC, and the distribution of the track momentum from the MDC. Black points with error bars illustrate data and red points are MC simulation. Note that the y-axis is in logarithmic scale and the MC is normalized to data by the number of events for each sub-plot. When drawing the distribution of one variable, the requirements on the other variables are applied.

5 Systematic uncertainty

The following sources of systematic uncertainties are considered: the uncertainty of the tracking efficiency, the uncertainty related to the requirements on the kinematic variables, the statistical uncertainty of the MC sample, the uncertainty of the beam energy measurement, the uncertainty of the trigger efficiency, and the systematic uncertainty of the event generator.

To estimate the systematic uncertainty related to the tracking efficiency, the Bhabha event sample is selected using information from the EMC only, without using the tracking information in the MDC. The selection criteria are: at least two clusters in the EMC for each candidate, and the two most energetic clusters are assumed to originate from the $e^+e^-$ pair; the deposited energies of the two clusters are required to be larger than $\sqrt{s}/4.26 \times 1.8$ (GeV). At CM energies above 4.420 GeV, the requirement is changed to $\sqrt{s}/4.26 \times 1.55$ (GeV). This adjustment allows us to avoid additional systematic uncertainties which would be introduced by the deviation of data and simulation in the deposited energy in the EMC, as discussed in Section 4. The polar angle of each cluster is required to be within $|\cos \theta_{\text{EMC}}| < 0.8$, where $\theta_{\text{EMC}}$ is the polar angle measured by the EMC; to remove the background from the di-photon process, $\Delta \phi = |\phi_1 - \phi_2| - 180^\circ$ and $\phi_{1,2}$ are the azimuthal angles of the clusters in the EMC boosted to the CM frame. The efficiency that the selected Bhabha events pass through the track requirements applied in the nominal analysis is calculated for both data and MC sample, and the difference between them is taken as the systematic uncertainty connected to the tracking efficiency.

The systematic uncertainty in the polar angle acceptance is estimated by changing the requirement from $|\cos \theta| < 0.8$ to $|\cos \theta| < 0.7$. The difference between the resulting and nominal luminosity is taken as the associ-
ated systematic uncertainty. The systematic uncertainty caused by the requirement on the energy deposited in the EMC is estimated by changing the requirement from $\frac{\sqrt{s}}{4.26} \times 1.55$ (GeV) to $\frac{\sqrt{s}}{4.26} \times 1.71$ (GeV). The systematic uncertainty caused by the requirement on the momentum is estimated by changing the requirement from $\frac{\sqrt{s}}{4.26} \times 2$ (GeV/c) to $\frac{\sqrt{s}}{4.26} \times 2.06$ (GeV/c). The ranges are picked as these cause the largest deviations from the nominal luminosity result near the requirements applied.

The statistical uncertainty of the efficiency determined from MC simulations is 0.25%. The CM energy is determined using e${}^+${}e${}^−$→(γ)μ${}^+$μ${}^−$ events. The invariant mass of the di-muon system is calculated taking into account ISR and FSR effects. The difference between the CM energy listed in Table 1 and that measured with the di-muon process is about 2 MeV, and the corresponding systematic uncertainty is estimated by changing the CM energy by 2 MeV in the MC simulation. The trigger efficiency for the Bhabha process is 100% with an uncertainty of less than 0.1% [8]. The theoretical uncertainty caused by the requirement on the momentum requirement is 0.43, the energy requirement is 0.09, the polar angle requirement is 0.38, the beam energy is 0.42, the trigger efficiency is 0.10, and the generator is 0.50, giving a total uncertainty of 0.97.

The same systematic uncertainty estimation method is applied to all the sub-samples. The largest relative uncertainty among them is taken as the associated uncertainty for all the sub-samples. The systematic uncertainties considered in this work are summarized in Table 3.

By assuming the sources of the systematic uncertainties to be uncorrelated, the total uncertainty is calculated as 0.97% by adding the contributions in quadrature.

6 Cross check

To verify the result, a cross check with di-gamma events is performed. The event selection criteria are the same as those used in estimating the systematic uncertainty caused by the tracking efficiency, except for the requirement on Δφ. In order to reduce the Bhabha background, the Δφ is required to be in the range of $[-0.8^\circ, 0.8^\circ]$, since photons are not deflected in the magnetic field.

The luminosity results of this cross check ($L_{ck}$) are shown in Table, together with the relative differences to the nominal ones. Both results have good consistency for all individual measurements, indicating the robustness of the result.

7 Summary

The integrated luminosity of the data samples taken at BESIII for studying the charmonium-like states and higher excited charmonium states is measured to an accuracy of 1% with Bhabha events. The total uncertainty is dominated by the systematic uncertainty. A cross check with di-gamma events is performed and the results are consistent with each other. The result presented here is essential for future measurements of cross sections with these data, and has already been used in the discovery of charged charmonium-like states [9–12].

The BESIII collaboration would like to thank the staff of BEPCII and the IHEP computing center for their dedicated support.

References


1) M Ablikim et al. (BESIII collaboration). Measurement of the center-of-mass energy of the data taken by BESIII at center-of-mass energies between 3.81 GeV and 4.60 GeV. Paper in preparation.