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# K2-280 b - a low density warm sub-Saturn around a mildly evolved star

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## ABSTRACT

We present an independent discovery and detailed characterisation of K2-280 b, a transiting low density warm sub-Saturn in a 19.9-day moderately eccentric orbit ( $e = 0.35^{+0.05}_{-0.04}$ ) from *K2* campaign 7. A joint analysis of high precision HARPS, HARPS-N, and FIES radial velocity measurements and *K2* photometric data indicates that K2-280 b has a radius of  $R_b = 7.50 \pm 0.44 R_\oplus$  and a mass of  $M_b = 37.1 \pm 5.6 M_\oplus$ , yielding a mean density of  $\rho_b = 0.48^{+0.13}_{-0.10} \text{ g cm}^{-3}$ . The host star is a mildly evolved G7 star with an effective temperature of  $T_{\text{eff}} = 5500 \pm 100 \text{ K}$ , a surface gravity of  $\log g_\star = 4.21 \pm 0.05$  (cgs), and an iron abundance of  $[\text{Fe}/\text{H}] = 0.33 \pm 0.08$  dex, and with an inferred mass of  $M_\star = 1.03 \pm 0.03 M_\odot$  and a radius of  $R_\star = 1.28 \pm 0.07 R_\odot$ . We discuss the importance of K2-280 b for testing formation scenarios of sub-Saturn planets and the current sample of this intriguing group of planets that are absent in the Solar System.

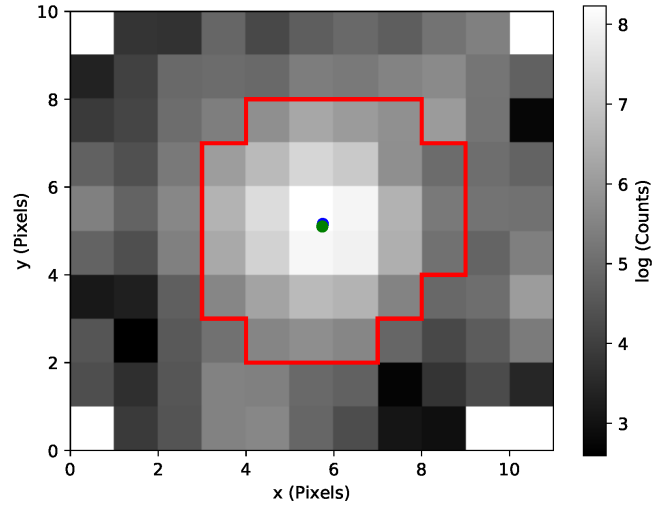
**Key words:** planets and satellites: detection – stars: individual (EPIC 216494238, K2-280) – techniques: photometric – techniques: radial velocities – techniques: spectroscopic

## 1 INTRODUCTION

The main advantage of the extended NASA’s *Kepler* mission (Borucki et al. 2010), known as *K2* (Howell et al. 2014), was a much larger number of bright stars in its fields of view located along the ecliptic. A significant number of planets transiting bright stars have been discovered in all *K2* campaigns (e.g., Montet et al. 2015; Crossfield et al. 2016; Vanderburg et al. 2016; Dressing et al. 2017; Petigura et al. 2018; Mayo et al. 2018; Yu et al. 2018a; Crossfield et al. 2018; Livingston et al. 2018). Some of these planets were only validated, but many were characterized by means of high precision radial velocity (RV) measurements that enabled mass determination with precision better than 20% (e.g., Vanderburg et al. 2015; Gandolfi et al. 2017; Christiansen et al. 2017; Prieto-Arranz et al. 2018; Rodriguez et al. 2018; Barragán et al. 2018b; Malavolta et al. 2018). However much higher precision is needed to distinguish between various possible planetary compositions (see e.g. Dorn et al. 2015, and references therein). Mass determination with such precision for small planets ( $R_p = 1\text{--}4 R_\oplus$ ) were possible only for short period ( $P_{\text{orb}} \lesssim 10$  days) sub-Neptunes and super-Earths, which induces RV semi-amplitudes on the parent stars of a few  $\text{m s}^{-1}$ , and for stars hosting ultra-short period planets, for which the Doppler reflex motion is enhanced by the extremely short orbital period ( $P_{\text{orb}} < 1$  day; Winn et al. 2018, and references therein). Most of the small *K2* planets with precise mass determination orbit bright stars, i.e., stars brighter than  $V = 11.5$ , which is the current limit for  $\sim 1 \text{ m s}^{-1}$  precision with spectrographs mounted at 3–4-m class telescopes (Pepe et al. 2013). The constraints for precise determination of planetary masses are naturally more relaxed for higher-mass planets, enabling us to study super-Neptune/sub-Saturn planets ( $R_p = 4\text{--}8 R_\oplus$ ) with longer orbital periods around fainter stars.

Sub-Saturns form a very intriguing group of planets that have no counterpart in the Solar System. Their main characteristic is a significant contribution of both heavy metal cores and low density gaseous envelopes to the total planet mass (Petigura et al. 2016). They are thus important laboratories to study envelope accretion. As shown by Petigura et al. (2017), the population of sub-Saturns has a very uniform distribution of the planetary mass between  $\sim 6$  and  $60 M_\oplus$ . Similar to gas-giant planets, they are also found to orbit mainly metal-rich stars. Finally, the most massive sub-Saturns are often the only detected planet in the system and orbit their parent stars on eccentric orbits, which suggests that dynamical instability might have played an important role in their formation.

Here we present an independent discovery and characterisation of a low mass, sub-Saturn planet on a 20-d orbit around a relatively faint ( $V = 12.5$ ), metal rich ( $[\text{Fe}/\text{H}] = 0.33 \pm 0.08$  dex), slightly evolved *K2* star that was proposed as a planet candidate by Petigura et al. (2018) and Mayo et al. (2018), and statistically validated as a planet by Livingston et al. (2018). This kind of planets was usually avoided by RV follow-up of *K2* candidates because of the faintness of their host stars. Besides, slightly evolved stars are typically avoided in RV follow-up projects, because of their higher expected stellar jitter (see e.g. Hekker et al. 2006, 2008; Tayar et al. 2019, and references therein). Both these effects may bias the statistical analysis of warm-giant



**Figure 1.** *K2* image of K2-280. Red lines shows the aperture defined by the amount of light of each pixel and level of background light. The electron count is indicated by the intensity of shading (light grey for high and dark grey for low count). The green circle indicates the current position of the target in the EPIC catalog, and the blue circle is the centre of the flux distribution. The scale of the image is the *Kepler* pixel scale of 3.98 arcsec/pix. The *K2* image is not a rectangle, but it is irregularly shaped. The white pixels in the corners contain no data.

planets. K2-280 b joins a sample of 30 sub-Saturns with mean densities determined with precision better than 50% discovered mainly by *Kepler* and *K2* (see Petigura et al. 2017, and references therein for first 23 planets).

This work was done as a part of the KESPRINT collaboration<sup>1</sup>, which aims to confirm and characterize *K2* and *TESS* planets. In Section 2 we describe the observations of K2-280, specifically the *K2* photometry, the NOT/FIES, ESO/HARPS, and TNG/HARPS-N high-resolution spectroscopy follow-up, and the high-contrast imaging. In Section 3 and 4 we present the properties of the host star K2-280 and the global analysis of photometric and Doppler data, respectively. In Section 5 we finally summarise and discuss the characteristics of K2-280 b in the context of the properties of the known population of sub-Saturn planets with mean densities determined with precision better than 50%.

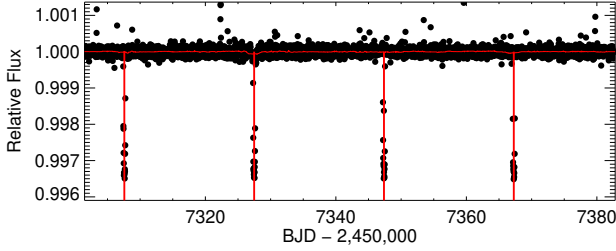
## 2 OBSERVATIONS AND DATA REDUCTIONS

### 2.1 *K2* Photometry

K2-280 was one of 13 469 long cadence targets observed from October 4<sup>th</sup> to December 26<sup>th</sup> 2015 (UT) during *K2* campaign 7. It was proposed as a target by GO programmes 7030 (PI Howard), and 7085 (PI Burke). We downloaded K2-280 images from the MAST archive<sup>2</sup> and used them to produce a de-trended *K2* light curve as described in detail in Dai et al. (2017). Figure 1 shows the pixel mask used to perform simple aperture photometry. We used the box

<sup>1</sup> <https://www.iac.es/proyecto/kesprint/>.

<sup>2</sup> [https://archive.stsci.edu/k2/data\\_search/search.php](https://archive.stsci.edu/k2/data_search/search.php).



**Figure 2.** Detrended *K2* light curve of K2-280. The 4 transits observed by *K2* are marked by vertical solid red lines. The horizontal red line is the local median flux level with a window of 0.5-day.

fitting least-square (BLS) routine (Kovács et al. 2002; Jenkins et al. 2010), improved by implementation of the optimal frequency sampling described in Ofir (2014) to search for transiting planet candidates in all Field 7 targets light curves. We detected transits of K2-280 b with a signal-to-noise ratio (SNR) of 24.5, depth of  $\sim 3.5 \times 10^{-3}$ , period of  $P = 19.89518 \pm 0.00028$  days, and a mid-time of the first transit  $T_0 = 2457307.58101 \pm 0.00059$  days in Barycentric Julian Date in the Barycentric Dynamical Time (BJD<sub>TDB</sub>; see, e.g., Eastman et al. 2010). The de-trended light curve of K2-280 with the correction for baseline flux variations and centroid motions is presented in Figure 2 with the 4 transits observed by *K2* highlighted with red lines. We removed the baseline flux variation by fitting a spline function with a width of 3 days. In Table 1 we report the main identifiers of K2-280, along with its equatorial coordinates, space motion, distance, and optical and near-infrared magnitudes.

## 2.2 High-dispersion spectroscopy

High-dispersion spectroscopic observations of K2-280 were obtained between April 30<sup>th</sup> 2016 (UT) and May 7<sup>th</sup> 2019 (UT) using ESO/HARPS, TNG/HARPS-N, and NOT/FIES spectrographs. We collected a total of 18 HARPS, 14 HARPS-N and 6 FIES spectra. The details of these observations are given in the subsections below. Table 2 gives the time stamps of the spectra in BJD<sub>TDB</sub>, the RVs along with their  $1\sigma$  error bars, as well as the bisector inverse slope (BIS) and full-width at half maximum (FWHM) of the cross-correlation function (CCF).

### 2.2.1 ESO/HARPS

We started the RV follow-up of K2-280 using the High Accuracy Radial velocity Planet Searcher (HARPS) spectrograph (Mayor et al. 2003, R $\approx$ 115 000) mounted at the ESO 3.57-m telescope of La Silla Observatory in Chile. We acquired 18 spectra between April 30<sup>th</sup> 2016 (UT) and April 27<sup>th</sup> 2018 (UT) under the observing programmes 097.C-0571(B), 097.C-0948(A), 098.C-0860(A), 099.C-0491(A), 0101.C-0407(A), and 60.A-9700(G), setting the exposure times to 1200–3600 seconds. The dedicated on-line HARPS Data Reduction Software (DRS) was used to reduce the spectra, and extract the Doppler measurements and spectral activity indicators. The SNR per pixel at 5500 Å is in

**Table 1.** Properties of K2-280.

Parameter	Value	Source
<i>Coordinates and Main Identifiers</i>		
RA 2000.0 (hours)	19:26:22.881	<i>Gaia</i> DR2
Dec 2000.0 (deg)	-22:14:51.552	<i>Gaia</i> DR2
<i>Gaia</i> DR2 Identifier	6772454416893148928	<i>Gaia</i> DR2
2MASS Identifier	19262288-2214514	2MASS PSC
UCAC Identifier	339-184113	UCAC4
EPIC Identifier	216494238	EPIC
TIC Identifier	119605900	TIC
<i>Optical and Near-Infrared Magnitudes</i>		
$K_p$ (mag)	12.302	<i>K2</i> EPIC
$B$ (mag)	$13.269 \pm 0.010$	UCAC4
$V$ (mag)	$12.536 \pm 0.040$	UCAC4
$R$ (mag)	$12.41 \pm 0.07$	UCAC4
$G$ (mag)	$12.3604 \pm 0.0002$	<i>Gaia</i> DR2
$g$ (mag)	$12.850 \pm 0.020$	UCAC4
$r$ (mag)	$12.320 \pm 0.020$	UCAC4
$i$ (mag)	$12.067 \pm 0.040$	UCAC4
$J$ (mag)	$11.141 \pm 0.021$	2MASS
$H$ (mag)	$10.854 \pm 0.024$	2MASS
$K$ (mag)	$10.765 \pm 0.019$	2MASS
<i>Space Motion and Distance</i>		
PM <sub>RA</sub> (mas yr <sup>-1</sup> )	$4.44 \pm 0.08$	<i>Gaia</i> DR2
PM <sub>DEC</sub> (mas yr <sup>-1</sup> )	$-12.50 \pm 0.07$	<i>Gaia</i> DR2
RV <sub><math>\gamma</math>,HARPS</sub> (km s <sup>-1</sup> )	$-1.1934^{+0.0011}_{-0.0011}$	This work
RV <sub><math>\gamma</math>,HARPS-N</sub> (km s <sup>-1</sup> )	$-1.1907^{+0.0011}_{-0.0011}$	This work
RV <sub><math>\gamma</math>,FIES</sub> (km s <sup>-1</sup> )	$-1.2349^{+0.0038}_{-0.0039}$	This work
$\pi_*$ (mas)	$2.526 \pm 0.111$	<i>Gaia</i> DR2
$d_*$ (pc)	$391.5^{+7.5}_{-7.2}$	Bailer-Jones et al. (2018)
$U$ (km s <sup>-1</sup> )	$-7.78 \pm 0.06$	This work
$V$ (km s <sup>-1</sup> )	$-5.98 \pm 0.84$	This work
$W$ (km s <sup>-1</sup> )	$-9.18 \pm 0.71$	This work
<i>Photospheric Parameters</i>		
$T_{\text{eff}}$ (K)	$5500 \pm 100$	This work
$\log g_*$ <sup>(a)</sup> (dex)	$4.00 \pm 0.10$	This work
$\log g_*$ <sup>(b)</sup> (dex)	$4.21 \pm 0.05$	This work
[Fe/H] (dex)	$0.33 \pm 0.08$	This work
<i>Derived Physical Parameters</i>		
$M_*$ (M <sub>⊙</sub> )	$1.03 \pm 0.03$	This work
$R_*$ (R <sub>⊙</sub> )	$1.28 \pm 0.07$	This work
$\rho_*$ (g cm <sup>-3</sup> )	$0.8^{+0.16}_{-0.13}$	This work
Age (Gyr)	$8.96 \pm 1.70$	This work
<i>Stellar Rotation</i>		
$v_{\text{rot}} \sin i_*$ (km s <sup>-1</sup> )	$3.0 \pm 1.0$	This work

Notes – <sup>(a)</sup> From spectroscopy. <sup>(b)</sup> From stellar mass and radius.

the range 22–46. Radial velocities were measured by cross-correlating the extracted spectra with a G2 numerical mask (Baranne et al. 1996). The uncertainties of the measured RVs are in the range 2.1–8.1 ms<sup>-1</sup> with a mean value of 4.2 ms<sup>-1</sup>.

### 2.2.2 TNG/HARPS-N

Between July 16<sup>th</sup> 2016 (UT) and May 7<sup>th</sup> 2019 (UT) we collected 14 spectra with the HARPS-N spectrograph (Cosentino et al. 2012, R $\approx$ 115 000) mounted at the 3.58-m Telescopio Nazionale Galileo (TNG) of Roque de los Muchachos Observatory in La Palma, Spain, under the observing programmes A33TAC\_15, OPT17A\_64, CAT17A\_91, and CAT19A\_97. The exposure time was set to 1200–3600, based on weather conditions and scheduling constraints, leading to a SNR per pixel of 15–47 at 5500 Å. The spectra were extracted using the off-line version of the HARPS-N DRS pipeline. Doppler measurements and spectral activity indicators were measured using an on-line version of the DRS, the YABI tool<sup>3</sup>, by cross-correlating the extracted spectra with a G2 mask (Baranne et al. 1996). The uncertainty of the measured RVs are in the range 2.0–8.6 ms<sup>-1</sup>, with a mean value of 4.5 ms<sup>-1</sup>.

### 2.2.3 NOT/FIES

We acquired 6 additional spectra using the Fibre-fed Échelle Spectrograph (FIES; Frandsen & Lindberg 1999; Telting et al. 2014) mounted at the 2.56-m Nordic Optical Telescope (NOT) of Roque de los Muchachos Observatory (La Palma, Spain). The observations were carried out between June 26<sup>th</sup> and September 6<sup>th</sup> 2016 (UT), as part of the OPTICON observing programme 53-109. We used the FIES high-resolution mode, which provides a resolving power of  $R=67\,000$  in the spectral range 3700–7300 Å. Following the observing strategy described in Buchhave et al. (2010) and Gandolfi et al. (2015), we traced the RV drift of the instrument by acquiring long-exposed ThAr spectra ( $T_{\text{exp}} \approx 35$  sec) immediately before and after each science exposure. The exposure time was set to 2700–3600 seconds, according to the sky conditions and scheduling constraints. The data reduction follows standard IRAF and IDL routines, which includes bias subtraction, flat fielding, order tracing and extraction, and wavelength calibration. Radial velocity measurements were computed via multi-order cross-correlations with the RV standard star HD 50692 (Udry et al. 1999), observed with the same instrument set-up as K2-280. The SNR per pixel at 5500 Å of the extracted spectra is in the range 15–35. The uncertainties are in the range 6.8–13.6 ms<sup>-1</sup> with a mean value of 9.7 ms<sup>-1</sup>.

## 2.3 High contrast imaging

To search for nearby stars and estimate a potential contamination factor from such sources we used a high contrast image of K2-280 publicly available on the ExoFOP-K2 website<sup>4</sup>. The image was acquired on June 19<sup>th</sup> 2016 (UT) using the Frederick C. Gillett Gemini North telescope and its adaptive optics (AO) system facility, ALTAIR with a natural guide star along with a Near InfraRed Imager and spectrograph (NIRI) (Hodapp et al. 2003) using the Brackett Gamma (Br $\gamma$ ) filter (Gemini-North ID

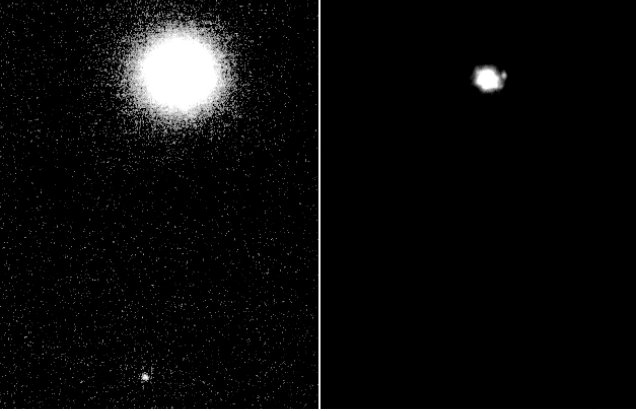
**Table 2.** HARPS, HARPS-N, and FIES radial velocities (RVs), bisector inverse slope (BIS), and full-width at half maximum (FWHM) of the cross-correlation function.

BJD <sub>TDB</sub>	RV	$\sigma_{\text{RV}}$	BIS	$\sigma_{\text{BIS}}$	FWHM
-2 450 000	(m s <sup>-1</sup> )	(m s <sup>-1</sup> )	(m s <sup>-1</sup> )	(m s <sup>-1</sup> )	(km s <sup>-1</sup> )
FIES					
7566.606106	-1236.447	10.145	-25.604	8.240	12.240
7570.575852	-1248.492	8.236	-19.745	5.424	12.237
7579.594362	-1221.025	13.568	-15.325	10.160	12.257
7583.604507	-1227.357	11.149	-1.138	9.710	12.232
7600.515019	-1233.204	8.028	-36.207	7.614	12.283
7638.387601	-1225.100	6.816	-23.315	4.940	12.220
HARPS-N					
7585.595242	-1192.158	6.794	-15.933	9.608	7.354
7603.514891	-1191.265	5.126	-13.175	7.249	7.363
7611.572851	-1201.759	8.653	-15.272	12.238	7.371
7892.699273	-1195.302	3.475	-5.393	4.915	7.348
7921.702566	-1191.619	3.014	-7.642	4.263	7.359
7958.534636	-1194.254	7.553	-39.517	10.682	7.360
7965.540153	-1197.765	2.357	-14.982	3.334	7.360
8013.363503	-1186.028	1.991	-18.747	2.816	7.366
8013.404784	-1185.686	2.117	-30.544	2.994	7.362
8014.362082	-1185.942	2.722	-28.539	3.850	7.352
8014.402379	-1189.116	2.756	-24.429	3.898	7.365
8015.359376	-1178.088	5.620	-11.056	7.948	7.349
8015.402150	-1187.664	5.605	-13.530	7.927	7.368
8610.724187	-1193.898	4.651	-17.515	6.577	7.359
HARPS					
7508.854151	-1195.313	3.185	-5.317	4.504	7.447
7511.892058	-1194.114	4.462	-2.236	6.310	7.409
7609.625616	-1195.440	2.732	19.011	3.865	7.415
7619.566531	-1184.749	4.551	16.474	6.437	7.413
7637.587335	-1181.617	6.366	-13.161	9.004	7.400
7638.596015	-1172.262	7.600	-21.324	10.748	7.447
7639.571271	-1181.987	6.990	-34.526	9.885	7.389
7645.538147	-1194.348	4.671	11.465	6.607	7.433
7682.520265	-1194.190	5.191	17.641	7.342	7.420
7984.617196	-1190.758	2.370	-5.857	3.352	7.422
7986.590931	-1203.107	3.223	-1.852	4.559	7.421
7987.587443	-1197.553	2.285	13.255	3.232	7.415
7990.553698	-1195.977	2.729	24.528	3.859	7.418
7990.593359	-1207.601	2.817	-20.535	3.983	7.396
7992.576732	-1189.811	2.093	16.287	2.950	7.419
7992.609762	-1187.701	2.759	-9.069	3.902	7.408
8003.594697	-1191.639	3.620	-0.841	5.110	7.401
8235.876739	-1189.777	8.100	-22.636	11.455	7.409

G0218) centered at 2.17  $\mu\text{m}$ , under Gemini Science Program GN-2016A-LP-5. Two faint stars are visible on the Gemini-North/NIRI+ALTAIR AO image of K2-280 (Figure 3): a very close-in companion at  $\sim 0.4''$  west-northwest (W–NW), and a distant source at  $\sim 6.6''$  south-southeast (S–SE) of K2-280. We carefully analyzed the Gemini-North/NIRI+ALTAIR AO image of K2-280. Table 3 reports separations, position angles, the magnitude difference  $\Delta m_{\text{Br}\gamma}$ , and the  $\Delta F_{\text{Br}\gamma}$  flux-ratio of these two objects relative to K2-280. Their brightness ratio at 2.17  $\mu\text{m}$  is comparable to the observed K2 transit depth (3500 ppm), which requires their consideration as sources of false positives (see Section 3.5).

<sup>3</sup> Available at <http://ia2-harps.oats.inaf.it:8000>.

<sup>4</sup> See [https://exofop.ipac.caltech.edu/k2/edit\\_target.php?id=216494238](https://exofop.ipac.caltech.edu/k2/edit_target.php?id=216494238).



**Figure 3.** AO image of the surroundings of K2-280 obtained with the Gemini-North/NIRI+ALTAIR instrument. Both panels show the same image, with a FOV of  $8.7''$  in the N-S and  $6.7''$  in the E-W direction (north to the top and east to the left), but with different brightness scales. The left panel shows the star at  $6.6''$  in the S-SE direction and the right one the close neighbor at  $0.4''$  W-NW of K2-280.

### 3 PROPERTIES OF THE HOST STAR

#### 3.1 *Gaia* measurements

K2-280 is among a small sub-sample of ESA’s *Gaia* mission (Gaia Collaboration et al. 2016) targets for which the *Gaia* DR2 (Gaia Collaboration et al. 2018)<sup>5</sup> – a first *Gaia*-only catalogue – provides not only astrometric measurements, but also astrophysical parameters (radii, luminosities, extinctions, and reddening) and median radial velocities. *Gaia* DR2 astrometric parameters of K2-280 are included in Table 1. *Gaia* DR2 values of stellar radius and median RV of K2-280 agree with the values determined in the subsections below.

#### 3.2 Photospheric parameters and stellar rotation velocity measurements using SME

We followed the procedure described in Fridlund et al. (2017) and Persson et al. (2018) and analysed the co-added spectra from HARPS, HARPS-N, and FIES with the spectral analysis package Spectroscopy Made Easy (SME; Valenti & Piskunov 1996; Valenti & Fischer 2005; Piskunov & Valenti 2017) to derive the effective temperature  $T_{\text{eff}}$ , surface gravity  $\log g_{\star}$ , iron abundance  $[\text{Fe}/\text{H}]$ , and projected rotational velocity  $v_{\text{rot}} \sin i_{\star}$ . Spectroscopy Made Easy uses grids of atmosphere models to calculate synthetic stellar spectra, which are fitted to the observed spectra using a  $\chi^2$ -minimising procedure. We use the line wings of  $\text{H}\alpha$ , which is rather insensitive to  $\log g_{\star}$  for this spectral type, to model  $T_{\text{eff}}$  (with a fixed  $\log g_{\star}$ ), and the line wings of the Ca I triplet to model  $\log g_{\star}$  (with a fixed  $T_{\text{eff}}$ ). We used the latest version of the software (5.2.2) and line lists from the Vienna atomic line database<sup>6</sup>. The model spectra were taken from ATLAS12 (Kurucz 2013). The calibration equations for Sun-like stars from Bruntt et al. (2010) and Doyle et al. (2014) were

<sup>5</sup> Released on April 25<sup>th</sup>, 2018.

<sup>6</sup> <http://vald.astro.uu.se>.

**Table 3.** Relative properties of the two nearby stars to K2-280 detected with the Gemini-North/NIRI+ALTAIR.

Parameter	W-NW close-in star	S-SE distant star
Separation ( $''$ )	$0.38 \pm 0.011$	$6.598 \pm 0.011$
Position Angle (deg)	$286.2 \pm 1.5$	$173.3 \pm 1.5$
$\Delta m_{B_{\text{r}\gamma}}$ (mag)	$4.72 \pm 0.15$	$6.65 \pm 0.15$
$\Delta F_{B_{\text{r}\gamma}}$ relative flux	$(1.3 \pm 0.2) \times 10^{-2}$	$(2.2 \pm 0.4) \times 10^{-3}$

adopted to fix the micro- and macroturbulent velocities,  $v_{\text{mic}}$  and  $v_{\text{mac}}$  to 0.5 and 1.0  $\text{km s}^{-1}$ , respectively. The spectroscopic parameters derived from the HARPS, HARPS-N, and FIES co-added spectra agree well within their nominal error bars. The final adopted values are  $T_{\text{eff}} = 5500 \pm 100 \text{ K}$ ,  $\log g_{\star} = 4.00 \pm 0.10$  (cgs), and  $[\text{Fe}/\text{H}] = 0.33 \pm 0.08$  dex (Table 1). They are defined as the weighted mean of the individual parameters derived from the HARPS, HARPS-N, and FIES co-added spectra.

#### 3.3 Photospheric parameters and radius measurements using SpecMatch-emp

As a sanity check, we also analysed the co-added HARPS and HARPS-N spectra using the SpecMatch-emp software package (Yee et al. 2017). SpecMatch-emp estimates the stellar effective temperature  $T_{\text{eff}}$ , radius  $R_{\star}$ , and iron abundance  $[\text{Fe}/\text{H}]$  by fitting the spectral region between 5000 and 5900  $\text{\AA}$  to hundreds of library spectra gathered by the California Planet Search programme. Following the procedure described in Hirano et al. (2018), we reformatted the co-added HARPS and HARPS-N spectra so that they can be read by SpecMatch-emp. We found  $T_{\text{eff}} = 5597 \pm 110 \text{ K}$  and  $[\text{Fe}/\text{H}] = 0.33 \pm 0.08$  dex, which agree with the effective temperature and iron abundance determined with SME (Table 1) within  $1\sigma$ . We found also that K2-280 is a slightly evolved star with a stellar radius of  $R_{\star} = 1.33 \pm 0.21 R_{\oplus}$ . We finally obtained a first estimate of the stellar mass ( $M_{\star} = 1.16 \pm 0.08 M_{\oplus}$ ) via Monte Carlo simulations using the empirical equations by Torres et al. (2010) alongside  $T_{\text{eff}}$ ,  $[\text{Fe}/\text{H}]$ , and  $R_{\star}$ .

#### 3.4 Physical parameters

We refined the fundamental parameters of K2-280 utilising the web interface<sup>7</sup> PARAM 1.3 along with PARSEC isochrones (Bressan et al. 2012). Following the method described in Gandolfi et al. (2008), we found that the interstellar extinction along the line of sight to the star is  $A_{\text{v}} = 0.10 \pm 0.05$ . Using the effective temperature and iron abundance derived in Section 3.2, alongside the extinction-corrected visual magnitude and the *Gaia* parallax<sup>8</sup> (Table 1), we determined a mass of  $M_{\star} = 1.03 \pm 0.03 M_{\odot}$  and a radius of

<sup>7</sup> Available at [http://stev.oapd.inaf.it/cgi-bin/param\\_1.3](http://stev.oapd.inaf.it/cgi-bin/param_1.3).

<sup>8</sup> We accounted for *Gaia* systematic uncertainties adding quadratically 0.1 mas to the nominal uncertainty of parallax (Luri et al. 2018).

$R_{\star} = 1.28 \pm 0.07 R_{\odot}$ , which agree with the values derived in Section 3.3. Stellar mass and radius implies a surface gravity of  $\log g_{\star} = 4.21 \pm 0.05$  (cgs), which is higher than our spectroscopic value of  $4.0 \pm 0.1$  (cgs), but within its  $2\sigma$  error bars. The age of the star was constrained to be  $8.9 \pm 1.7$  Gyr, further confirming the evolved status of K2-280. The values of stellar radius and mass agree within  $3\text{-}\sigma$  with the ones determined by Petigura et al. (2018) ( $R_{\star} = 1.45^{+0.20}_{-0.18} R_{\odot}$ ,  $M_{\star} = 1.17^{+0.10}_{-0.08} M_{\odot}$ ), Mayo et al. (2018) ( $R_{\star} = 1.064^{+0.069}_{-0.047} R_{\odot}$ ,  $M_{\star} = 1.101^{+0.025}_{-0.028} M_{\odot}$ ), and Livingston et al. (2018) ( $R_{\star} = 1.28 \pm 0.03 R_{\odot}$ ,  $M_{\star} = 1.11 \pm 0.04 M_{\odot}$ ). We stress that the parameter estimates determined in the three works listed above are based on spectra with relatively low SNR, in contrast to our co-added, high SNR, HARPS, HARPS-N and FIES spectra. Petigura et al. (2018) and Livingston et al. (2018) used the same spectra collected with the High Resolution Echelle Spectrometer (HIRES Vogt et al. 1994) mounted on 10-m Keck I telescope, with typical SNR = 45 for stars with  $V < 13.0$ . Mayo et al. (2018) used spectra collected with Tillinghast Reflector Echelle Spectrograph (TRES) mounted on the 1.5-m Tillinghast telescope at the Whipple Observatory on Mt. Hopkins in Arizona with even lower SNR.

We also calculated the UVW space velocities of K2-280 using the IDL code `gal_uvw`<sup>9</sup> (based upon Johnson & Soderblom 1987), using the *Gaia* DR2 proper motions and parallax, and the average of the HARPS and HARPS-N systemic velocities  $\gamma$  (Table 1). Our calculated values of UVW are listed in Table 1; we quote values in the local standard of rest using the solar motion of Coşkunoğlu et al. (2011). We then used the methodology of Reddy et al. (2006) to determine the Galactic population membership of K2-280. We found that K2-280 has a  $> 99\%$  probability of belonging to the Galactic thin disk, and less than 1% of belonging to either the thick disk or the halo. This is consistent with K2-280's high metallicity of  $[\text{Fe}/\text{H}] = 0.33 \pm 0.08$  dex.

The final adopted stellar parameters are listed in Table 1. The effective temperature and surface gravity translate into a G7V spectral type (Gray & Corbally 2009).

### 3.5 Faint AO companions

From the two faint companions to K2-280 identified in the Gemini-North/NIRI+ALTAIR AO image (Section 2.3), the one located  $6.6''$  S-SE of K2-280 was identified in the *Gaia* DR2 as the source 6772454206445987712. Based on its very small proper motion ( $\text{PM}_{\text{RA}} = 0.29 \pm 0.52 \text{ mas yr}^{-1}$  and  $\text{PM}_{\text{DEC}} = -0.92 \pm 0.45 \text{ mas yr}^{-1}$ ) and distance found by Bailer-Jones et al. (2018) ( $d = 5.103^{+3.435}_{-2.094}$  kpc), we concluded that it is a background star. Using the *Gaia* G-band magnitude ( $G = 18.765 \pm 0.010$ ), we derived a G-band brightness ratio relative to K2-280 of  $0.0027 \pm 0.0001$ . Considering the close similarity between the *Gaia* G-band and the *Kepler* passband, this companion is too faint to be the source of the transit signal detected in the *K2* data.

For the close-in W-NW companion we cannot determine whether it is physically bound or unbound to K2-280. Yet, its very small angular separation of  $0.4''$  supports the binary scenario for K2-280. Based on the Besançon Galactic

population model<sup>10</sup> (Robin et al. 2003) and following the procedure described in Hjorth et al. (2019) we calculated the probability of a chance alignment to be 0.04%. Assuming that the W-NW companion is physically bound to K2-280, we can then obtain further information about it.

The central wavelength of  $2.19 \mu\text{m}$  of the Bry filter is nearly identical to that of the near-infrared *K* band. Therefore we used the apparent *K* magnitude of K2-280 from Table 1 ( $m_K = 10.765 \pm 0.019$ ) and the magnitude difference from Table 3 to calculate absolute *K* magnitudes of both stars. They are equal to  $M_K = 2.778 \pm 0.090$  for K2-280 and  $M_K = 7.50 \pm 0.22$  for the nearby companion. Making use of the Dartmouth isochrone table (Dotter et al. 2008) for metallicity  $[\text{Fe}/\text{H}] = 0.36$  and ages between 9 and 11 Gyr, we estimated that the nearby companion is a M3.5–M4 red-dwarf with a mass between  $0.21$  and  $0.28 M_{\odot}$ . Using its angular separation from Table 3 and the DR2 parallax of K2-280 we calculated a lateral separation from K2-280 of  $150.4^{+8.2}_{-7.7}$  au. We note that current models of planetary formations in wide binary stellar systems predict a shortage of giant planets in binaries with separations of  $\leq 100$  au (e.g., Nelson 2000; Mayer et al. 2005; Thébault et al. 2006); the nearby companion should therefore not have affected the formation of the K2-280 planetary system.

Based on the Dartmouth isochrone table for metallicity  $[\text{Fe}/\text{H}] = 0.36$  and ages between 9 and 11 Gyr, we also estimated the nearby star's absolute *Kepler* magnitude ( $M_{Kp}$ ) as 10.75–11.5 mag, and its apparent *Kepler* magnitude as 18.75–19.5 mag. That is, its *Kepler* brightness is  $0.0019 \pm 50\%$  of K2-280's brightness. However, a false-positive scenario with an equal mass eclipsing binary (eclipse depth equal to 50%) and a transit signal with a depth of  $3.5 \times 10^{-3}$  can only be caused by a binary that is brighter than 0.007 times the host's brightness. Therefore, assuming that the nearby W-NW star is physically bound with K2-280, we may exclude it as a source of a false positive.

## 4 GLOBAL ANALYSIS

We used the code `pyaneti` (Barragán et al. 2019) to perform the joint analysis of the RV and *K2* transit data. The code uses the limb-darkened quadratic model by Mandel & Agol (2002) to fit the transit light curves and a Keplerian model for the RV measurements. We integrated the light curve model over 10 steps to simulate the *Kepler* long-cadence integration (Kipping 2010). Fitted parameters, parametrizations and likelihood are similar to previous analysis performed with `pyaneti` (e.g. Barragán et al. 2016, 2018a).

The photometric data includes  $\sim 17$  hours (i.e, twice the transit duration) of data-points centered around each of the 4 transits observed by *K2*. We de-trended the photometric chunks using the program `exotrending`. (Barragán & Gandolfi 2017), fitting a second order polynomial to the out-of-transit data. The Doppler measurements include the 6 FIES, 14 HARPS-N, and 18 HARPS RVs presented in Section 2.2.

We adopted uniform priors for all the parameters; details are given in Table 5. We started 500 Markov chains randomly distributed inside the prior ranges. Once all chains

<sup>9</sup> Available at [https://idlastro.gsfc.nasa.gov/ftp/pro/astro/gal\\_uvw.pro](https://idlastro.gsfc.nasa.gov/ftp/pro/astro/gal_uvw.pro).

<sup>10</sup> Available at <http://modele2016.obs-besancon.fr>.

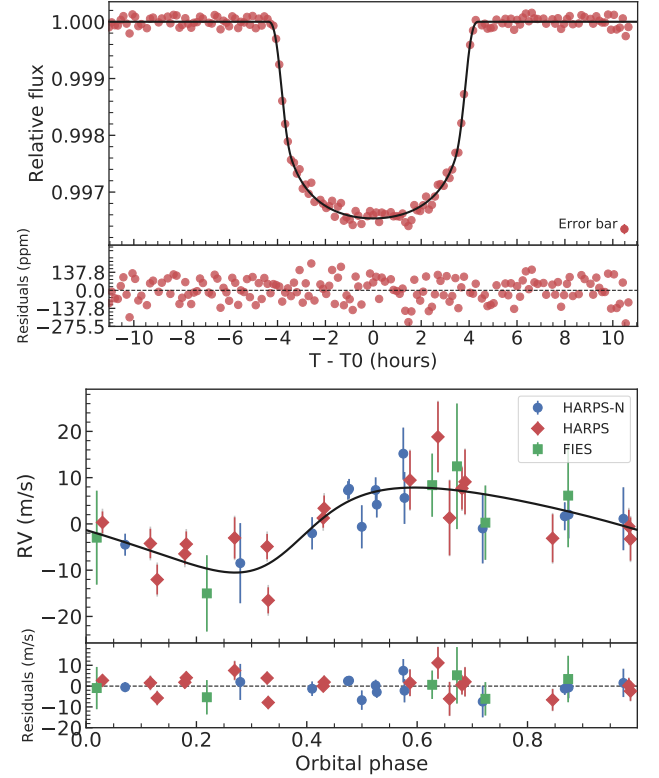
converged<sup>11</sup>, we ran 5000 additional iterations. We used a thin factor of 10 to generate a posterior distribution of 250,000 independent points for each parameter.

We first explored the properties of the Doppler signal by fitting the RV data alone. We tested different models: one model assumes there is no Doppler reflex motion; one model assumes the presence of a planet on a circular orbit; another model assumes the presence of a planet on an eccentric orbit. These three models were run with and without a jitter term for each instrument. This generates a set of 6 different models. The main statistical properties of each model are listed in Table 4. From this Table we can draw the following conclusions: 1) the models including a planet signal are strongly preferred over the models without it; 2) the eccentric model is preferred, as suggested also by the transit fit (see the following paragraph); 3) the model does not require to add a jitter term for each spectrograph, suggesting that any extra signal (stellar variability, other planets, etc.) are below the instrumental precision. This supports our RV analysis assuming only a Keplerian orbit. We note that we still fit for a jitter term for each instrument to allow more flexibility to our modelling and to mitigate the effects of the relatively sparse sampling of our data on the accuracy of the semi-amplitude estimate.

We used Kepler’s third law to check if the stellar density derived from the modeling of the transit light curves is consistent with an eccentric orbit (see, e.g., Van Eylen & Albrecht 2015). We first ran an MCMC analysis assuming the orbit is circular. The derived stellar density is  $0.32^{+0.02}_{-0.06} \text{ g cm}^{-3}$ . This density disagrees with the stellar density of  $0.8^{+0.16}_{-0.13} \text{ g cm}^{-3}$  obtained from the spectroscopic parameters derived in Section 3. We then performed a joint analysis allowing for an eccentric solution. We derived a stellar density of  $0.82^{+0.38}_{-0.35} \text{ g cm}^{-3}$ , which is consistent with the spectroscopically derived stellar density. This provides further evidence that the planetary orbit is eccentric. For the final analysis, we decided to set a Gaussian prior on  $a/R_{\star}$  using Kepler’s third law and the stellar mass and radius derived in Section 3 and listed in Table 1.

The median and 68.3% percent credible intervals of the marginalized posterior distributions are reported in Table 5. Figure 4 displays the RV and transit data together with the best fitting model. We show a corner plot of the fitted parameters in Figure A1.

The HARPS, HARPS-N, and FIES Doppler measurements show an RV variation in phase with the transit ephemeris (Figure 4, lower panel). However, as described by Cunha et al. (2013), contaminant stars that are within the sky-projected angular size of the spectrograph fibre ( $1''$  for HARPS and HARPS-N,  $1.3''$  for FIES) may affect the radial velocity measurements of the target star. If the radial velocity of the contaminant star is changing, i.e., its spectrum is shifting across the spectrum of target star, it can distort the spectral line profile of the target (and hence its CCF), mimicking the presence of an orbiting planet. As presented by Cunha et al. (2013) in their Table 8, for magnitude differences of  $\sim 5$ -6 mag, the impact of F2 V–K5 V contami-



**Figure 4.** *Top panel:* Transit light curve folded to the orbital period of K2-280 b and residuals. The thick black line is the rebinned best-fitting transit model. The red points are the K2 data. *Bottom panel:* The RV curve of K2-280 b phase-folded to the orbital period of the planet. The best fitting solution is marked with a solid black line. HARPS-N, HARPS, and FIES data are shown with blue circles, red diamonds, and green squares, respectively. The lower panel shows the residuals to the best-fitting model.

nant star on a G8 V target star can be as high as  $10 \text{ m s}^{-1}$ . If the nearby N–NW star, which has an angular separation  $0.38 \pm 0.011''$  from K2-280, is an F or G background eclipsing binary, it may not only generate a transit-like signal in the light curve of K2-280 every 19.9 days, but also a low-amplitude radial velocity signal at this period. We carefully checked the the FWHM and BIS of the HARPS, HARPS-N, and FIES cross-correlation functions to search for potential line profile variation induced by the blend companion. The generalized Lomb-Scargle periodograms (Zechmeister & Kürster 2009) of these indicators show no significant signal neither at the 19.9-day period and its harmonics, nor at any other period. We also found no correlation between the FWHM and BIS, and the RV measurements (Figure 5). In particular, the Spearman correlation coefficient between the HARPS RV measurements and the BIS of CCFs is equal to  $r_{RV-BIS,HARPS} = -0.45$  and between HARPS RVs and FWHM is equal to  $r_{RV-FWHM,HARPS} = -0.18$ . The Spearman correlation coefficient between the HARPS-N RVs and BIS is equal to  $r_{RV-BIS,HARPS-N} = -0.19$  and between the HARPS-N RVs and FWHM is equal to  $r_{RV-FWHM,HARPS-N} = -0.06$ . In the case of FIES measurements, the Spearman correlation coefficient between RVs

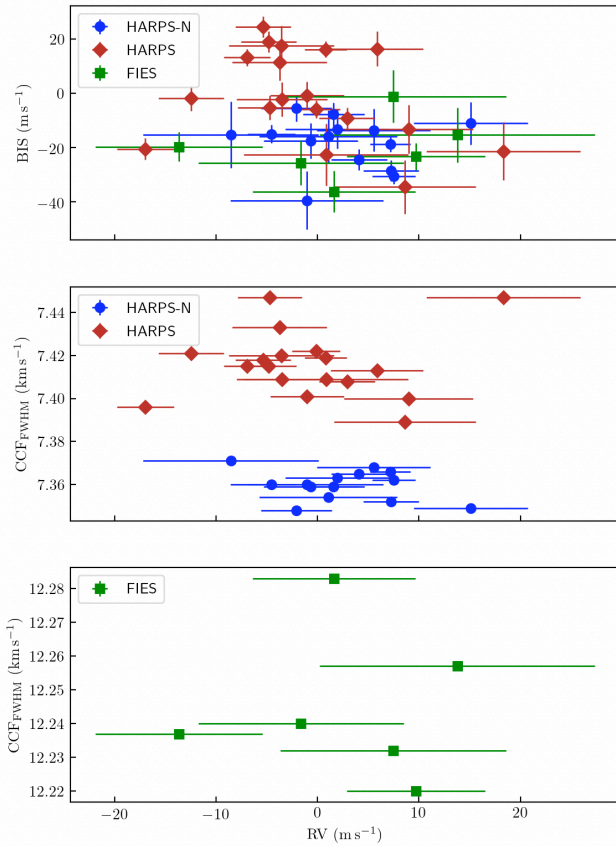
<sup>11</sup> We define convergence as when chains have a scaled potential factor  $< 1.02$  for all the parameters (see Gelman & Rubin 1992, for more details).



**Table 4.** Model comparison for our RV fits.

Test	Npars	log likelihood	BIC	K(m/s)
No planet - no jitter	3	118	-226	0
No planet - jitter	6	135	-248	0
planet - circular orbit- no jitter	6	136	-249	$7.20 \pm 1.15$
planet - circular orbit - jitter	9	143	-254	$7.18 \pm 1.60$
planet - eccentric orbit - no jitter	8	154	-280	$9.31 \pm 1.20$
planet - eccentric orbit - jitter	11	154	-269	$9.27 \pm 1.30$

*Note* – Further details about the Bayesian Information Criterion (BIC) are given in, e.g., [Burnham & Anderson \(2002\)](#).



**Figure 5.** *Top panel:* The BIS versus RVs from HARPS-N, HARPS, and FIES. *Middle panel:* The  $CCF_{FWHM}$  versus RVs from HARPS-N and HARPS. *Bottom panel:* The  $CCF_{FWHM}$  versus RVs from FIES. RVs from all instruments have been subtracted by the systemic velocities listed in Table 5 and derived by our joint analysis.

and BIS is equal to  $r_{RV-BIS,FIES} = 0.37$  and between RVs and FWHM is equal to  $r_{RV-FWHM,FIES} = -0.08$ .

We note that we could not measure the stellar rotation period from the *K2* light curve. Using the stellar radius determined in Section 3.4 and the projected rotation velocity determined in Section 3.2, we found the upper limit of the stellar rotation period to be  $P_{rot} = 21.6^{+12.6}_{-6.3}$  days. This means that the stellar rotation period of K2-280 is shorter than 34.2 days. Following the prescription given by [Aigrain et al. \(2012\)](#), the photometric variation found in the *K2* light

curve ( $\sim 600$  ppm) implies an activity induced RV signal of about  $2 \text{ m s}^{-1}$  ( $1.9 \text{ m s}^{-1}$  for stellar rotation period equal to orbital period of K2-280 b (19.9 days) or  $1.1 \text{ m s}^{-1}$  for  $P_{rot}$  equal to 34.2 days). The probability that stellar rotation modulation may generate RV variations of K2-280 is therefore very low. We conclude that most likely the Doppler shift of K2-280 is induced by the orbital motion of a planet transiting K2-280 rather than a blended eclipsing binary or stellar rotation modulation. We stress however that the activity-induced RV signal at a level of  $\sim 2 \text{ m s}^{-1}$  is larger than the jitter terms listed in Table 5 and larger than the precision on our estimate of the Doppler semi-amplitude variation induced by the planet ( $K = 9.18 \pm 1.27 \text{ m s}^{-1}$ ). Therefore, we warn the reader that our semi-amplitude estimate might be affected by unaccounted for stellar activity.

## 5 DISCUSSION AND SUMMARY

### 5.1 K2-280 b and the current sample of sub-Saturn planets

With a mass of  $M_b = 37.1 \pm 5.6 M_\oplus$  and a radius of  $R_b = 7.50 \pm 0.44 R_\oplus$ , K2-280 b joins the group of sub-Saturns planets – defined as planets having radii between 4 and  $8 R_\oplus$  ([Petigura et al. 2017](#)) – whose masses and radii have been measured. The basic physical parameters of a sample of 23 sub-Saturns with densities measured with a precision better than 50% have been presented and discussed by [Petigura et al. \(2017\)](#). We here extend this sample by adding K2-280 b alongside 6 additional sub-Saturns that have densities measured with a precision better than 50%, as described below. WASP-156 b ([Demangeon et al. 2018](#)), a  $\sim 0.5 R_{Jup}$  planet with a Jupiter-like density was discovered by the ground-based SuperWASP transit survey ([Pollacco et al. 2006](#); [Smith & WASP Consortium 2014](#)). Kepler-1656 b, a dense sub-Saturn with a high eccentricity of  $e = 0.84$  transiting a relatively bright ( $V = 11.6$  mag) solar-type star, was recently reported by [Brady et al. \(2018\)](#). Three sub-Saturns were discovered and characterised by the KESPRINT consortium, two of them in *K2* campaign 3 (K2-60 b, [Eigmüller et al. 2017](#)) and campaign 14 (HD 89345 b, aka K2-234 b, [Van Eylen et al. 2018](#); [Yu et al. 2018b](#)), and HD 219666 b ([Esposito et al. 2019](#)) in *TESS* Sector 1. One sub-Saturn, GJ 3470 b ([Bonfils et al. 2012](#)), orbiting an M1.5 dwarf was not included by [Petigura et al. \(2017\)](#), but we add it to the current sample, adopting the parameters from [Awiphan et al. \(2016\)](#). All of these new sub-Saturns, including K2-280 b, reside in apparently single systems. Figure 6 shows the mass–radius and mass–

**Table 5.** K2-280 Stellar and Planetary Parameters

Parameter	Prior <sup>(a)</sup>	Inferred value <sup>(b)</sup>
<i>Model Parameters</i>		
Orbital period $P_{\text{orb}}$ (days)	$\mathcal{U}[19.89, 19.90]$	$19.89526 \pm 0.00028$
Transit epoch $T_0$ (BJD <sub>TDB</sub> −2 450 000)	$\mathcal{U}[7307.55, 7307.65]$	$7307.58114 \pm 0.00056$
Scaled semi-major axis $a/R_{\star}$	$\mathcal{N}[25.58, 0.90]$	$25.79^{+0.87}_{-0.90}$
Scaled planet radius $R_p/R_{\star}$	$\mathcal{U}[0, 0.2]$	$0.05354^{+0.00094}_{-0.00056}$
Impact parameter, $b$	$\mathcal{U}[0, 1]$	$0.27^{+0.16}_{-0.17}$
$\sqrt{e} \sin \omega_{\star}$	$\mathcal{U}[-1, 1]$	$-0.547^{+0.047}_{-0.05}$
$\sqrt{e} \cos \omega_{\star}$	$\mathcal{U}[-1, 1]$	$-0.235^{+0.038}_{-0.043}$
Radial velocity semi-amplitude variation $K$ (m s <sup>−1</sup> )	$\mathcal{U}[0, 50]$	$9.18 \pm 1.27$
Parameterized limb-darkening coefficient $q_1$	$\mathcal{U}[0, 1]$	$0.54^{+0.14}_{-0.11}$
Parameterized limb-darkening coefficient $q_2$	$\mathcal{U}[0, 1]$	$0.249^{+0.082}_{-0.071}$
Systemic velocity $\gamma_{\text{HARPS-N}}$ (km s <sup>−1</sup> )	$\mathcal{U}[-2.2, -0.2]$	$-1.1934^{+0.0011}_{-0.0011}$
Systemic velocity $\gamma_{\text{HARPS}}$ (km s <sup>−1</sup> )	$\mathcal{U}[-2.2, -0.2]$	$-1.1907^{+0.0011}_{-0.0011}$
Systemic velocity $\gamma_{\text{FIES}}$ (km s <sup>−1</sup> )	$\mathcal{U}[-2.2, -0.2]$	$-1.2349^{+0.0038}_{-0.0039}$
Jitter term $\sigma_{\text{HARPS-N}}$ (m s <sup>−1</sup> )	$\mathcal{U}[0, 100]$	$0.36^{+0.97}_{-0.3}$
Jitter term $\sigma_{\text{HARPS}}$ (m s <sup>−1</sup> )	$\mathcal{U}[0, 100]$	$1.28^{+1.76}_{-1.13}$
Jitter term $\sigma_{\text{FIES}}$ (m s <sup>−1</sup> )	$\mathcal{U}[0, 100]$	$0.67^{+2.58}_{-0.58}$
<i>Derived Parameters planet b</i>		
Planet mass $M_p$ ( $M_{\oplus}$ )	...	$37.1 \pm 5.6$
Planet radius $R_p$ ( $R_{\oplus}$ )	...	$7.50 \pm 0.44$
Planet density $\rho_p$ (g cm <sup>−3</sup> )	...	$0.48^{+0.13}_{-0.10}$
Semi-major axis of the planetary orbit $a$ (au)	...	$0.1461^{+0.0099}_{-0.0097}$
Orbital eccentricity, $e$	...	$0.35^{+0.05}_{-0.04}$
Angle of Periastron, $\omega_{\star}$ (deg)	...	$246.77^{+4.34}_{-5.28}$
Time of periastron $T_p$ (BJD <sub>TDB</sub> −2 450 000)	...	$7315.06^{+0.238}_{-0.44}$
Transit duration $\tau_{14}$ (hours)	...	$8.267^{+0.063}_{-0.054}$
Equilibrium temperature <sup>(c)</sup> $T_{\text{eq}}$ (K)	...	$787 \pm 17$
Linear limb-darkening coefficient $u_1$	...	$0.367^{+0.074}_{-0.079}$
Quadratic limb-darkening coefficient $u_2$	...	$0.37^{+0.15}_{-0.15}$
Planet surface gravity <sup>(d)</sup> (cm s <sup>−2</sup> )	...	$648.0^{+107.0}_{-102.0}$
Planet surface gravity (cm s <sup>−2</sup> )	...	$647.0^{+133.0}_{-117.0}$

Note – <sup>(a)</sup>  $\mathcal{U}[a, b]$  refers to uniform priors between  $a$  and  $b$ , and  $\mathcal{N}[a, b]$  to Gaussian priors with median  $a$  and standard deviation  $b$ . <sup>(b)</sup> The inferred parameter value and its uncertainty are defined as the median and 68.3 percent credible interval of the posterior distribution. <sup>(c)</sup> Assuming albedo = 0. <sup>(d)</sup> Calculated from the scaled-parameters as suggested by Southworth et al. (2007).

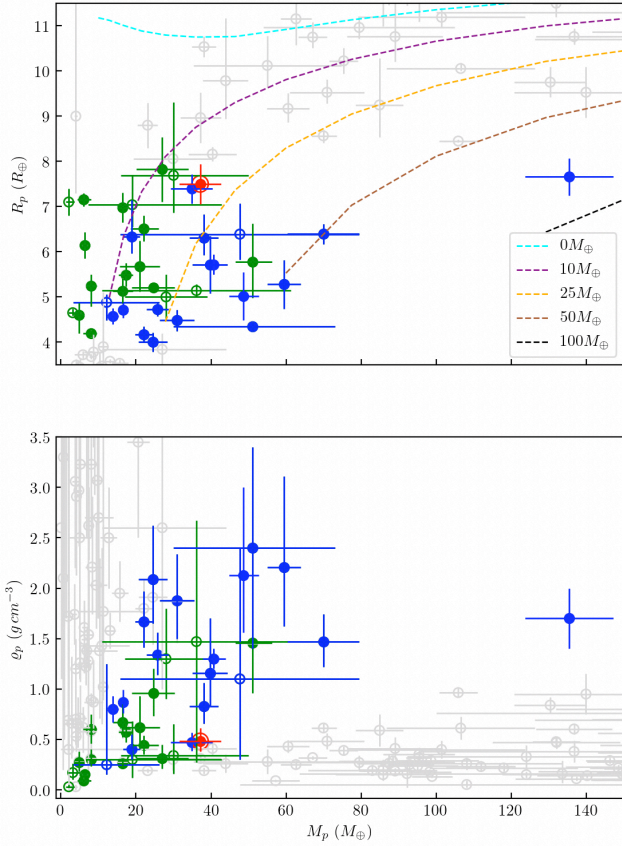
density diagrams for this extended sample of 30 planets. Sub-Saturns found to be in multi-planet systems are marked with green filled circles, whereas those in single systems are marked with blue filled circles. The position of K2-280 b is indicated with a red-rimmed circle. Sub-Saturns whose density has been measured with a precision slightly worse than 50% are marked with green and blue open circles. All the remaining transiting planets with measured radii and masses are marked with open gray circles<sup>12</sup>. According to the Fortney et al. (2007)’s models – also shown in the mass-radius diagram (Figure 6, upper panel) – K2-280 b has a core of about 10–25  $M_{\oplus}$ , accounting for ~25–65% of its total mass.

The diagrams in Figure 6 confirm the main characteristics found by Petigura et al. (2017) for the population of sub-Saturns. One of the main property is the uniform distribution of masses in the range ~5–75  $M_{\oplus}$ . With a mass of  $135 \pm 12 M_{\oplus}$  and radius of  $7.66 \pm 0.41 R_{\oplus}$ , K2-60 b (Eigmüller et al. 2017) is close to the lower envelope of giant planets on the mass–radius diagram and is the only sub-Saturn-sized

planet with a mass higher than Saturn ( $95.16 M_{\oplus}$ ). With a mean density of  $1.7 \pm 0.3 \text{ g cm}^{-3}$  (i.e. Neptune’s density), K2-60 b is also the most dense planet in the mass range ~75–250  $M_{\oplus}$ . As stressed by Eigmüller et al. (2017), K2-60 b with radius smaller than expected from the models of Laughlin et al. (2011) is more dense than expected and close to the sub-Jovian desert characterised by scarcity of planets with orbital periods below 4 days and masses lower than ~300  $R_{\oplus}$  (Szabó & Kiss 2011; Beaugé & Nesvorný 2013; Mazeh et al. 2016). The underestimation of its radius was excluded based on adaptive optics imaging (Schmitt et al. 2016). Only radial accelerations lower than  $2 \text{ m s}^{-1} \text{ day}^{-1}$  that can not be excluded based on RVs collected by Eigmüller et al. (2017) suggest that mass of K2-60 b may be lower than current determination. Nevertheless, this intriguing planet may help with a study of sub-Jovian desert and its borders.

Although the mass distribution of sub-Saturns is quite uniform, the most massive ones have radii close to and below ~6  $R_{\oplus}$ , visible as a correlation on the mass–density diagram (Figure 6, lower panel). The Spearman correlation coefficient between mass and density for the current sample of 30 sub-Saturns (excluding K2-60 b) is equal to  $r = 0.72$ . This

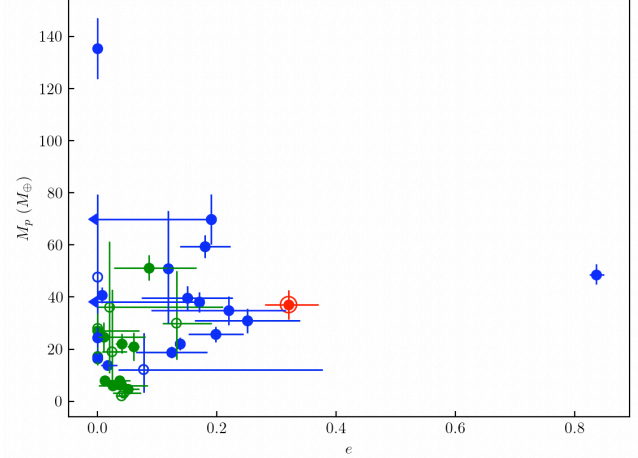
<sup>12</sup> As retrieved from the NASA Exoplanet Archive (Akeson et al. 2013) – July 2019.



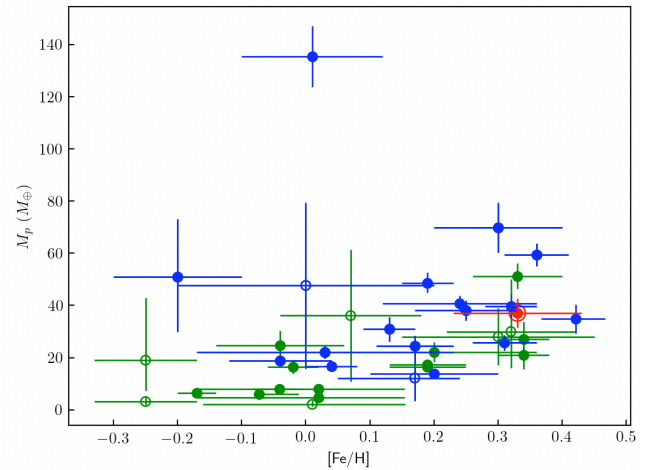
**Figure 6.** Mass–radius (upper panel) and mass–density (lower panel) diagrams for a sample of sub-Saturns ( $R_p = 4 - 8 R_{\oplus}$ ). Sub-Saturns whose mean densities have been measured with a precision better than 50% located in multi-planet systems are marked with green filled circles, whereas those in single systems are marked with blue filled circles. The position of K2-280 b is indicated as a red-rimmed circle. Sub-Saturns with densities measured with a precision slightly worse than 50% are marked with green and blue open circles. The remaining planets with measured radii, masses and mean densities (NASA Exoplanet Archive (Akesson et al. 2013), as of July 2019) are marked with open gray circles. The dashed lines on the mass–radius diagram (upper panel) correspond to the Fortney et al. (2007) models for planet core masses of 0, 10, 25, 50, and  $100 M_{\oplus}$  and age 10 Gyrs.

correlation is comparable to the one for the sample of 23 planets discussed in Petigura et al. (2017) that is equal to  $r = 0.79$ . Notably, almost all of the most massive sub-Saturns from the current sample of 30 planets reside in apparently single-planet systems (blue circles in Figure 6). Sub-Saturns in single-planet systems have also often moderate eccentricities, higher than their counterparts in multi-planet systems, as shown in Figure 7. As suggested by Petigura et al. (2017), the moderate eccentricities of more massive sub-Saturns in apparently single systems and the lack of high-eccentricity, high-mass objects in multi-planet systems may be explained by scattering and merging events during the formation process.

Petigura et al. (2017) found a marginal correlation between the stellar metallicity and the mass of sub-Saturn planets (the Spearman correlation coefficient  $r = 0.57$ ),



**Figure 7.** Mass of sub-Saturn planets as a function of the eccentricity. Samples and point symbols as in Figure 6.



**Figure 8.** Mass of sub-Saturn planets as a function of iron content of their host stars. Samples and point symbols as in Figure 6.

with the massive sub-Saturns found to orbit metal-rich stars. We confirm this for the current sample of 30 sub-Saturns (excluding K2-60 b) with exactly the same value of the Spearman correlation coefficient. This is consistent with the results of Buchhave et al. (2012) who, based on the sample of Kepler planets, found that planets larger than  $\sim 4 R_{\oplus}$  orbit stars with relatively high metal content ( $-0.2 < [\text{Fe}/\text{H}] < 0.5$  dex). For the sake of consistency with Figures 6 and 7, we included in Figure 8 the sub-Saturns orbiting binary stars, namely, Kepler-47 (AB) c and d ( $[\text{Fe}/\text{H}] = -0.25 \pm 0.08$  dex) and Kepler-413 (AB) b ( $[\text{Fe}/\text{H}] = -0.2 \pm 0.1$  dex), which were omitted by Petigura et al. (2017). Given its mass of  $M_p = 37.1 \pm 5.6 M_{\oplus}$  and the iron content of its host star ( $[\text{Fe}/\text{H}] = 0.33 \pm 0.08$  dex), K2-280 b follows this trend, being a relatively massive sub-Saturn orbiting a metal rich star.

K2-280 b has a relatively long orbital period of  $\sim 19.9$  days and transits a slightly evolved star in an ap-

parently single-planet system. With an eccentricity of  $e = 0.35^{+0.05}_{-0.04}$ , K2-280 b is exactly within the range of eccentricities found by Van Eylen et al. (2019) for *Kepler* systems with single transiting giant planets ( $R_p > 6 R_\oplus$ ). After Kepler-1656 b (Brady et al. 2018), K2-280 b is the second most eccentric sub-Saturn known to date. In contrast to Kepler-1656 b, the mass of  $M_b = 37.1 \pm 5.6 M_\oplus$ , radius of  $R_b = 7.50 \pm 0.44 R_\oplus$ , and mean density of  $\rho_b = 0.48^{+0.13}_{-0.10} \text{ g cm}^{-3}$ , make K2-280 b more similar to HD 89345 b (aka K2-234 b; Van Eylen et al. 2018; Yu et al. 2018b). The moderate eccentricity of K2-280 b suggests a formation pathway involving planet-planet gravitational interactions, and make this sub-Saturn planet a member of a relatively rare group of exoplanets and an interesting object for possible future follow-up.

## 5.2 Prospects for atmospheric characterisation and Rossiter-McLaughlin effect measurements

Although K2-280 b is a quite puffy planet, the relatively large radius of its host star ( $R_\star = 1.28 \pm 0.07 R_\odot$ ) results in a quite low transmission signal per scale height ( $H$ ) of the planetary atmosphere (55 ppm). This makes it a difficult target for atmospheric characterisation with current ground- and space-based facilities. The transmission spectroscopy metric (TSM) defined by Kempton et al. (2018) for JWST/NIRISS is  $\sim 45$  for K2-280 b, i.e. two times lower than the threshold TSM for planets with radii  $R_p \in (1.5 - 10.0) R_\oplus$  to be selected as high-quality atmospheric characterization targets. The long transit duration ( $\sim 8$  hours) further complicates ground-based follow-up observations.

Still, there is a possibility of Rossiter-McLaughlin (RM) effect measurements, for which the overall amplitude is expected to be  $\sim 6 \text{ m s}^{-1}$ , depending on the real values of stellar projected rotation velocity and planetary and stellar radii. With an impact parameter of  $b = 0.27^{+0.16}_{-0.17}$ , the transit of K2-280 b is close to being central. In such a case the shape of the RM effect would not change significantly with the sky-projected spin-orbit angle  $\lambda$ , but mainly the RM amplitude, leading to a strong correlation between  $\lambda$  and  $v_{\text{rot}} \sin i_\star$  (see, e.g., Albrecht et al. 2011). Therefore more precise determination of  $v_{\text{rot}} \sin i_\star$  of this slow rotator, based for instance on the Fourier transform technique (e.g. Smith & Gray 1976; Dravins et al. 1990; Gray 2008, and references therein) applied to single very high resolution and high SNR line profiles, would be needed. Measurements of the sky-projected spin-orbit angle through RM observations may help to test formation scenarios of warm sub-Saturn planets. This gives additional arguments for attempting RM observations, as the probability of a misalignment between the planet's orbital angular momentum vector and its host star's spin axis should be higher if caused by a perturber than by primordial misalignment of the protoplanetary disk. Two full transits of K2-280 b observable from the Chilean observatories will occur on July 7<sup>th</sup>/8<sup>th</sup> 2020 and August 9<sup>th</sup>/10<sup>th</sup> 2021.

## 6 CONCLUSIONS

We report here detailed characterisation of a low-density ( $\rho_b = 0.48^{+0.13}_{-0.10} \text{ g cm}^{-3}$ ) sub-Saturn transiting a mildly evolved, metal rich G7 star K2-280. With a mass of

$M_b = 37.1 \pm 5.6 M_\oplus$ , a radius of  $R_b = 7.50 \pm 0.44 R_\oplus$ , and an eccentricity of  $e = 0.35^{+0.05}_{-0.04}$ , K2-280 b joins the group of sub-Saturns planets in apparently single-planet systems. This second most eccentric sub-Saturn known to date is an interesting object for possible future follow-up observations that may help to test formation scenarios of this intriguing group of planets that are absent in the Solar System.

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*Facility:* *Kepler*, *Gaia*, Gemini-North/NIRI+ALTAIR, ESO/HARPS, TNG/HARPS-N, NOT/FIES.

*Software:* BLS, IRAF, PARAM 1.3, *pyaneti*, SME, *SpecMatch-emp*.

*Data availability:* The data underlying this article are available in the article and in its online supplementary material.

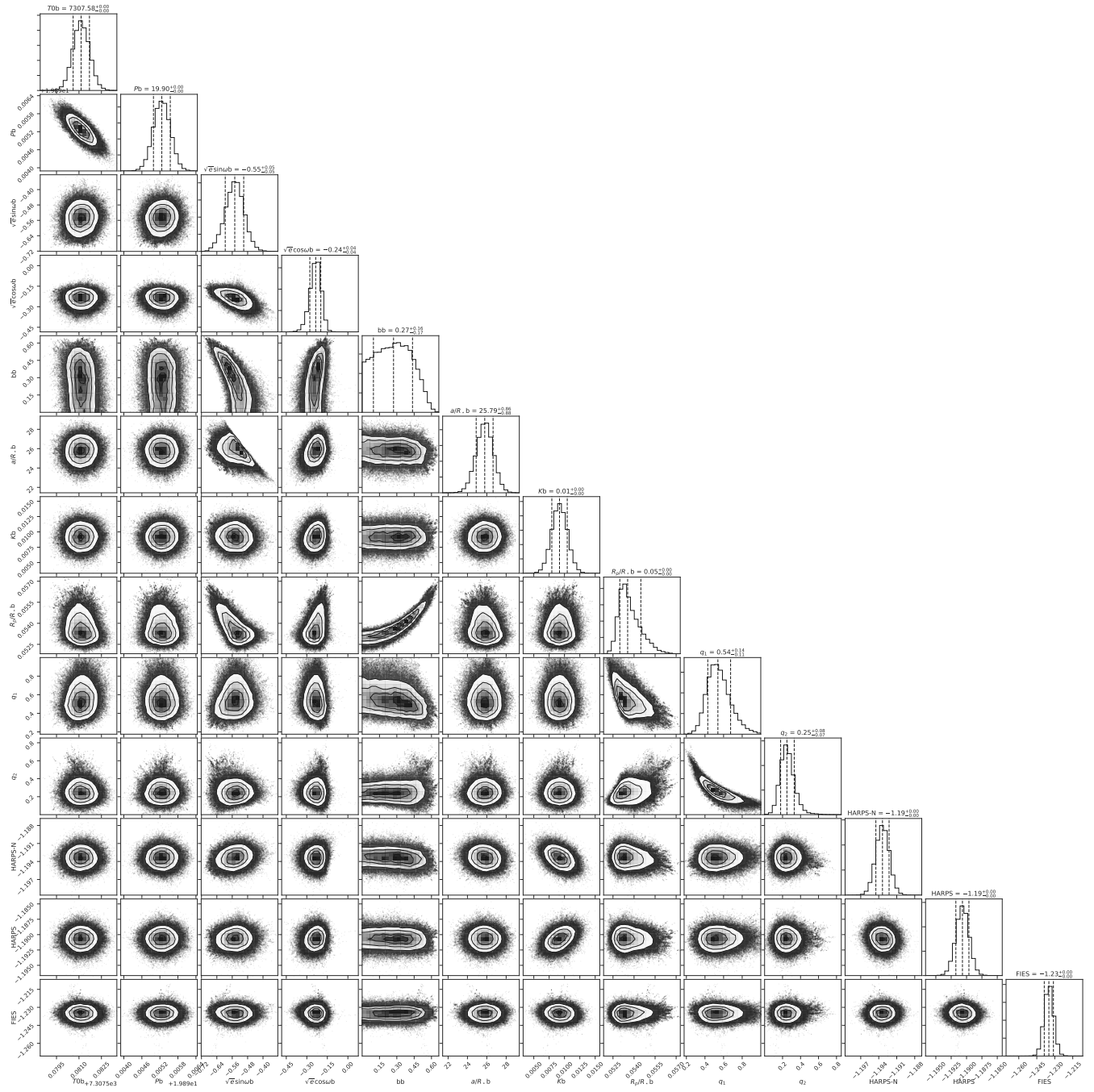
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## APPENDIX A: CORNER PLOT FOR FITTED PARAMETERS

This paper has been typeset from a  $\text{\TeX}/\text{\LaTeX}$  file prepared by the author.



**Figure A1.** Corner plot for the fitted parameters of the K2-280 system. This figure was created using `corner.py` (Foreman-Mackey 2016).